# Burn control in fusion reactors via isotopic fuel tailoring

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# The Need for Burn Control

- For tokamak fusion power to become an economical source of energy, reactors must operate for extended periods of time with a high power gain
- The gain is determined by the burn condition, i.e. the plasma density and temperature
- Some of the more economically and technologically attractive operating points may be unstable
  - Small decreases in temperature can lead to quenching
  - Small increases can cause the system to move to a stable, but undesired, high-temperature equilibrium
- Even when operating at stable equilibria, system performance during transients and disturbances could be undesirable without control
- Active burn control could enable operation at unstable operating points and could improve overall reactor performance.

## An Unstable Equilibrium



# Approaches Used in Previous Studies



Many previous works **utilize only one of the available actuators**. Most also use **linearized models** to generate control laws. The stable operating range of these controllers is therefore limited.

# Approach Used in Our Previous Work



Our previous work **combined actuation methods** and used a **nonlinear control design** to extend the region of attraction of the control scheme. However, the use of impurity injection could cause accumulation of impurities

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# Approach Used in This Work



In this work we continue the use of **nonlinear control design** techniques, while **avoiding impurity injection** by adjusting the deuterium and tritium fuel sources separately, a technique called **isotopic fueling**.

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# Isotopic Fueling



By adjusting the relative density of tritium and deuterium in the plasma, denoted as  $\gamma = n_T/(n_D + n_T)$ , the fusion power output can be modulated.

## Model Used for Controller Design

Alpha particles: 
$$\frac{dn_{\alpha}}{dt} = -\frac{n_{\alpha}}{\tau_{\alpha}} + S_{\alpha}$$
 (1)

Deuterium: 
$$\frac{dn_D}{dt} = -\frac{n_D}{\tau_D} - S_{\alpha} + S_D$$
 (2)

Tritium: 
$$\frac{dn_T}{dt} = -\frac{n_T}{\tau_T} - S_{\alpha} + S_T$$
 (3)

Energy: 
$$\frac{dE}{dt} = -\frac{E}{\tau_E} + Q_\alpha S_\alpha - P_{rad} + P_{aux}$$
 (4)

Reaction rate: 
$$S_{\alpha} = \gamma (1 - \gamma) n_{H}^{2} \langle \sigma \nu \rangle$$
 (5)  
Radiation losses:  $P_{rad} = A_{b} Z_{eff} n_{e}^{2} \sqrt{T}$  (6)

# **Desired Equilibrium**

An equilibrium can be obtained by choosing three parameters, for example:  $\bar{T}, \bar{n}, \bar{\gamma}$ , and solving:

Alpha particles: 0 = 
$$-\frac{\bar{n}_{\alpha}}{\tau_{\alpha}} + \bar{S}_{\alpha}$$
 (7)

Deuterium: 0 = 
$$-\frac{\bar{n}_D}{\tau_D} - \bar{S}_{\alpha} + \bar{S}_D$$
 (8)

Tritium: 
$$0 = -\frac{\bar{n}_T}{\tau_T} - \bar{S}_{\alpha} + \bar{S}_T$$
 (9)

Energy: 
$$0 = -\frac{\bar{E}}{\tau_E} + Q_\alpha \bar{S}_\alpha - \bar{P}_{rad} + \bar{P}_{aux}$$
 (10)

# Stabilization of Energy Subsystem

Defining  $\tilde{E} = E - \bar{E}$ , we can obtain

$$\frac{d\tilde{E}}{dt} = -\frac{\tilde{E}}{\tau_E} - \frac{\bar{E}}{\tau_E} + Q_\alpha S_\alpha - P_{rad} + P_{aux}$$
(11)

we can cancel the terms in red and stabilize the subsystem by satisfying the following condition:

$$Q_{\alpha}S_{\alpha} - P_{rad} + P_{aux} = \frac{E}{\tau_{E}}$$
(12)

This condition can be satisfied by modulating

- Paux directly
- *P<sub>rad</sub>* through impurity injection
- $S_{\alpha}$  through the tritium fraction

In this work we avoid using impurity injection and instead use modulation of the tritium fraction.

#### Desired tritium fraction

We first set  $P_{aux} = 0$  and seek a tritium fraction  $\gamma = \gamma^*$  that satisfies condition (12).

$$\mathbf{Q}_{\alpha}\gamma^{*}(1-\gamma^{*})\mathbf{n}_{H}^{2}\langle\sigma\nu\rangle-\mathbf{P}_{rad}=\frac{E}{\tau_{E}}$$
(13)

- Either two real solutions or two complex solutions
- If real, we take the solution satisfying  $\gamma^* \leq$  0.5
- If complex, auxiliary heating is needed to satisfy the condition, so we set  $\gamma^* = 0.5$  and move on

### Heating control law

We use the desired tritium ratio to calculate

$$P_{aux} = \frac{\bar{E}}{\tau_E} - Q_\alpha \gamma^* (1 - \gamma^*) n_H^2 \langle \sigma \nu \rangle + P_{rad}$$
(14)

 $\gamma$  is a virtual input, so it has dynamics. We define the error

$$\hat{\gamma} = \gamma - \gamma^* \tag{15}$$

which is governed by

$$\dot{\hat{\gamma}} = \frac{1}{n_H} \left[ -\frac{n_T}{\tau_T} - S_\alpha + S_T - n_H \dot{\gamma^*} - \gamma \left( -\frac{n_T}{\tau_T} - \frac{n_D}{\tau_D} - 2S_\alpha + S_D + S_T \right) \right]$$
(16)

The error between the actual and desired tritium ratio enters into the energy dynamics:

$$\frac{d\tilde{E}}{dt} = -\frac{\tilde{E}}{\tau_E} + \hat{\gamma}\phi \tag{17}$$

Finally, the dynamics of the plasma density deviations can be written as

$$\dot{\tilde{n}} = -3\frac{n_{\alpha}}{\tau_{\alpha}} - 2\frac{n_{\tau}}{\tau_{\tau}} - 2\frac{n_{D}}{\tau_{D}} - (Z_{I} + 1)\frac{n_{I}}{\tau_{I}} - S_{\alpha} + 2S_{D} + 2S_{\tau} + (Z_{I} + 1)S_{I}$$
(18)

We take the Lyapunov functional candidate

$$V = \frac{k_1^2 \tilde{E}^2 + k_2^2 \hat{\gamma}^2 + \tilde{n}^2}{2}$$

where  $k_1 \approx 10^{15}$  and  $k_2 \approx 10^{20}$  (recall that  $\tilde{E} \approx 10^5$ ,  $\hat{\gamma} \approx 10^{-1}$ , and  $\tilde{n} \approx 10^{20}$ ). We then compute the time derivative of *V* as

$$\dot{V} = k_{1}^{2}\tilde{E}\dot{E} + k_{2}^{2}\hat{\gamma}\dot{\gamma} + \tilde{n}_{H}\dot{\tilde{n}}_{H}$$

$$= \frac{k_{2}^{2}\hat{\gamma}}{n_{H}} \left[ \frac{k_{1}^{2}n_{H}\tilde{E}\phi}{k_{2}^{2}} - \frac{n_{T}}{\tau_{T}} - S_{\alpha} + S_{T} - n_{H}\dot{\gamma}^{*} - \gamma \left( -\frac{n_{T}}{\tau_{T}} - \frac{n_{D}}{\tau_{D}} - 2S_{\alpha} + S_{D} + S_{T} \right) \right]$$

$$+ \tilde{n} \left[ -3\frac{n_{\alpha}}{\tau_{\alpha}} - 2\frac{n_{T}}{\tau_{T}} - 2\frac{n_{D}}{\tau_{D}} - (Z_{I} + 1)\frac{n_{I}}{\tau_{I}} - S_{\alpha} + 2S_{D} + 2S_{T} + (Z_{I} + 1)S_{I} \right] - \frac{k_{1}^{2}\tilde{E}^{2}}{\tau_{E}}$$
(19)

#### Tritium fueling control law

$$S_{T} = -\frac{k_{1}^{2}n_{H}\tilde{E}\phi}{k_{2}^{2}} + \frac{n_{T}}{\tau_{T}} + S_{\alpha} + n_{H}\dot{\gamma}^{*} - K_{\gamma}\hat{\gamma} + \gamma \left(-\frac{3}{2}S_{\alpha} + \frac{3}{2}\frac{n_{\alpha}}{\tau_{\alpha}} + \frac{(Z_{I}+1)}{2}\frac{n_{I}}{\tau_{I}} - \frac{(Z_{I}+1)}{2}S_{I} - \frac{1}{2}K_{n}\tilde{n}\right)$$
(20)

### Deuterium fueling control law

$$S_{D} = \frac{1}{2} \left[ 3 \frac{n_{\alpha}}{\tau_{\alpha}} + 2 \frac{n_{D}}{\tau_{D}} + 2 \frac{k_{1}^{2} n_{H} \tilde{E} \phi}{k_{2}^{2}} - S_{\alpha} - K_{n} \tilde{n} + 2K_{\gamma} \hat{\gamma} - 2n_{H} \dot{\gamma}^{*} - \gamma \left( -3S_{\alpha} + 3 \frac{n_{\alpha}}{\tau_{\alpha}} + (Z_{l} + 1) \left( \frac{n_{l}}{\tau_{l}} - S_{l} \right) - K_{n} \tilde{n} \right) \right]$$
(21)

where  $K_{\gamma} > 0$  and  $K_n > 0$ 

The choice of control laws (20) and (21) reduces (19) to

$$\dot{V} = -rac{k_1^2 ilde{E}^2}{ au_E} - K_n ilde{n}^2 - rac{k_2^2}{n_H} K_\gamma \hat{\gamma}^2$$

- This shows that the control laws guarantee asymptotic stability of the energy, plasma density and requested tritium ratio.
- Since no linearization of the model was used, the control laws globally stabilize the model used for design
- The scheme works for both ignited and sub-ignited operating points
- The controller depends parametrically on the equilibrium, so it can be used to go from one operating point to another

















# **Conclusions and Future Work**

- A multi-input nonlinear fusion burn controller has been proposed
- Impurity injection is avoided by use of isotopic fuel tailoring
- The controller depends parametrically on the operating point and works for ignited and sub-ignited equilibria
- Simulation results show good performance
  - Stabilization of unstable equilibria
  - Improved transient performance when moving between equilibria
  - Successful transition between ignited and sub-ignited points
- Future work will focus on dealing with model uncertainty and spatially distributed effects