

First results from Electron Energy Distribution Function (EEDF) Analysis of Langmuir Probe Signals

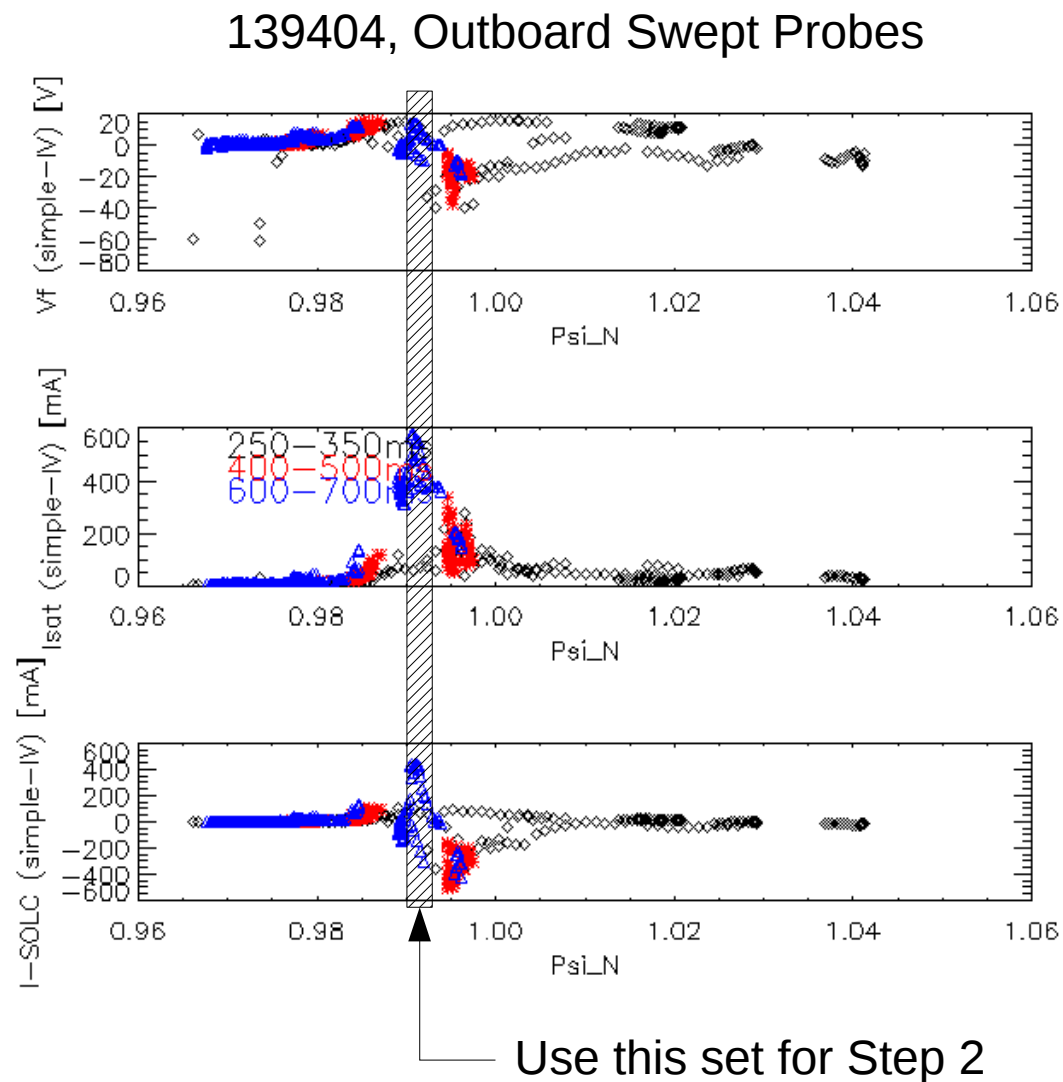
M.A. Jaworski, et al.

Motivation

- LLD discharge observations of reduced V_f (coming up)
- Recent research trends suggesting single Maxwellian fluid approach posing difficulties (e.g. Chankin, 2007, “Discrepancy between modelled and measured radial electric fields in the scrape-off layer of divertor tokamaks: a challenge for 2D fluid codes?” **NF 47**)
- Previous work indicating kinetic effects in the SOL should be expected (Chodura 1992 and Batischev 1997)

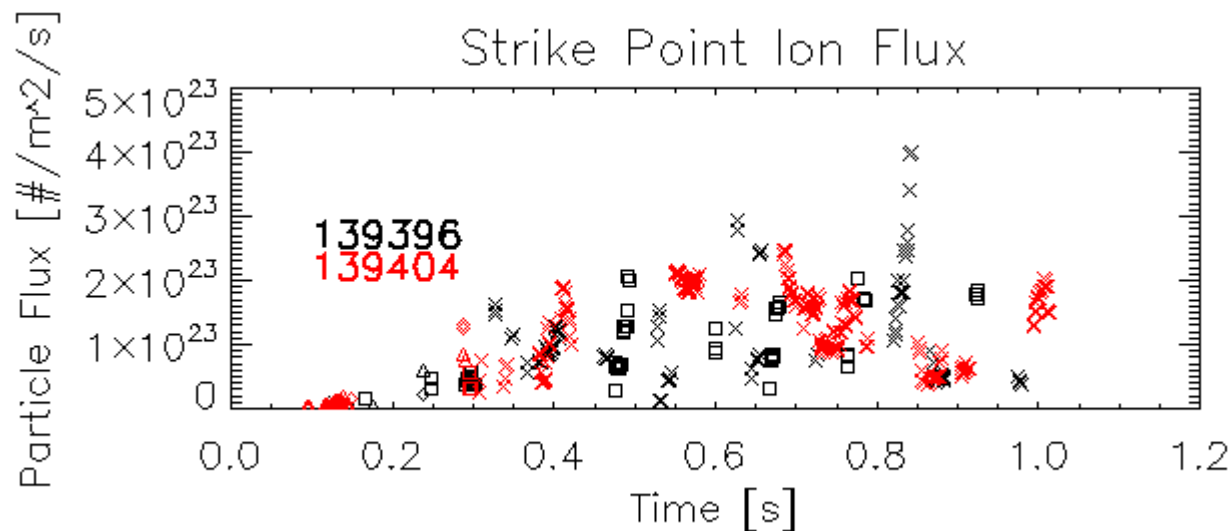
Discharge comparison process

- General problem of high-gradient region moving about over spatially localized measurement points
- Step 1 makes use of EFIT to generate a psi value for each probe in time and then construct strike-point sweeps (but indicates time evolution)
- Step 2 Filter resulting data set along specific psi value to generate a time-base plot of the value of a quantity on a specific magnetic surface



Can Now Compare Apples-to-Apples

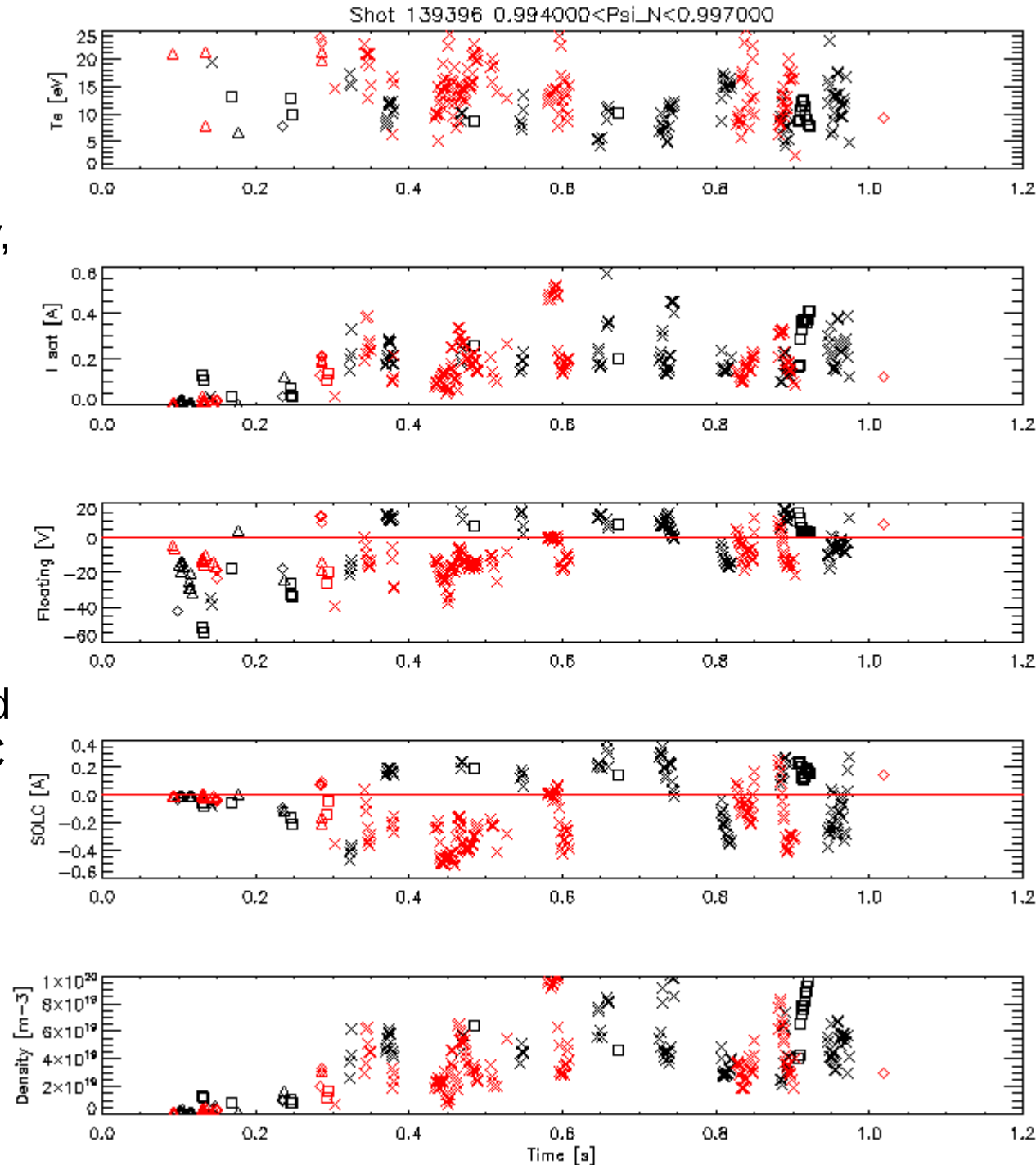
- Use same psi-surfaces where peak ion current occurs to provide strike-point flux values
- Colors denote shot
- Symbols indicate different probes falling within desired psi range



$$\Gamma = \frac{I_{scat}}{e A_{probe}}$$

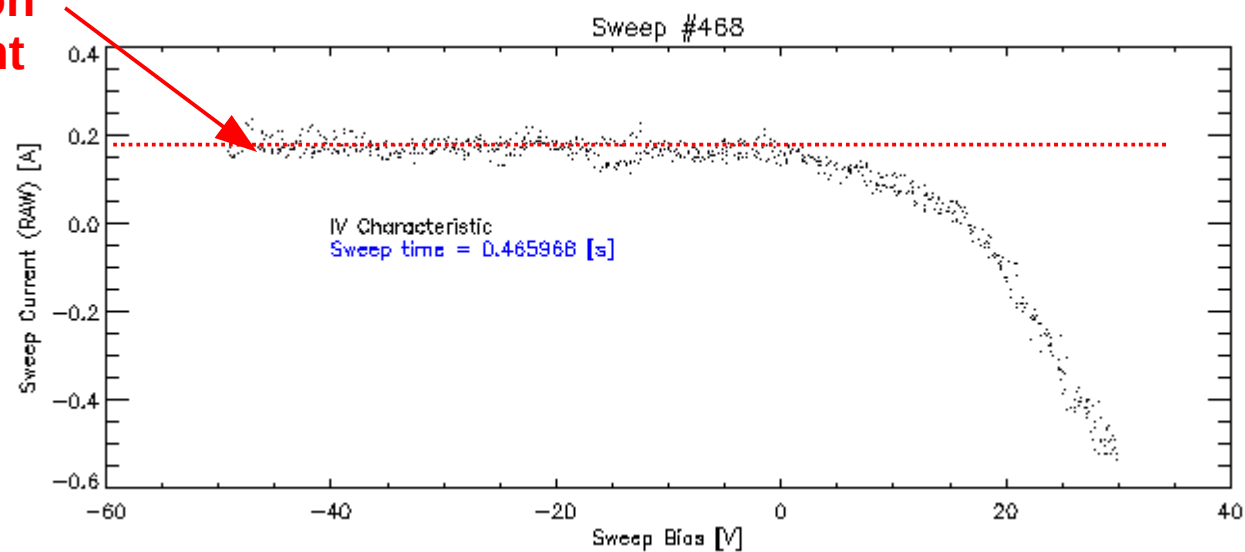
Discharge comparison during XP1041A

- 139396 compared with 139404
 - More gas, but not much change in electron density, nor in deuteron content (**where did it go?**)
 - Type V ELMs return by end of sequence – not present in 139404
- Tantalizingly higher electron temperatures in 139404
- Reduced floating potential and corresponding negative SOLC
- **But can I trust T_e and classical interpretation method?**
- **Motivation for EEDF**

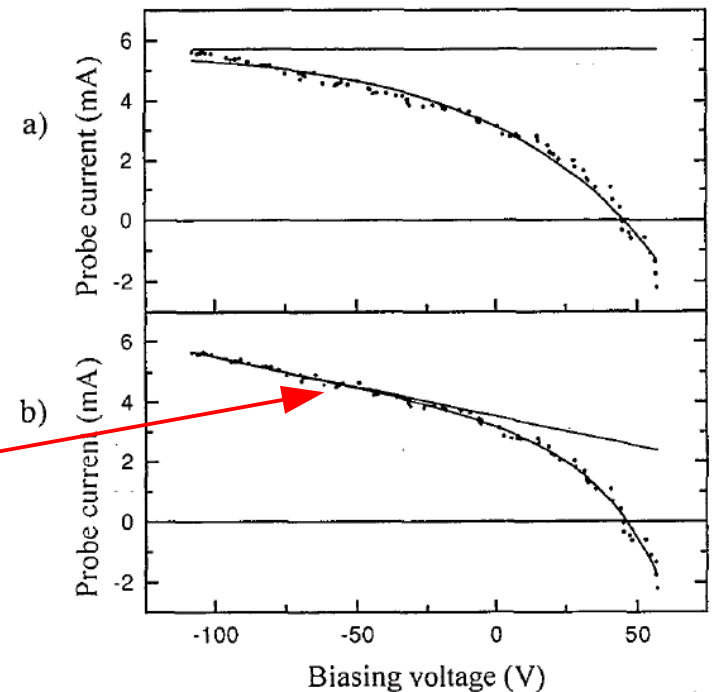


Probe Regime Enables EEDF Analysis

Clean Saturation
of probe current



- Outboard probes typically in “thin-sheath” regime (due to probe geometry)
- Simplifies analysis by removing sheath growth effects (Gunn probes, NSTX inboard probes)



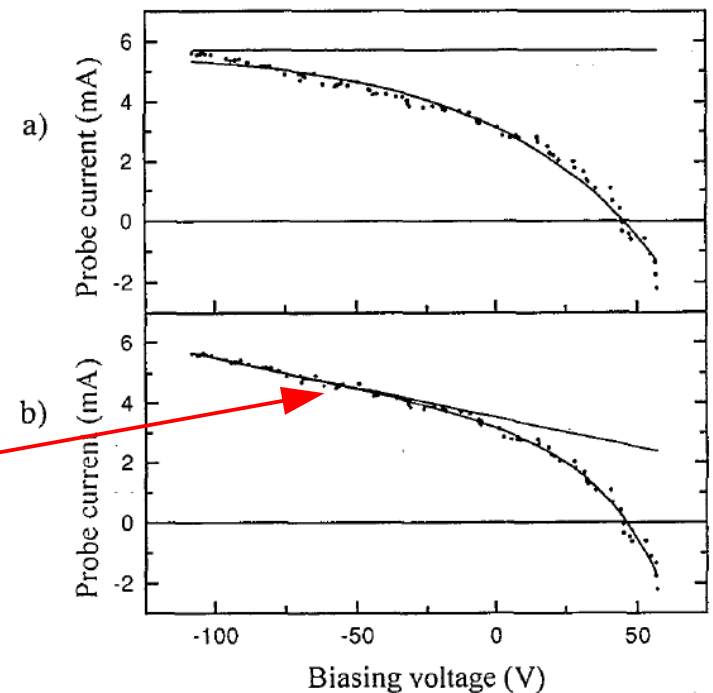
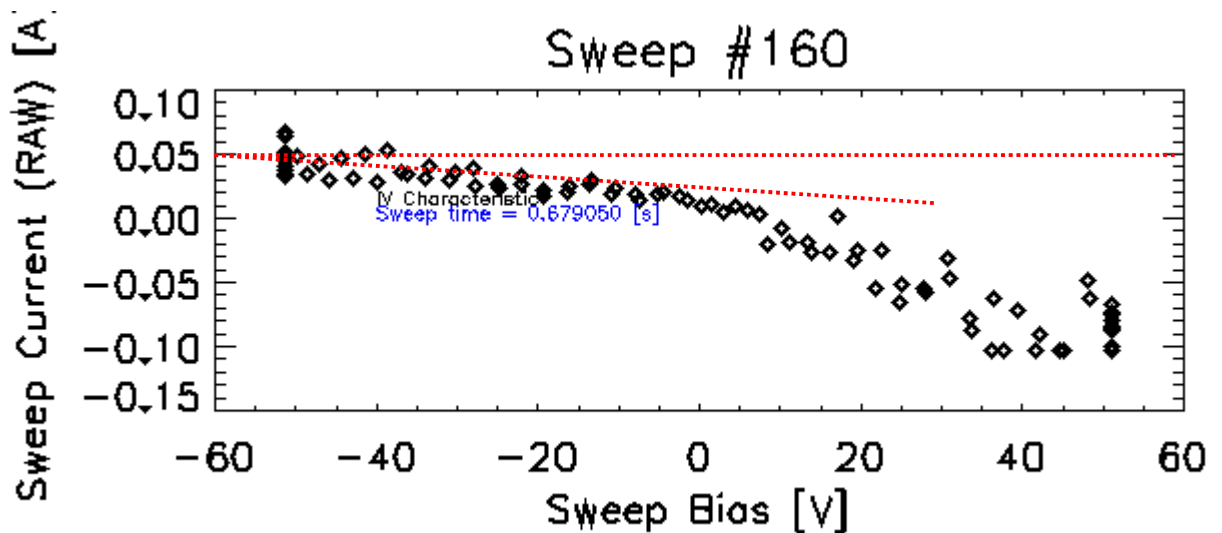
Gunn, RSI, 1995

Probe Regime Enables EEDF Analysis

Inboard probes do not always saturate due to sheath growth (probe 3, 139396) (else Te~28eV!)

Shift from I_{sat} is function of angle of incidence, density and temperature [Gunn, 1997, RSI]

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Gunn, RSI, 1995

Recent Work Demonstrates EEDF Extraction in Magnetized Plasmas

- Requires energy scale length to be long compared to probe perturbation scale length

Arslanbekov, **PSST**, 1995
-High pressure and magnetic fields

- Magnetic field makes distribution function $f(e)$ proportional to first derivative in some cases, avoiding noise-prone second derivatives

Demidov, **PoP**, 1999
-Magnetic fields

Popov, **PPCF**, 2009
-CASTOR Tokamak, cylindrical probes

$$\tau_{\varepsilon}^{-1} \approx \nu_{ee} + \delta \nu_{e-n} + \nu_{e-i}$$

$$\lambda_{\varepsilon} = \sqrt{4 D_e \tau_{\varepsilon}}$$

~6cm vs. 1-2 mm

This will be the first usage of this EEDF method in a high-power tokamak and with divertor probes

Non-Local Probe Theory

- When energy scale-length is much larger than probe scale, then leading terms in kinetic equation is diffusion in spatial coordinates [Arslanbekov, 1994, Golubovskii, 1981]
- Solution to current drawn by probe given as:

$$j(V) = \frac{8\pi e}{3m^2 \gamma_0} \int_{eV}^{\infty} \frac{(\varepsilon - eV) f_{\infty}(\varepsilon) d\varepsilon}{1 + \frac{\varepsilon - eV}{\varepsilon} \psi(\varepsilon, V)}$$

- Where psi is a diffusion parameter depending on probe geometry and system type (i.e. e-n collisions, classical (Demidov) or Bohm (Popov) cross-field diffusion)

$$\psi(\varepsilon, V) = \frac{1}{\gamma_0 \lambda(\varepsilon)} \int_a^{\infty} \frac{D(\varepsilon) dr}{(r/a)^n D(\varepsilon - eV(r))}$$

Limit cases for Non-Local Theory

- In the limit of very small psi parameter (i.e. collisionless or non-magnetized) resulting form is identical to Langmuir formula:

$$j(V) \propto \int_{eV}^{\infty} (\varepsilon - eV) f(\varepsilon) d\varepsilon$$

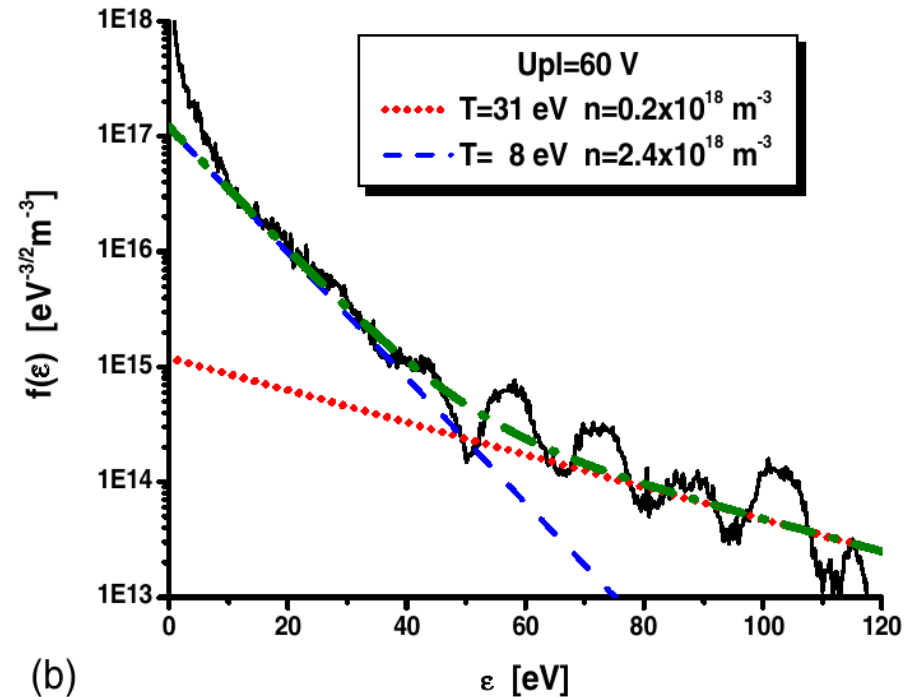
- Which results in the usage of second derivative for EEDF
- In the limit of very large psi parameter (i.e. high neutral pressure, high magnetic field), results in first-derivative method:

$$j'(V) \propto -\frac{eV \sqrt{eV}}{\psi_0} f(eV)$$

- Popov, 2009, categorized according to psi_0 value, which is value of psi parameter for a 1eV electron
- For intermediate values, mixture of methods.

CASTOR Usage of EEDF Method

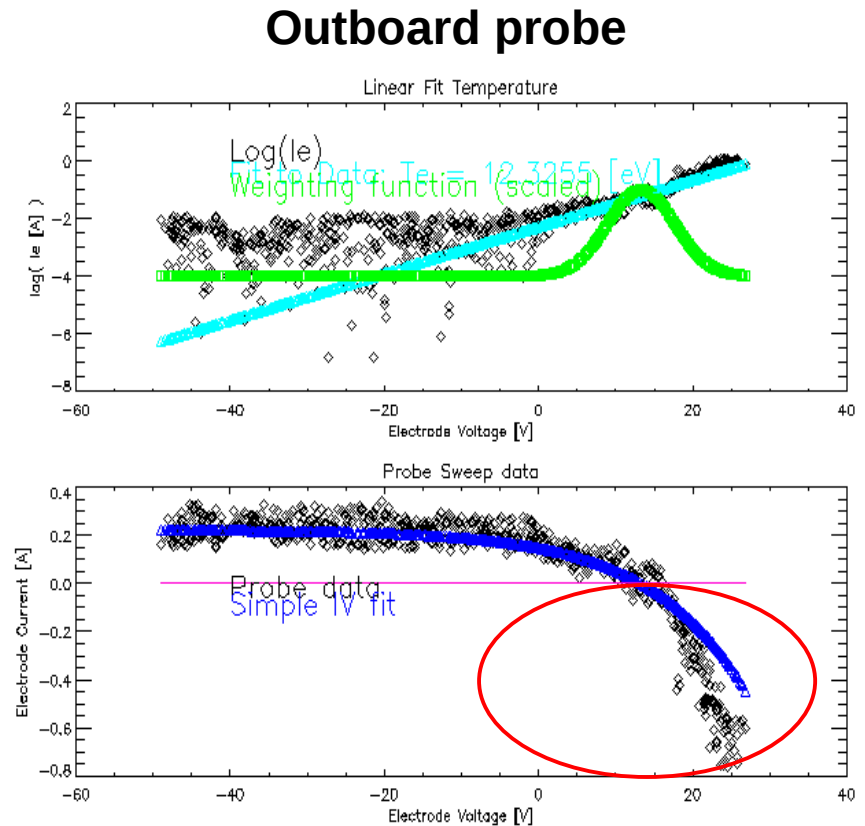
- Found that for values of $\psi_0 > 100$, little error (2-5%) in using first derivative method alone
- For values less than 100, recommended using iterative method due to integral equation for first derivative
- Found bi-modal distribution in edge of their plasmas
- **Initial estimate of NSTX outboard probe diffusion parameter was $\Psi_0 \sim 75$**



CASTOR tokamak EEDF
obtained via first-derivative
method (Popov, 2009, **PPCF**)

NSTX Data Example

- Simple IV method (“Classical”) measures sub-section of data
 - Initial analyses used Gaussian weighting around floating potential (Jaworski, **RSI**, 2010)
 - Typical usage is for $V < V_f$ (Stangeby, **PPCF**, 1995; Matthews, **PPCF**, 1994)
- Discards most of the electron current, samples only the tail (or less)



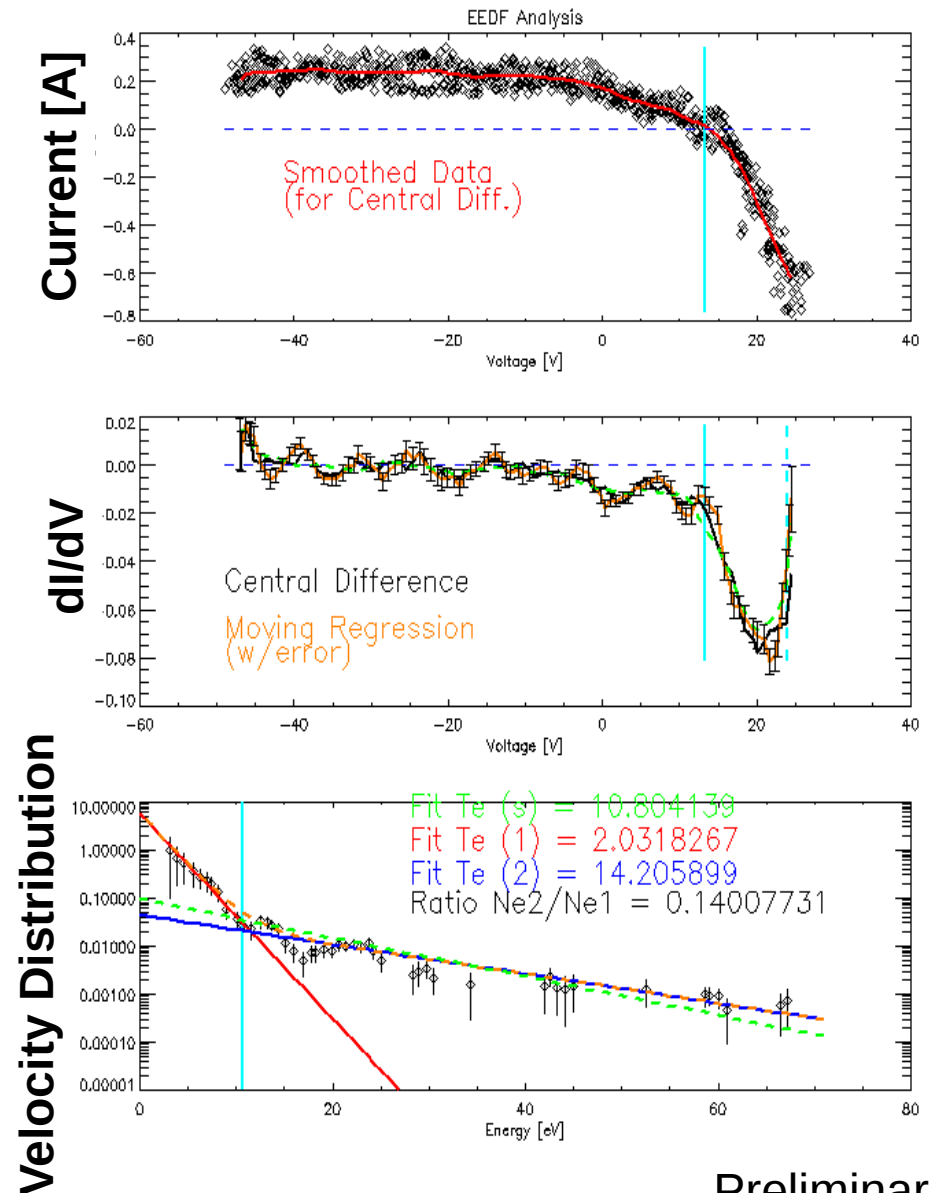
Dark blue is the resulting fit – Classical method assumes single Maxwellian and only measures the tail of the distribution (below V_f).

Red ellipse indicates unused data in the classical method.

Bi-modal EEDF in Divertor

- First derivative method requires some processing (two methods shown for calculating, non-iterative method used initially, hence “preliminary”)
- Reveals bi-modal distribution (also seen in CASTOR) and relative fractions of hot and cool populations
 - Will attempt to get absolute densities from EEDF
 - Chi²-minimization algorithm developed to determine best fit (less subjective input from user)

Same data set shown previously

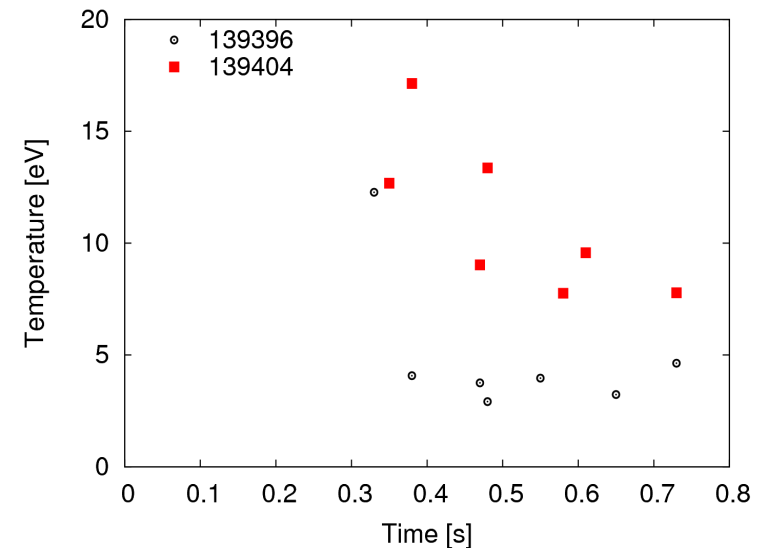
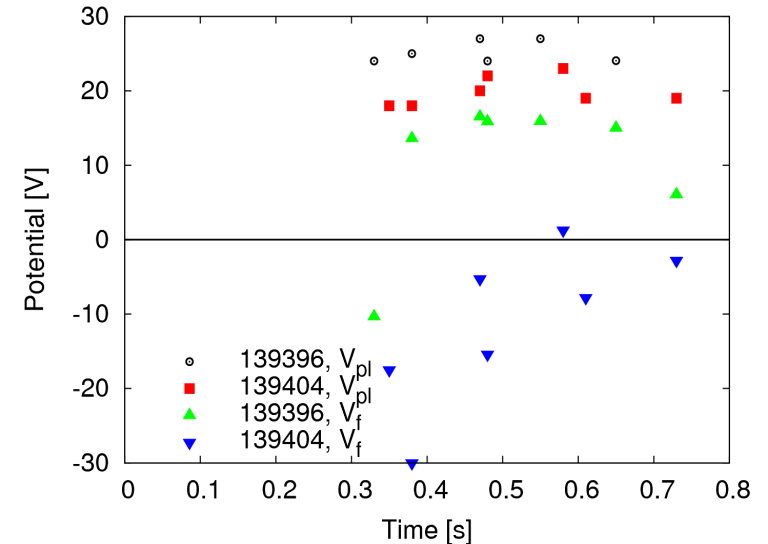


Preliminary

Application of Analysis to LLD

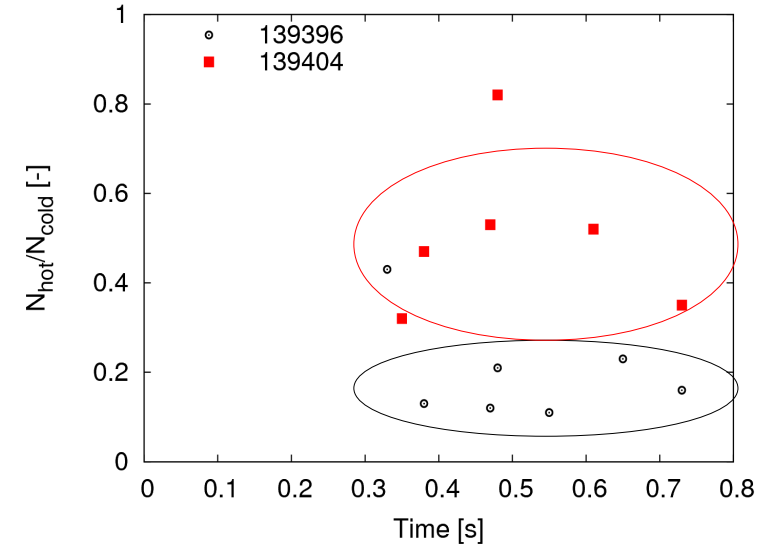
- Return to XP1041A discharges
- Can the floating potential reduction be attributed to plasma biasing of plasma potential downward?
 - Floating potential difference is about 20V
 - Plasma potential difference is only about 6V
 - Cannot remove with only V_p reduction (although there is a reduction in this quantity)
- Increase in $V_p - V_f$ is often attributed to an increase in $T_e \sim (V_p - V_f)/2.8$ (for pure D plasma)

NB: Sparse data set because EEDF still a “by-hand” process. Time points correspond to points in shot shown earlier.

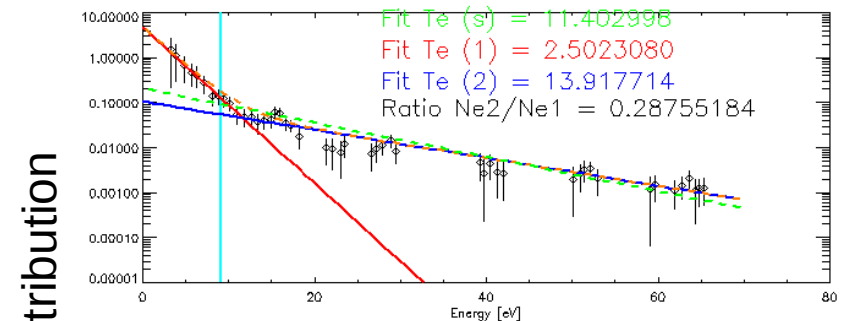


Hot Electron Fraction Larger in 139404

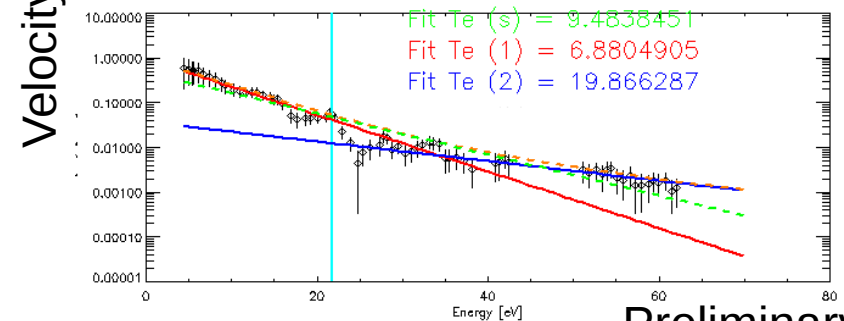
- Bi-modal temperatures compared
 - Still determining reliability of low-temperature solution, but similar Te's within scatter but...
 - Fraction of higher temperature population larger in 139404 discharge
- Hot fraction typically less than 20% in 139396 vs. >40% in 139404
- Have since added functionality to determine when single Maxwellian better than bi-modal via Goodness-of-Fit statistic



139396, 0.65s Bimodal is better



139404, 0.58s Single is better



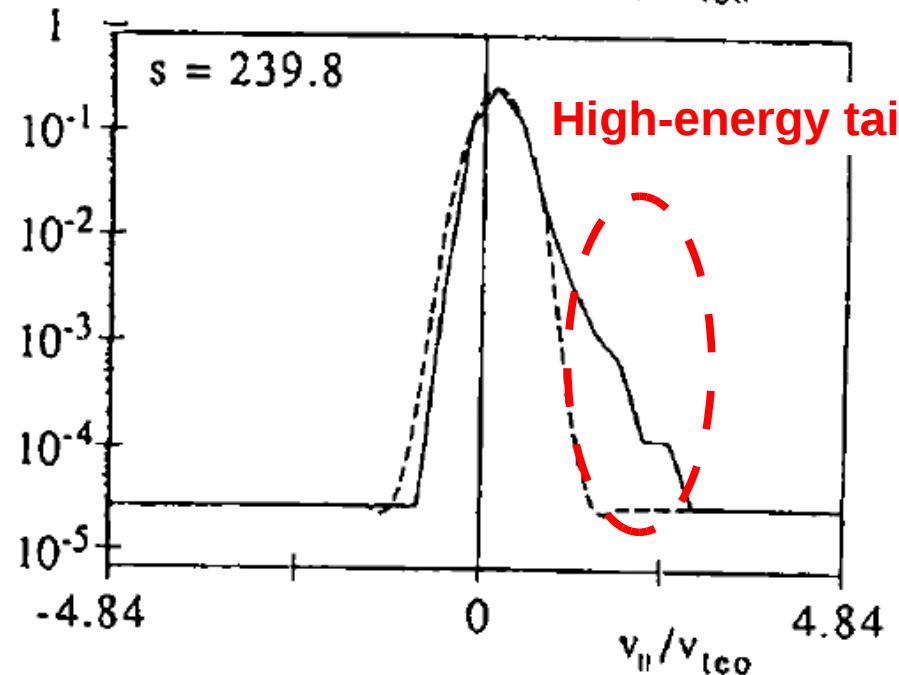
Preliminary

Larger hot fraction consistent with lower R

- Chodura kinetic effects in the SOL paper from 1992 simulated low/high recycling with Fokker-Planck code (Chodura, **Contrib. Plasma Phys.** 1992)
 - Demonstrated high-energy tail effects (Also seen in Batischev paper of 1997)
 - High energy tail predicted to be in recycling plasmas
 - Low recycling looks more like “mid-plane” distribution (i.e. single, hot Maxwellian)
- **The increase in hot fraction during the LLD experiment toward a value of 1 is consistent with reduced recycling (i.e. evolving toward single Maxwellian with higher temperature)**

Chodura states in body of paper that: “[in the case of no recycling] the velocity distribution stays nearly Maxwellian if the source is Maxwellian.” I.e. electrons free-stream from midplane to target. On the other hand, recycling plasma is cool, but tail can penetrate to PFC.

Recycling boundary, high collisionality



Improvements in the current method

- Add in iterative solution methods for the full integral equation to handle arbitrary ψ values
- More rigorous calculation of a divertor probe diffusion parameter
- Careful examination needed to determine if Plasma-Electron Spectroscopy (PLES) is possible for impurity identification/tracking (Demidov, 2010, **CPP**)
- Analyze more shots and start determining relevant quantities for comparison with other diagnostics (e.g. energy flux from EEDF for comparison with IR)

Summary

- Outboard probe geometry seems compatible with EEDF analysis
 - First usage of this method in a high-power tokamak and with divertor mounted probes – encouraging results so far
 - Still some theory development required for more confidence
- EEDF method reveals far more information about the plasma than classical interpretation
 - Plasma potential, actual distribution function
 - Initial survey indicates it may be possible to perform PLES with the Langmuir probes (impurity diagnostic)
- Comparison of LLD discharges with EEDF method reveals:
 - Depression in floating potential (direct measure/observation) in discharge with suspected “actively pumping” LLD
 - Larger fraction of hot electrons in the suspected “actively pumping” discharge – would explain floating potential depression, increase in temperature
 - Increase in hot fraction consistent with kinetic modeling of the SOL comparing non-recycling/recycling boundary condition

BACKUP

Other Kinetic Modeling

- Batishchev also predicted high-energy tail at divertor plate
- Utilized ALLA Fokker-Planck code to simulate attached and detached divertors
- High-gradient regions do not have enough interactions to slow high-energy electrons
- Consistent with Chodura modeling

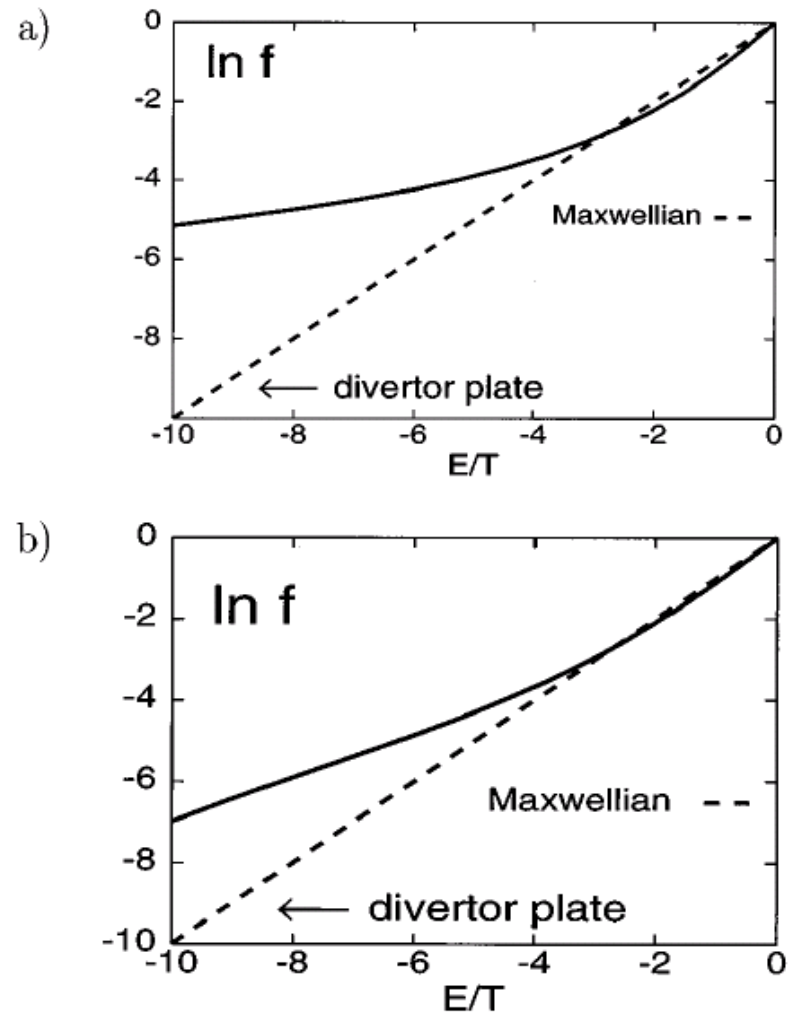


FIG. 12. Normalized distribution function at the plate for detached (a) and attached (b) cases compared to a Maxwellian.

Shots:
139396
139404

