# Nonlinear Simulations of NBI-driven GAE modes in NSTX

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#### **Motivation:**

- Multiple sub-cyclotron frequency modes are observed in NSTX during NBI injection.
- CAE and GAE modes are predicted to be driven unstable by super Alfvenic NBI ions with V<sub>b</sub>~3V<sub>A</sub> (80 keV) through the Doppler shifted cyclotron resonance.
- Strong anisotropy in the fast-ion pitch-angle distribution provides the energy source for these instabilities.
- Both CAE and GAE modes are observed.
- <u>Numerical simulations are needed to include</u>: self-consistent anisotropic equilibrium, FLR effects, thermal ion and nonlinear effects.

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# HYM – Parallel Hybrid/MHD Code

HYM code developed at PPPL and used to investigate kinetic effects on MHD modes in toroidal geometry (FRCs and NSTX)

- 3-D nonlinear, parallel (3D domain decomposition, MPI).
- Physics model: MHD/particle (one fluid thermal plasma + energetic particle ions)
- Full-orbit kinetic ions using delta-f / full-f numerical scheme.

# Model

Grad-Shafranov equation for two-component plasma: MHD plasma (bulk) and fast ions

$$\frac{\partial^2 \psi}{\partial z^2} + R \frac{\partial}{\partial R} \left( \frac{1}{R} \frac{\partial \psi}{\partial R} \right) = -R^2 p' - HH' - GH' + RJ_{i\phi} \qquad \mathbf{B} = \nabla \phi \times \nabla \psi + h \nabla \phi \\ h(R, z) = H(\psi) + G(R, z) \\ \mathbf{J}_{ip} = \nabla G \times \nabla \phi \qquad , G - \text{poloidal stream function}$$

Self-consistent MHD + fast ions coupling scheme

Background plasma - fluid: 
$$\rho \frac{d\mathbf{V}}{dt} = -\nabla p + (\mathbf{j} - \mathbf{j}_i) \times \mathbf{B} - n_i (\mathbf{E} - \eta \mathbf{j})$$

**Fast ions** – delta-f scheme:  $F_0 = F_0(\epsilon, \mu, p_{\phi})$ , where  $\mu$  is calculated up to 1<sup>st</sup> order in  $\rho_i / L$ ; realistic  $F_0$  to match the distribution functions computed from the TRANSP code.

#### Plasma parameters and profiles are matched to NSTX shot #114147 (TRANSP)



## 3D simulations of energetic ion-driven instabilities in NSTX

- Growth rates of unstable modes are very sensitive to details of distribution function (pitch-angle); equilibrium profiles are not.
- Most unstable mode toroidal number shifts to larger n for larger q<sub>0</sub>: n=5, m=-1



$$\begin{array}{c} \gamma_4 = 0.005\omega_{ci} \quad \text{and} \ \omega = 0.3\omega_{ci} \\ \gamma_5 = 0.014\omega_{ci} \quad \omega = 0.3\omega_{ci} \end{array} \qquad k_{\parallel} = \frac{\omega_{ci} - \mid \omega \mid}{v_{\parallel}}$$

Observed features agree with that of GAE mode, which exists just below the lower edge of the Alfven continuum:

- For each *n*, several *m* are unstable with large  $k_{||}$  and nm < 0.
- Localized near magnetic axis.
- Large  $\delta B_{\perp}$  component.

Dependence of the linear growth rate on viscosity.

Main damping mechanism for GAE is continuum damping (modeled in HYM with artificial viscosity):  $\gamma_d/\omega \sim (r/r_{res})^{2m+\delta}$ 

### Mode structure (equilibrium #114147)





### Magnetic mode structure



At peak amplitude  $\delta B_{\parallel} \approx 1/3 \ \delta B_{\perp}$ , but at the edge there is significant compressional component  $\delta B_{\parallel} > \delta B_{\perp}$ .