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Results from MHD XP's 701, 702, 703

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w/ contributions from
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NSTX Results Review

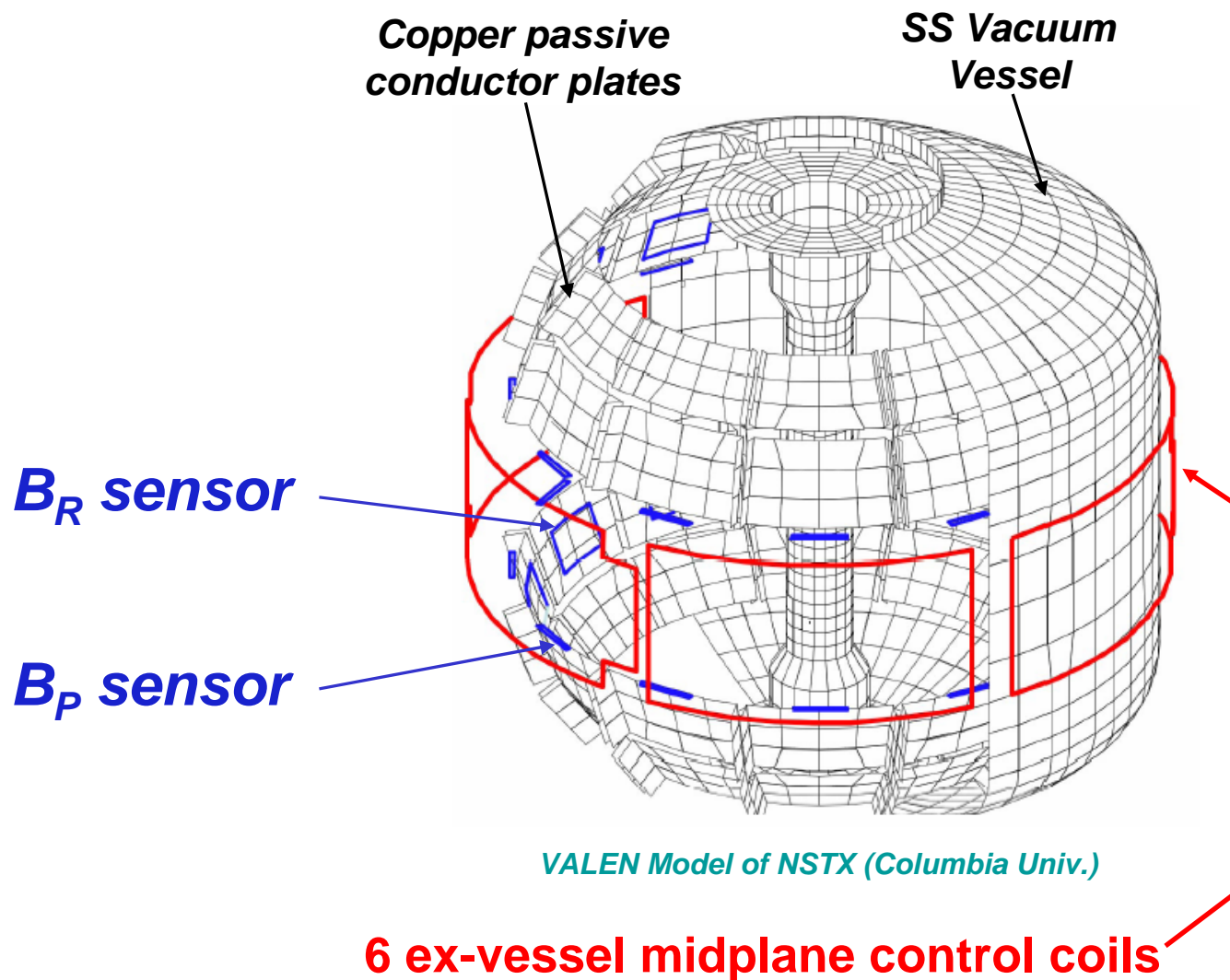
July 23-24, 2007

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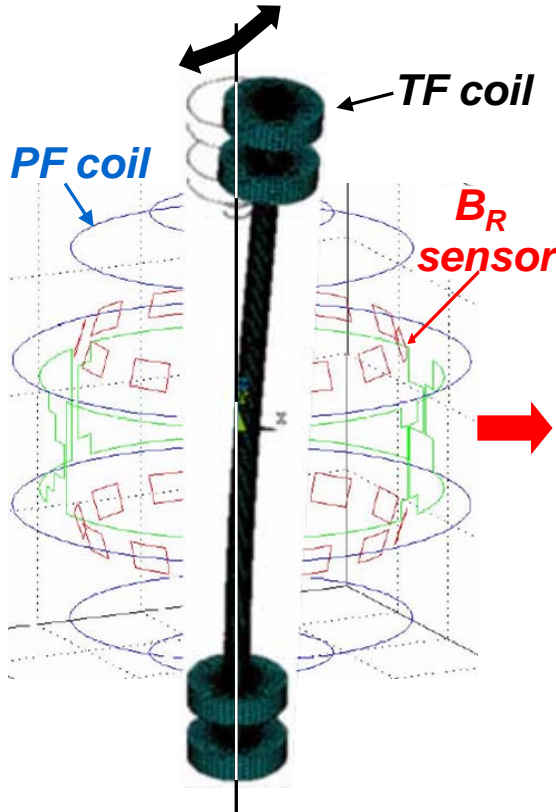
*Culham Sci Ctr
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The NSTX RWM/EF coil and sensor system

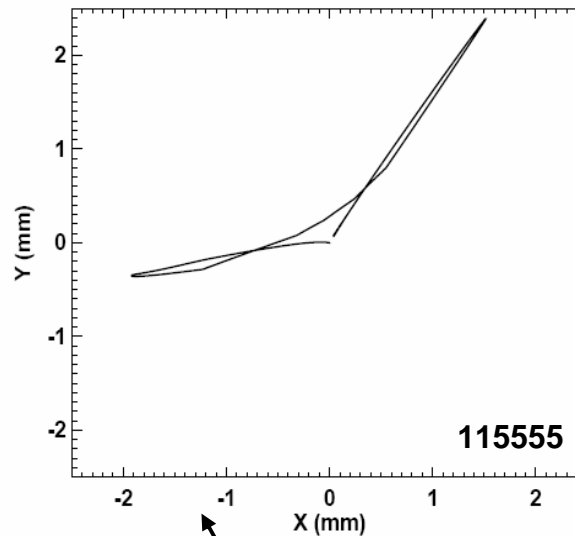


Error field source identification and compensation in NSTX

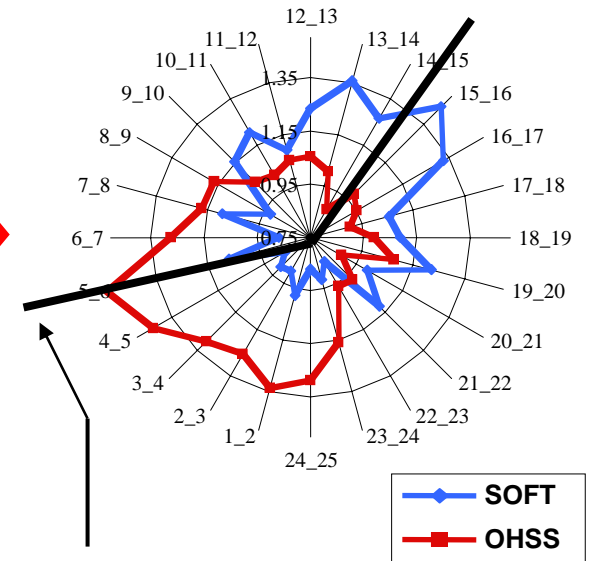
- Some scenarios with lower κ , δ exhibit Ω_ϕ and β_N collapse when $\beta_N > \beta_N$ (no-wall)
- Measure 2-3 Gauss $B_\perp^{2/1}$ EF in LM experiments... what is EF source?
- **Present picture of EF Source:** EF from/near OH leads at top of machine induces TF coil motion relative to B_R sensors (plates, vessel) and thus **the PF coils**



TF coil shift at mid-plane
inferred from B_R sensors
during OH+TF vacuum shot



Normalized TF bottom joint voltage drop from OH+TF test



TF flag-joint voltage variation direction
consistent with magnetics

Shim between TF bundle and OH tension tube added before 2007 run to reduce motion

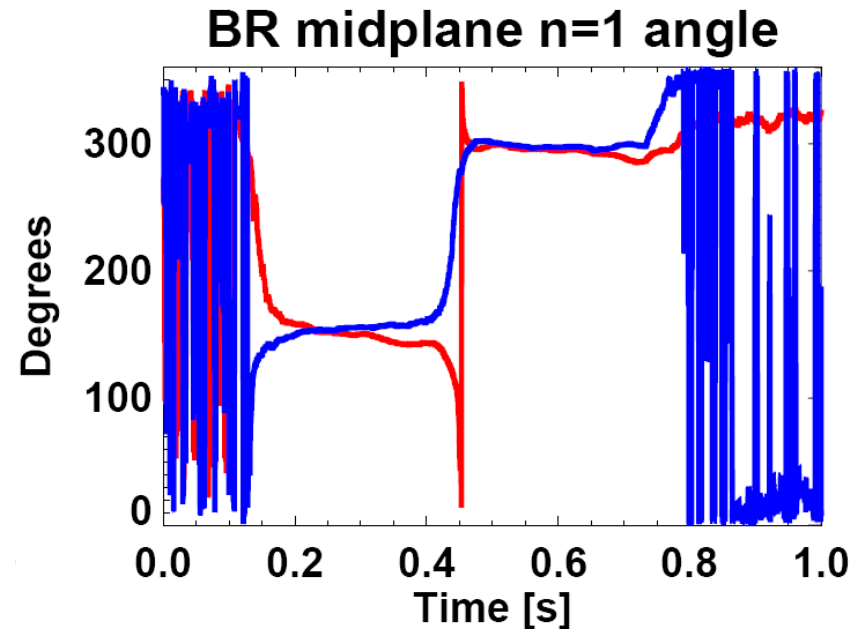
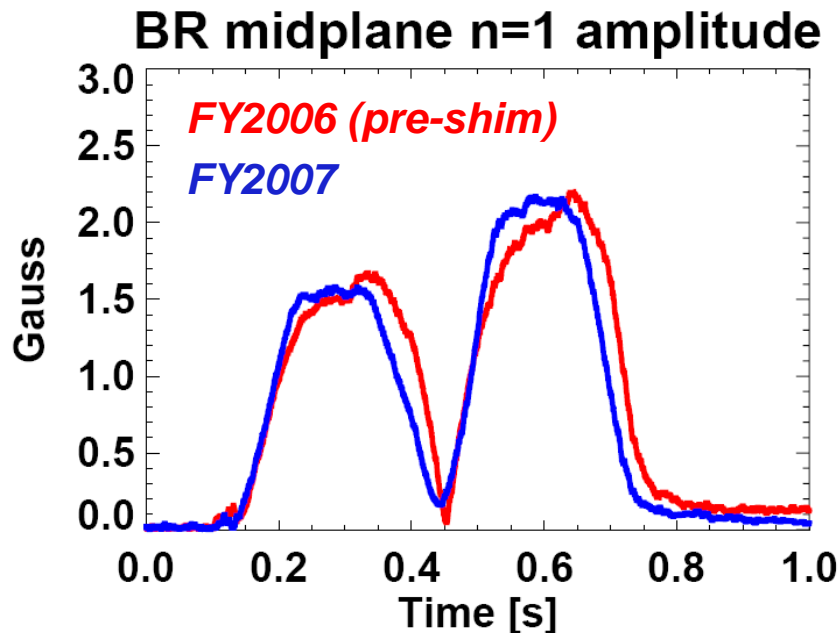
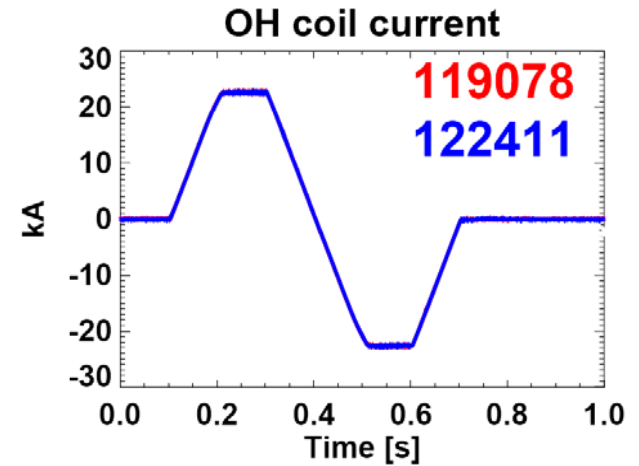
XP701 - Assessment of intrinsic error fields after TF centering



- Assessed modifications to TF coil centering w.r.t. intrinsic EF
 - Intrinsic EF very similar to 2006 for LM ohmic shots
 - Larger difference for long-pulse – lower EF after $I_{OH}=0$ crossing
- Verified rotation response asymmetry to $n=1$ applied field
- Could not reproduce 2006 reference discharge for OHxTF EFC algorithm optimization (rotation collapse not observed)
 - Used externally applied $n=1$ pulses instead in XP702
- NEW: Measured plasma response asymmetry to $n=3$
 - Pulse-length increases with “corrective” $n=3$
 - Rotation increases with “corrective” $n=3$

Error field from fast OH variation largely unchanged

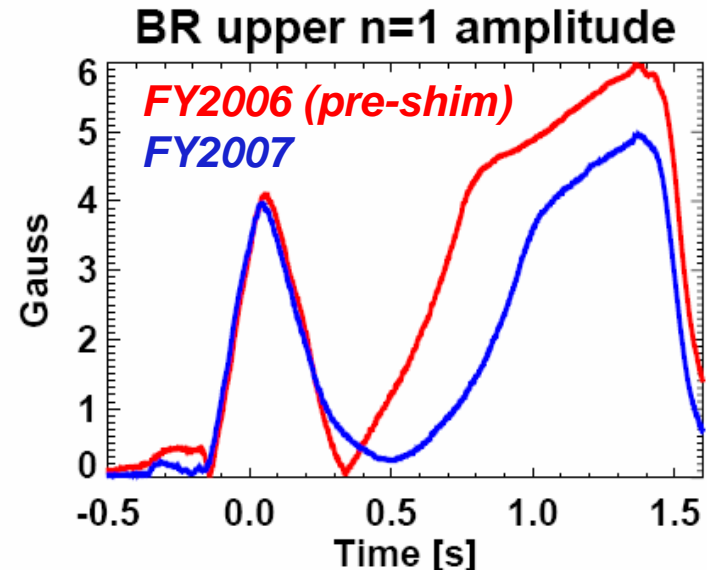
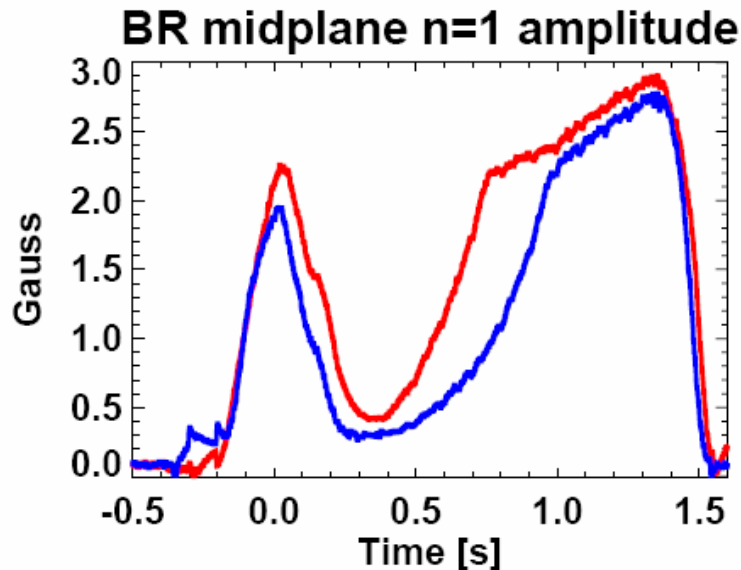
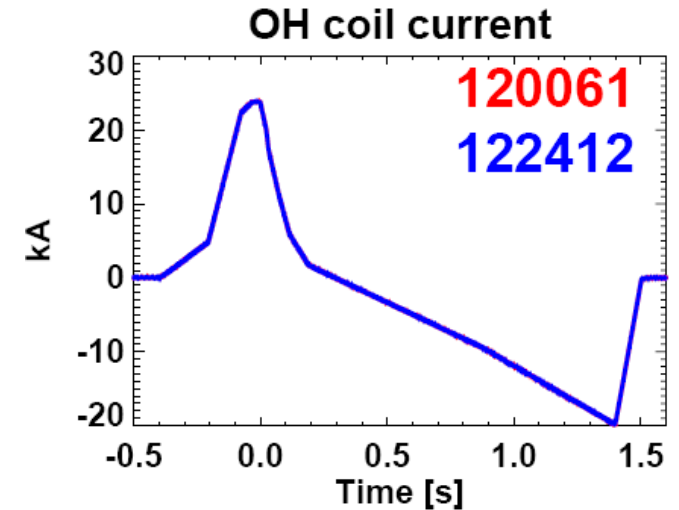
- OH waveform similar to that in ohmic discharges used for locked-mode experiments



EF from slow OH variation different after $I_{OH}=0$ crossing



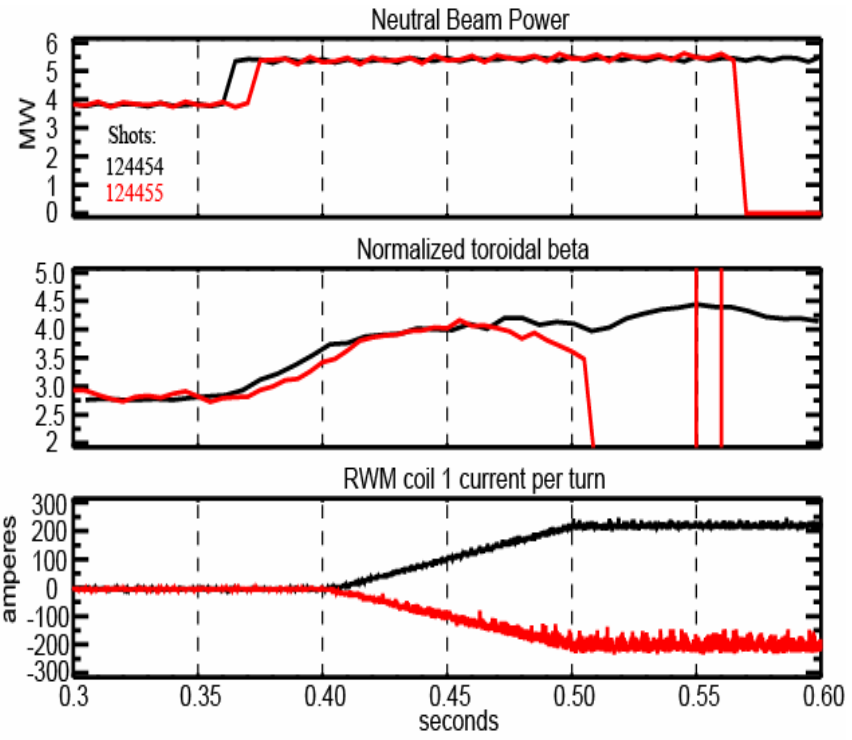
- OH waveform similar to that for NBI H-mode long-pulse
- EF increase is delayed after I_{OH} crossing, but eventually reaches similar amplitude and slope
- Midplane EF amplitude similar, upper EF amplitude reduced late in shot



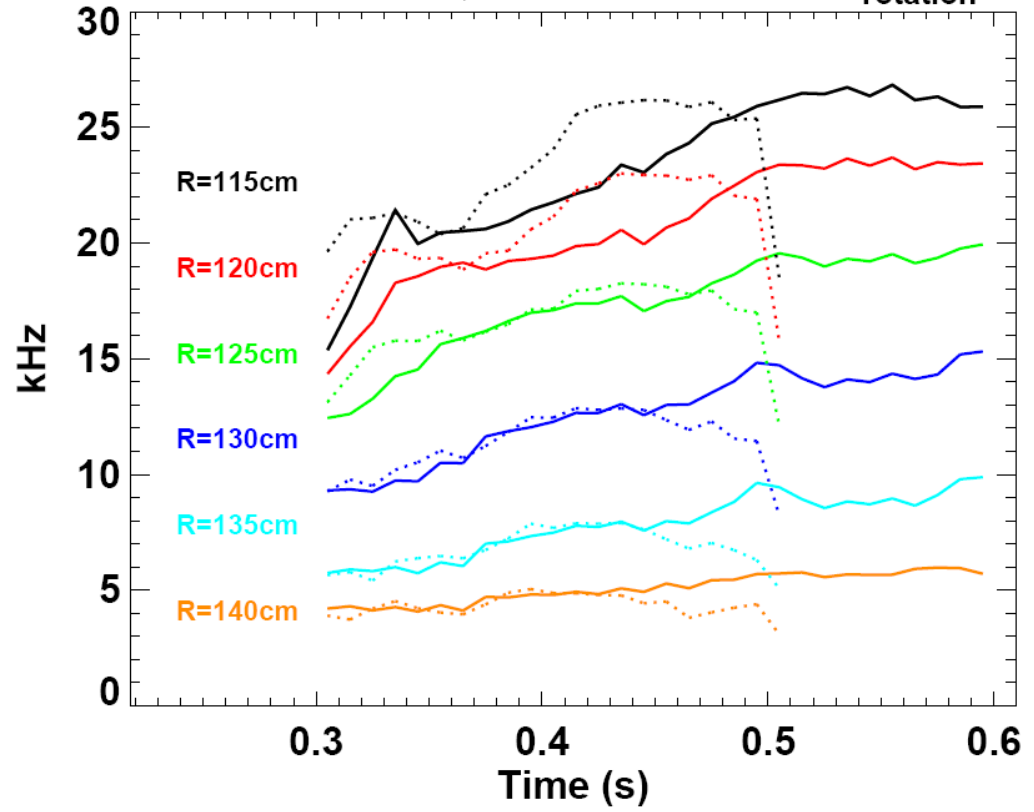
Plasma still exhibits asymmetric response to phase of applied $n=1$ field



- Rotation collapse begins at $t=0.45\text{s}$, and is most clearly evident for radial positions $R > 1.25\text{m}$



Shot 124454, 124455 carbon f_{rotation}

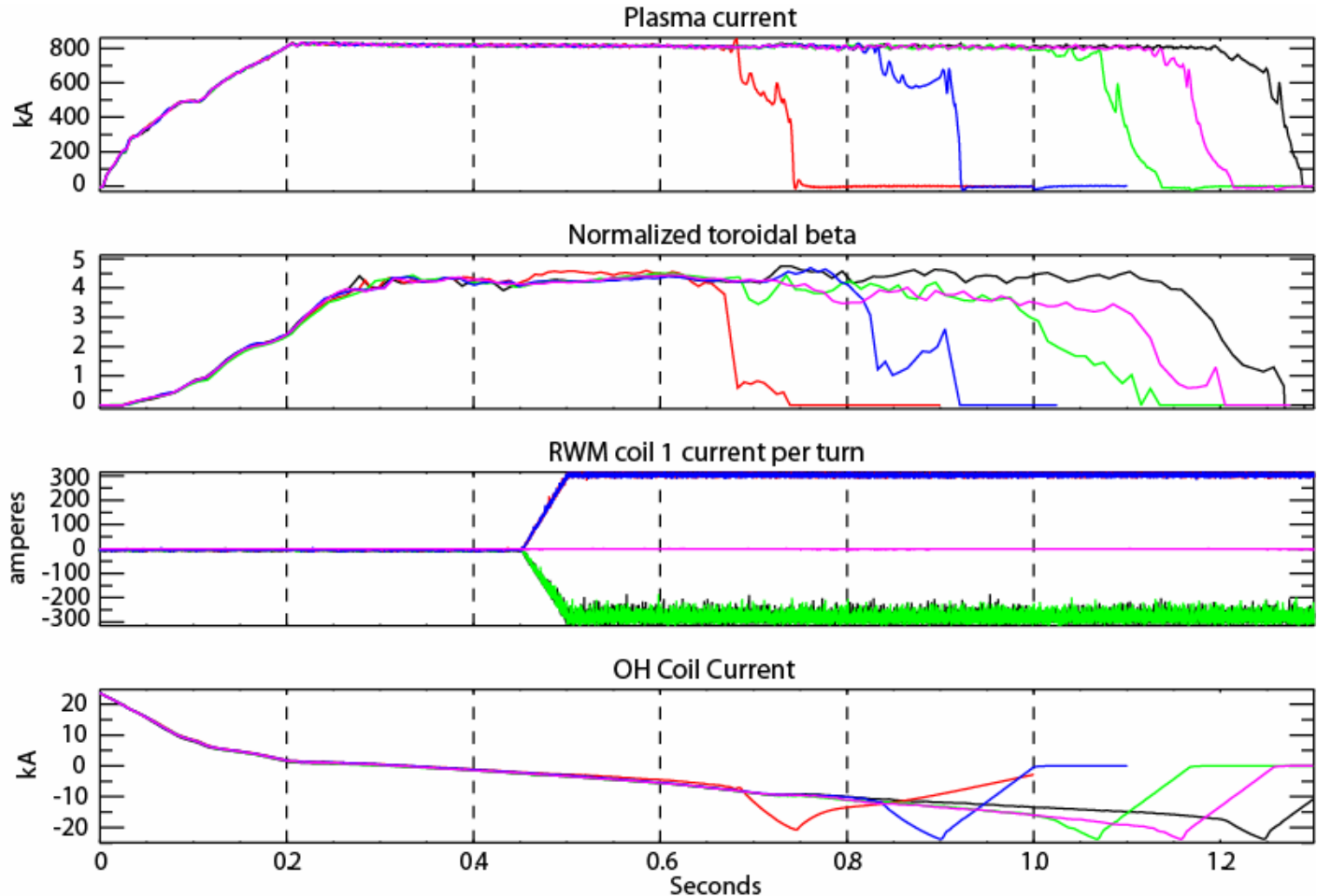


NEW: Plasma exhibits asymmetric response to applied n=3 field



- Pulse-length depends on polarity of applied n=3 field

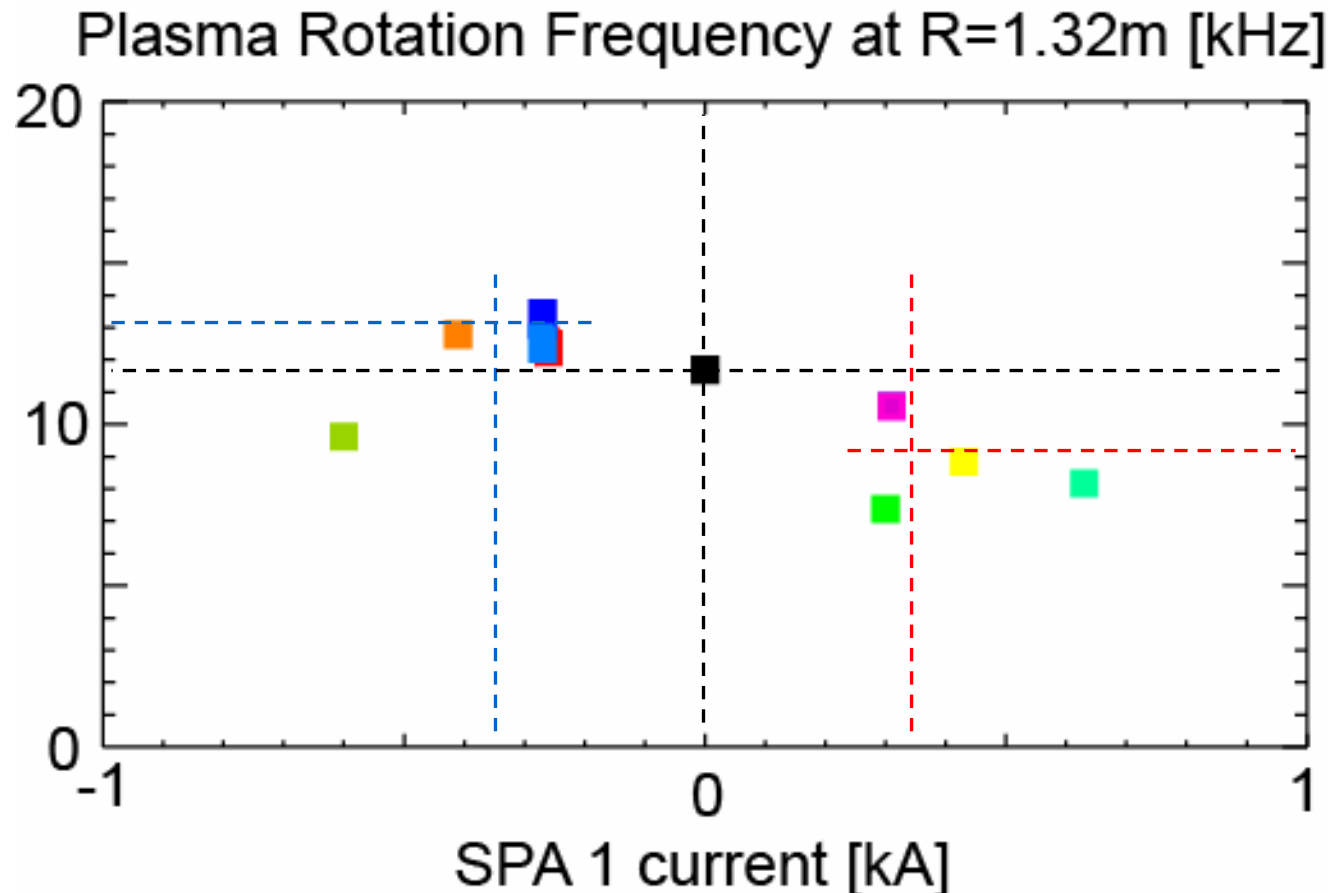
Shots:
124428
124429
124430
124432
124411



Outboard Ω_ϕ changes by 30-40% with n=3 polarity flip



- Optimal n=3 current magnitude = 300-400A
- PF5 coil shape data from 2004 → PF5 is source of n=3 EF
 - Need to assess if this is consistent with empirical correction below



XP702 - RFA detection optimization during dynamic error field correction



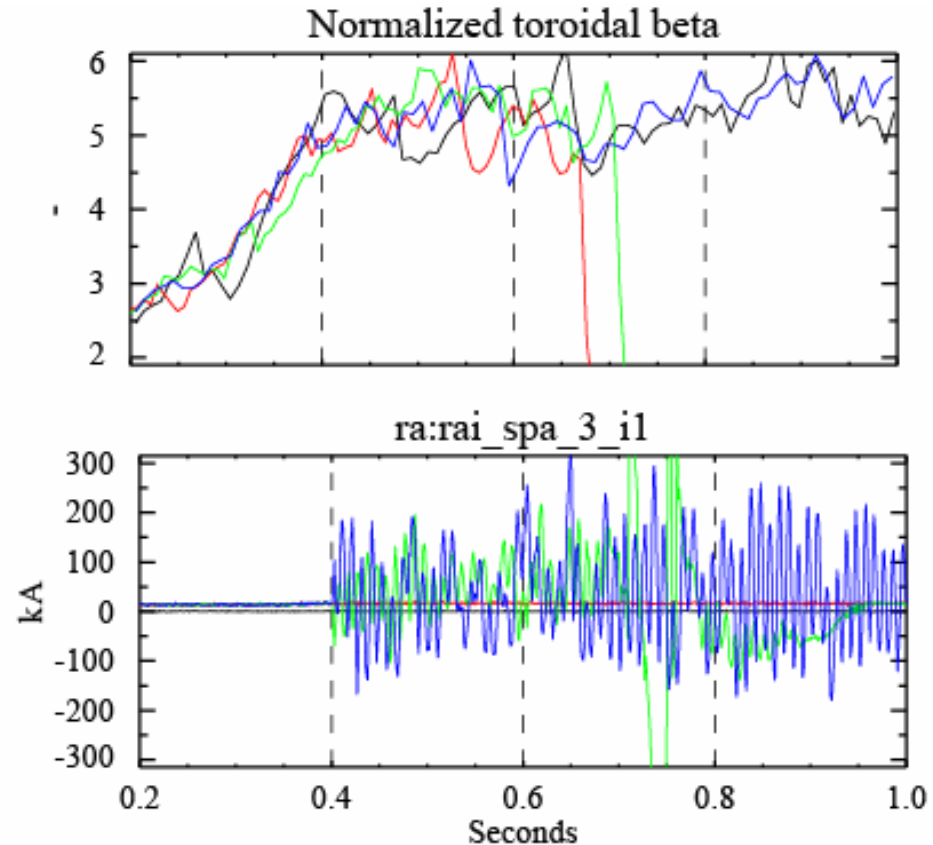
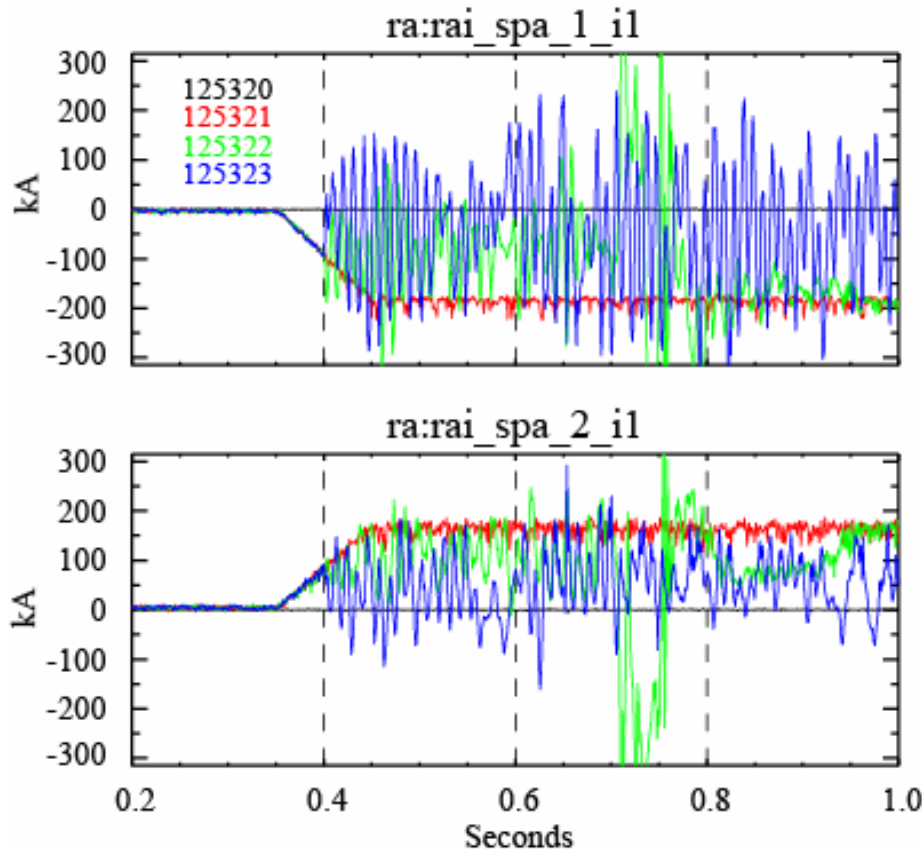
- Implemented real-time mode-ID using U+L B_p and B_R sensors
- Compared DEFC response using upper and lower B_p sensors vs. just using upper B_p sensors (as was used in 2006)
 - More robust mode-ID achieved (higher signal / baseline offset)
 - Higher proportional gain possible (0.7-0.8 vs. 0.5)
- Could not reproduce 2006 reference discharge which previously exhibited intrinsic rotation collapse →
 - Instead, applied $n=1$ EF pulse to induce collapse when OHxTF small
 - Scanned DEFC phase and gain until applied currents were nulled
 - **feedback system “trained” to eliminate EFA of known source**
- Combination of “trained” $n=1$ DEFC + $n=3$ EFC →
 - longest pulse of all shots in XP702
 - sustained plasma rotation

DEFC system trained to null RFA from externally applied/known n=1 error field source



- Optimal phase difference ($\delta=270^\circ$) between measured B_p (U/L avg) and applied B_R required to null n=1 EF pulse
- Sufficient feedback gain also required:

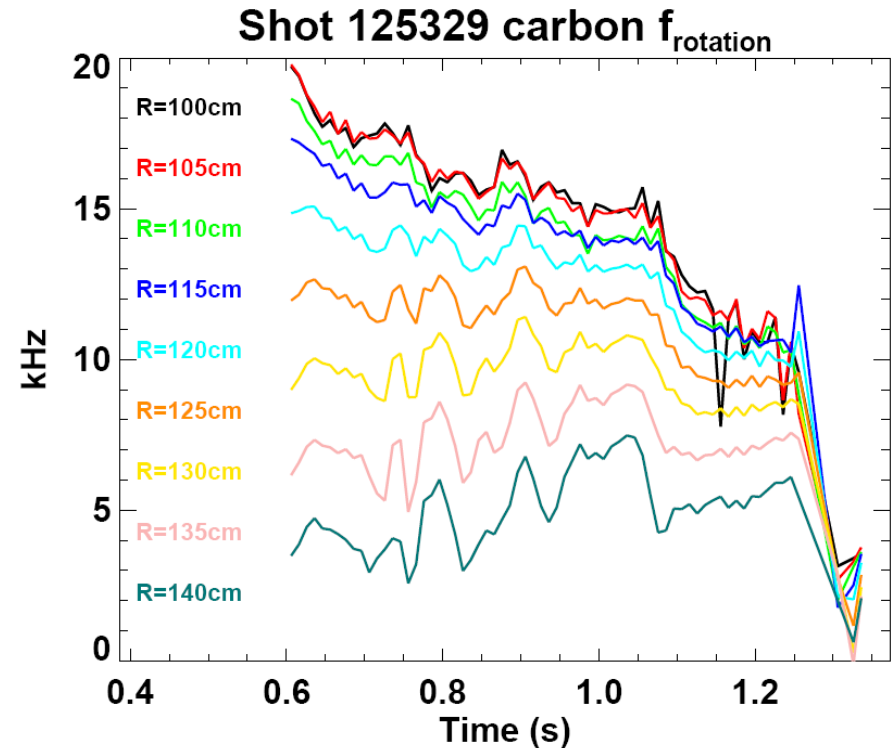
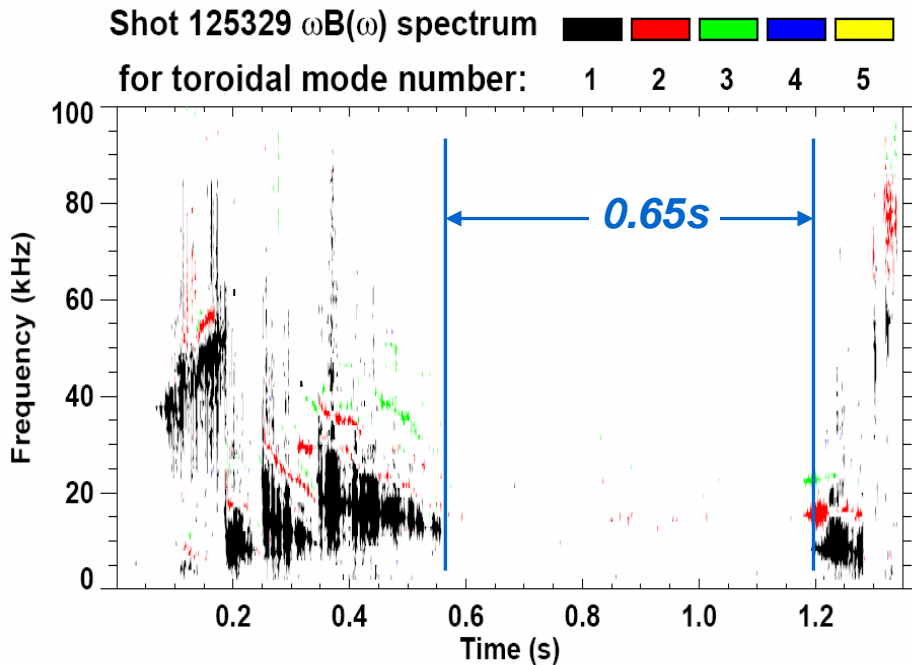
$G_p=0.0$ $G_p=0.5$ $G_p=0.7$



n=1 RFA feedback + n=3 EFC improves performance



- Long period free of core low-f MHD activity
 - Plasma rotation sustained over same period
 - Core rotation decreases with increasing $f_{GW} \rightarrow 0.75$, $P_{RAD} \rightarrow 3.5\text{MW}$
 - $R > 1.2\text{m}$ rotation slowly increases until $t=1.1\text{s}$ (large ELM?)
- Longest pulse at $I_p=900\text{kA}$



XP703 - B and q scaling of low-density locked-mode threshold at low-A



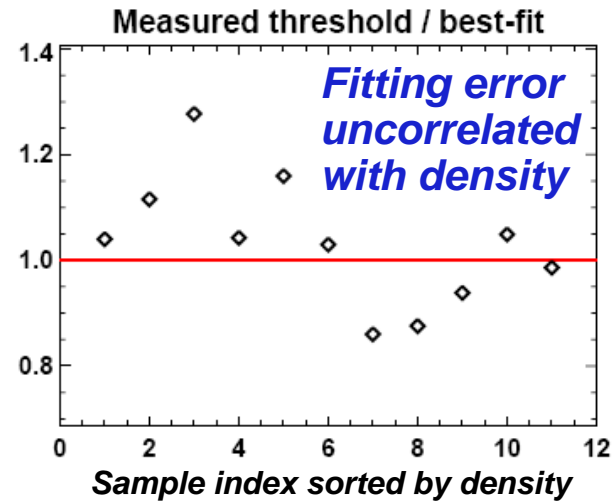
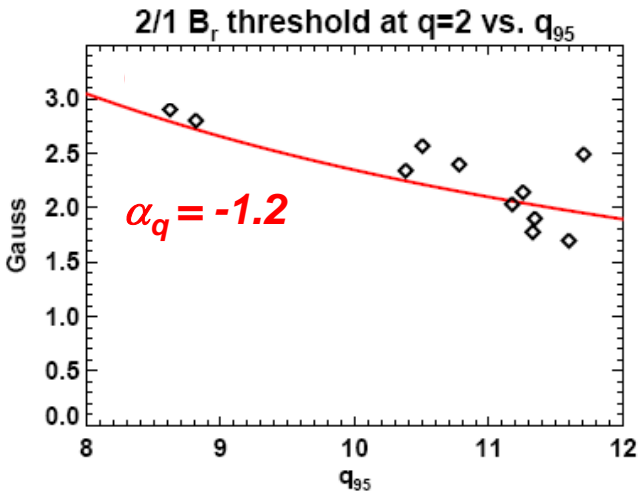
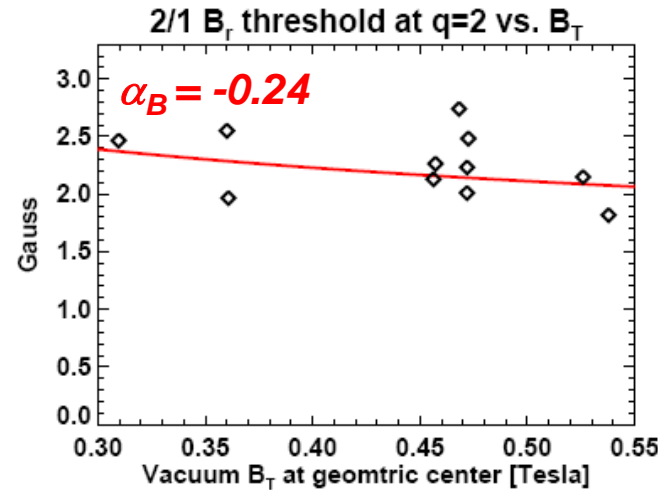
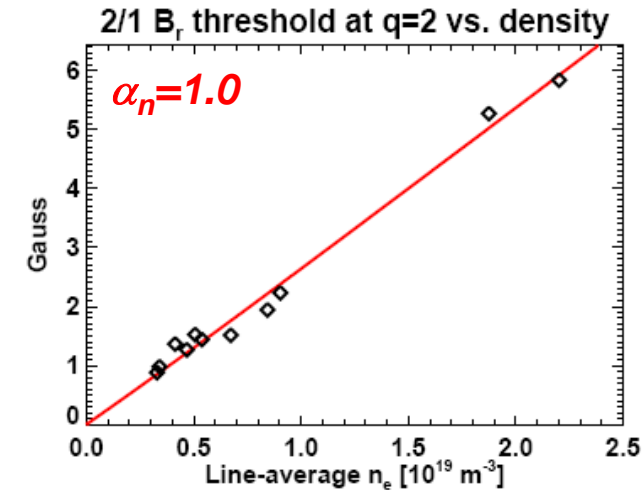
- Extended density range for threshold – now have factor of 4
- Performed B_T scan from 3kG to 5.5kG
- q_{95} scan difficult because high-q does not have $q=2$ in plasma
 - Found this out after MSE data was obtained
- Threshold increases with increased edge q-shear (w/o MSE)
 - Similarly - also increases with internal inductance
- Obtained MSE data for 4 scenarios of interest
 - $q=2$ surface is in plasma at time of locking, but NO $q=1$ surface
 - Core shear is often weakly reversed
 - Measured q profiles not yet included in analysis shown below!
- Locking threshold scaling favorable for ITER

NOTE: Scaling form used below: $B_{21}(\text{lock}) \propto n^{\alpha_n} B_T^{\alpha_B} q^{\alpha_q} R^{\alpha_R}$

NSTX locking data shows linear density scaling and weak B_T dependence, but unexpected inverse scaling with q_{95}



$2/1 B_r$

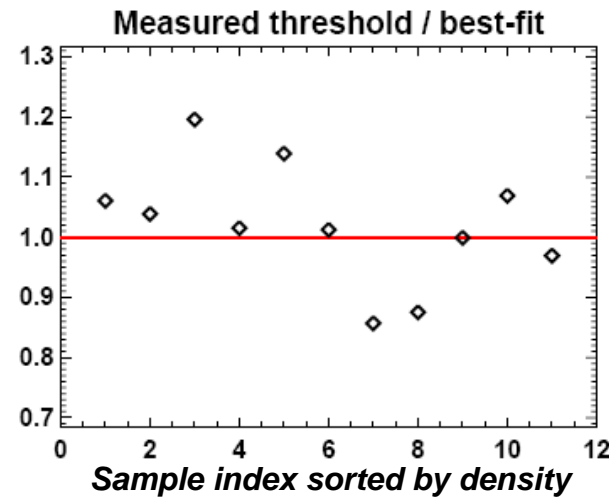
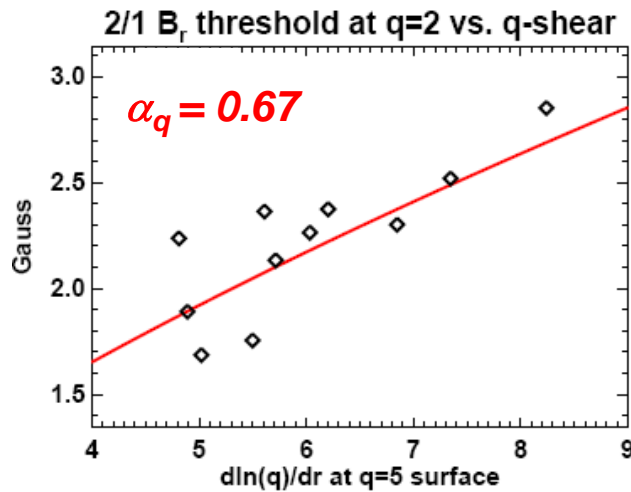
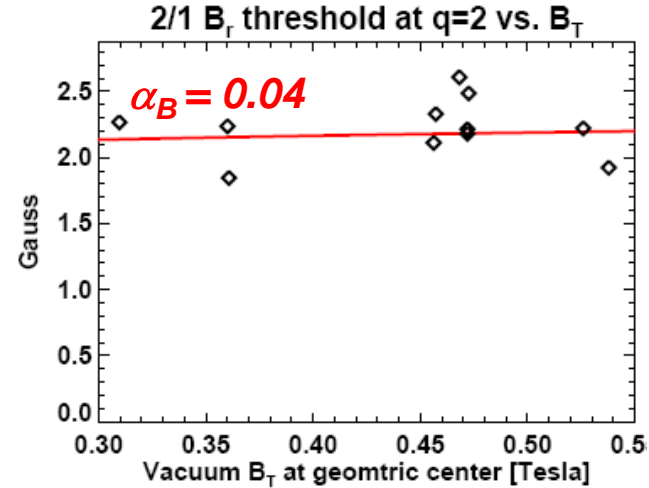
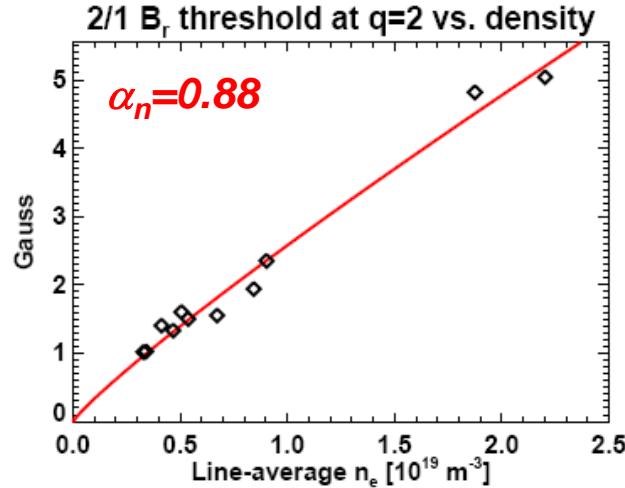


MSE data (not used for scaling above) shows variation in q_{min} when q_{95} is varied.
 $q(0)$ does not = 1 as for other experiments \rightarrow q_{95} not good proxy for q -shear for NSTX.

NSTX locking data shows nearly linear density scaling and very weak B_T dependence, and expected positive scaling with edge q-shear



$2/1 B_r$

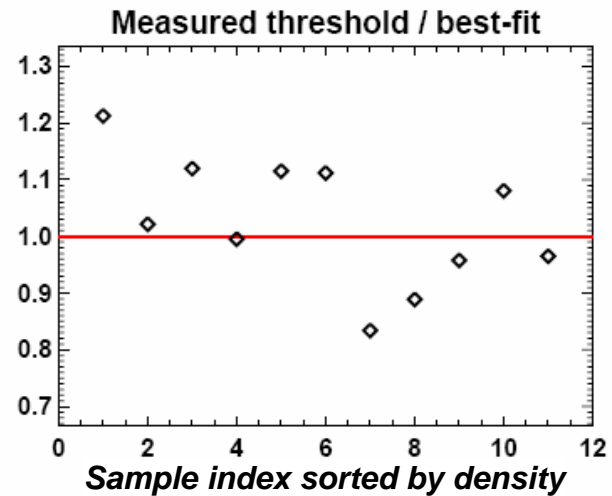
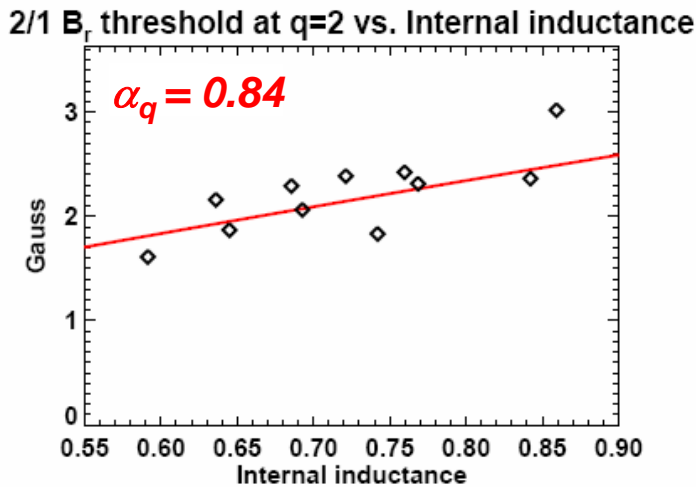
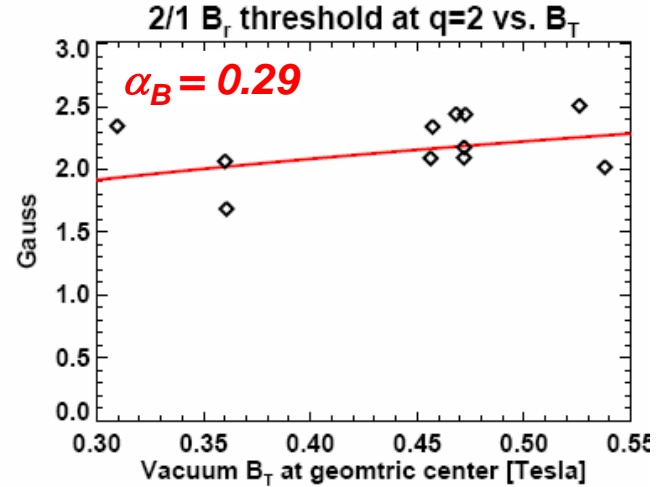
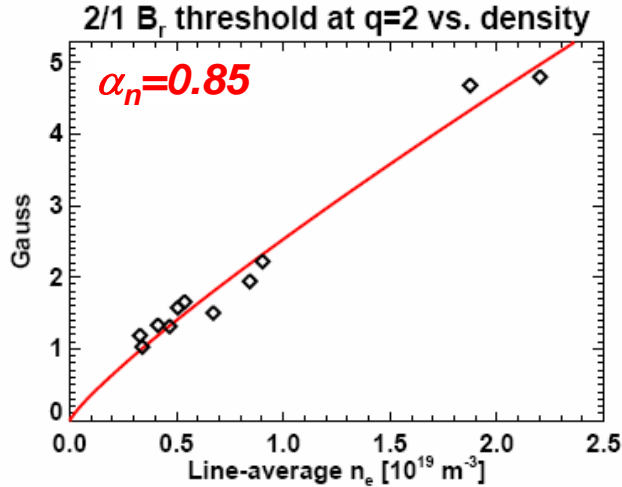


Density and B-field scalings are sensitive to choice of q-scaling variable, but, assuming size scaling coefficient $\alpha_R = 2\alpha_n + 1.25(\alpha_B - 1) \rightarrow$ NSTX $\alpha_R = 0.45$ to 0.56

NSTX locking data shows nearly linear density scaling, weak positive B_T dependence, and nearly linear scaling with internal inductance



$2/1 B_r$



Density and B-field scalings are sensitive to choice of q -scaling variable:
 Range in NSTX: $\alpha_R = 0.45$ (q_{95}), $\alpha_R = 0.56$ (q' at $q=5$), $\alpha_R = 0.8$ (I_i)

Large extrapolation from NSTX to ITER, but here it is...



For NSTX q_{95} scaling with:

$$n_e=10^{19}\text{m}^{-3}, B_T=5.7\text{T}, R_0=6\text{m}, \alpha_R = 0.45, q_{95}=3$$

$$\text{ITER } B_{21} \rightarrow 14\text{G}, \text{ or } B_{21}/B_T = 2.5 \times 10^{-4}$$

For NSTX I_i scaling with:

$$n_e=10^{19}\text{m}^{-3}, B_T=5.7\text{T}, R_0=6\text{m}, \alpha_R = 0.8, I_i=1.0$$

$$\text{ITER } B_{21} \rightarrow 30\text{G}, \text{ or } B_{21}/B_T = 5.2 \times 10^{-4}$$

Caution – no $q=1$ surface in NSTX plasmas which could lower thresholds in ITER (and NSTX)

Also need to propagate uncertainties through analysis properly...

IPEC analysis not yet systematically included in scalings



- IPEC analysis shows that total B_{\perp}^{mn} (including plasma response) on resonant surfaces differs significantly from vacuum external B_{\perp}^{mn}
 - So how did we account for intrinsic error field in NSTX if plasma response is important, and we didn't include it?
- Empirically find that different normalization for B_{\perp}^{mn} can improve accounting for EF using only vacuum fields:

– Instead use perturbed helical flux:
$$\delta\psi_{mn}^h = \sigma \left[R^3 \left| \frac{q}{F} \right| (B_Z \delta B_R - B_R \delta B_Z) \right]_{mn}$$

Where σ is a polarity and phase factor and $F = B_{\phi} R$

– $\delta\psi^h$ scales as $R^2 |\nabla\psi| q/F B_{\perp}^{mn}$

R^2 scaling of $\delta\psi^h$ apparently provides better geometric representation of fields that generate singular currents, torques, and mode locking

ALL NSTX DATA SHOWN ABOVE USES RENORMALIZED VACUUM $\delta\psi^h$ to compute $\delta B_{2/1}$

