

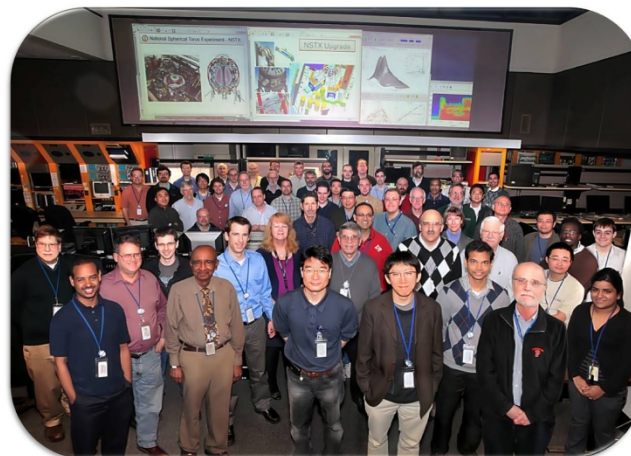
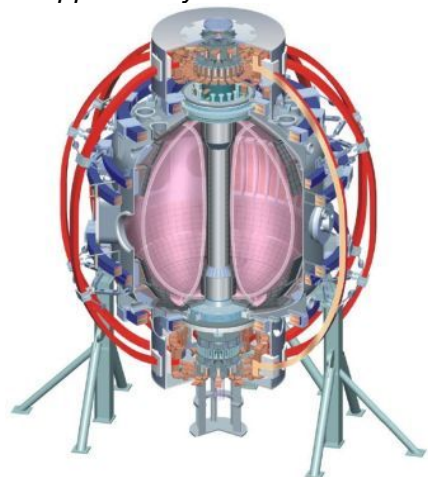
Novel Alfvén eigenmode structure measurements in NSTX via reflectometry*

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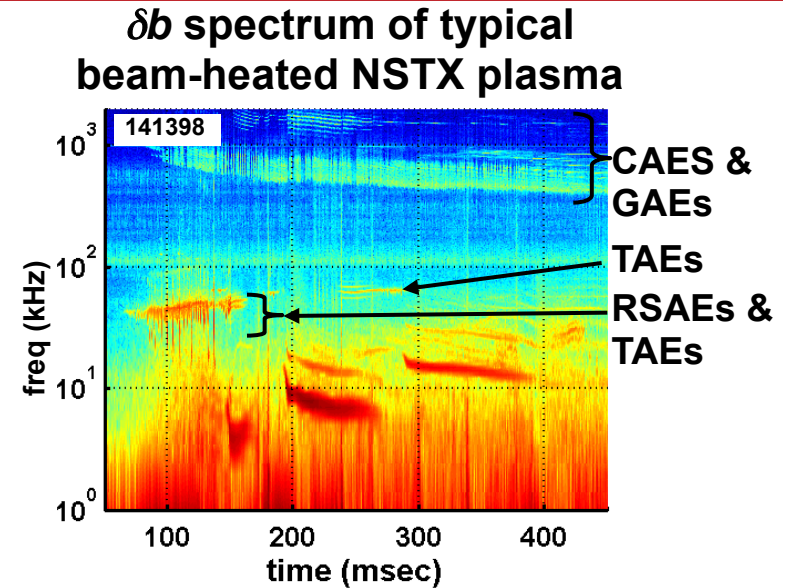
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Summary

- Motivation: Alfvén eigenmodes (AE) can impact plasma performance
 - Low Frequency AEs (TAE & RSAE), $f < \sim 200$ kHz \Rightarrow fast-ion transport
 - High Frequency AEs (CAE & GAE), $f > \sim 400$ kHz \Rightarrow correlated with enhanced χ_e
- AE structure measured via 16-channel reflectometer array,
 $n_0 \sim 1 - 7 \times 10^{19} \text{ m}^{-3}$ (30 – 75 GHz)
- High frequency AE structure measured in core of beam heat-heated H-mode plasmas \Rightarrow identified as CAEs and GAEs
- **CAEs strongly core localized \Rightarrow may cause thermal electron transport**
 - GAE activity correlates with enhanced electron thermal transport; proposed as cause [D. Stutman *et al.*, PRL **102** 115002 (2009); K. Tritz, APS DPP 2010 Invited PI2.2]
 - CAEs share key GAE characteristics: core localization & frequency
- **TAE structure** measured (amplitude and phase) with high resolution \Rightarrow **strong radial phase variation across plasma ($\Delta\phi \sim \pi/2$)**
 - indicates radial propagation
- TAE phase variation indicates non-ideal physics
 - inconsistent with ideal MHD
 - M3D-K simulation implicates coupling with fast-ions

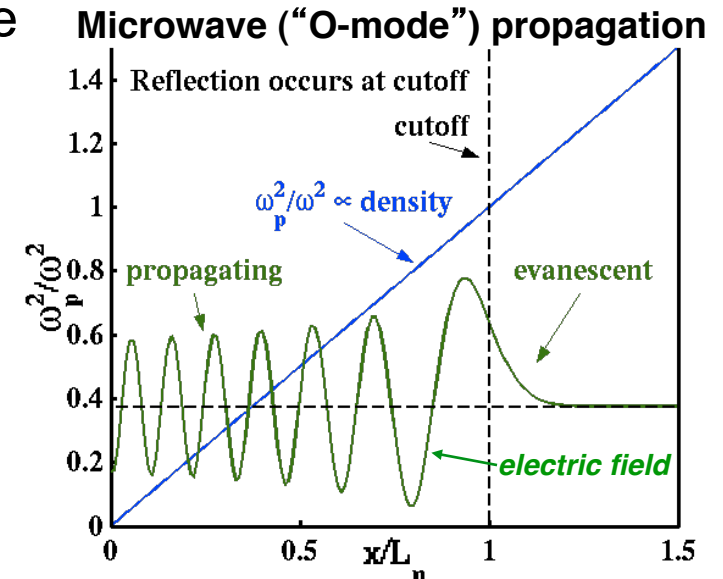
Motivation: Alfvén eigenmode structure measurement promotes better understanding of plasma performance

- NSTX plasmas feature rich spectrum of Alfvén eigenmodes
 - Low frequency AEs ($f < \sim 200$ kHz): Toroidicity-induced (TAE) & Reversed shear (RSAE)
 - High frequency AEs ($f > \sim 400$ kHz): Compressional (CAE) & Global (GAE)
- Alfvén eigenmodes (AE) play critical role in many aspects of plasma performance
 - Low frequency AEs cause fast-ion transport and loss:
 - change equilibrium sources (momentum, energy ...)
 - damage plasma facing components
 - High frequency AE activity correlates with enhanced χ_e in core of H-mode beam heated plasmas
 - D. Stutman *et al.*, PRL **102** 115002 (2009); K. Tritz, APS DPP 2010 Invited PI2.2
- Mode δn structure routinely measured in NSTX via fixed-frequency reflectometer radial array
 - 16 channels, $n_0 \sim 1 - 7 \times 10^{19} \text{ m}^{-3}$ (30 – 75 GHz)



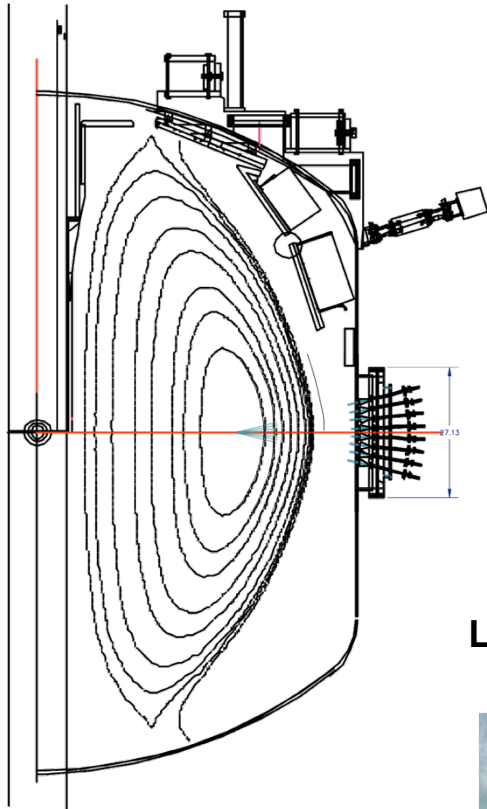
Reflectometers measure local density fluctuation in plasma

- Microwaves propagate to “cutoff” layer, where density high enough for reflection ($\omega_p = \omega$)
 - Dispersion relation of “ordinary mode” microwaves: $\omega^2 = \omega_p^2 + c^2k^2$,
 ω_p^2 proportional to density ($\omega_p^2 = e^2n_0/\epsilon_0m_e$)
 - $k \rightarrow 0$ as $\omega \rightarrow \omega_p$,
microwaves reflect at $k = 0$
- Reflectometer measures path length changes of microwaves reflected from plasma
 - phase between reflected and launched waves changes ($\delta\phi$)
- Wave propagation controlled by density
 - for large scale modes $\delta n/n_0 \sim \delta\phi/(2k_{\text{vac}}L_n)$, $L_n = n_0/|\nabla n_0|$



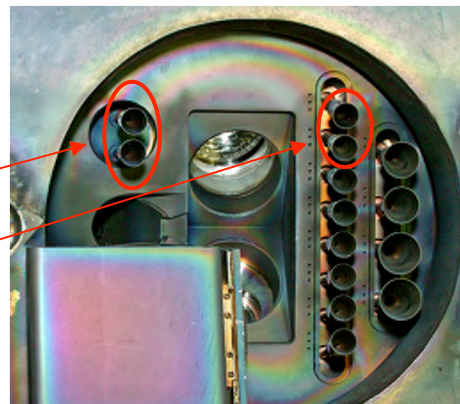
Reflectometers provide radial array of measurements

NSTX cross-section



- Two arrays: “Q-band” & “V-band”
 - Q-band: 30, 32.5, 35, 37.5, 42.5, 45, 47.5 & 50 GHz
 - V-band: 55, 57.5, 60, 62.5, 67.5, 70, 72.5 & 75 GHz
- Arrays closely spaced (separated $\sim 10^\circ$ toroidal)
- Single launch and receive horn for each array
- **Horns oriented perpendicular to flux surfaces \Rightarrow frequency array = radial array**
- Cutoffs span large radial range in high density plasmas ($n_0 \sim 1 - 7 \times 10^{19} \text{ m}^{-3}$)

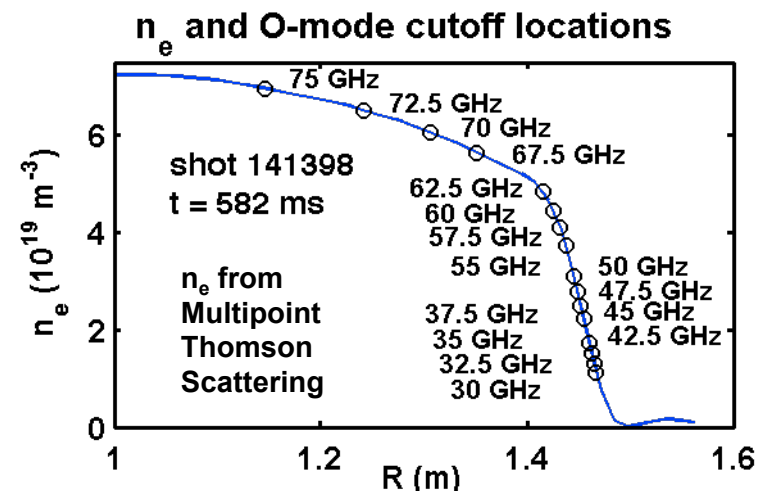
Launch and Receive Horns
(Interior View)



30-50 GHz

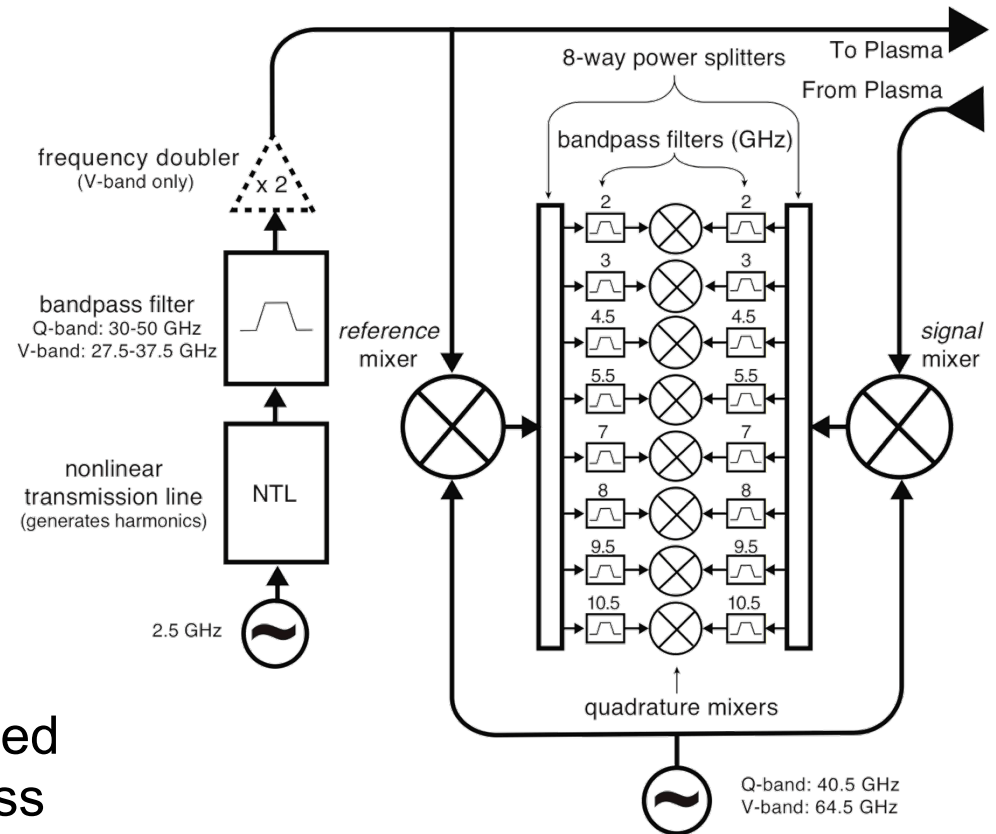
55-75 GHz

(not shown: horns modified to optimize for frequency range)



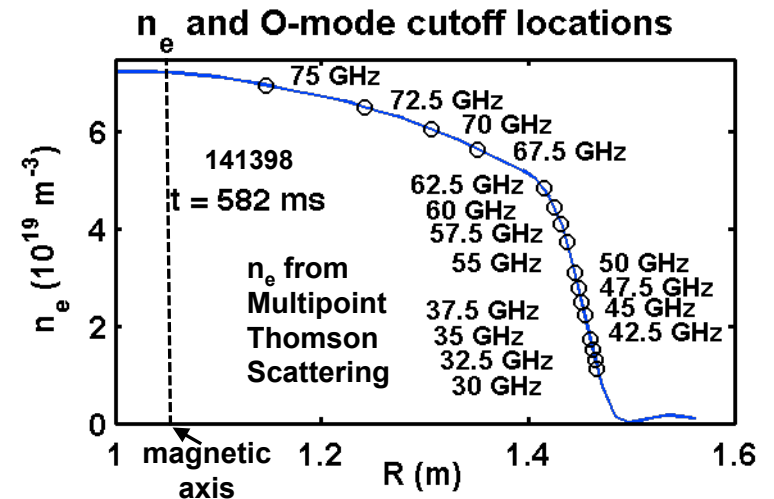
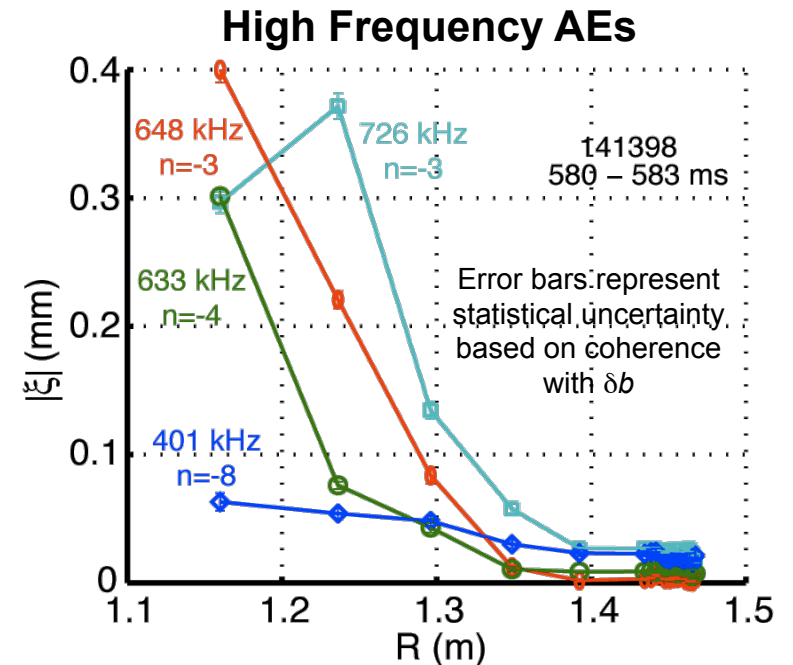
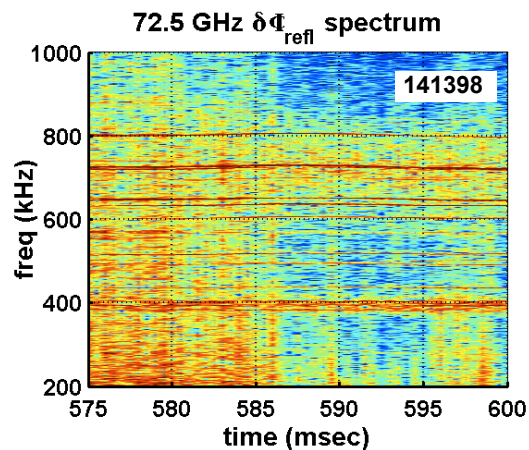
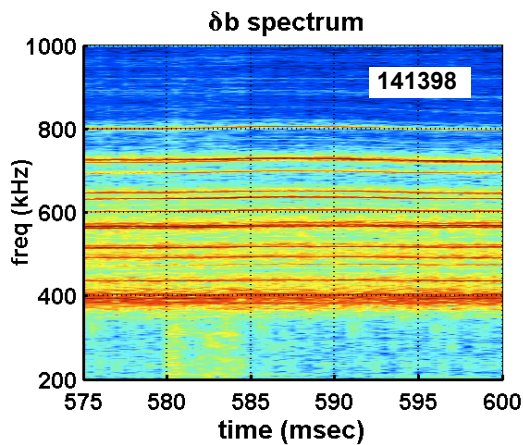
Microwave circuit exploits nonlinear transmission line

- 2.5 GHz source + nonlinear transmission line \Rightarrow harmonics up to 50 GHz
- Bandpass filter selects frequency range
 - Q-band: 30 – 50 GHz
 - V-band: 27.5 – 37.5 GHz (frequencies doubled after filter)
- Frequencies downshifted to ≤ 10.5 GHz for processing
- Individual freq. isolated & processed via 8-way power splitters, bandpass filters & quadrature mixers



High Frequency AE structure measured in core of H-mode beam heated plasma*

- High frequency AE structures measured with 16 channels in 6 MW beam-heated H-mode plasma
- Observed modes fall in **two categories**:
 - Small amplitude + broad structure, $f < \sim 600$ kHz, $n = -7$ & -8 (e.g. $f = 401$ kHz on right)
 - Large amplitude + strongly core localized, $f > \sim 600$ kHz, $n = -3, -4$ & -5 (e.g. $f = 633, 648$ and 726 kHz on right)



*N A Crocker et al. PPCF 53 105001 (2011)

Modes can be identified as CAEs or GAEs via mode number and frequency evolution

- GAE and CAE frequency have distinctly different sensitivity to equilibrium

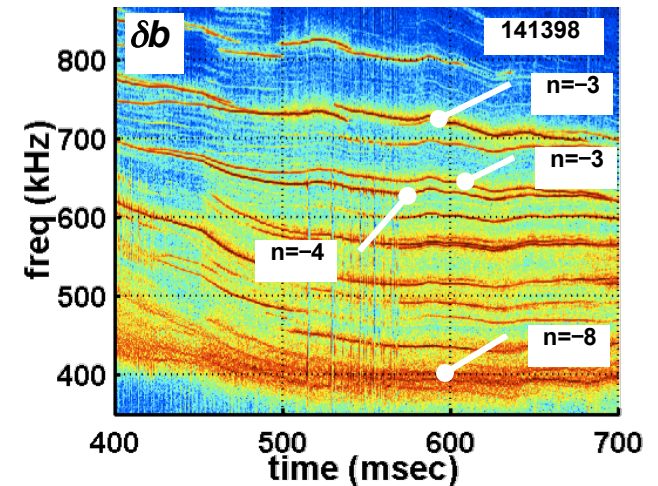
- Mode frequency evolution can be compared to local dispersion relations:

$$f_{CAE} \sim \frac{k v_A}{2\pi} + n f_{ROT}, \quad k \sim \left[\left(\frac{n}{R} \right)^2 + \left(\frac{m}{r} \right)^2 + \delta_r^{-2} \right]^{1/2}$$

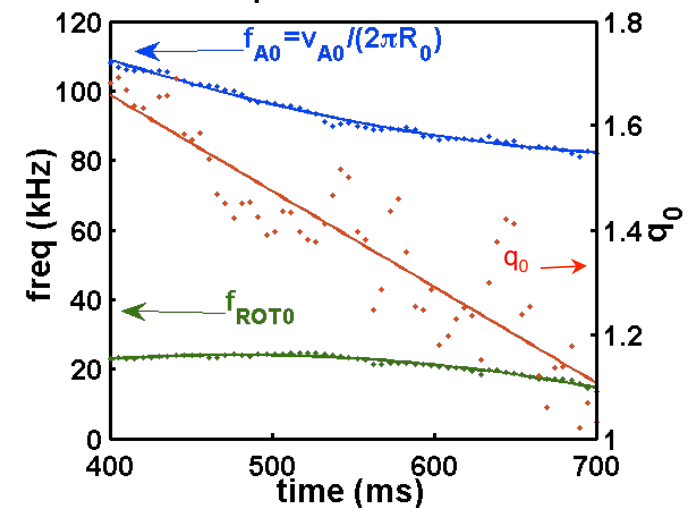
$$f_{GAE} = \frac{k_{\parallel} v_A}{2\pi} + n f_{ROT}, \quad k_{\parallel} \approx \frac{1}{R} \left| \frac{m}{q} - n \right|$$

- Toroidal mode numbers (n) measured via 10 B-dot coils in toroidally distributed array
- Equilibrium trends known from measurement:
 - q_0 and B_0 from equilibrium reconstruction using magnetic field pitch from Motional Stark Effect
 - n_{e0} measured via Multipoint Thomson Scattering
 - Alfvén velocity, $v_{A0} = B_0 / (\mu_0 \rho_0)^{1/2}$
 - $\rho_0 = m_D n_{e0}$, $m_D = \text{Deuterium mass}$
 - Toroidal rotation frequency, f_{ROT0} , from Charge Exchange Recombination Spectroscopy

AE frequency evolution

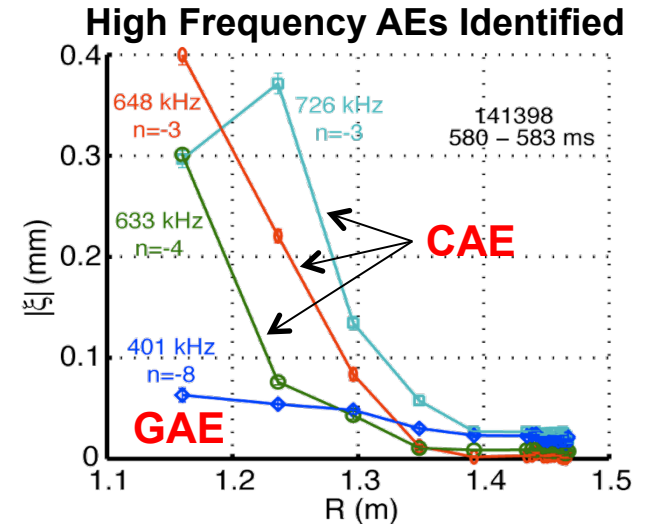


Equilibrium trends

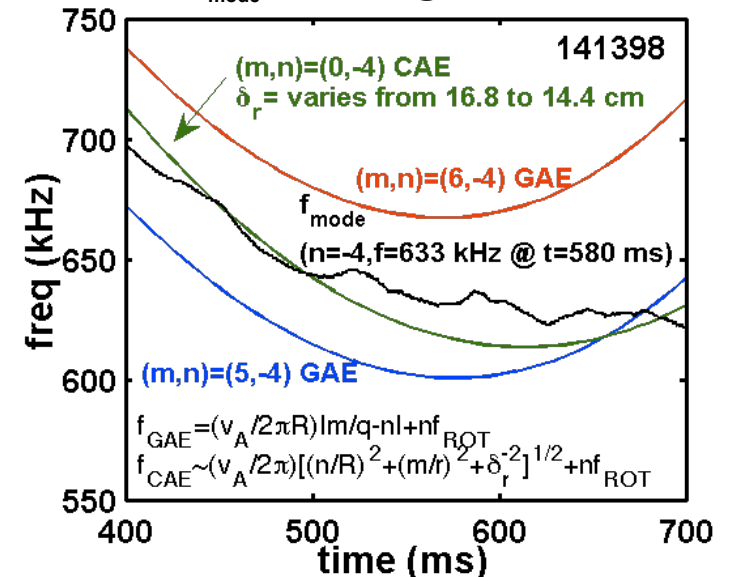


Frequency sensitivity to q_0 key in distinguishing CAEs from GAEs

- Expected GAE frequency evolution sensitive to m/q_0 if $|m| \gg 1$
 - q_0 varies substantially (1.7 – 1.1) over $t = 400 - 700$ ms
- Modes with $f < \sim 600$ kHz, $n = -7$ & -8 Identified as GAEs
 - $|n| \gg 1 \Rightarrow f(t) \sim f_{\text{GAE}}(t)$ if low $|m|$
- Modes with $f > \sim 600$ kHz, $n = -5, -4$ & -3 Identified as CAEs
 - $f(t)$ consistent with $f_{\text{CAE}}(t)$
 - Assume low $|m|$: $m = 0, 1$
 - Assume slowly varying plausible radial scale (from mode structure): $\delta_r \sim 15 - 25$ cm
 - $f(t)$ NOT consistent with $f_{\text{GAE}}(t)$: low $|n|$, high $f \Rightarrow$ high $|m| \Rightarrow$ strong q_0 sensitivity



Example of identification analysis:
 $f_{\text{mode}} = 633$ kHz @ $t = 580$ ms



Identification of core-localized CAEs motivates future investigation of CAEs as possible of cause enhanced χ_e

- GAEs proposed as cause enhanced χ_e in core of NSTX beam-heated H-mode plasmas
 - D. Stutman *et al.*, PRL **102** 115002 (2009)
 - K. Tritz, APS DPP 2010 Invited PI2.2
 - High frequency AE activity, identified as GAEs, correlates with enhanced χ_e in core
 - GAE frequency \sim electron orbit frequencies \Rightarrow resonant orbit modification
 - GAE core localization expected \Rightarrow active in region of enhanced χ_e
- Modes identified as CAEs can also modify electron orbits in core
 - CAEs shown to be core localized, with frequencies similar to GAEs
 - CAEs stronger & more core localized than GAEs \Rightarrow more impact on electrons than GAEs?

D. Stutman *et al.*, PRL 102 115002 (2009)

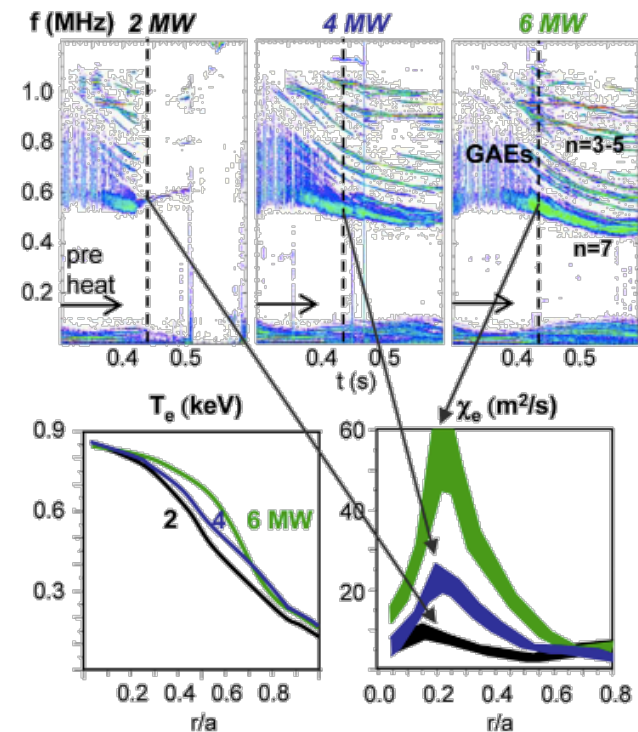
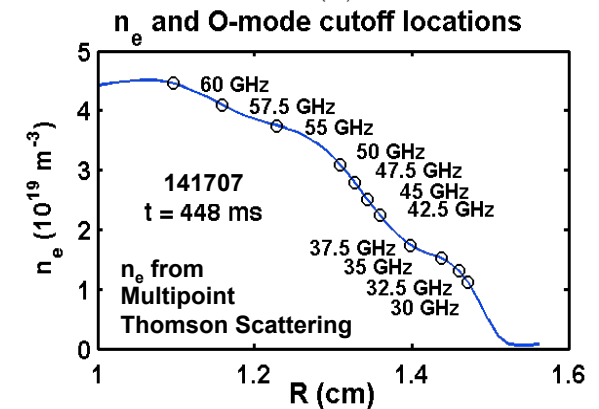
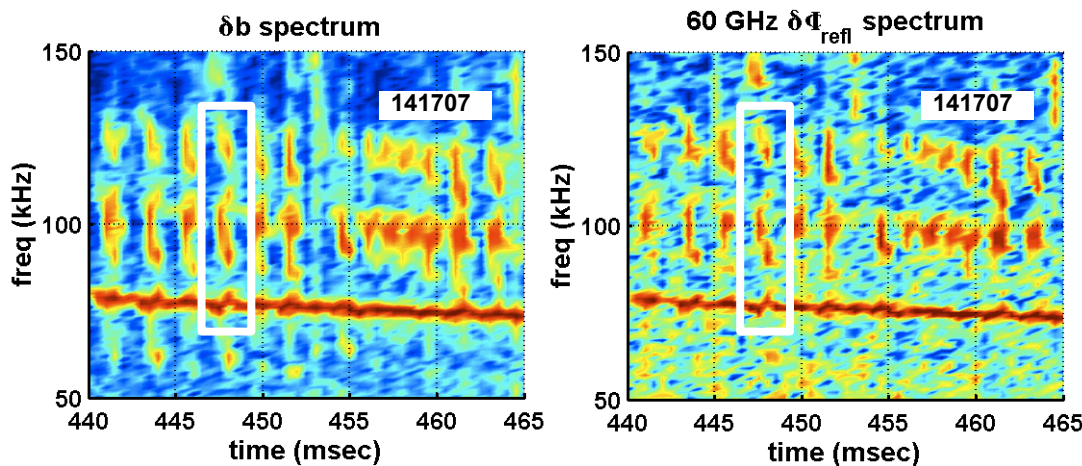
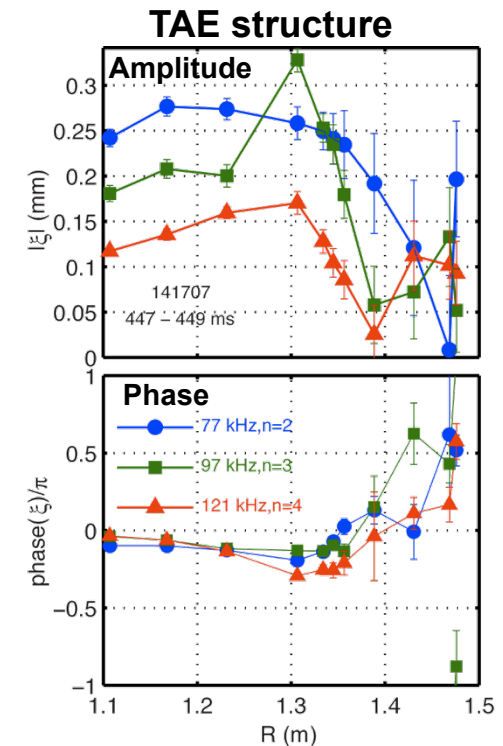


FIG. 3 (color online). Correlation between GAE activity, T_e flattening, and central χ_e increase in NSTX H modes heated by 2, 4, and 6 MW neutral beam, at $t \sim 0.44$ s. Within the uncertainties, the q , n_e , and $\omega_{E \times B}$ profiles are the same in all discharges at the time of the transport correlation [13].

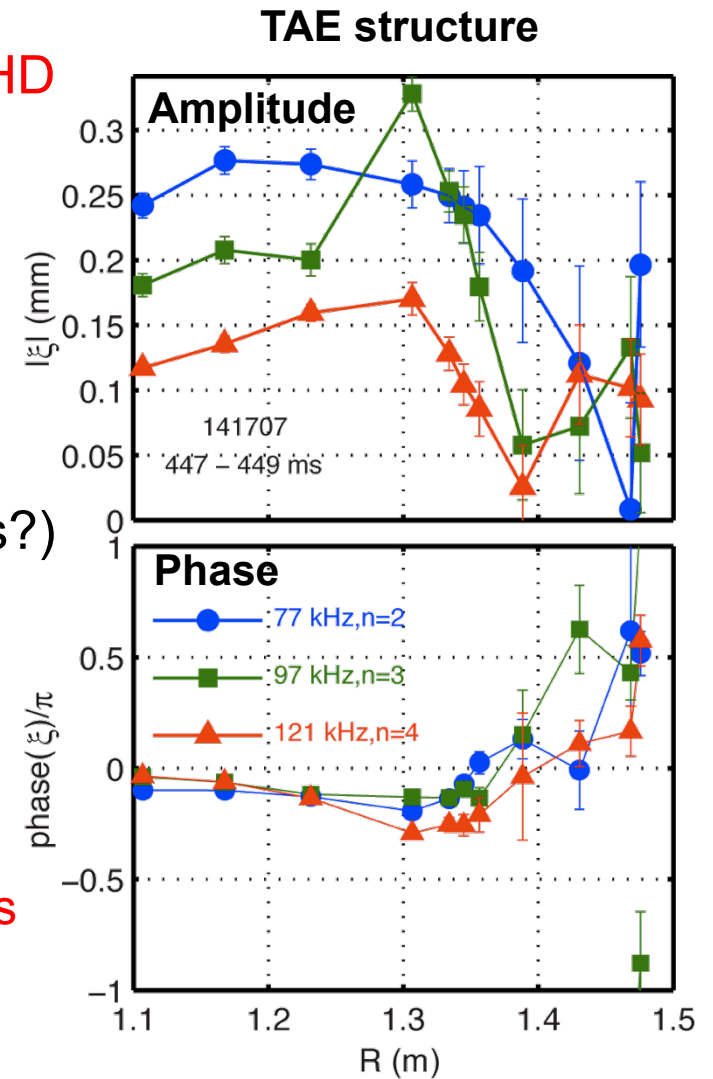
TAE structure measured with high spatial resolution

- TAE structure (amplitude and phase) measured across plasma at 12 locations near midplane
N A Crocker et al. PPCF 53 105001 (2011)
- Strong radial phase variation ($\Delta\Phi \sim \pi/2$) observed
 - Indicates radial propagation
- Measurement advances campaign to validate M3D-K code
 - High resolution strengthens comparison



Strong radial phase variation of TAEs indicates non-ideal MHD physics*

- **Radial phase variation inconsistent with ideal MHD**
 - Measurements approx. in midplane
 - Plasma up-down symmetric \Rightarrow ideal MHD predicts real-valued eigenmode in midplane (phase = 0 or 180°)
 - Caveat: effect of plasma rotation not clear from theory
- Coupling to fast-ions (and other non-ideal effects?) may explain phase variation
 - Radial phase variation observed for RSAEs in DIII-D [BJ Tobias *et al.* PRL **106** 075003 (2011)]
 - Comparison with TAE/FL code implicates coupling with fast-ions
 - M3D-K simulation NSTX included fast-ion & shows similar phase variation

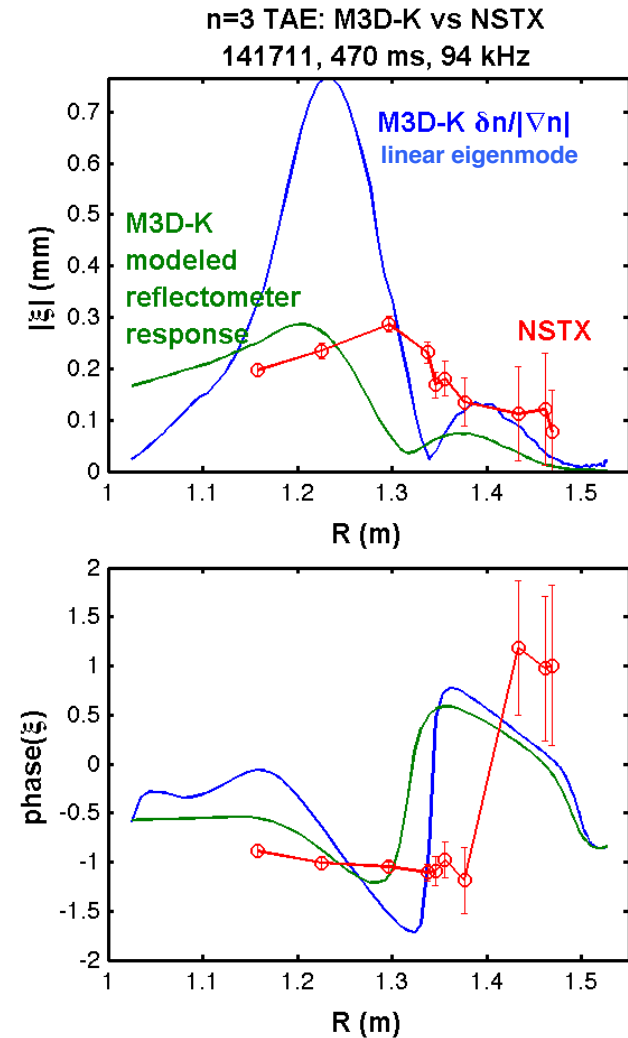


Observed TAE structure similar to M3D-K linear eigenmode prediction

- M3D-K solves for TAE eigenmode in NSTX plasma
 - initial value code
 - MHD plasma coupled to kinetically treated energetic ions
- Reflectometer response (ξ) modeled for M3D-K δn (i.e. “synthetic diagnostic”)
 - WKB approximation for path length (L) used:

$$L = L_0 + \xi = \int_{edge}^{\omega_p^2(R)=\omega^2} \sqrt{1 - \omega_p^2(R)/\omega^2} dR$$

- M3D-K structure [$|\xi|$ & phase(ξ)] similar to NSTX, but shifted radially inward ~ 7 cm
 - further work needed to understand shift

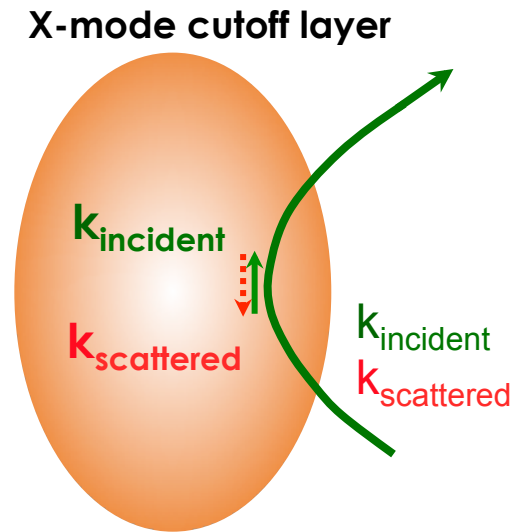


Conclusions

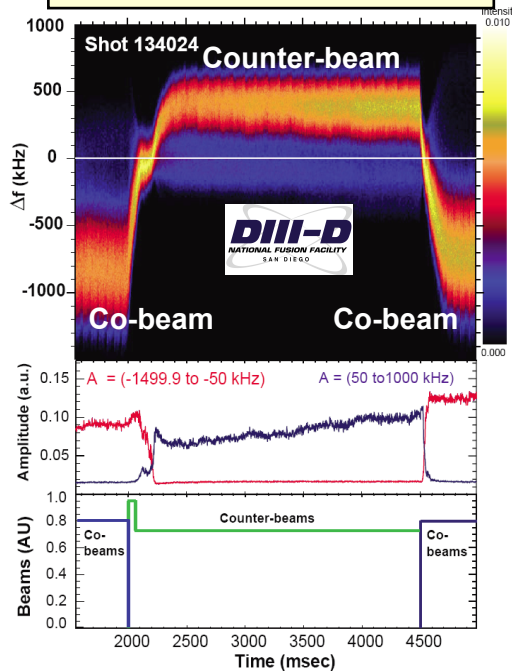
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Doppler backscattering (DBS): Determines equilibrium E_r & oscillating flows associated with GAMs & Alfvén modes

- DBS locally measures scattering from intermediate scale turbulence
 - Doppler shift provides information on local turbulent flow. Equilibrium & oscillating flows are imprinted onto Doppler shift
 - Allows determination of radial profile of oscillating electric fields of GAMs, Alfvén waves, etc.



Quadrature data for co- and counter beams



DBS phase analysis of GAMs & Alfvén modes

