

# Studying High-k Turbulence With Microwave Scattering on NSTX

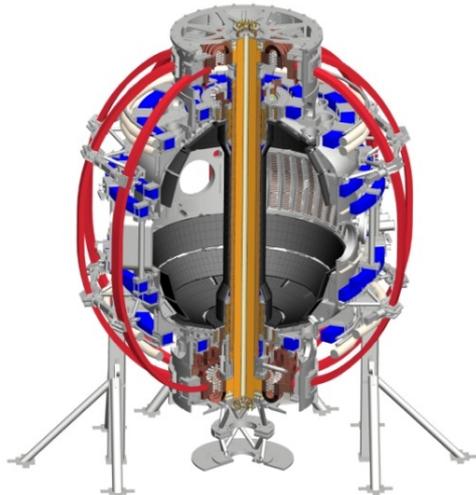
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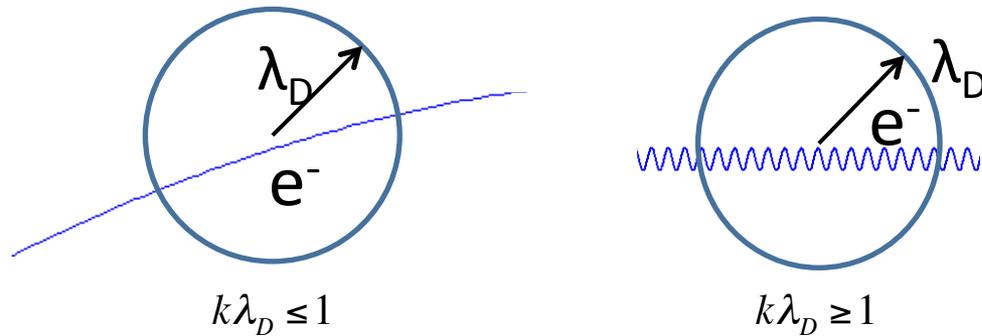
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# Anomalous electron thermal transport is observed in all NSTX confinement regimes

- Theory and experiments suggest that toroidal **ETG** turbulence is a candidate for anomalous electron thermal transport.
- A ***microwave collective scattering diagnostic*** at NSTX measures electron density fluctuations indicative of high-k turbulence on the electron gyro-scale  $k_{\perp}\rho_e \leq 1$ .
- A correlation between **electron density gradient** and electron scale high-k fluctuations measured by the high-k scattering diagnostic is established on a set of ***current ramp-down***, NBI heated H-mode plasmas.
- Linear stability analysis using GS2 code confirms the stabilizing effect of electron density gradient.
- GS2 linear runs show that wavenumbers at maximum linear growth rates shift to even higher k values ( $k_{\perp}\rho_s > 40$ ) as electron density gradient is increased.
- Linear growth rates at high wavenumbers are less sensitive to electron density gradient than linear growth rates at lower wavenumbers.

# Collective Thomson Scattering Theory is used to measure ETG-scale turbulence

- Collective/coherent and incoherent scattering



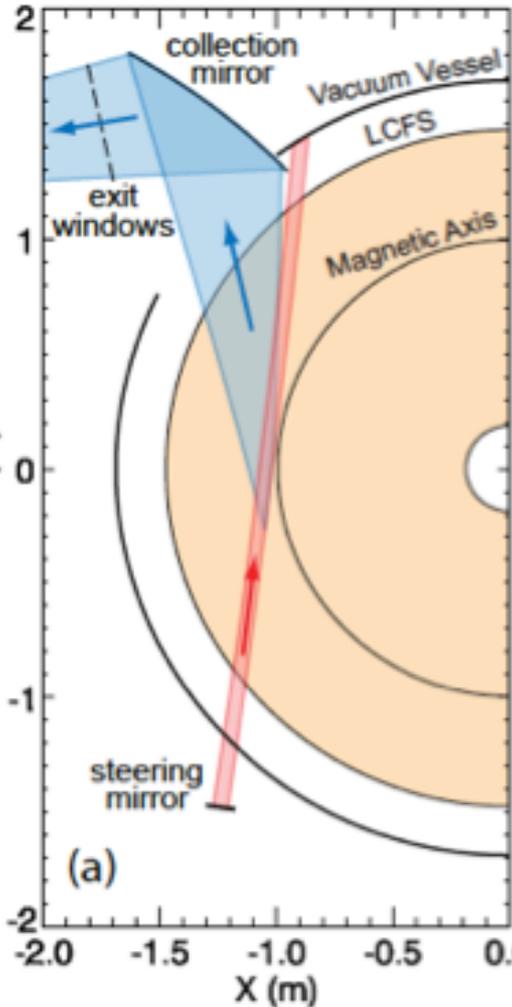
- Typical values at NSTX  $\lambda_D \sim 10^{-5} \text{ m}$ ,  $k \sim k_{\perp} < 10^4 \text{ m}^{-1}$   $\Rightarrow k\lambda_D < 1$  (**collective scattering**)

- ETG/High-k** fluctuations:  $k_{\perp} \rho_e \leq 1$

- Scattered power density  $\frac{d^2 P}{d\Omega d\nu} = P_i r_e^2 L_z |\Pi \cdot \hat{e}|^2 \frac{|\tilde{n}_e(k, \omega)|^2}{VT}$

$r_e$  classical electron radius  
 $V, L_z$  volume and length of scattering volume  
 $\Pi$  polarization tensor  
 $\hat{e}$  direction of incident electric field  
 $T$  observation time

# High-k Microwave Scattering Diagnostic at NSTX



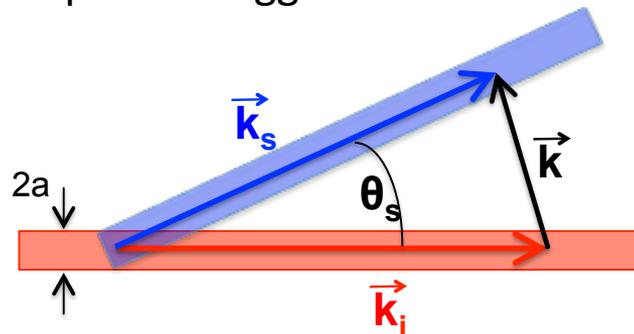
View from top of NSTX  
(D.R. Smith PhD thesis 2009)

- Gaussian Probe beam: 200 mW, 280 GHz,  $\lambda \sim 1.07$  mm,  $a = 3$  cm ( $1/e^2$  radius).
- Propagation close to midplane  $\Rightarrow k_r$  spectrum.
- 5 detection channels  $\Rightarrow$  Wavenumber range  $k_r \sim 5$ -30  $\text{cm}^{-1}$ .
- Wavenumber resolution  $\Delta k = \pm 0.7$   $\text{cm}^{-1}$ .
- Radial coverage:  $R = 106$ -144 cm.
- Radial resolution:  $\Delta R = \pm 2$  cm.
- **Three wave-coupling** between incident probe beam ( $k_i, \omega_i$ ) and plasma ( $k, \omega$ ).

$$\vec{k}_s = \vec{k} + \vec{k}_i \quad \omega_s = \omega + \omega_i$$

- $\omega_i, \omega_s \gg \omega$  imposes Bragg condition

$$k = 2k_i \sin(\theta_s/2)$$



# Spatial Localization and Wavenumber Resolution

- Volume overlap of incident and scattered beams leads to poor spatial localization.
- Theory [cf. Horton *Rev. Mod. Phys.* 1999] predicts  $k_{||} \sim 1/qR \ll k_{\perp} \Rightarrow \vec{k} \cdot \vec{B} \approx 0$
- Plasma fluctuations must satisfy:
 

$k \cdot B \approx 0$	(1) Perpendicular fluctuations.
$k = 2k_i \sin(\theta_s / 2)$	(2) Bragg Condition
- When incident beam forms a small angle with  $\vec{B}$ , (1) and (2) become highly dependent on **toroidal curvature** of magnetic field (cf. scattered beams at  $P_1$  and  $P_2$  in the figure). **Oblique propagation** (outside the midplane) of incident beam exploits this phenomenon and enhances *longitudinal localization* of fluctuations [cf. Mazucatto *Phys. Plasmas* 2003].

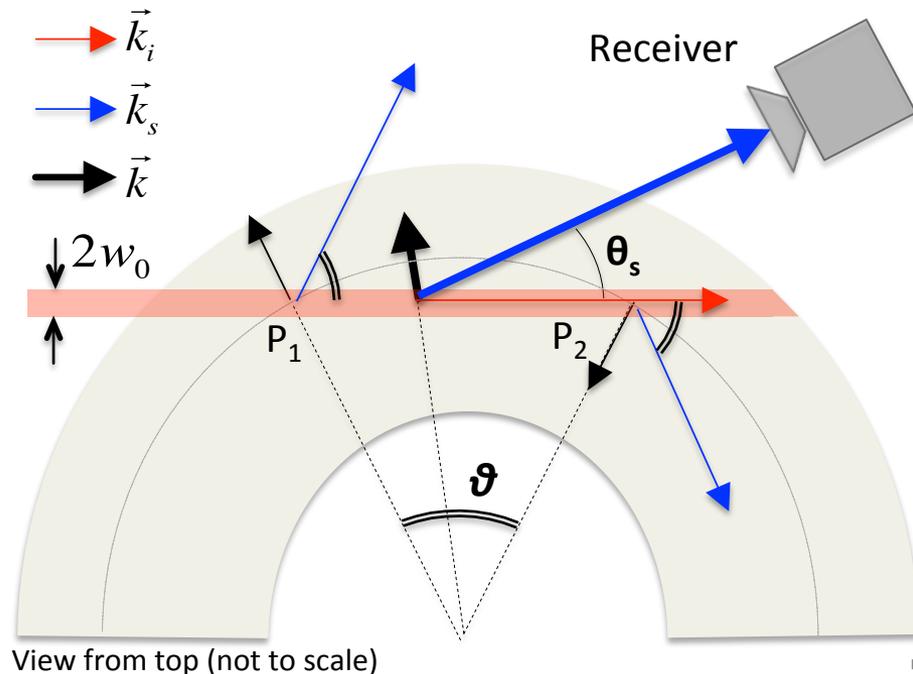
- For midplane propagation, (1) and (2) are only satisfied at  $P_1$  and  $P_2$  and fluctuation wavenumber is purely in the **radial direction**.
- In practice, beam propagation is out of midplane, but oblique angle is small ( $\sim 5^\circ$ ).  $k$  is *mostly* radial.

- Gaussian beam width dictates  $k$  and R-resolution

$$A(r_{\perp}) = \exp(-r_{\perp}^2 / w_0^2)$$

$$G(k_{\perp}) = \exp(-k_{\perp}^2 / \Delta k^2)$$

$$\Delta k = 2 / w_0$$

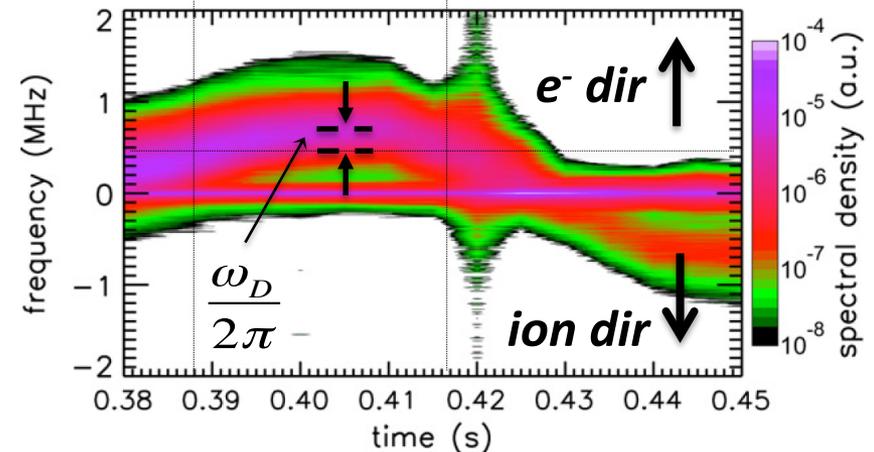
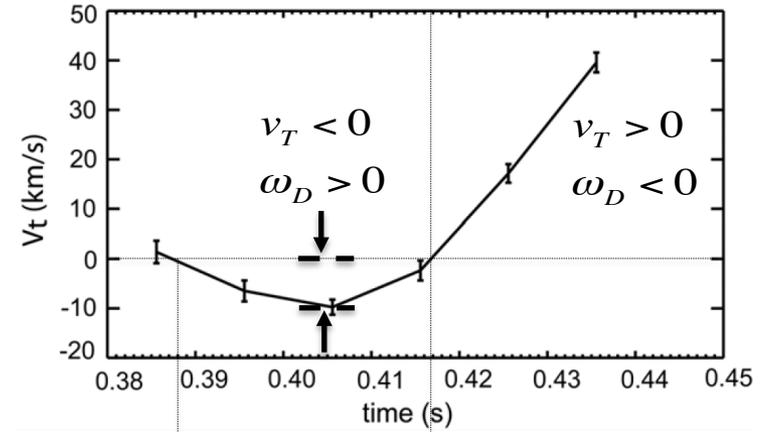
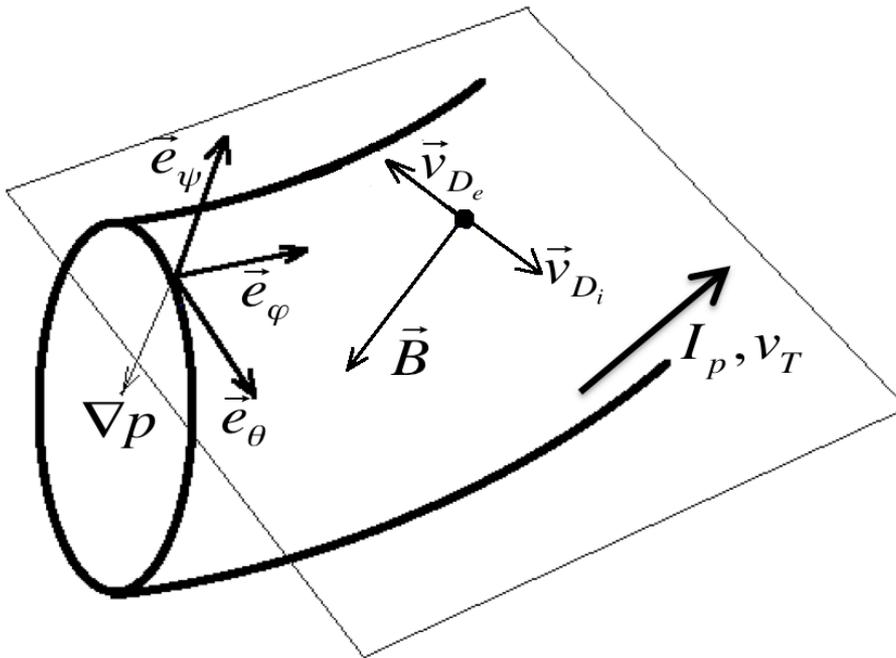


# Toroidal Rotation has an Effect on Measured Fluctuations: Doppler Shift

Outboard measurements ( $k_r < 0$ )

$$\text{if } \vec{v}_{D_e} \cdot \vec{e}_\varphi < 0 \quad (k_T < 0)$$

$$\Rightarrow \omega_D \approx k_T v_T < 0 \quad (\text{for } v_T > 0)$$

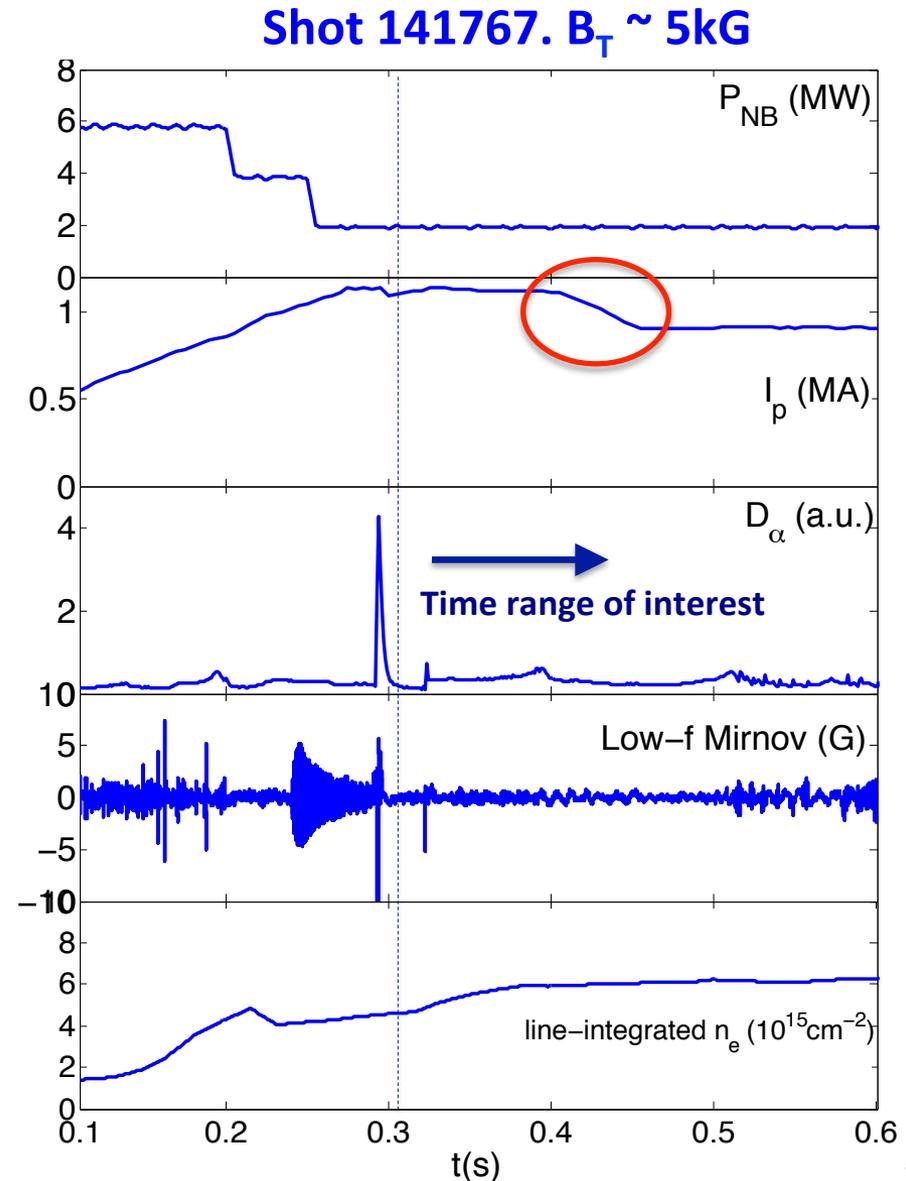


Adapted from Mazzucato Nuc. Fusion 2009.  
Here  $f > 0 \Rightarrow e^-$  direction.

Doppler shift is mainly affected by *toroidal* rotation component ( $\omega_D \sim k_T v_T$ ). **Doppler shift** can make fluctuations *appear* to propagate in the **ion diamagnetic drift direction** (fig.  $t > 0.42$  s).

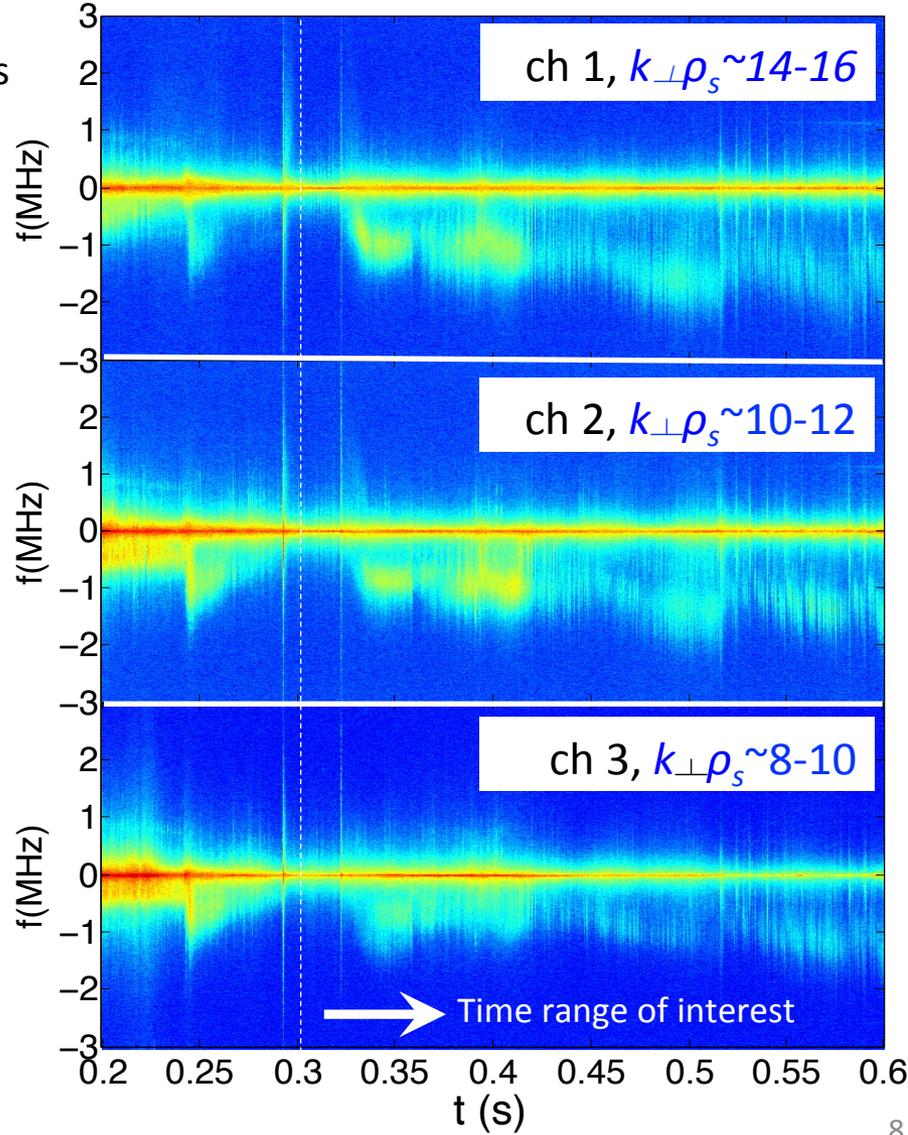
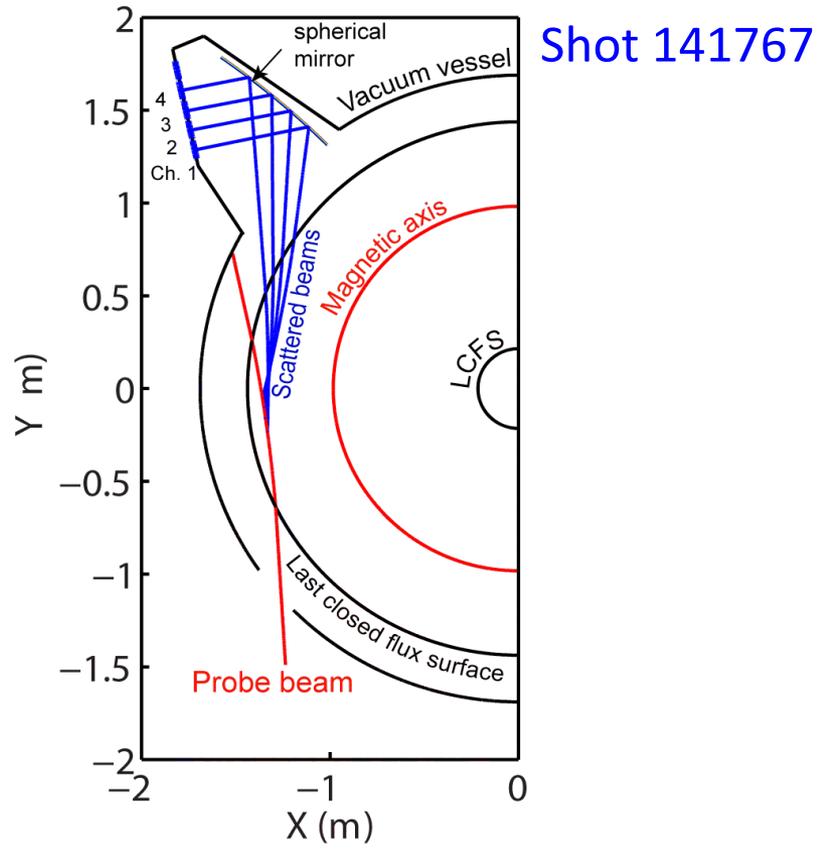
# A Set of NBI-heated H-mode Plasmas is Used to Study High-k Turbulence during Current Ramp-down

- **NBI heated**, HHFW heating is absent during the run.
- **Current ramp down** between  $t = 400$  ms and  $t = 450$  ms (from LRDFIT).
- Time range of interest is  $t > \sim 300$  ms, covering current ramp-down phase, and after ELM event at  $t \sim 290$  ms.
- **MHD activity is quiet during that time.** Before  $t \sim 290$  ms MHD activity is high (*cf.* low-f Mirnov signal).
- Line integrated density is fairly constant during the time range of interest.

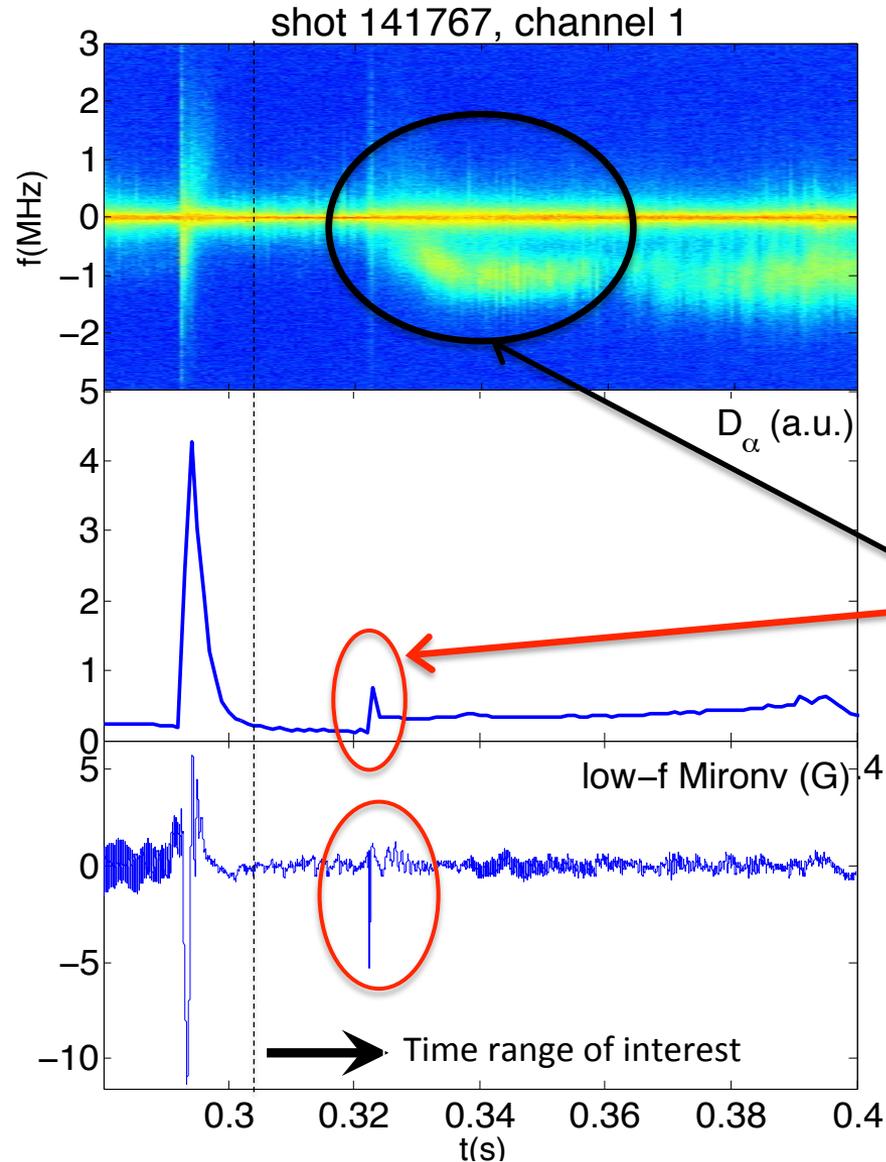


# Each Channel of the NSTX High-k Scattering System Detects a Fluctuation Wavenumber $k$

- Channel 1 detects highest  $k_{\perp}$  and  $k_T$ , Doppler shift is greatest.
- High peak at  $f \sim 0$  corresponds to stray radiation.



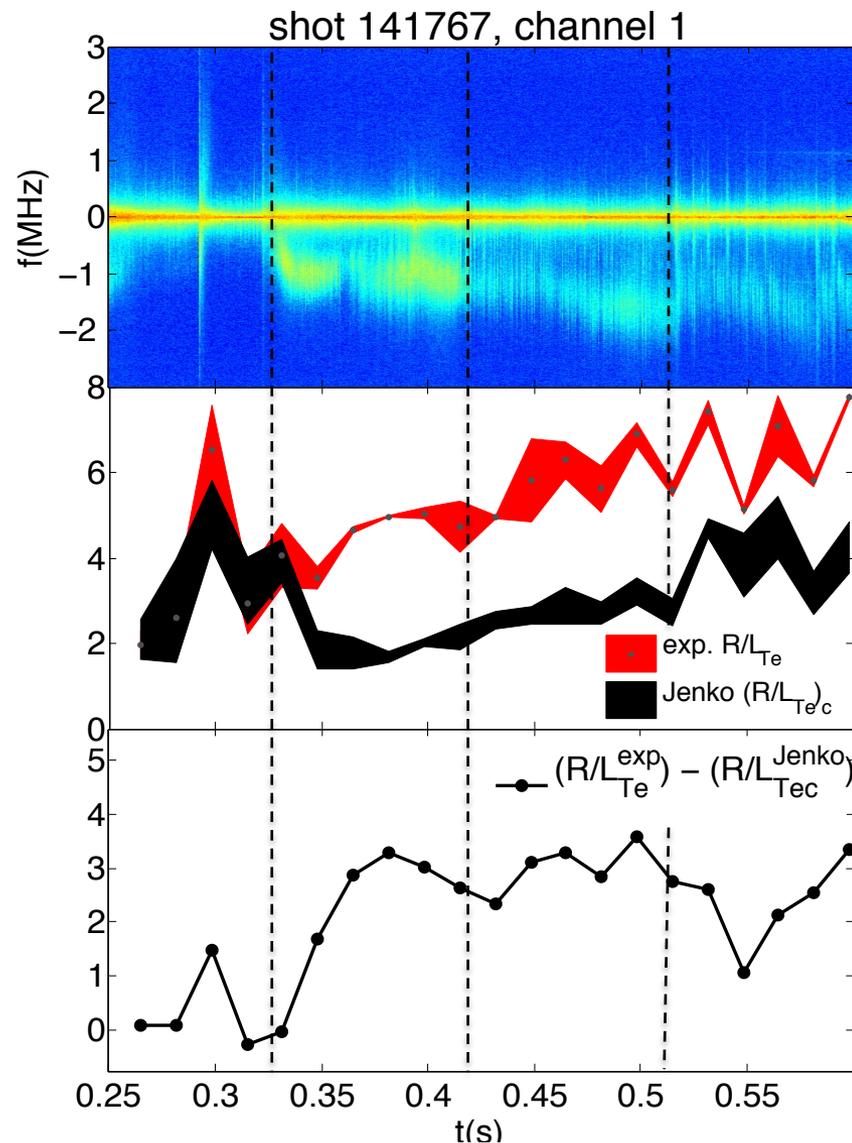
# High-k Fluctuations Start after Small Spike in $D_\alpha$ and Mirnov Signal



- Before  $t \sim 290$  ms, MHD activity is high. At  $\sim 290$  ms, an ELM event takes place and MHD activity quiets.
- Between  $t \sim 290$  ms and  $t \sim 320$  ms, high-k fluctuations are absent and MHD activity is quiet.
- **High-k fluctuations** start at  $t \sim 320$  ms, after small ELM event, detected in  **$D_\alpha$  and Mirnov signal**.

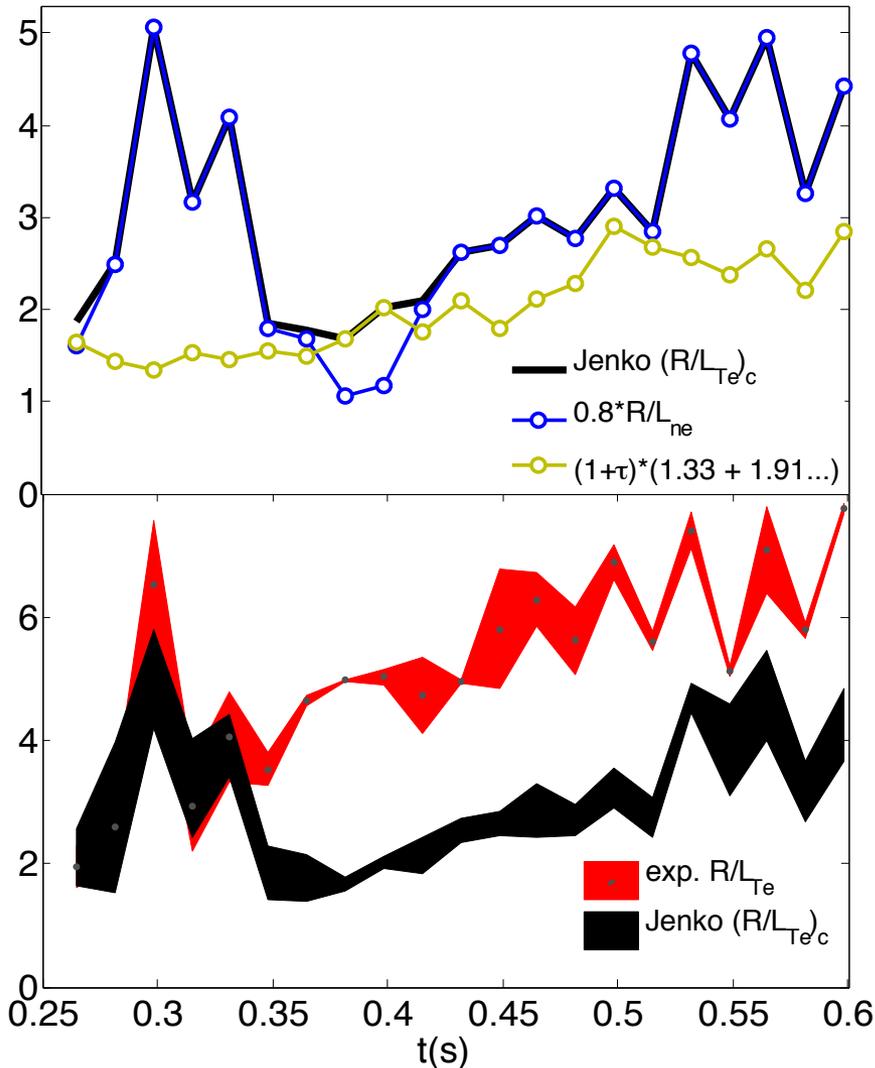
# Observed Fluctuations Correlate to Difference Between Critical and Experimental ETG

- **Normalized electron temperature gradient** ( $R/L_{Te}$ ) is calculated at the scattering location ( $R \sim 135$  cm), and plot along with the **Jenko critical gradient** [cf. *Jenko Phys. Plasmas 2001*].
- Prior to  $t \sim 320$  ms,  $R/L_{Te} \sim (R/L_{Tec})^{Jenko}$  and ETG turbulence is marginally stable. No observed fluctuations. After  $t \sim 320$  ms,  $R/L_{Te} > (R/L_{Te})_{crit}$  and fluctuations develop.
- As  $(R/L_{Te})_{crit}$  increases and gets closer to  $R/L_{Te}$  detected fluctuations decrease ( $t \sim 410$  ms and  $t \sim 510$  ms).
- During period  $\sim 350$  ms  $< t < \sim 500$  ms, a similar difference in  $R/L_{te} - (R/L_{Te})_{crit}$  produces very different fluctuation levels. This will be later explained by the *density gradient stabilization of lower numbers*.
- The applicability of Jenko's critical ETG is discussed in slides 14 and 15.



# Theory Predicts Electron Density Gradient Can Drive the Difference $(R/L_{Te})_c - R/L_{Te}$ and Stabilize Turbulence

Shot 141767



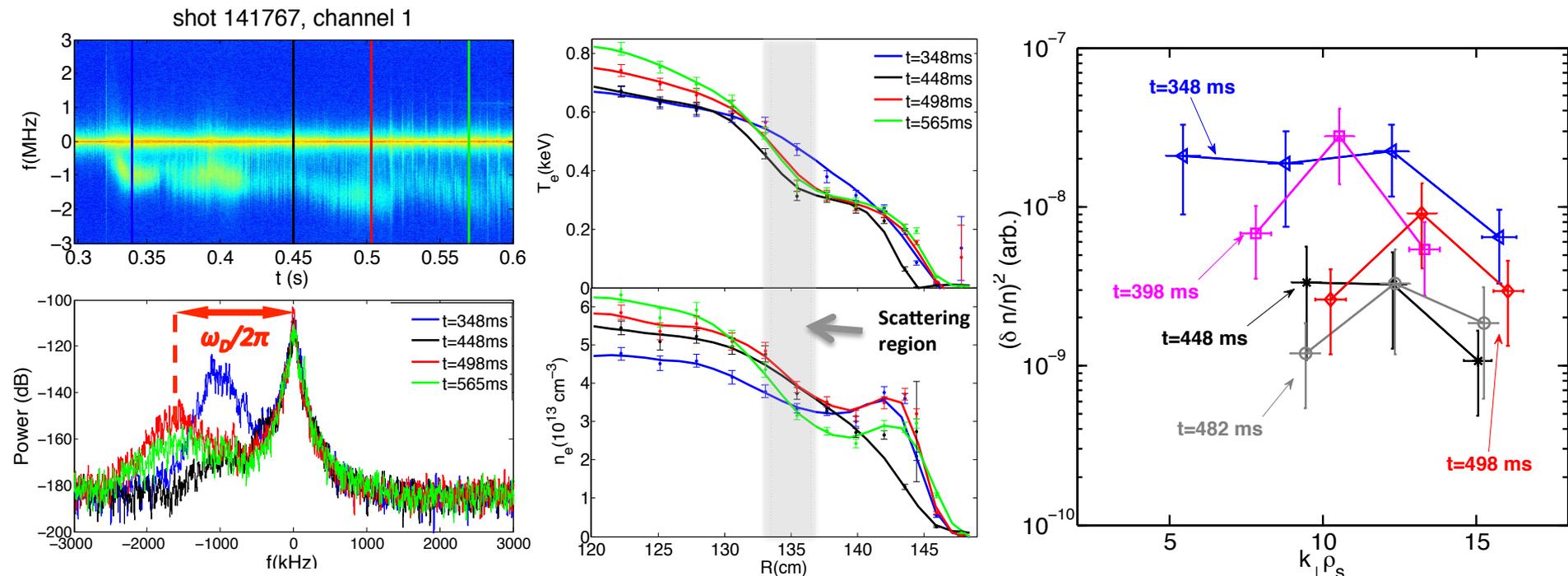
- Jenko critical gradient is a maximum of a  $R/L_{ne}$  term and an  $s/q$  term [cf. *Jenko Phys. Plasmas* 2001].

$$(R/L_{Te})_{crit} = \max \left\{ \begin{array}{l} 0.8R/L_{ne} \\ (1+\tau)(1.33+1.91\hat{s}/q)(1-1.5\epsilon) \end{array} \right.$$

with  $\tau = Z_{eff} T_e / T_i$

- High values of  $R/L_{ne}$  can make it the dominant term in Jenko's  $(R/L_{Te})_{crit}$  formula. According to this formula,  $R/L_{ne}$  can drive the difference between critical and experimental ETG.
- High enough values of  $R/L_{ne}$  could bring critical ETG to experimental ETG levels or even higher. This *should* have a **stabilizing** effect on turbulence.

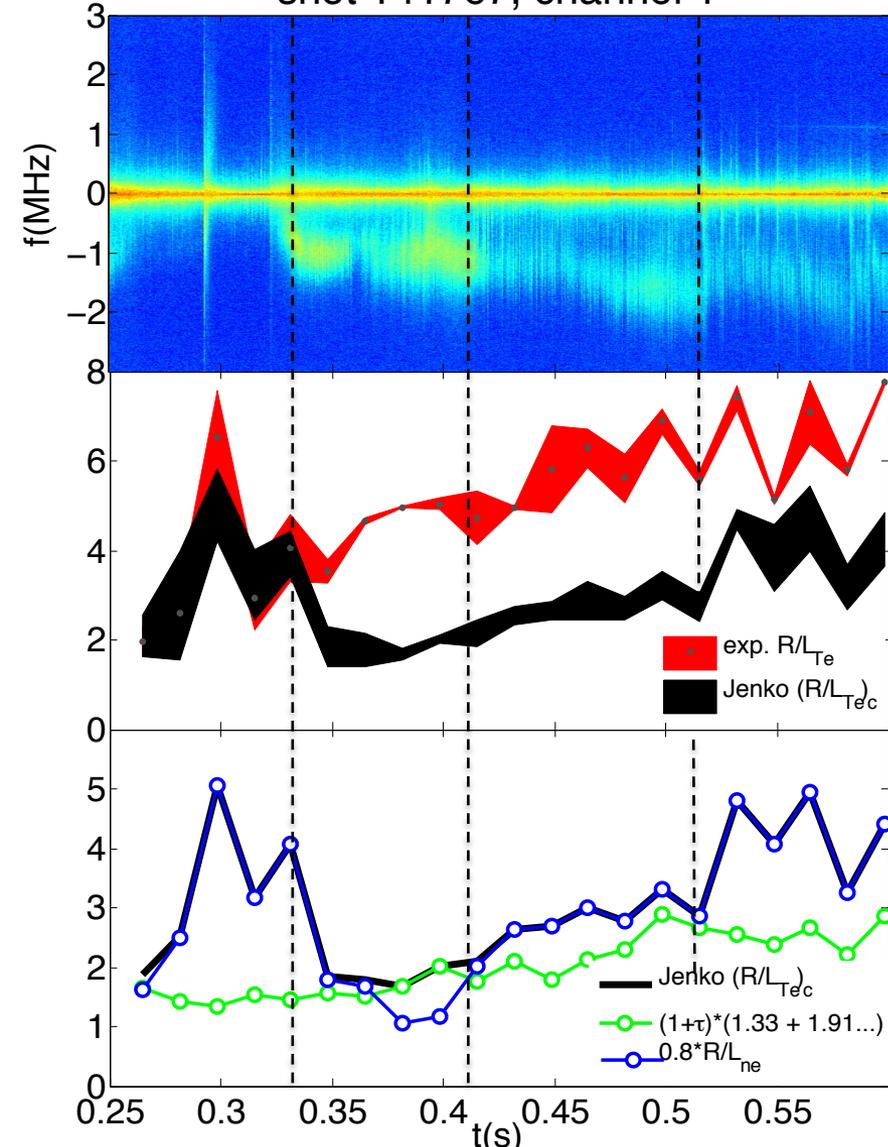
# Observed High-k Fluctuations Correlate to Local Electron Density Gradient



- Electron Density Gradient suffers the biggest change in the **scattering region**.
- Doppler shift is measured as distance from 0 to observed fluctuation frequency.
- A big change in the electron density fluctuation level (right figure) is observed between  $t = 348$  ms and  $t = 448$  ms. A notable change in  $R/L_{ne}$  takes place during that time.
- In general, fluctuation levels  $(\delta n_e/n_e)^2$  at low wavenumbers ( $k_{\perp} \rho_s < 10$ ) decrease in time.
- After  $t = 448$  ms,  $(\delta n_e/n_e)^2$  at high wavenumbers ( $k_{\perp} \rho_s \sim 15$ ) increases.

# Time Traces of Local Electron Density Gradient Confirm its Influence on Observed Fluctuations

shot 141767, channel 1



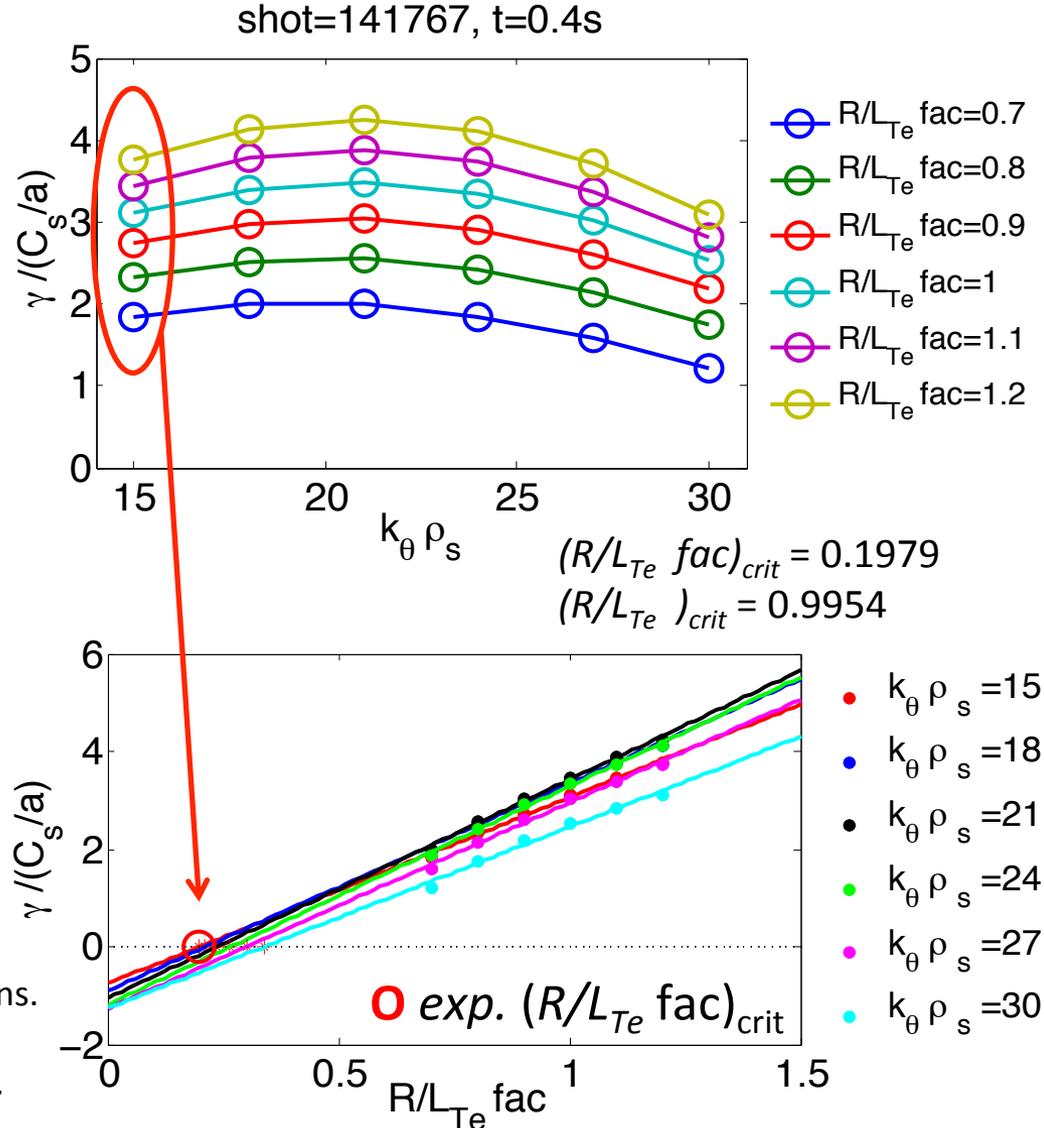
- As  $R/L_{ne}$  increases, it dominates in Jenko's formula  $(R/L_{Te})_{crit}$  ( $t < 340$  ms,  $t > 410$  ms &  $t > 515$  ms). Fluctuations decrease during that time.
- Previous to  $t \sim 320$  ms ETG is marginally stable with respect to Jenko critical gradient. No fluctuations are observed.
- $R/L_{ne}$  has a **stabilizing** effect when it dominates Jenko critical gradient.
- Electron density gradient stabilization of ETG turbulence was already observed by Y. Ren *et al*, Phys. Rev. Letters 2011.

# A Scan on $R/L_{Te}$ is Performed to Compute a Critical Gradient with GS2 Linear Runs

- Electron temperature gradient (ETG or  $R/L_{Te}$ ) is varied keeping all other quantities constant. The factor is called ( $R/L_{Te}$  fac)
- High-k linear growth rates saturate with increasing ( $R/L_{Te}$ ), for different wavenumbers.
- $(R/L_{Te})_{crit}$  is found to be the minimum  $R/L_{Te}$  to satisfy  $\gamma = 0$ .

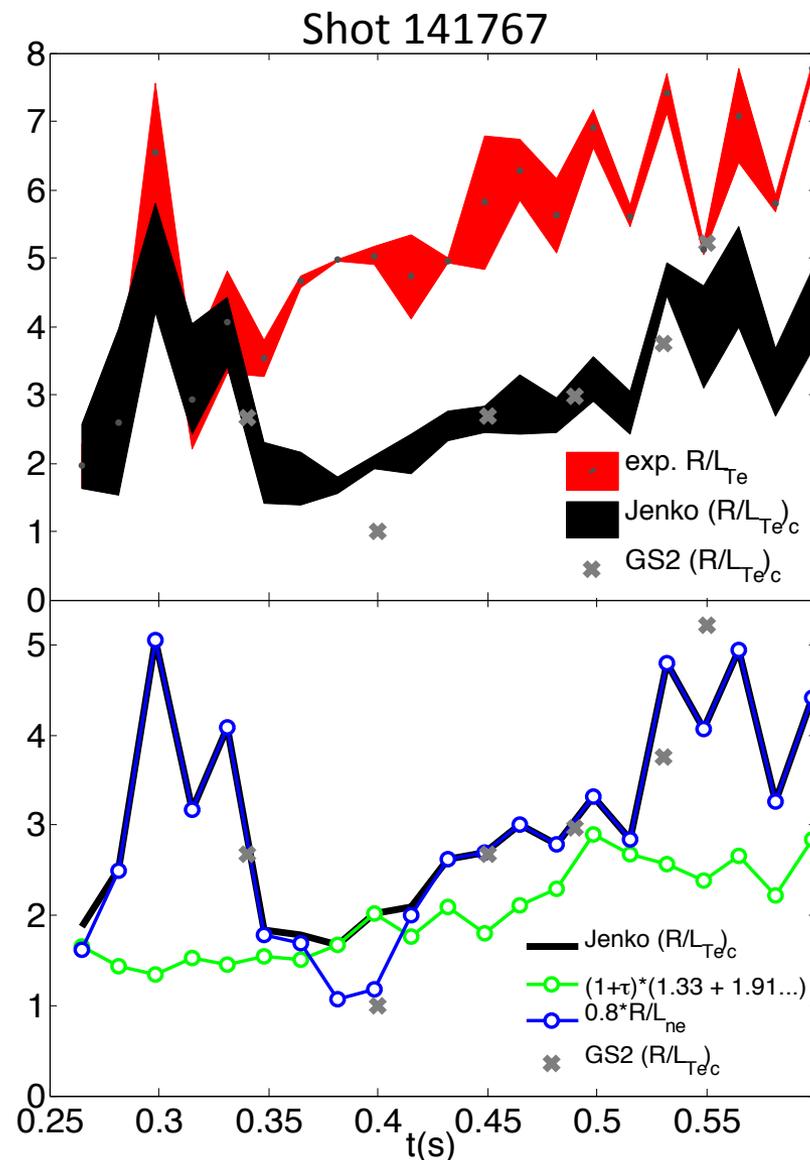
## Features of GS2

- GS2 is an initial value gyrokinetic code.
- Uses flux tube geometry.
- Includes electromagnetic effects in these linear runs.
- Tracks fastest growing modes.
- Cf. Kotschenreuther et al, Comp. Phys. Comm 1995.



# Critical Gradient Computed with GS2 Linear Runs Agrees with Jenko's Critical ETG

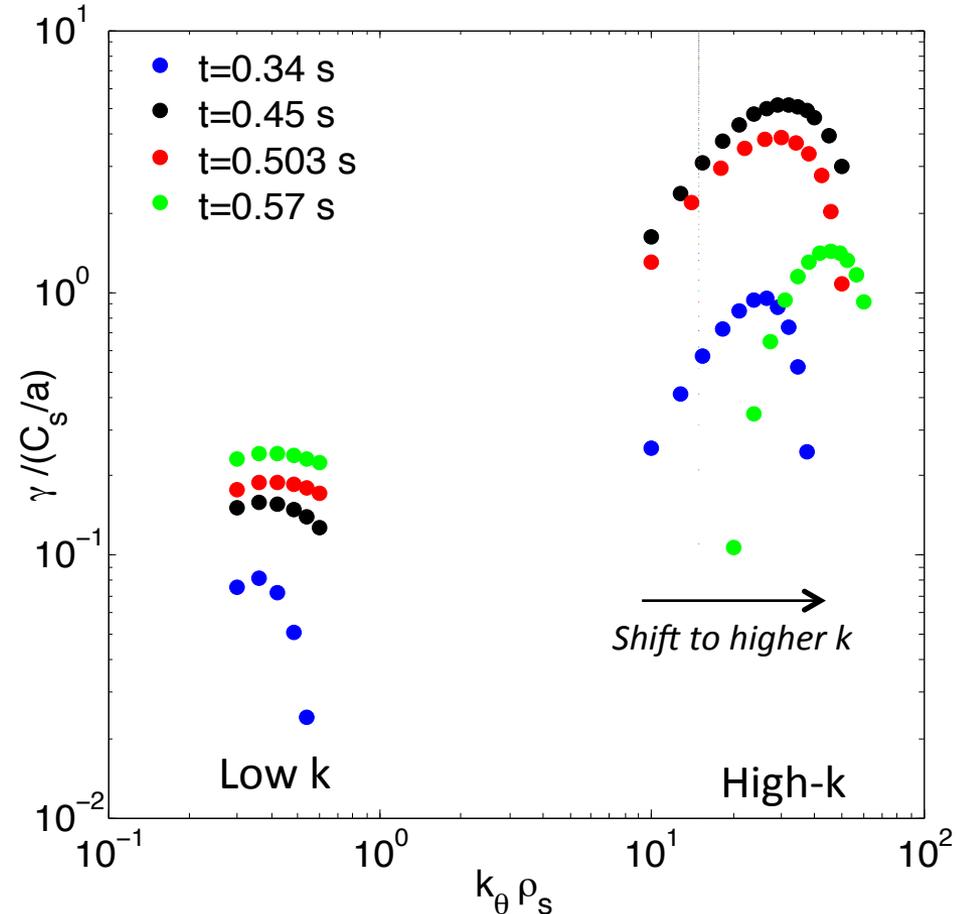
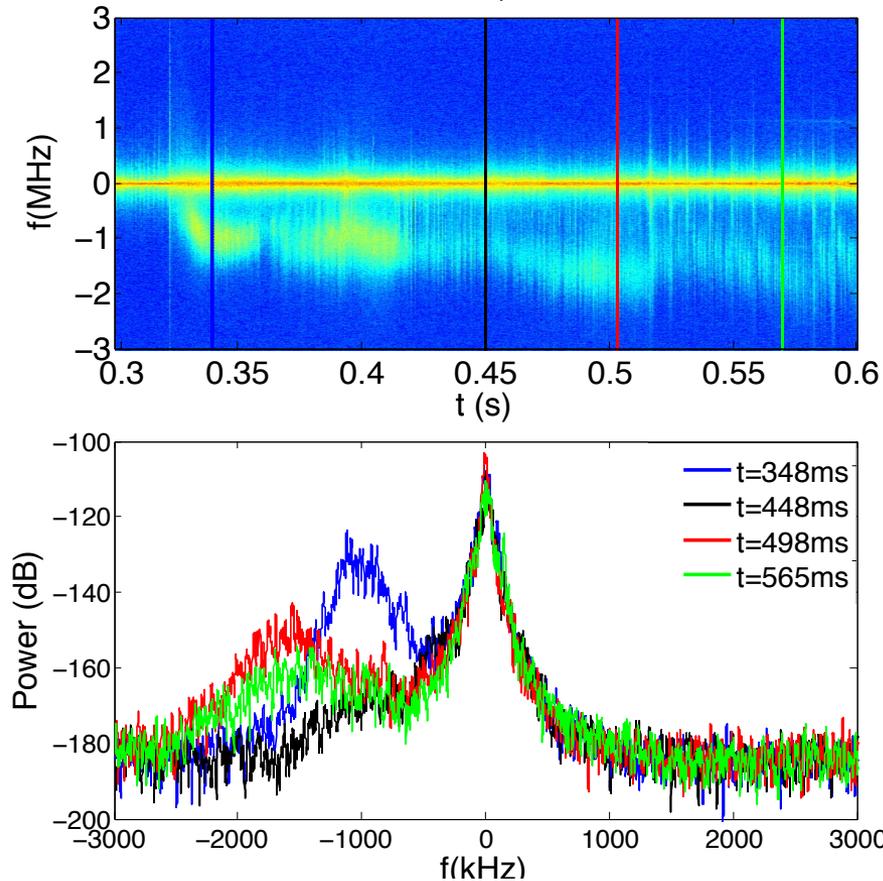
- A critical gradient computed with GS2 linear runs is plotted along with Jenko's critical gradient formula.
- Fairly good agreement is observed between GS2 calculations and Jenko's formula.
- This **validates** Jenko's critical ETG formula and previous comparisons with experimental ETG.



# GS2 Linear Simulations Show the Wavenumbers at Maximum Growth Rate Shift to Higher $k$ in Time

- Low- $k$  linear growth rates ( $k_{\perp} \rho_s \leq 1$ ) appear to increase with time, and are lower than high- $k$ .
- Wavenumbers corresponding to maximum linear growth rate shift towards higher- $k$ .
- As wavenumbers corresponding to maximum growth rate move to higher wavenumbers in time (e.g.  $t = 570$  ms), observed fluctuations decrease (*cf.* bottom left figure).

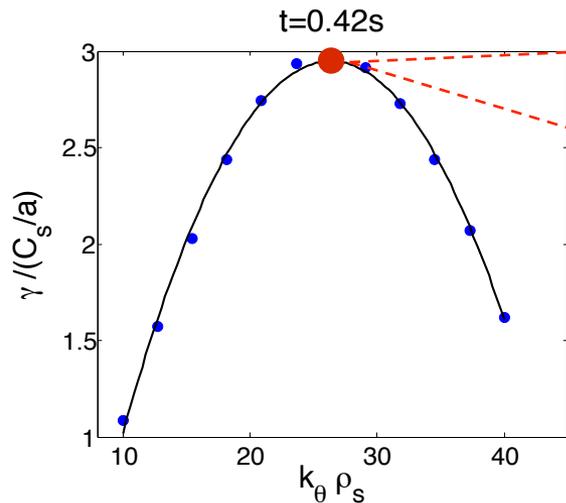
shot 141767, channel 1



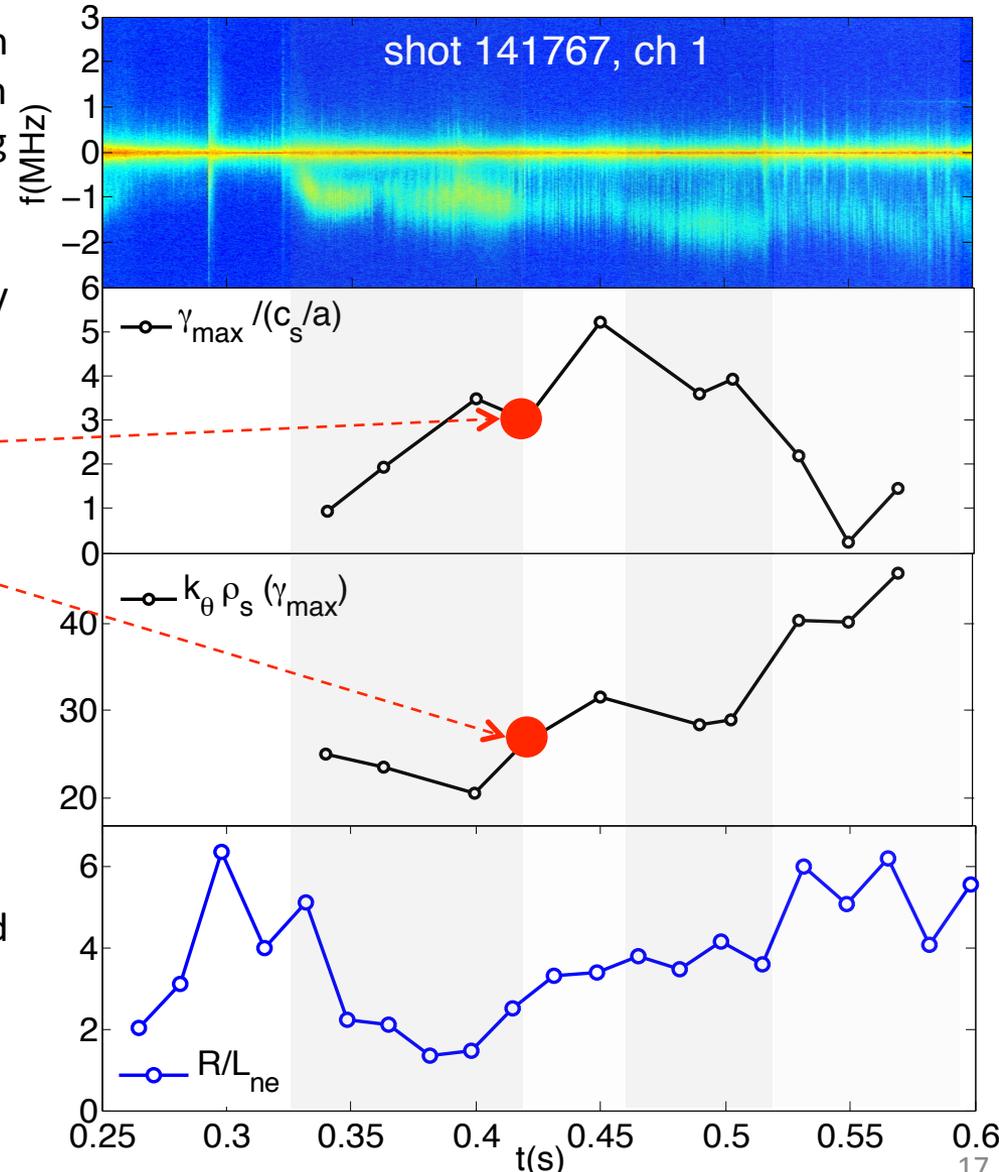
Linear growth rate  $\gamma / (c_s/a)$  is a weak function of  $T_e$  through normalization.  $T_e$  does not change significantly during time range of interest (*cf.* slide 12).

# Wavenumber at Maximum Linear Growth Rate Correlates to Electron Density Gradient and Observed Fluctuations

- Linear growth rates are plotted against  $k$  (at each time, *cf.* left figure) to determine the maximum growth rate  $\gamma_{\max}/(c_s/a)$  and the corresponding wavenumber  $k_{\theta}\rho_s(\gamma_{\max})$  (red dot).
- Note the evolution of  $k_{\perp}\rho_s(\gamma_{\max})$  in time is very similar to  $R/L_{ne}$  at the scattering location.

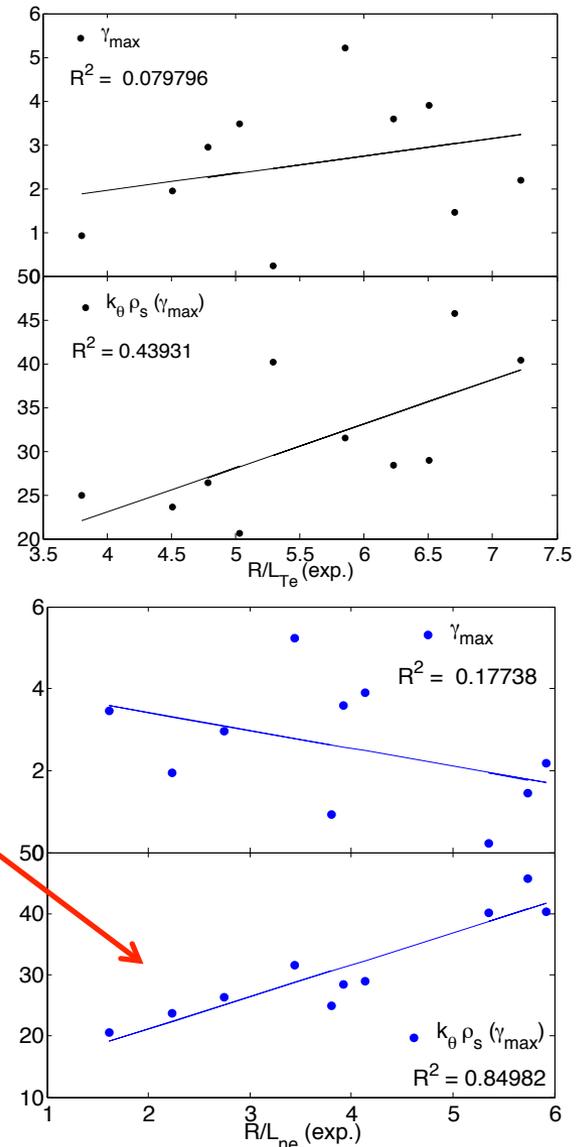


- $R/L_{ne}$  and  $k_{\perp}\rho_s(\gamma_{\max})$  correlate well with observed fluctuations (*cf.* evolution within time panels).
- No correlation between maximum growth rate  $\gamma_{\max}/(c_s/a)$  and fluctuations is observed.



# Correlation Between Wavenumber Values at Maximum Growth Rates and Electron Density Gradient

- **Maximum linear growth rates** ( $\gamma_{max}$ ) and the corresponding wavenumber values are plotted against local experimental  $R/L_{Te}$  (black) and  $R/L_{ne}$  (blue) along with linear fits.
- Low correlation is observed between  $\gamma_{max}$  and experimental  $R/L_{Te}$  and  $R/L_{ne}$ .
- **Wavenumber values at maximum linear growth rates** ( $k_{\theta} \rho_s(\gamma_{max})$  cf. slide 17) correlate better to  $R/L_{Te}$  and  $R/L_{ne}$  than linear growth rates.
- Best correlation is observed between  $k_{\theta} \rho_s(\gamma_{max})$  and  $R/L_{ne}$ . This could also be observed in the plots of slide 17.
- The conjecture is that  $R/L_{ne}$  is driving high-k turbulence to higher wavenumbers.
- Correlation is quantified by  $R^2$  coefficient of linear fit.



# Correlation Between Wavenumber Values at Maximum Growth Rate and Wavenumber Spectrum of Fluctuations

Fig. a

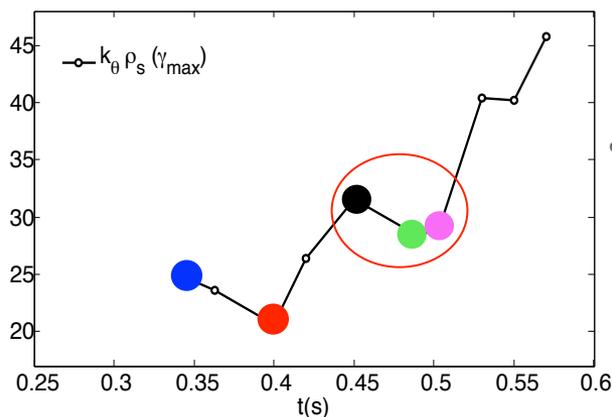


Fig. b

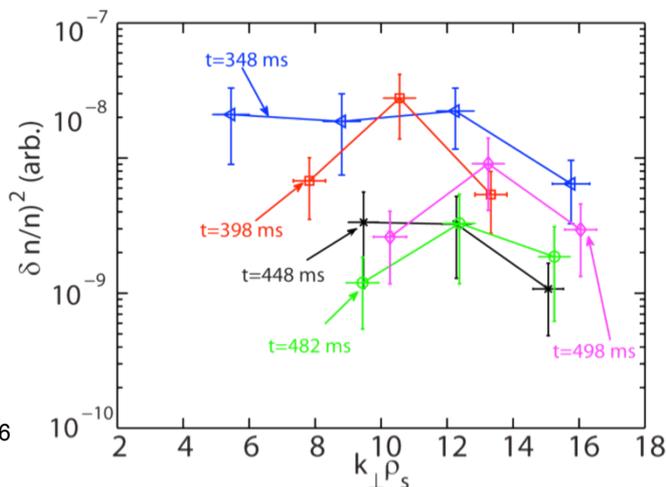
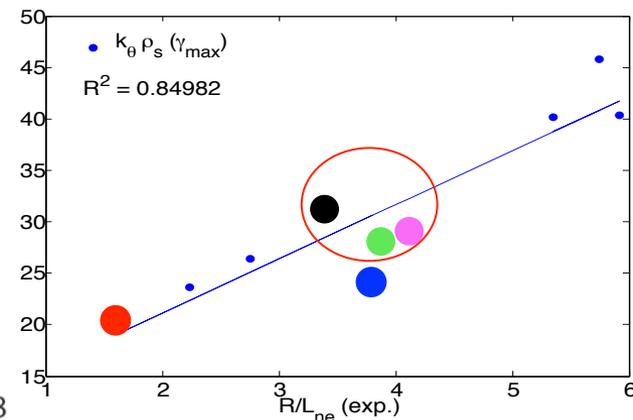


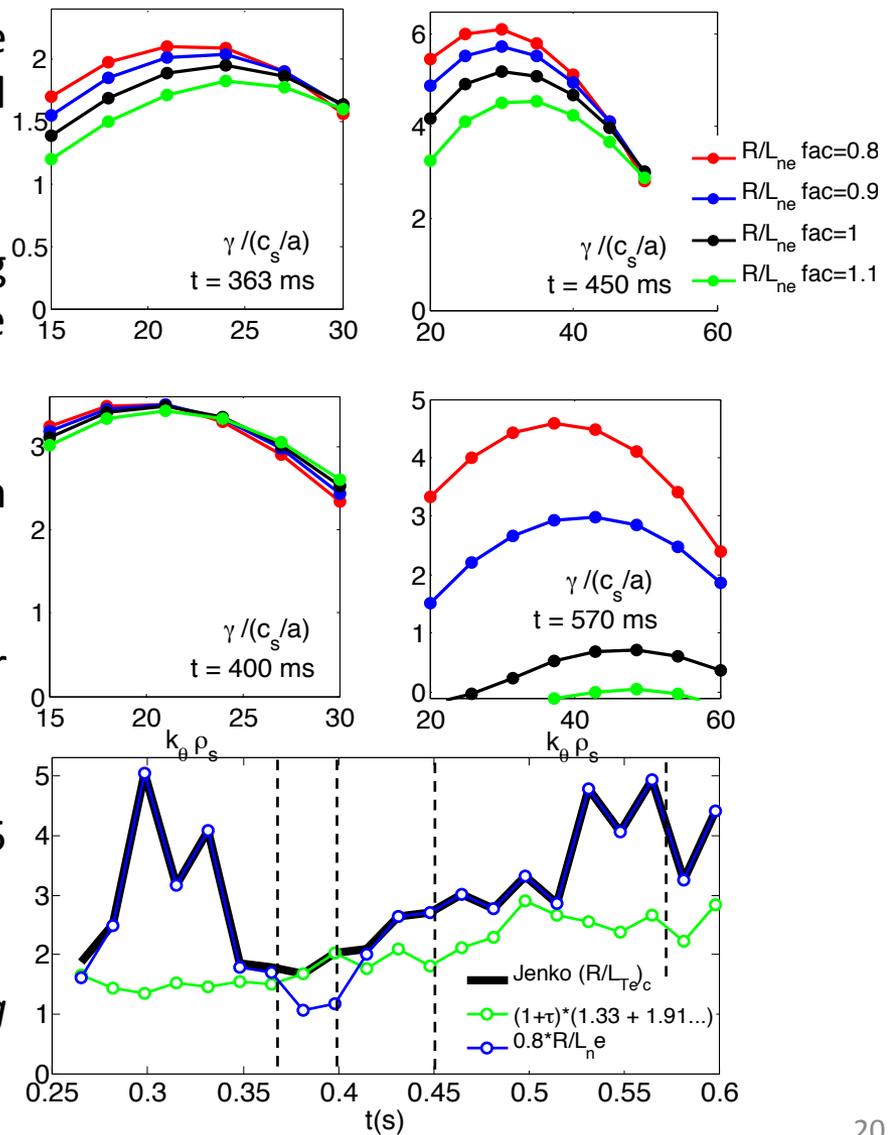
Fig. c



- A general increase of  $k_{\perp} \rho_s(\gamma_{max})$  in time (linear GS2 simulations, Fig. a) is in agreement with a global decrease of lower- $k$  ( $k_{\perp} \rho_s < 10$ ) fluctuation level ( $(\delta n_e/n_e)^2$  Fig. b).
- Big increase of  $k_{\perp} \rho_s(\gamma_{max})$  for  $398 \text{ ms} < t < 448 \text{ ms}$  agrees with decrease in observed fluctuation level for all  $k$ .
- After  $t \sim 448 \text{ ms}$ , higher  $k$  ( $k_{\perp} \rho_s \sim 12-15$ ) fluctuation levels increase. During that time,  $R/L_{ne}$  increases (*cf.* red circle in Fig. c corresponding to same time period).
- Nonlinear gyrokinetic simulations have previously shown the effect of electron density gradient stabilization of low- $k$  turbulence and the consequences on electron thermal transport [*cf.* Y. Ren *Phys. Plasmas* 2012].

# A Scan in Electron Density Gradient is Performed to Confirm its Effect on High-k Turbulence

- Linear growth rates with GS2 are most sensitive when  $0.8 \cdot R/L_{ne}$  term dominates Jenko's critical ETG (t = 450 ms, t = 570 ms).
- At t = 400 ms,  $0.8 \cdot R/L_{ne}$  term is not dominating Jenko's critical ETG, and linear growth rates are not sensitive to  $R/L_{ne}$ .
- Lower-k values are more sensitive to  $R/L_{ne}$  than higher-k values.
- Except at t = 400 ms,  $R/L_{ne}$  decreases linear growth rates => **stabilizing effect**.
- At t = 400 ms, linear growth rates at  $k_{\theta} \rho_s > 25$  increase with  $R/L_{ne}$ .
- $R/L_{ne}$  could be a responsible factor for driving turbulence to higher k values (cf. slide 18).



# Summary and Future Work

## Summary

- High-k **electron scale density fluctuations** are detected with the coherent microwave scattering diagnostic at NSTX.
- Difference between local ETG values and Jenko's critical ETG correlates to observed fluctuations.
- As local **electron density gradient** ( $R/L_{ne}$ ) increases, it dominates Jenko's critical ETG and is observed to have a stabilizing influence on observed fluctuations.
- Increasing  $R/L_{ne}$  produces a shift of high-k fluctuations to even higher k values.
- A scan on local  $R/L_{ne}$  with GS2 linear runs shows linear growth rates can be very sensitive to local  $R/L_{ne}$  when it is the dominant term in Jenko's **critical ETG**. In the opposite case, linear growth rates appear to be practically insensitive to local  $R/L_{ne}$ .
- Lower k-values in the high-k range appear more sensitive to  $R/L_{ne}$  than higher k-values.

## Future Work

- Carry out further studies for other NSTX shots with similar characteristics and compare the influence of local electron density gradient.
- Perform transport analysis to study influence of local electron density gradient in electron thermal transport.
- Carry out nonlinear gyrokinetic simulations to evaluate the effects of electron density gradient on turbulence and electron thermal transport.