

# Impurity transport studies in NSTX beam heated H-mode plasmas

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### Abstract

# Motivation

- i) Perform the first particle transport studies in high (Power &  $\beta$ ) performance NSTX H-modes, important for NSTX and the next-step ST.
  - a)  $\mathbf{D}_{\mathbf{Z}}$  (diffusion) &  $\mathbf{v}_{\mathbf{Z}}$  (convective velocity)
  - b) Identify regimes with possible impurity "screening" ( $v_z > 0$ ).
  - c)  $\rho^*$  scaling at *fixed q-profile* and independent scans of  $I_p$  and/or  $B_{\phi}$ .

# •Experiment

- i) **Neon** injected in **MHD** quiescent high- $\beta$ , high confinement H-modes.
- ii) The main diagnostic used was the **newly-developed "multi-color" optical SXR array**.
- iii) Atomic-based simulations constrained with the bolometer and poloidal USXR arrays.

### •Results

i)  $D_{Ne} \le 1m^2/s$  for  $(r/a \le 0.8)$  is in the neoclassical range.

ii) Low particle diffusivity that suggests anomalous transport driven by low-k electrostatic turbulence is suppressed.

#### Previous experimental work done in STs

# Neoclassical impurity transport in CDX-U ohmic plasmas





# Neoclassical impurity transport in NSTX L-mode plasmas



[2] D. Stutman, et. al., Phys. Plasmas, **10**, 4387, (2003).

# Main diagnostic used in impurity transport experiments



[3] L. F. Delgado-Aparicio, *et. al.*, to appear, PPCF, 2007.
[4] L. F. Delgado-Aparicio, *et. al.*, to appear, J. Appl. Phys., 2007

# Principle of the "optical" soft x-ray (OSXR) array

#### **Conversion of XUV emission to visible light**



#### To discrete channels and light detectors (PMT and/or APDs) + (RC/TIA) amplifiers

It's a system that uses a **fast** (~1  $\mu$ s) and **efficient** scintillator (CsI:Tl) in order to **convert soft x-ray photons** (0.1<E<sub>ph</sub><10 keV) to **visible green light** ( $\lambda$ ~550 nm).

# XUV energy discrimination & cut-off energies $(E_c)$



[3] L. F. Delgado-Aparicio, *et. al.*, to appear, PPCF, 2007.
[4] L. F. Delgado-Aparicio, *et. al.*, to appear, J. Appl. Phys., 2007

SXR metallic filters for the present tOSXR array

Metallic foil	E <sub>c</sub> (eV) for Transm.=10%	E <sub>c</sub> (eV) for Transm.=50%
Be 10 µm	780	1170
Be 100 µm	1690	2497
Be 300 µm	2416	3550

"Ideal" contrast factor (300/10 vs 100/10) ~ 1.81



### H-mode target (1MA, 4.5 kG) – background shots

- $I_p=1$  MA,  $B_0 \sim 4.5$  kG,  $W_T \sim 230$  kJ,  $\tau_E \sim 50$  ms (no large type I ELMs)
- MHD-free during time of interest.
- Constant boundary elongation ( $\delta$ ~2.25) & triangularity ( $\kappa$ ~0.6) for ~ 500 ms
- NBI modulation: 5.4 MW and 4.2 MW
- Gas puff: 1.5 torr·l/s,  $\Delta t$ ~50 ms, injection time: 0.35
- Target shots were reproducible; important for background subtraction (carbon accumulation)



#### The background shots are reproducible



# Small Neon perturbation of the background plasma



*X*-ray *E*xtreme *U*ltraviolet *S*pectrometer (integrated signal) (J. Lepson - UC – Berkeley, P. Beiersdorfer, LLNL)



# Neon plasmas don't vary significantly from background shot

Carbon Charge Exchange Recombination Spectroscopy (R. Bell, PPPL)



#### OSXR show edge Neon builds up quickly, core builds up slowly



- Good SNR of OSXR signals when Neon was injected.
- Fast edge vs slow core Neon build up.
- Strong peaking of impurities (~0.7 s) and flattening of  $T_e$  (consistent with MPTS)
- Late (1,1) MHD mode (~0.75 s).

#### No difference in flux surfaces or q-profile in neon seeded plasmas



0.2

0.4

0.6

0.8

R (m)

1.0

1.2

1.4

12

events and edge localized modes

#### Matrix-based 1-D Abel inversion



$$B_{X}(r_{T},t) = 2\int_{R}^{R_{\max}} E_{X}(R,t) \frac{RdR}{\sqrt{R^{2} - r_{T}^{2}}} \Longrightarrow E_{X}(R,t) = -\frac{1}{\pi} \int_{R}^{R_{\max}} \frac{dB_{X}(r_{T},t)}{dr_{T}} \frac{dr_{T}}{\sqrt{r_{T}^{2} - R^{2}}}$$

equivalent to

$$B_i = \sum_j L_{i,j} E_j \Longrightarrow E_j = \sum_i L_{j,i}^{-1} B_i$$
 13

# 1D (radial) impurity transport simulation (MIST) code

The time evolution of the neon emissivity after the *injection* was modeled using the time-dependent <u>M</u>ultiple Ionization <u>S</u>tage <u>T</u>ransport code (MIST).

- Atomic processes included in the model:
  i) Electron ionization ii) excitation iii) recombination
- Computes the evolution of all charge states through the experimental MPTS time history of  $n_e(R,t)$  &  $T_e(R,t)$  profiles assuming external profiles of diffusivity (D) and convective velocity (V).

#### SXR filter contribution

- 1. He- & H- like ions emit resonances lines between 0.9-1.0 keV  $\Rightarrow$  detected through the Be 10 µm.
- 2. The fully stripped ions (Ne+10) emit strong recombination continuum radiation above the 1.4 keV  $\Rightarrow$  detected through the Be 100  $\mu$ m.

#### MIST simulation constraints



#### b) Poloidal USXR arrays



# Background subtracted OSXR and MIST simulation





The impurity diffusivity estimated inside  $r/a \le 0.7$  is in good agreement with the NCLASS predicted [8] neoclassical transport coefficients.



The existence of a small to moderate inward pinch velocity depend strongly on the plasma profiles.

[8] W. Houlberg, et. al., Phys. Plasmas, 4, 3320, (1997).

#### $n_Z(R,t)$ & $T_e(R,t)$ decoupling during late impurity accumulation



by increasing the foil thickness from 10 to  $300 \,\mu m$ 

#### Summary

- The neon penetrates the core on the tens to hundred ms, with a final peaking of the neon density indicating a **diffusivity**  $\leq 1m^2/s$ , and the existence of a small to moderate inward pinch velocity depending strongly on the plasma profiles.
- Preliminary NCLASS simulations indicate that the **particle diffusivity is in the order of the values predicted by the neoclassical transport theory**.
- Thermal ion transport from TRANSP is also in the neoclassical range.
- Particle transport around the ion neoclassical values would suggest a very low level of turbulent ion transport.
- Electron transport in this discharges is close to two orders of magnitude higher (  ${\it O}~(m_i/m_e)^{1/2}$  )

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