



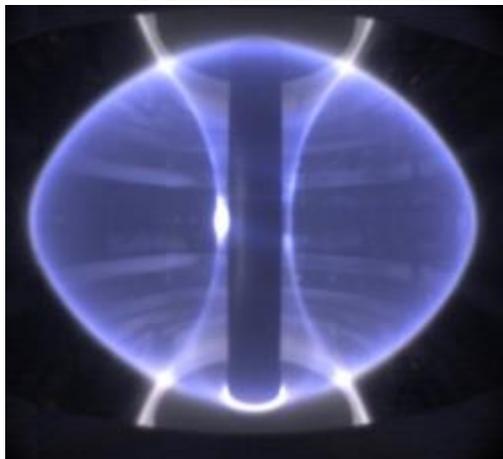
# TH/P3-14: Measurement and prediction of momentum transport in spherical tokamaks

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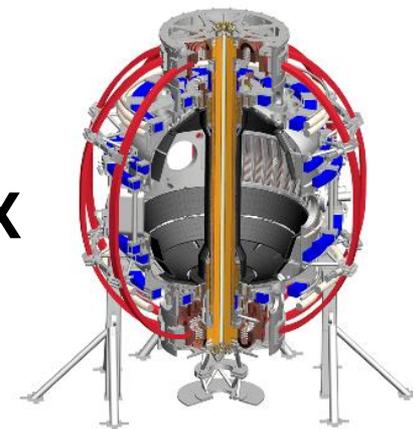


IAEA FEC 2016, Kyoto, Japan



**MAST**

**NSTX**



# OVERVIEW

- To predict rotation profile (important for macro- and micro-instabilities) need to understand torques sources/sinks and momentum transport [Ida & Rice NF (2014)]
- Here, investigating momentum pinch in low aspect ratio, high beta spherical tokamak plasmas (NSTX & MAST) as an additional constraint on theory (stems out of ITPA T&C activity)
- **Summary:**
  - Previous NSTX H-mode experiments inferred momentum pinch comparable to conventional tokamaks,  $(RV_{\phi}/\chi_{\phi})_{\text{exp}} \sim (-1) - (-10)$
  - However, local, quasi-linear GK theory predicts negligible pinch  $(RV_{\phi}/\chi_{\phi})_{\text{sim}} \sim 0 - (-1)$  in NSTX H-modes due to electromagnetic and low aspect ratio effects on mode-symmetry
  - MAST L-mode experiments (i.e. lower beta) were conducted, analysis not inconsistent with significant pinch,  $(RV_{\phi}/\chi_{\phi})_{\text{exp}} \sim (-1) - (-9)$
  - However, non-stationary conditions also allow for weaker or outward pinch,  $(RV_{\phi}/\chi_{\phi})_{\text{exp}} \sim (-2) - (5)$
  - Local, quasi-linear GK theory predicts weak pinch  $(RV_{\phi}/\chi_{\phi})_{\text{sim}} \sim (-1)$  similar to NSTX H-modes, too much uncertainty in experiments to constrain predictions

# Toroidal angular momentum transport considered as sum of diffusion ( $-\chi_\phi \nabla \Omega$ ), convection ( $V_\phi \Omega$ ) and residual stress ( $\Pi_{RS}$ )

- Transport equation: 
$$\frac{\partial}{\partial t} (n_i m_i \langle R^2 \rangle \Omega) + \nabla \cdot \Pi_\phi = T$$

- Assumed transport form: 
$$\Pi_\phi = nmR^2 \left( -\chi_\phi \nabla \Omega + V_\phi \Omega \right) + \cancel{\Pi_{RS}}$$

*Ignoring residual stress contributions throughout this work*

- Can identify different physical mechanisms by how they break symmetry of microinstability [Peeters, NF (2011); Angioni, PFR (2012); Diamond, NF (2013)]

- Pinch expected due to Coriolis effect (Peeters, 2007), or equivalently turbulent equipartition (Hahm, 2007) + thermoelectric force (Peeters, 2009)

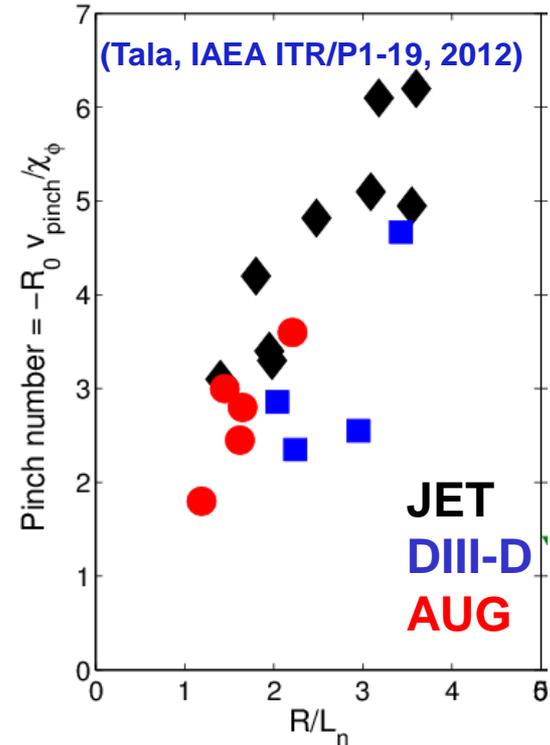
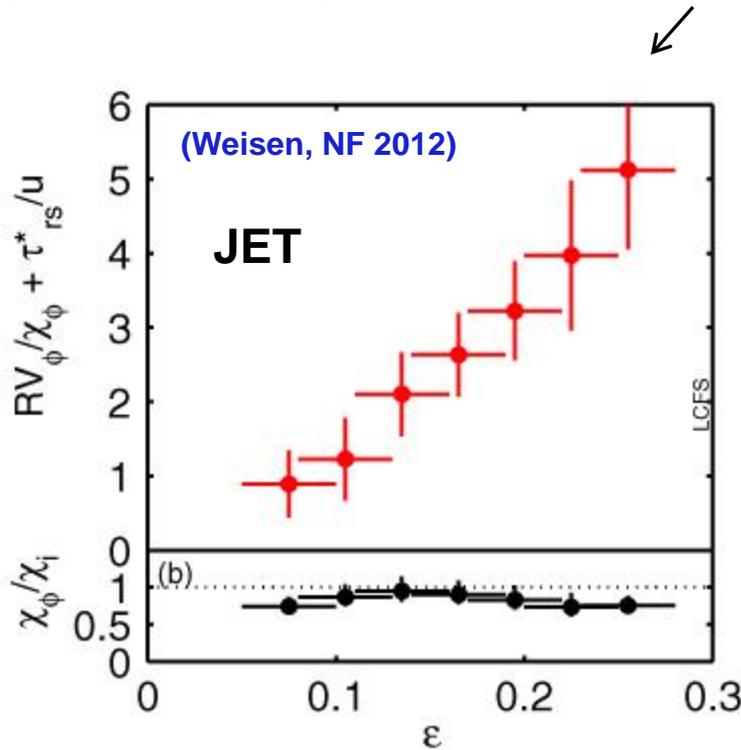
$$\hat{\Pi}_\phi = \hat{\chi}_\phi \left( \hat{u}' + \frac{RV_\phi}{\chi_\phi} \hat{u} \right)$$

$$\hat{u}' = \frac{-R^2 \nabla \Omega}{c_s}$$

$$\hat{u} = \frac{R\Omega}{c_s}$$

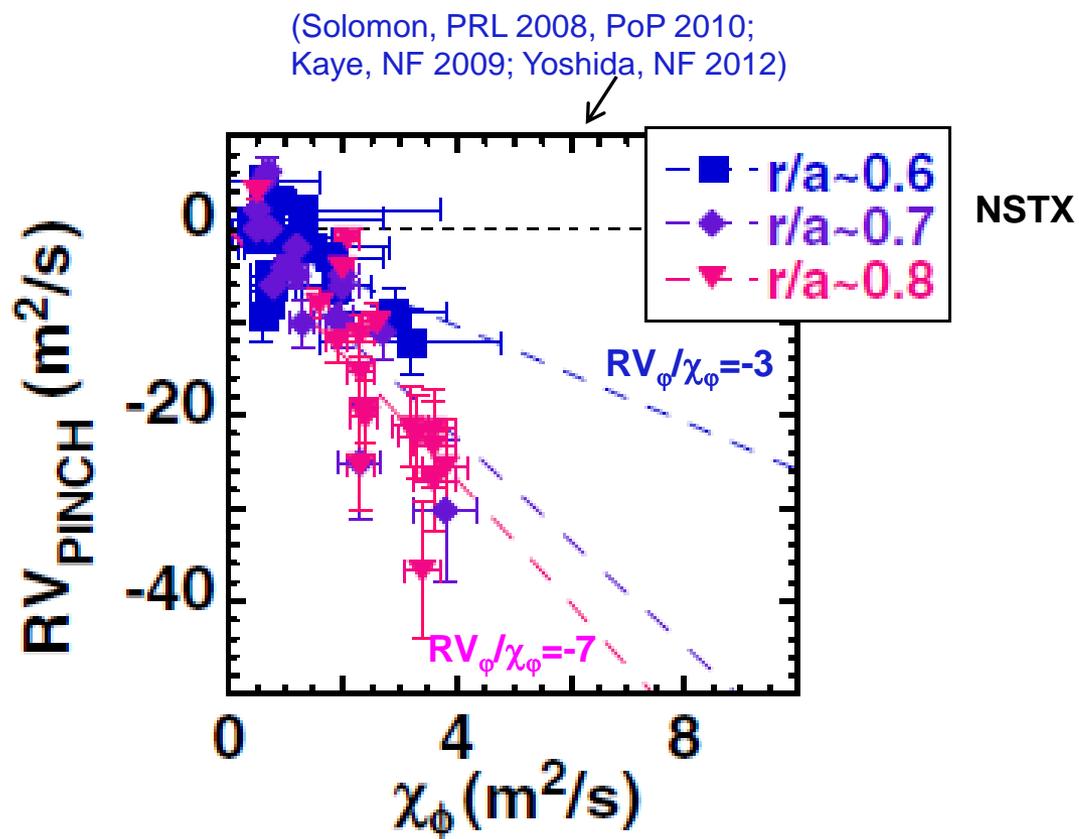
# Momentum pinch measured and predicted in conventional tokamaks

- Measurements in many machines from both perturbative experiments (NBI, 3D coils) and statistical regression analysis



- Increase in (inward) pinch observed with  $\varepsilon=r/R$  and  $R/L_n$ , also predicted by ITG theory (Peeters, PRL 2007; PoP 2009)

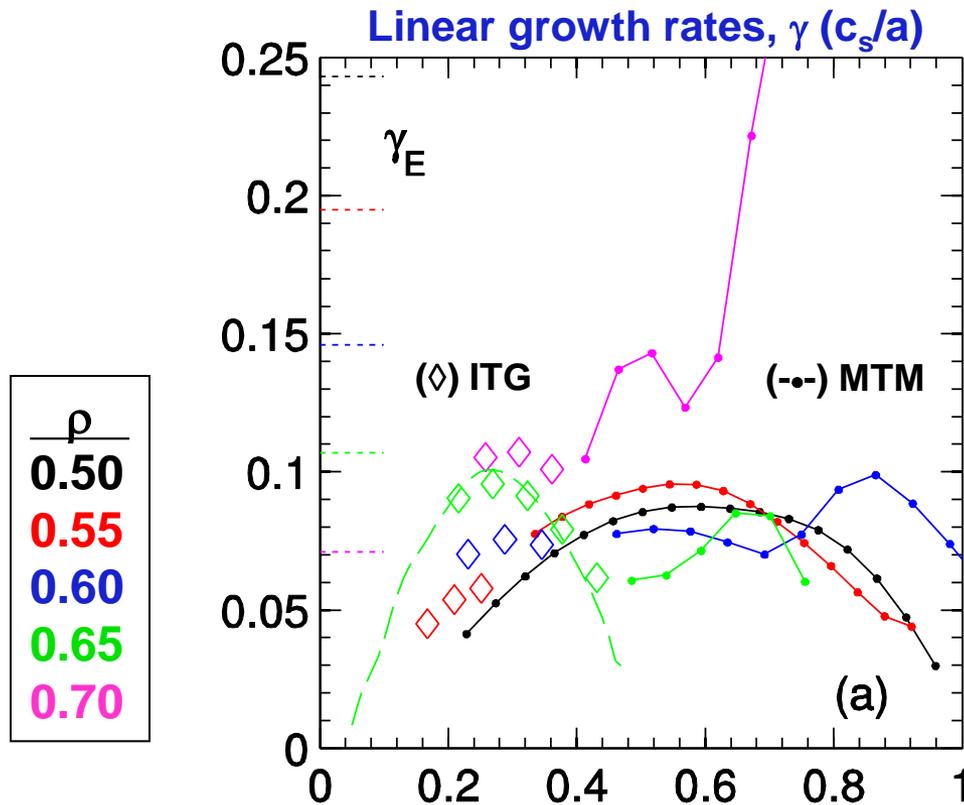
# Perturbative NSTX H-mode experiments indicate existence of an inward momentum pinch, $RV_{\phi}/\chi_{\phi} \approx -(1-7)$



- Local, linear gyrokinetic simulations of ITG turbulence describe pinch and scaling in conventional tokamaks  $\Rightarrow$  **does this hold for STs?**

# Higher beta NSTX H-modes often dominated by microtearing modes (MTM) with sub-dominant ballooning modes

- Most cases have  $\gamma_{\text{MTM}} > \gamma_{\text{ballooning}}$  ( $\Pi_{\phi}=0$  for MTM)
- Sub-dominant modes can be ITG, KBM or compressional ballooning modes – calculate pinch *assuming* they contribute to transport



**Linear GYRO simulations**  
(Candy, Waltz, 2003)

3 species: D, C, e

EM:  $\phi$ ,  $A_{\parallel}$ ,  $B_{\parallel}$

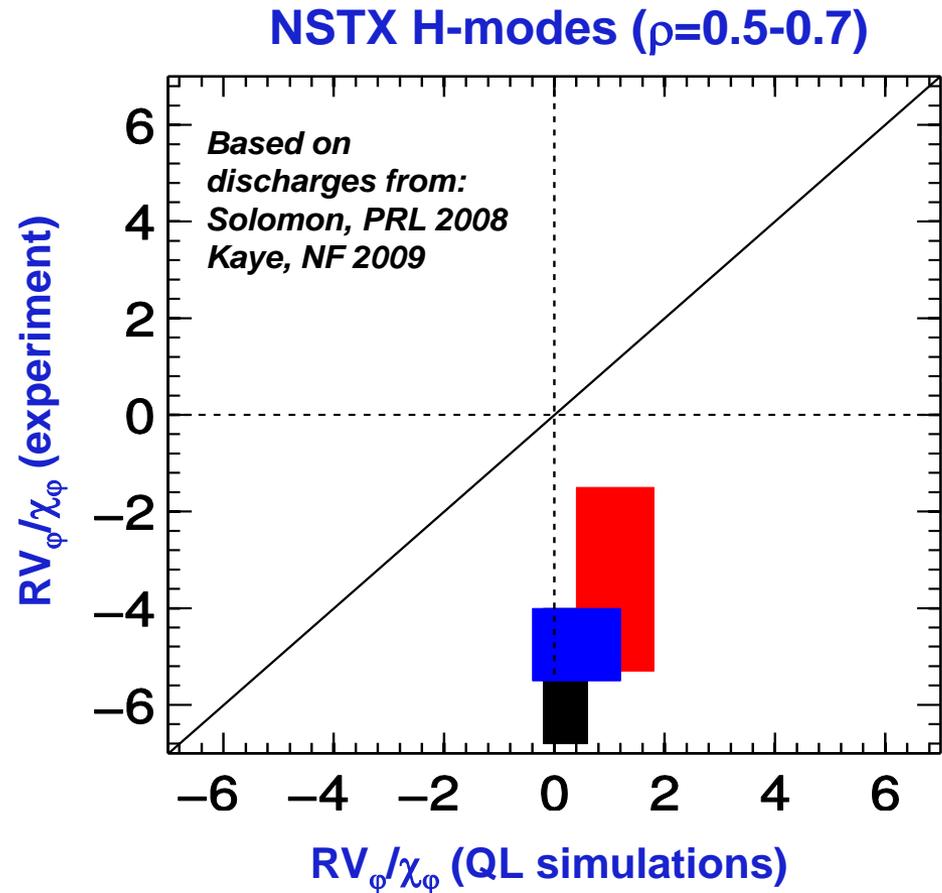
Equilibrium reconstruction

$\beta_T=12\%$ ,  $\beta_N=3.5$

Guttenfelder, 2016 (Phys. Plasmas, *in review*)

# Negligible or outward momentum convection predicted from ES and EM ballooning modes in NSTX

- Weak/outward pinch consequence of parallel mode structure response at high beta, low aspect ratio, see:
  - Peeters, PoP (2009)
  - Kluy, PoP (2009)
  - Hein, PoP (2011)
  - Guttenfelder, PoP (2016)



# Momentum pinch coupled to symmetry breaking in parallel mode structure

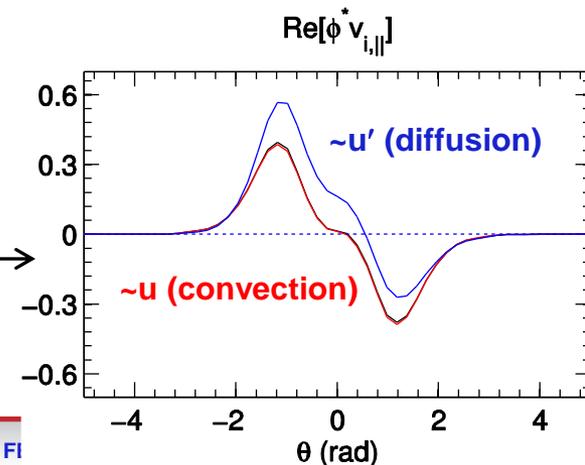
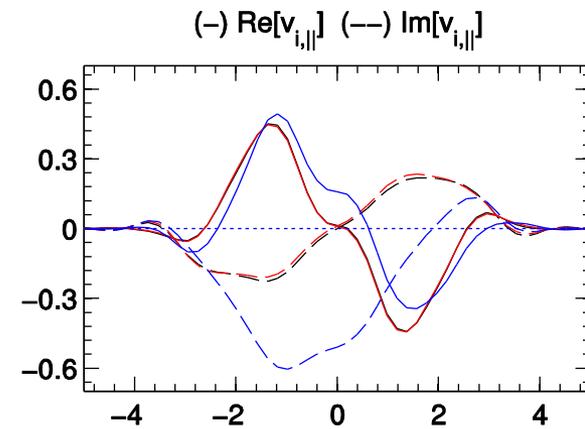
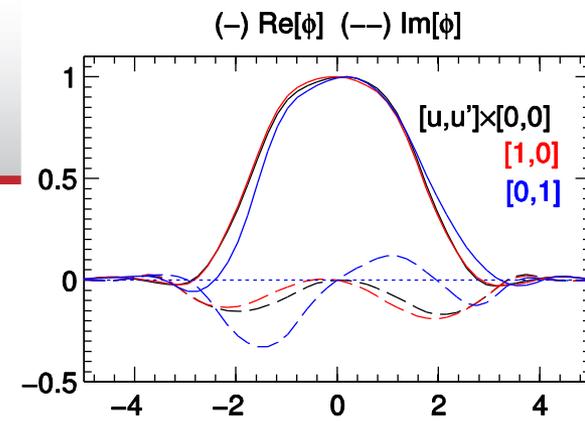
- Component of curvature drift in lab frame ( $M < 1$  smaller than curv. drift)

$$v_{\kappa} \approx \frac{mv_{\parallel}^2}{eBR} \rightarrow \frac{2m(v_{\parallel} + u_0)^2}{eBR} = \frac{mv_{\parallel}^2}{eBR} + \frac{2mv_{\parallel}u_0}{eBR} + \frac{mu_0^2}{eBR}$$

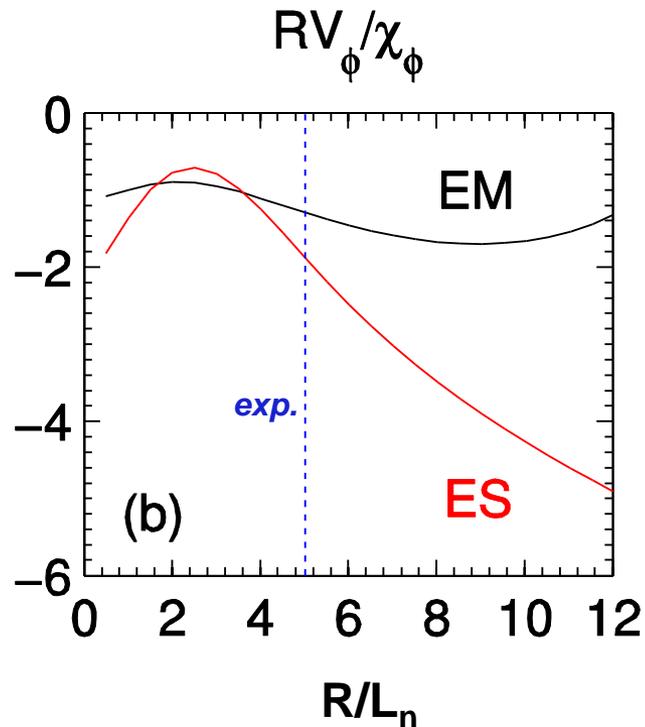
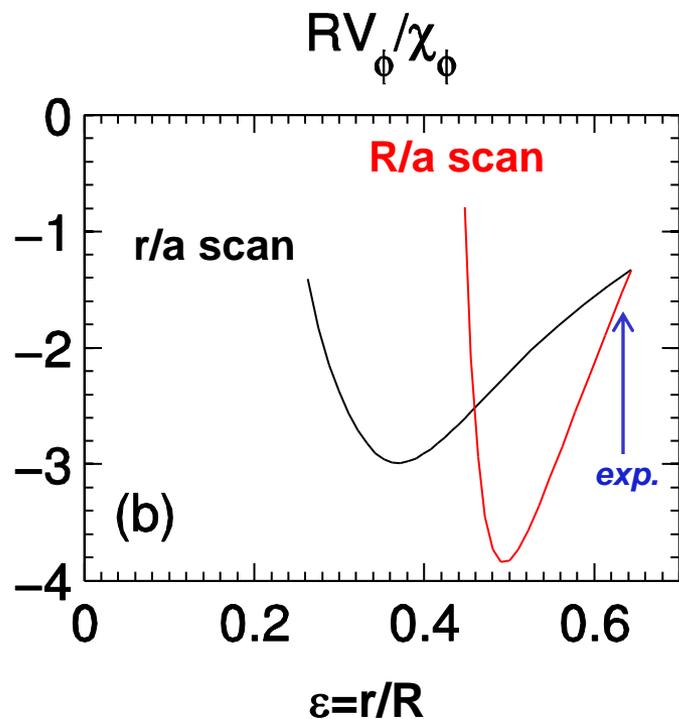
$\text{Curvature}$   
( $v_{\kappa}$ )
 $\text{Coriolis}$   
( $\sim M \cdot v_{\kappa}$ )
 $\text{Centrifugal}$   
( $\sim M^2 \cdot v_{\kappa}$ )

- Does not influence stability, but toroidal flow couples  $\delta n, \delta T$  with  $\delta u \rightarrow$  can cause momentum transport if eigenfunctions develop parallel asymmetry

- Parallel asymmetry from  $u > 0$  very small in NSTX due to strong particle trapping & toroidicity  $\rightarrow$  little convective transport**  $\longrightarrow$



# A larger (inward) pinch can be found: (i) at increased aspect ratio, (ii) in purely ES limit at high $\nabla n$

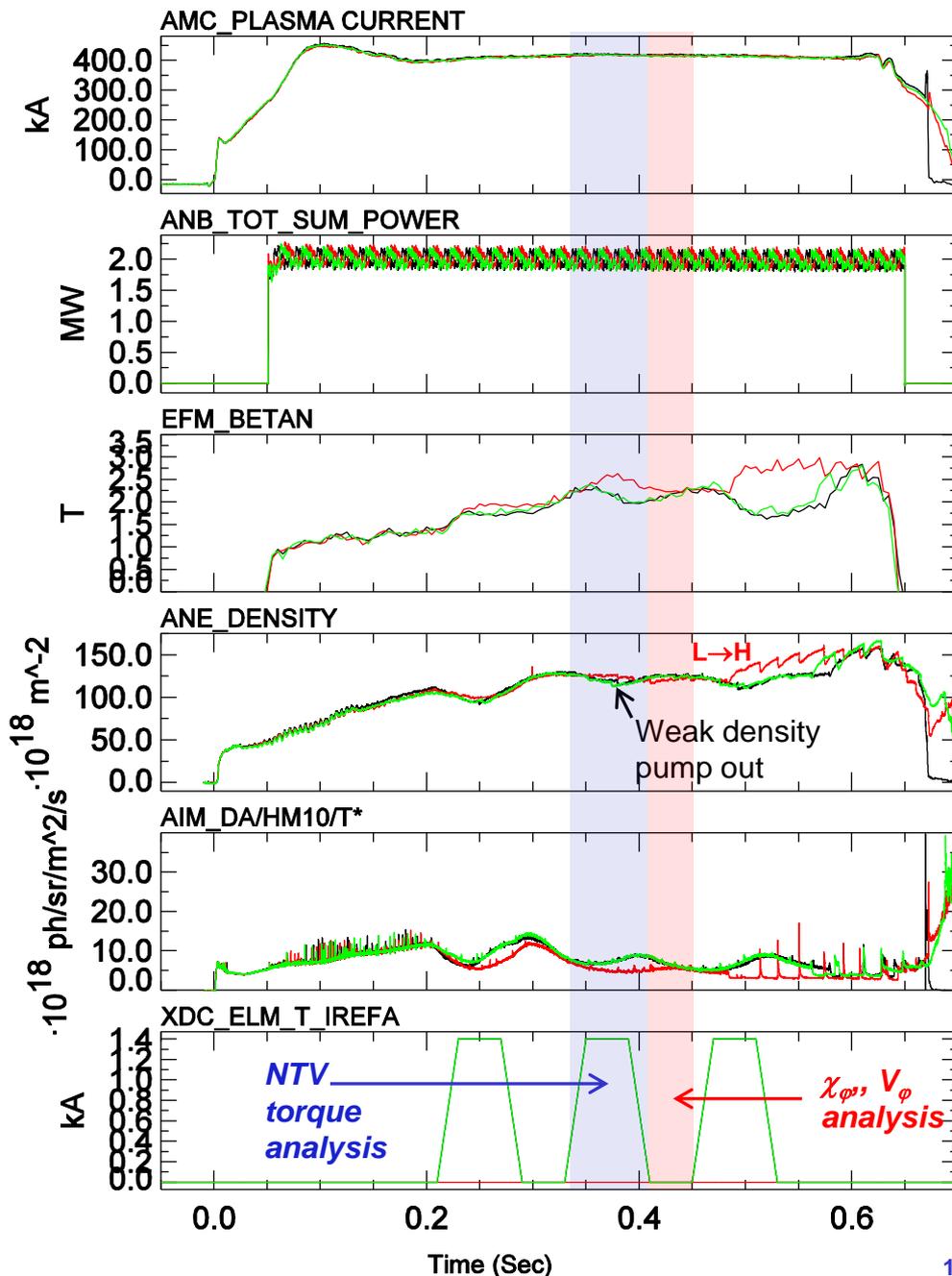


- Above simulations based on NSTX L-mode discharge (Ren, NF 2013)
- Variation in pinch related to changes in parallel mode structure and symmetry (Guttenfelder, PoP 2016)
- Can't do aspect ratio scan, can try to do similar analyses at lower beta...

# MAST L-mode experiment conducted in 2013

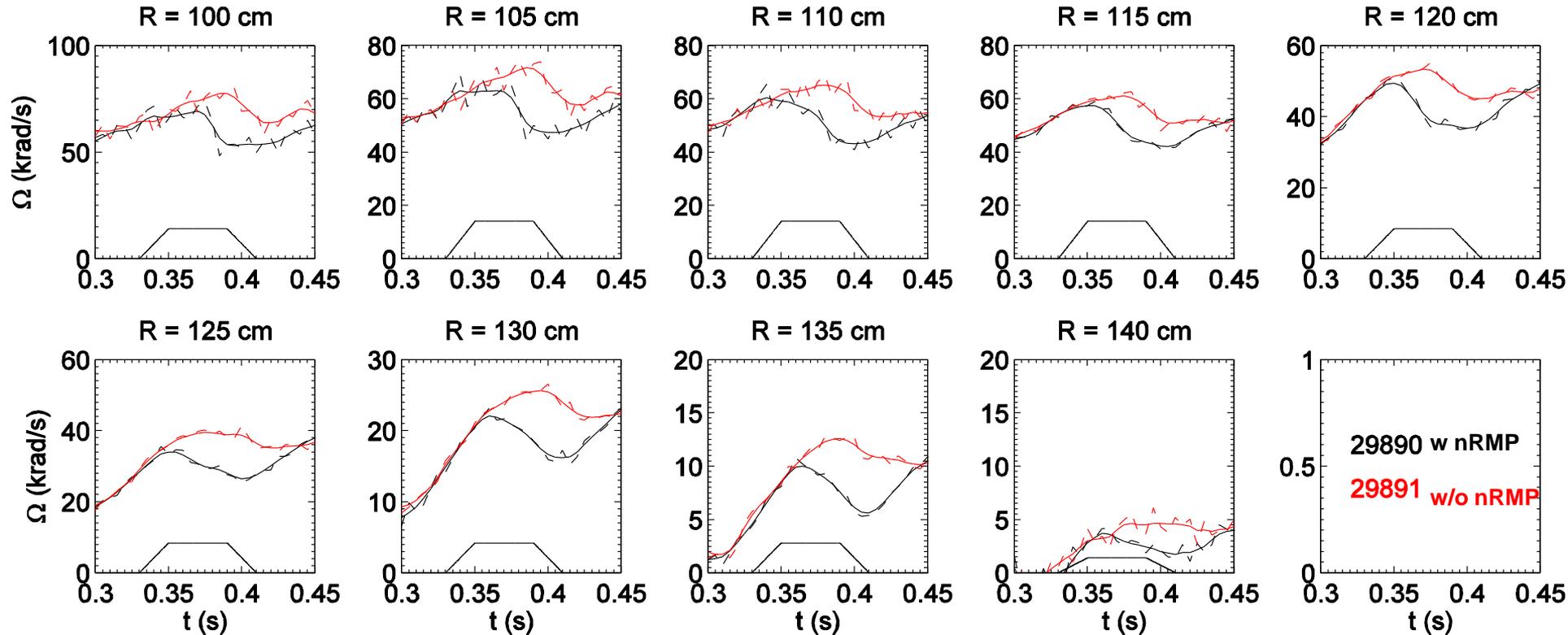
- 2 MW LSN L-mode
  - $\langle n_e \rangle = 2.3 \times 10^{19} \text{ m}^{-3}$
  - $B_T = 0.5 \text{ T}$ ,  $I_p = 0.4 \text{ MA}$  ( $q_{95} \approx 5$ )
  - $\beta_N \sim 2$ ,  $\beta_T \sim 4\%$
- Using applied 3D fields ( $n=3$ ) to perturb rotation
  - 29890/ 29892 – three  $n=3$  field pulses applied to brake rotation
  - 29891 – no nRMP pulses
- Weak density pump out w/ nRMP, drop in  $\beta_N$
- Without nRMP, eventual transition into H-mode ( $t \sim 0.47 \text{ s}$ )

Shot: — 29890 — 29891 — 29892



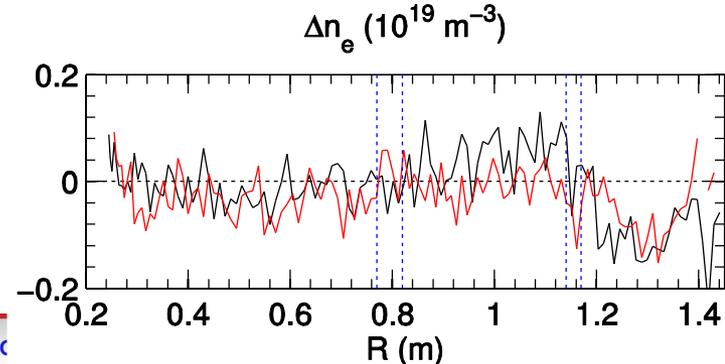
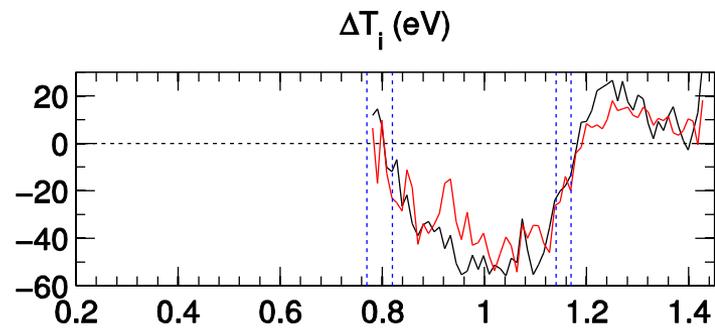
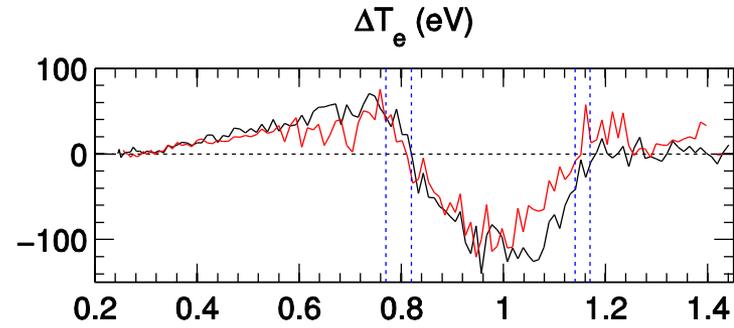
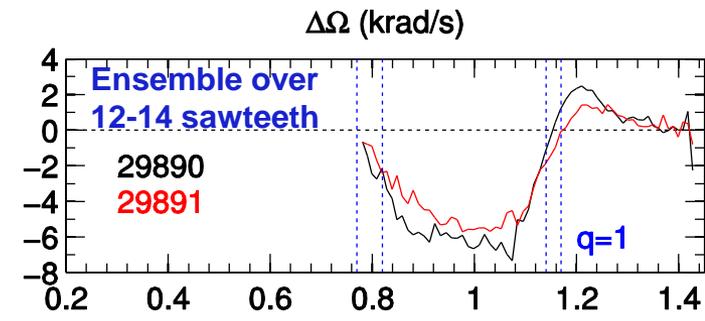
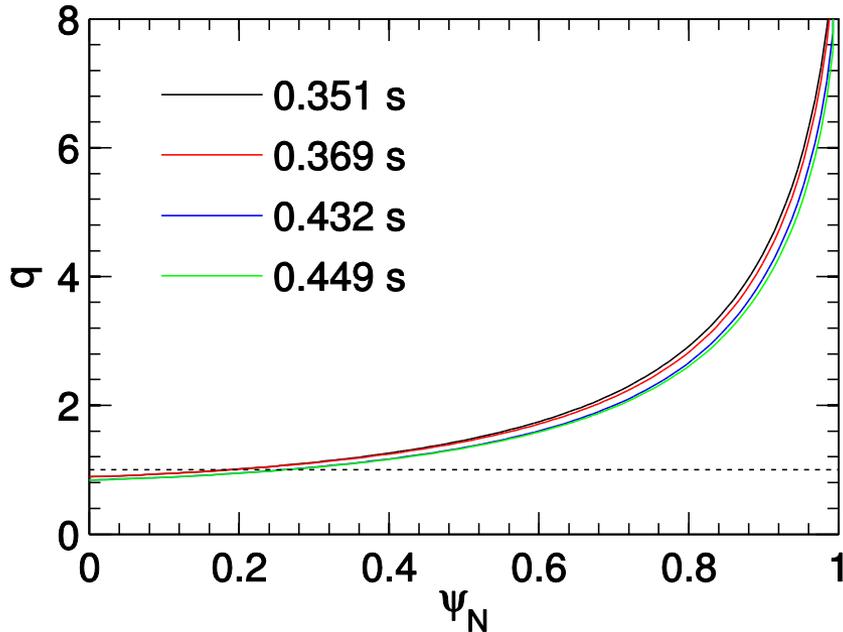
# Changes in toroidal rotation due to 3D fields clearly observed

- Non-stationary conditions -- control shot (29891) provides a baseline for analysis (will impact analysis discussed later)
- Filtering to remove faster sawteeth oscillations ( $\Delta t_{ST} \sim 6-12$  ms)
  - $\Delta \Omega_{ST} \sim 2-6$  krad/s  $<$   $\Delta \Omega_{3D} \sim 10-20$  krad/s



# Effect of sawteeth on rotation weaker than applied 3D fields

- Sawteeth cause  $\sim 6$  krad/s ( $\sim 8\%$ ) deceleration inside inversion radius
- $q=1$  surface  $\psi_N \sim 0.19-0.26$  ( $R_{\text{out}} \sim 114-118$  cm) consistent with  $\Delta T_e$  inversion
  - $\Delta T_e \sim 120$  eV ( $\sim 16\%$  of  $T_{e,0} \sim 750$ )
  - $\Delta T_i \sim 50$  eV ( $\sim 6\%$  of  $T_{i,0} \sim 800$ )



# Method to infer $\chi_\phi$ and $V_\phi$ from transient rotation response after RMP turn-off

- TRANSP solves for momentum flux,  $\Pi$ , using the flux-surface-averaged toroidal angular momentum transport equation (Goldston, Varenna 1985), plus NUBEAM calculations for torque sources & sinks:

$$\frac{\partial}{\partial t} \left( \sum_i n_i m_i \langle R^2 \rangle \Omega \right) + \frac{1}{V'} \frac{\partial}{\partial \rho} [V' \cdot \Pi] = \sum T_{\text{source}} - \sum T_{\text{sink}}$$

- Assuming momentum flux composed of only diffusive and convective contributions:

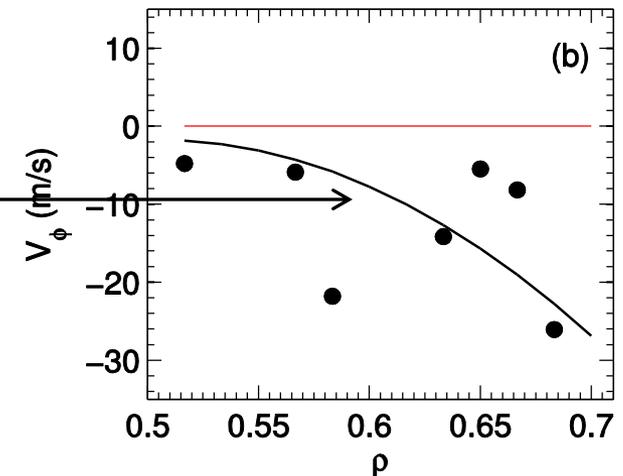
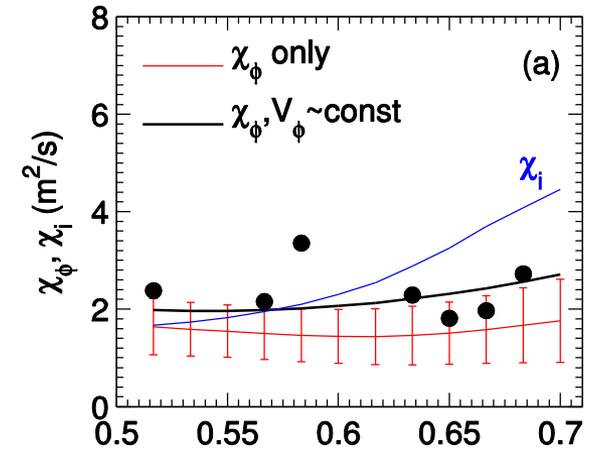
$$\Pi = \sum_i n_i m_i \left[ -\langle R^2 (\nabla \rho)^2 \rangle \chi_\phi \frac{\partial \Omega}{\partial \rho} + \langle R^2 \rangle \langle \nabla \rho \rangle V_\phi \Omega \right]$$

we can use  $\Pi(\rho, t)$ ,  $d\Omega/d\rho(\rho, t)$ , and  $\Omega(\rho, t)$  in a nonlinear least squares fit algorithm to determine best fit  $\chi_\phi(\rho)$ ,  $V_\phi(\rho)$  (assumed constant in time)

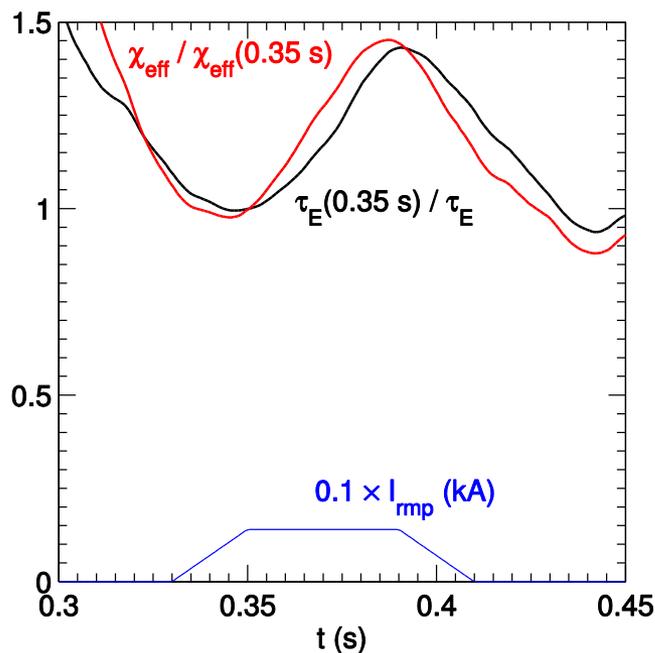
- Using time window (0.41-0.45 s) after removal of NTV torque ( $T_{\text{NTV}}=0$ )
- Note: method only valid if  $d\Omega/d\rho(t)$  and  $\Omega(t)$  are sufficiently decorrelated

# Inward momentum pinch inferred from transient recovery

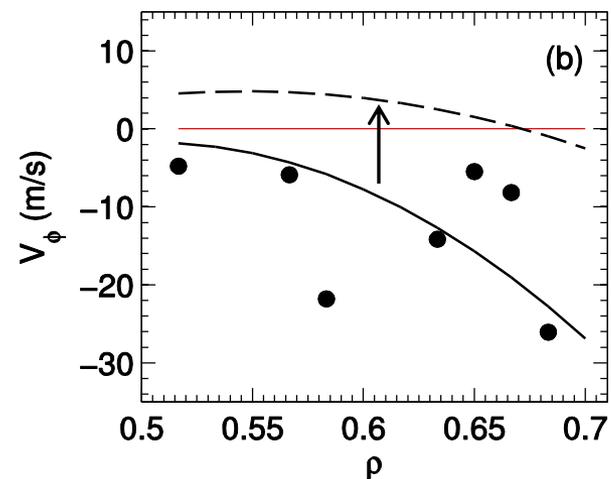
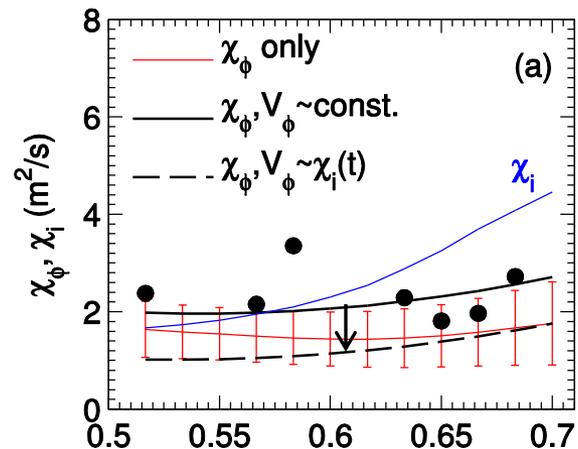
- $\chi_\phi$ ,  $V_\phi$  assumed constant in time
- Using both  $\chi_\phi$  and  $V_\phi$  improves the quality of fit ( $\chi_v^2$  smaller than  $\chi_\phi$ -only fit)
- At locations where there is a strong  $\Omega$ - $\nabla\Omega$  linear correlation, method is ill-posed  $\Rightarrow \chi_\phi$  &  $V_\phi$  tend to large values
- Can fit entire analysis region simultaneously using polynomial profiles
  - Best fit (lowest  $\chi_v^2$ ) using quadratic



# Energy confinement and local thermal transport is non-stationary during analysis window

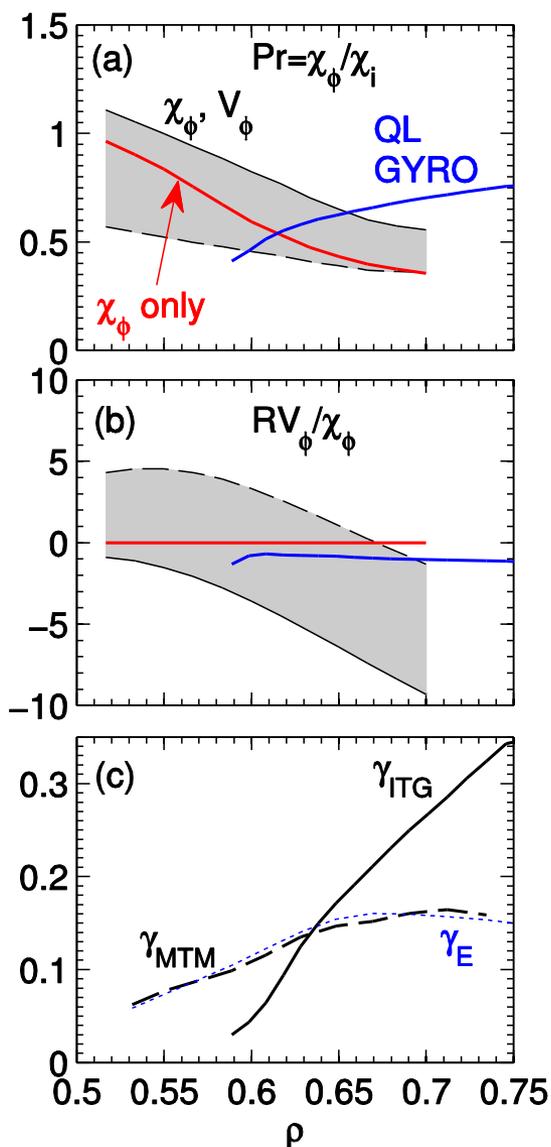


- $\tau_{E,\text{th}}$  drops  $\sim 40\%$ ,  $\chi_{\text{eff}}$  increases 40%
- Inferred pinch becomes smaller or outward assuming  $\chi_{\phi}, V_{\phi} \sim \chi_i(t) \sim 1/\tau_E(t)$



# Resulting pinch parameter covers broad range, too much uncertainty to constrain quasi-linear predictions

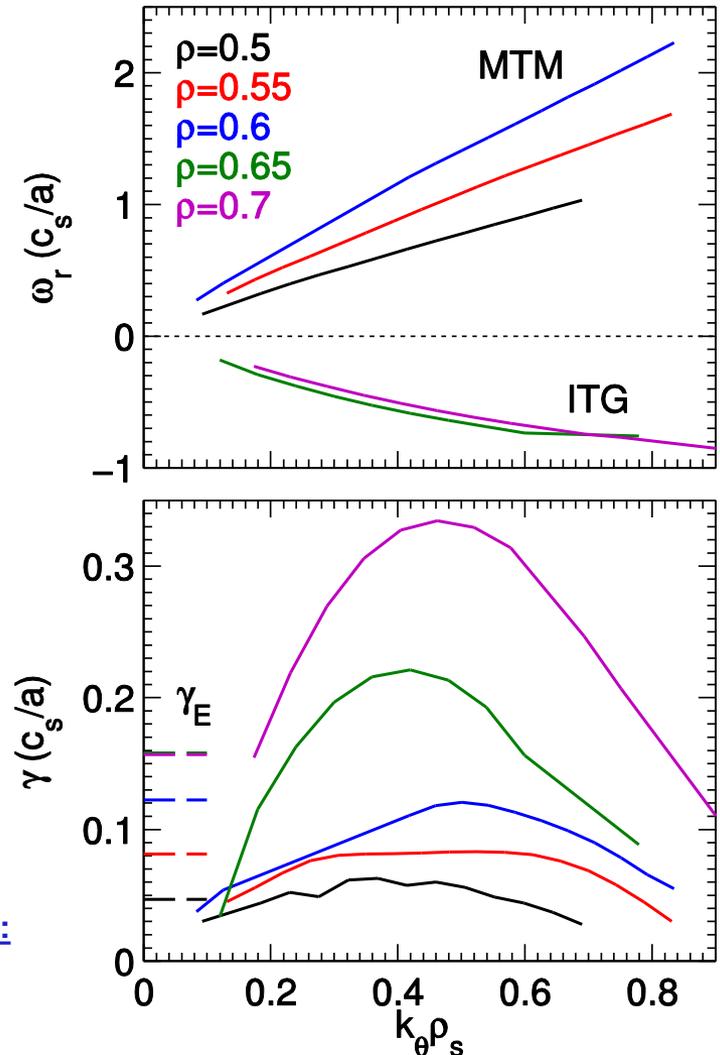
- Prandtl number varies between  $Pr_{\text{exp}}=0.4-1.1$ 
  - Quasilinear Prandtl number from unstable ITG  $Pr_{\text{ITG}} \sim 0.5-0.8$ , in range of experimental inference
- Pinch parameter ranges between  $(RV_{\phi}/\chi_{\phi})_{\text{exp}}=(-1)-(-9)$  assuming fixed coefficients, or  $(RV_{\phi}/\chi_{\phi})_{\text{exp}}=(-1)-(5)$  for time-varying coefficients
  - Quasilinear pinch parameter is very small,  $(RV_{\phi}/\chi_{\phi})_{\text{ITG}} \sim -1$
  - Similar to weak pinch predicted in NSTX L and H-modes
- **Too much experimental uncertainty to constrain pinch predictions**



# Linear GYRO\* simulations used to predict unstable modes and corresponding momentum pinch

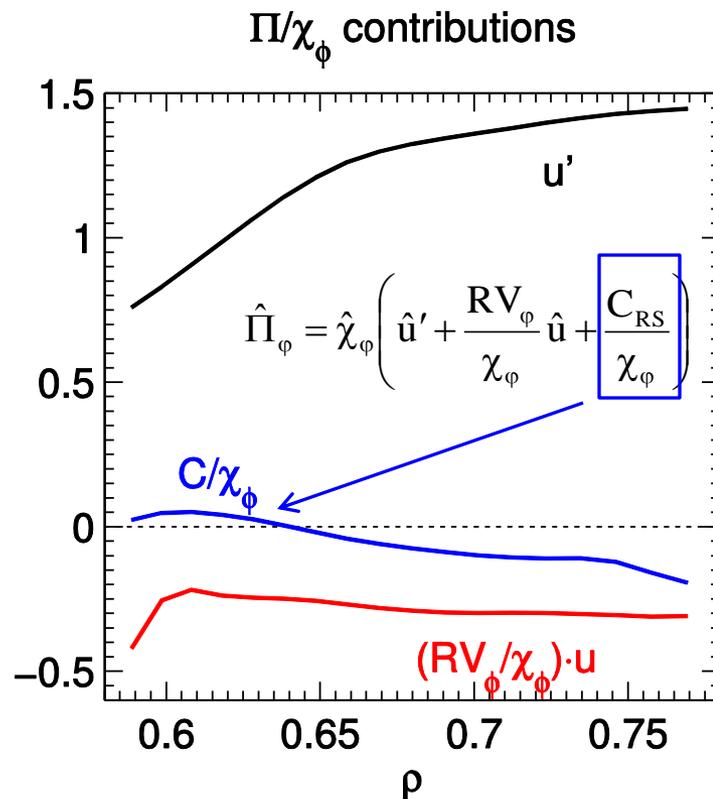
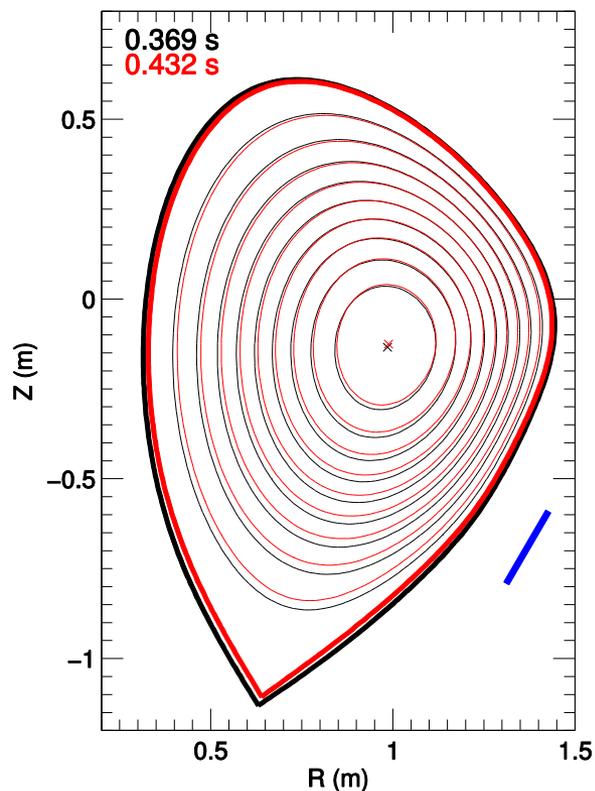
- Broad spectrum of microtearing modes (MTM) predicted  $\rho=0.5-0.6$ 
  - Even in L-mode, beta relatively high ( $\beta_N=2$ ,  $\beta_T=4\%$ )
- However, no momentum transport predicted from MTM
- Electrostatic ITG becomes dominant  $\rho>0.6$
- Can compute quasi-linear Prandtl number and momentum pinch for ITG mode (previous slide)

\*Linear GYRO simulations  
(Candy, Waltz, 2003) using:  
3 species: D,C,e  
EM:  $\phi$ ,  $A_{\parallel}$ ,  $B_{\parallel}$   
Equilibrium reconstruction



# Have also begun investigating theoretical residual stress contributions

- Quasilinear residual stress from strong up-down asymmetry (Camenen, PRL 2009) predicted to be smaller than diffusive or convective contributions



- Also investigating residual stress sources due to finite  $\rho_* \sim 1/100$  with global GTS sims, e.g. profile shear  $\sim \omega_r' \cdot \rho_*$  (Camenen, NF 2011), intensity shear  $\sim d(\gamma_{ITG} - \gamma_E)/dr \cdot \rho_*$  (Gurcan, PoP 2010), zonal flow shear (Wang, FEC2016 TH/P3-12; PoP 2010)

# Summary & future work

- Previous NSTX H-mode experiments inferred momentum pinch comparable to conventional tokamaks,  $(RV_\phi/\chi_\phi)_{\text{exp}} \sim (-1) - (-10)$
- However, local, quasi-linear GK theory predicts negligible pinch  $(RV_\phi/\chi_\phi)_{\text{sim}} \sim 0 - (-1)$  in NSTX H-modes due to electromagnetic and low-aspect-ratio effects on mode-symmetry
- MAST L-mode experiments (i.e. lower beta) were conducted, analysis not inconsistent with significant pinch,  $(RV_\phi/\chi_\phi)_{\text{exp}} \sim (-1) - (-9)$
- However, non-stationary conditions ( $\chi_\phi, V_\phi \sim \chi_i(t)$ ) also allow for weaker or outward pinch,  $(RV_\phi/\chi_\phi)_{\text{exp}} \sim (-2) - (5)$
- Local, quasi-linear GK theory predicts weak pinch  $(RV_\phi/\chi_\phi)_{\text{sim}} \sim (-1)$  similar to NSTX H-modes - **too much uncertainty in experiments to constrain MAST L-mode predictions**
- Future NSTX-U L-mode experiments are planned to continue investigation of momentum pinch and residual stress contributions