

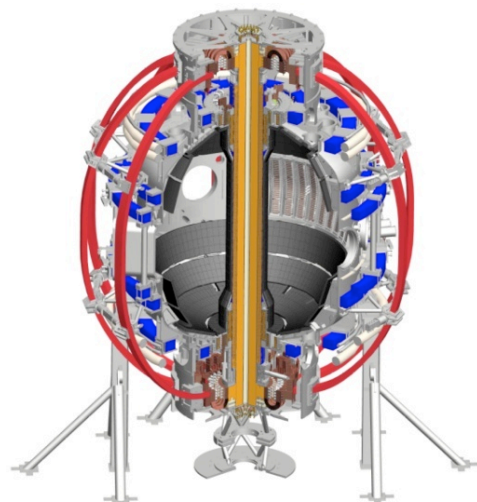
# Pedestal Structure, Fluctuations, and Transport Analysis during an ELM cycle on NSTX

**Ahmed Diallo<sup>1</sup>**

*R. Maingi<sup>2</sup>, D. Smith<sup>3</sup>, G. Kramer<sup>1</sup>, S. Kubota<sup>4</sup>, T. Osborne<sup>5</sup>, J. Canik<sup>2</sup>, S-H Ku<sup>1</sup>, M. Podesta<sup>1</sup>, W. Guttenfelder<sup>1</sup>, R. Bell<sup>1</sup>, B. LeBlanc<sup>1</sup>, C-S. Chang<sup>1</sup> and the NSTX Research Team*

- <sup>1</sup> PPPL
- <sup>2</sup> ORNL
- <sup>3</sup> University Wisconsin
- <sup>4</sup> UCLA
- <sup>5</sup> GA

**TTF-Workshop  
 Annapolis, MD  
 April 2012**



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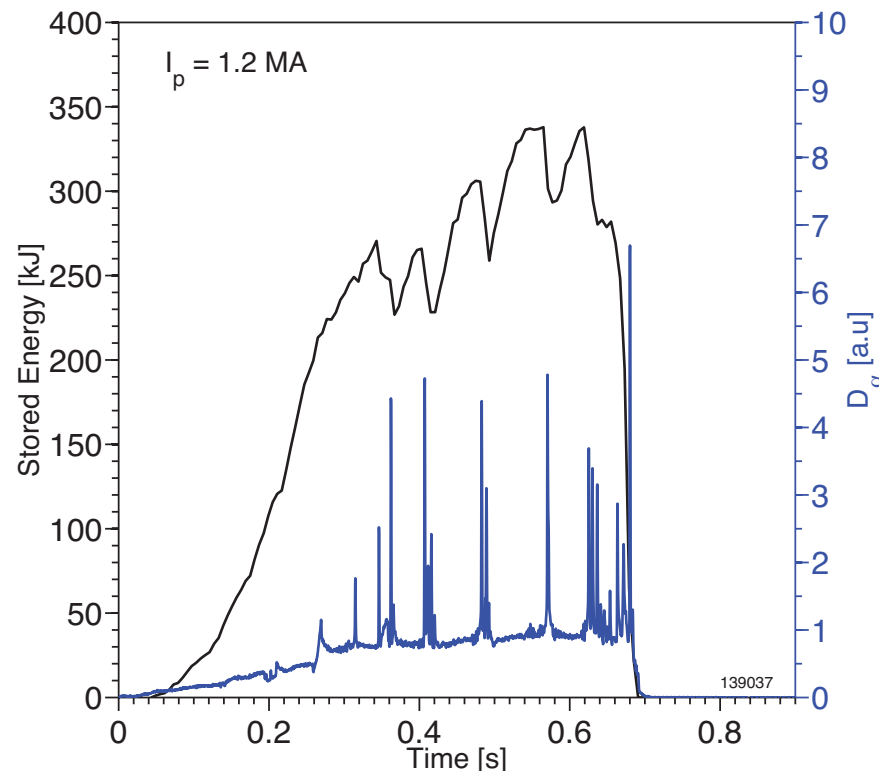
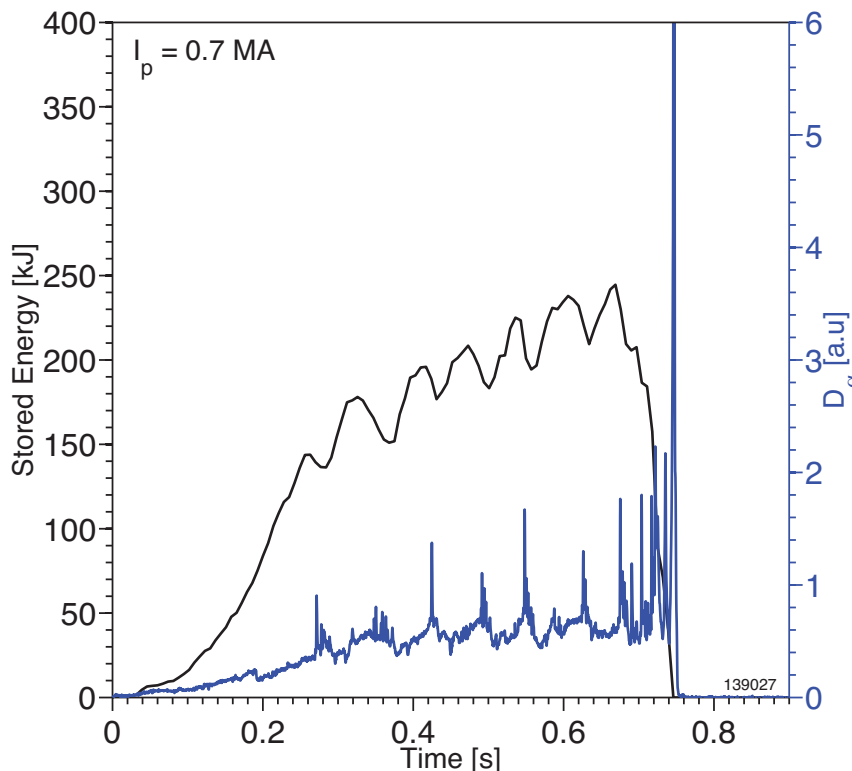
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## Characterize the pedestal structure evolutions and pedestal region turbulence to elucidate the role of transport in pedestal constraints

- Summary of the pedestal structure characterization in ELMy H-modes
- Interpretive transport analysis
  - SOLPS calculations
  - Comparison with XGC0 - neoclassical simulations
- Fluctuations characterizations during the inter-ELM phase
- Preliminary  $\delta f$  XGC1 calculations and comparison with observations
- Analysis of lithium effects on transport and turbulence will be shown in this afternoon session by **D. Boyle; R. Maingi**

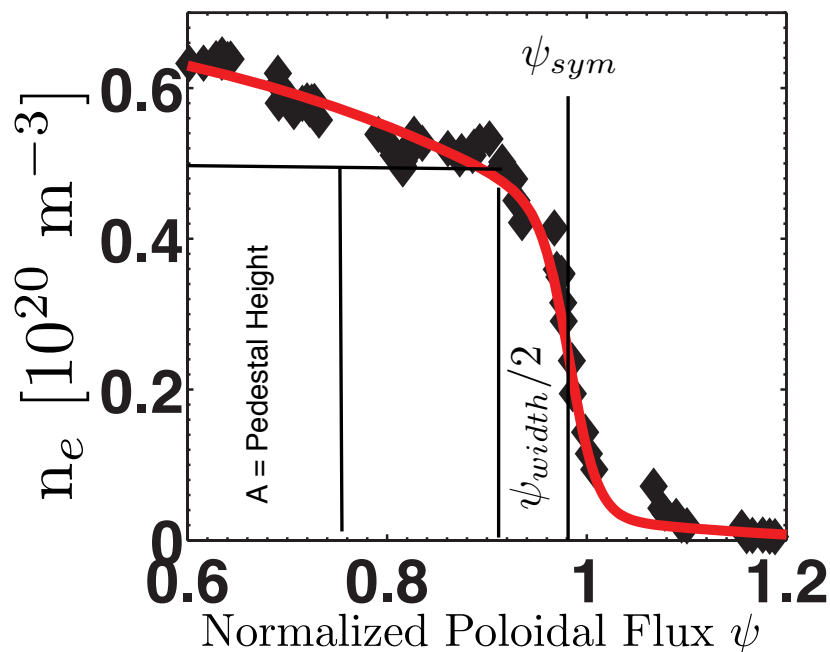
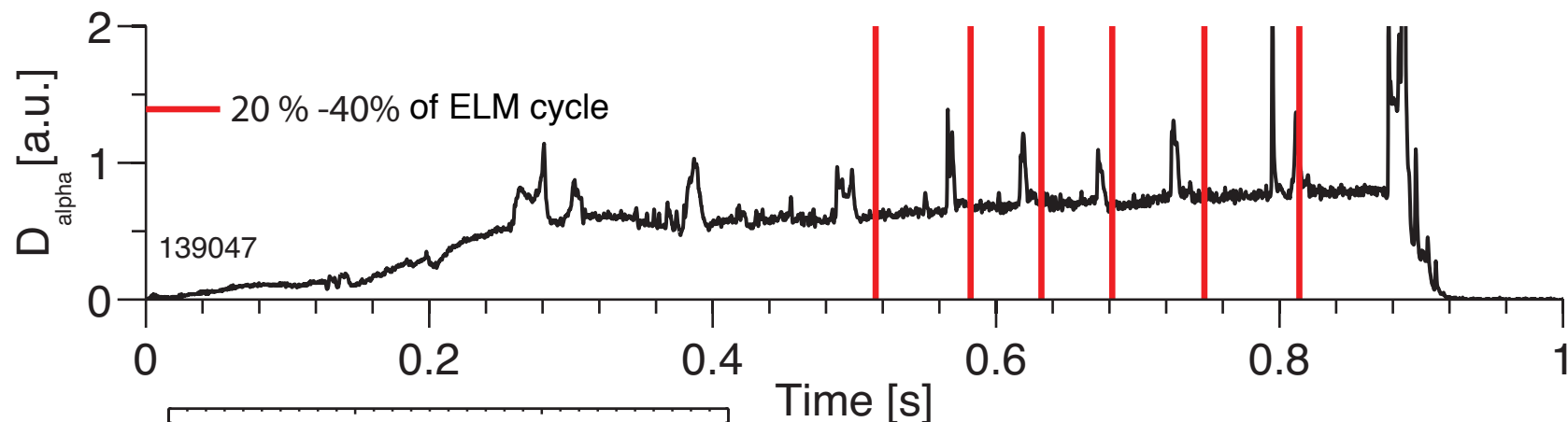
# Dedicated experiments to vary the pedestal pressure height and width through $I_p$ scans were performed on NSTX

- Constant injected power ( $P_{\text{NBI}}$ ) and magnetic field ( $B_T$ )
- Lower single null slightly downward and fixed high triangularity shaping.



- Large drop (up to 15%) of stored energy ( $W_{\text{mhd}}$ ) after each ELM crash.
  - Pedestal stored energy  $\sim 25\% - 40\%$  of  $W_{\text{mhd}}$
- Implicitly generating scans of the pedestal structure.

# Radial profiles of density, temperature and pressure are composite of times between multiple fraction of ELMs



- Electrons edge radial profiles are fitted using a modified tanh function

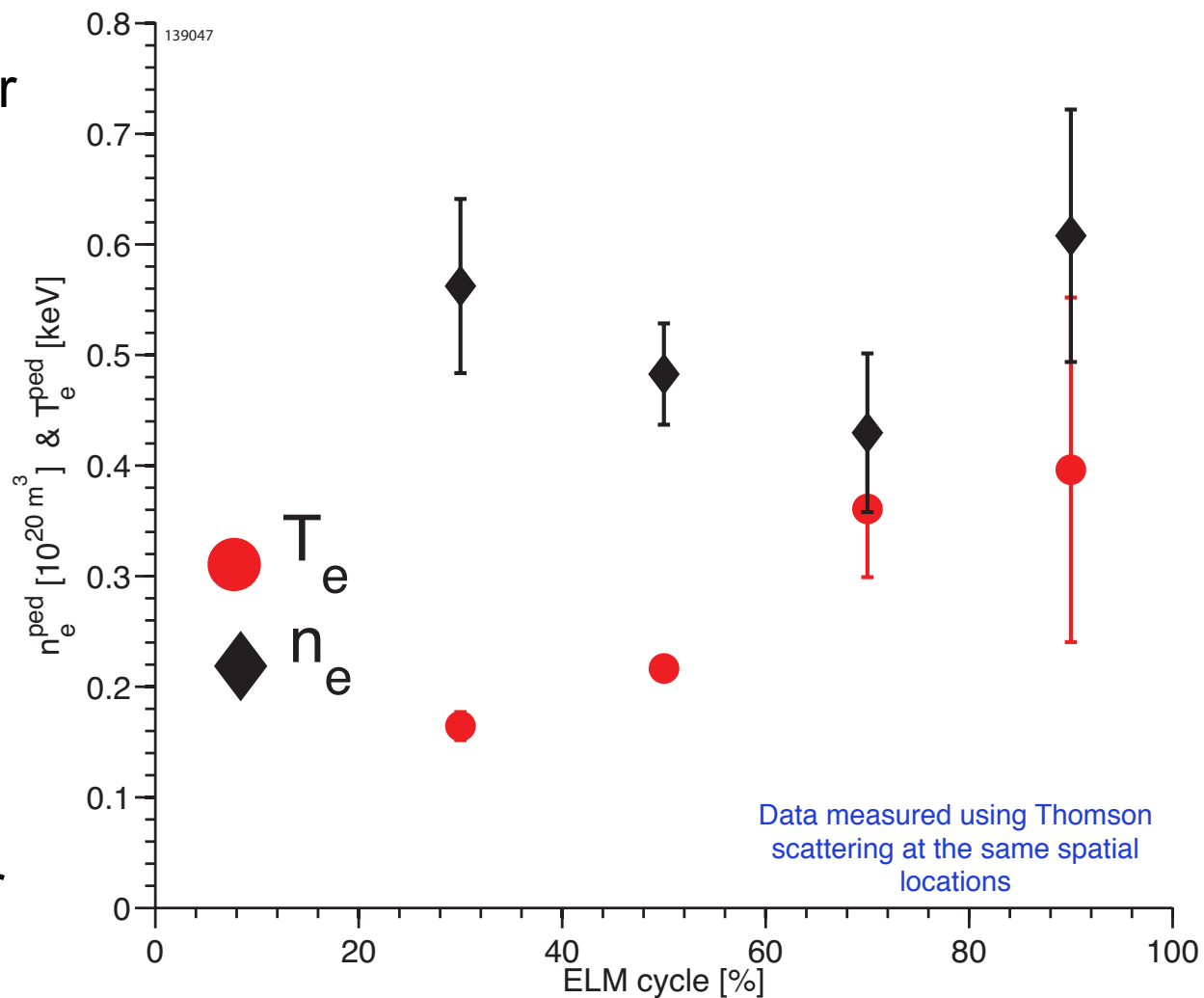
$$N(\psi) = A \tanh\left(\frac{\psi_{sym} - \psi}{\psi_{width}}\right) + offset$$

R. Groebner and T. Osborne PoP (1998)

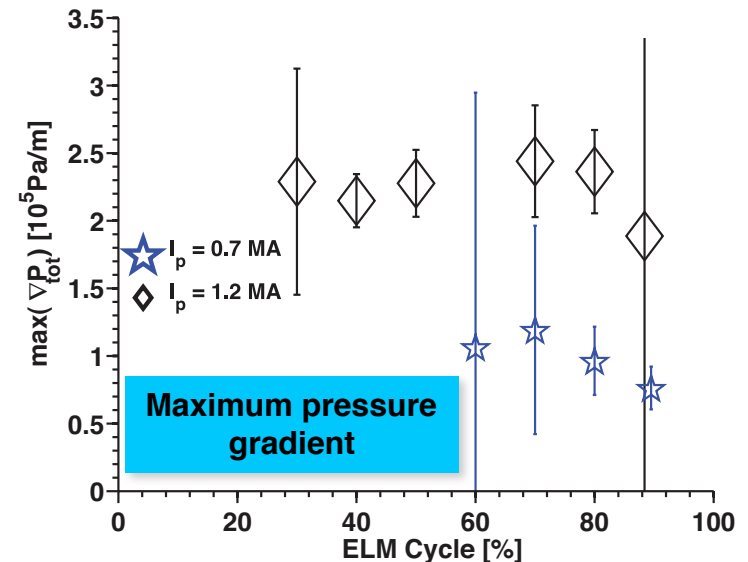
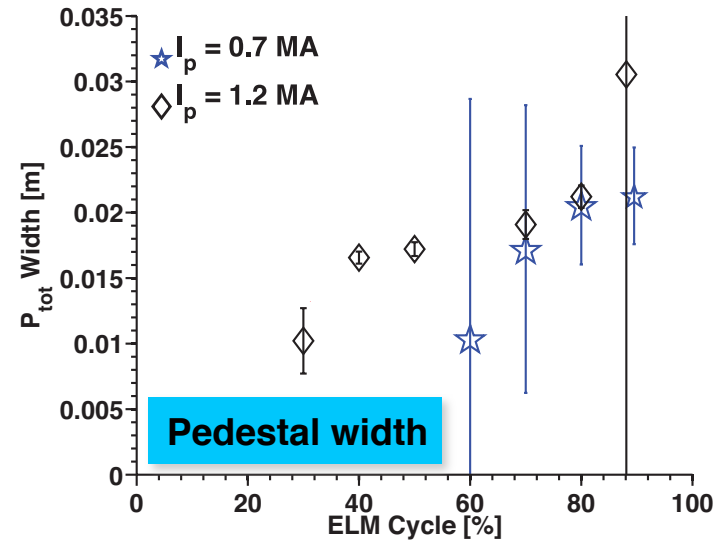
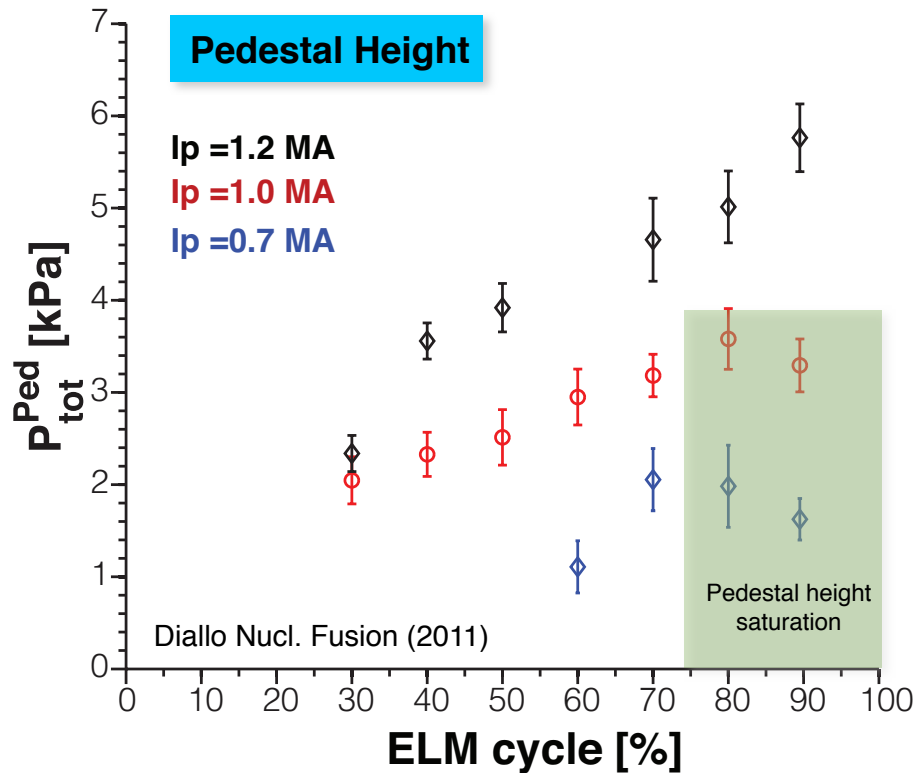
- Ion profiles are spline fitted

# Temperature pedestal height increases during the ELM cycle while the density pedestal show no convincing trend

- More than a factor of two increase in pedestal temperature
- Density pedestal is much less sensitive to the ELM cycle
- Heat and particle evolutions appear to be decoupled

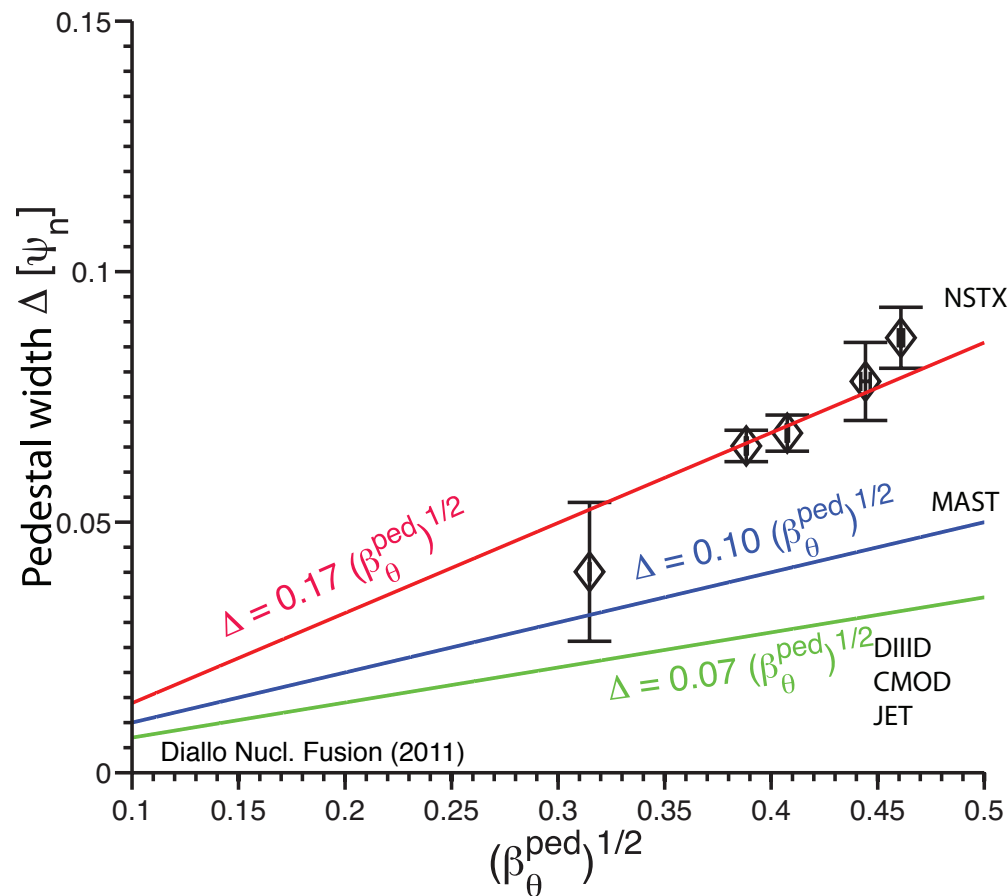


# Pedestal width and height progressively increase during ELM cycle but the peak pressure gradient remains clamped



- Pedestal height increases by a factor  $\leq 3$ 
  - Height scales with  $I_p$
- Pedestal width increases independently of  $I_p$
- Gradient is clamped early in ELM cycle

# Measured pedestal pressure width scales with $\sqrt{\beta_\theta}$



$$\beta_\theta^{ped} = 2\mu_0 P_{ped} / B_\theta^2$$

- Good description of the width scaling over multiple machines (DIII-D, CMOD, JET, MAST)

Groebner, NF, (2009)

Kirk, PPCF, (2009)

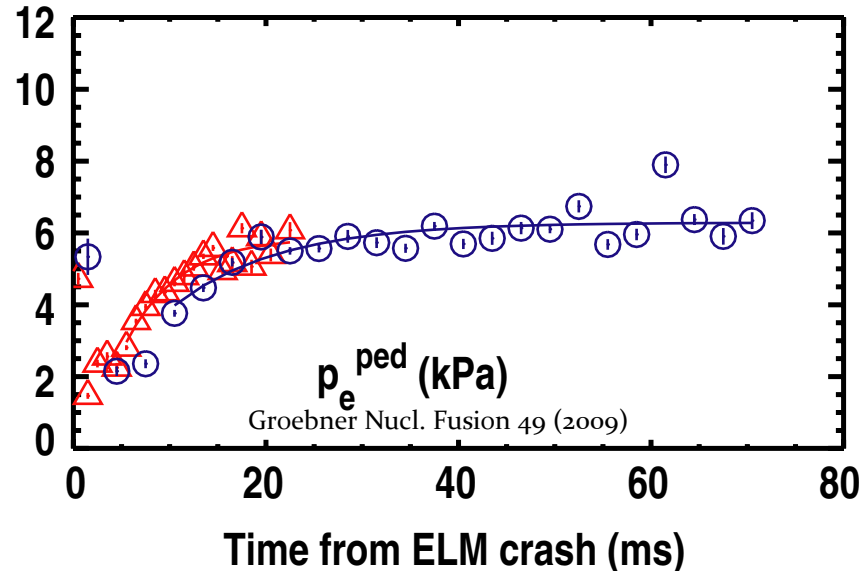
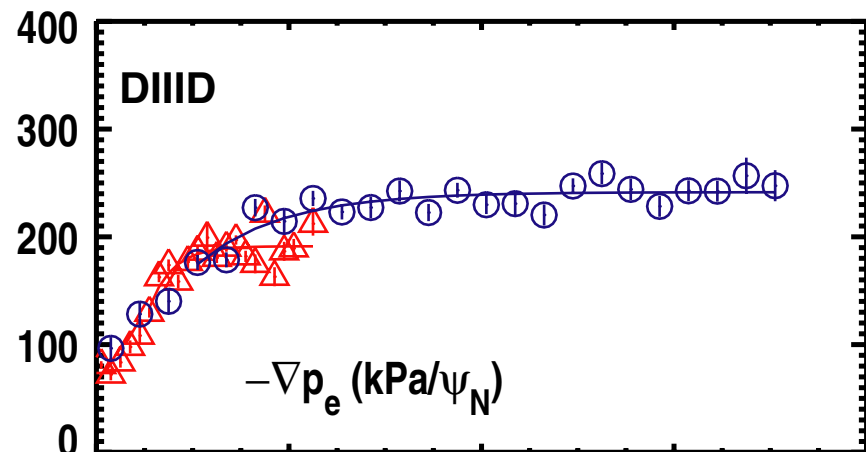
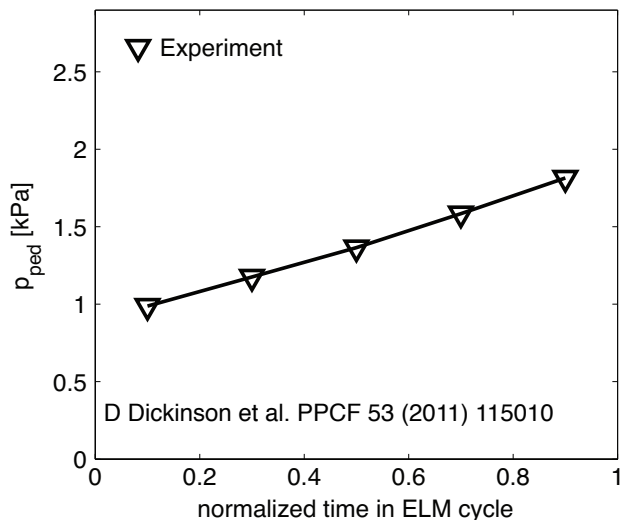
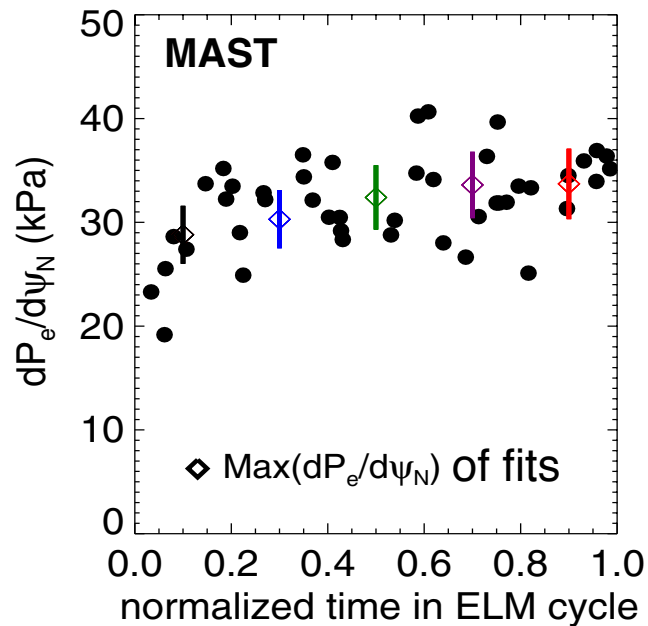
Beurkens, PoP, (2011)

Walk, to be published, Nucl. Fusion (2012)

- In NSTX, the observed width is larger than conventional tokamaks
  - NSTX pedestal width is 1.7 and 2.4 larger than MAST and DIII-D respectively

- Larger width in NSTX could sustain significant edge current

# Saturation of the gradient is ubiquitous across devices, but different trends in pedestal height evolution are observed

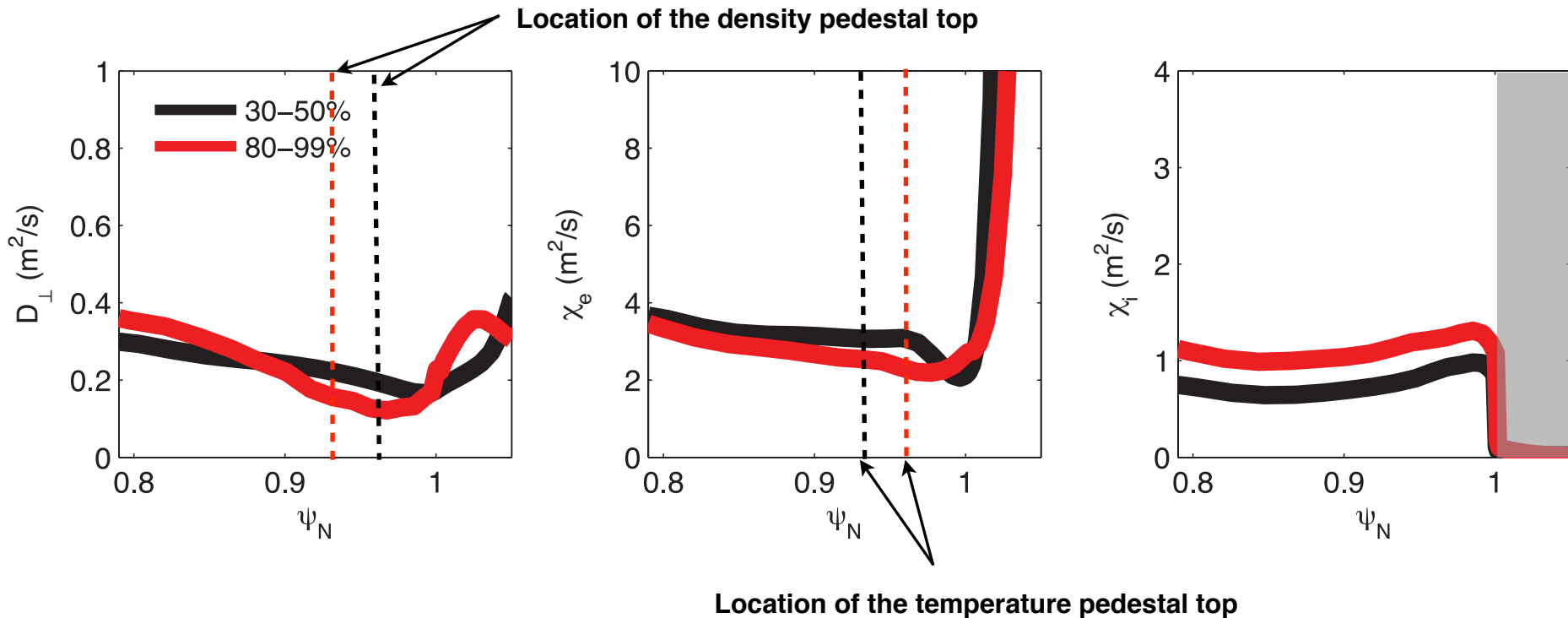




# Characterize the pedestal structure evolutions and pedestal turbulence to elucidate the role of transport setting the pedestal structure

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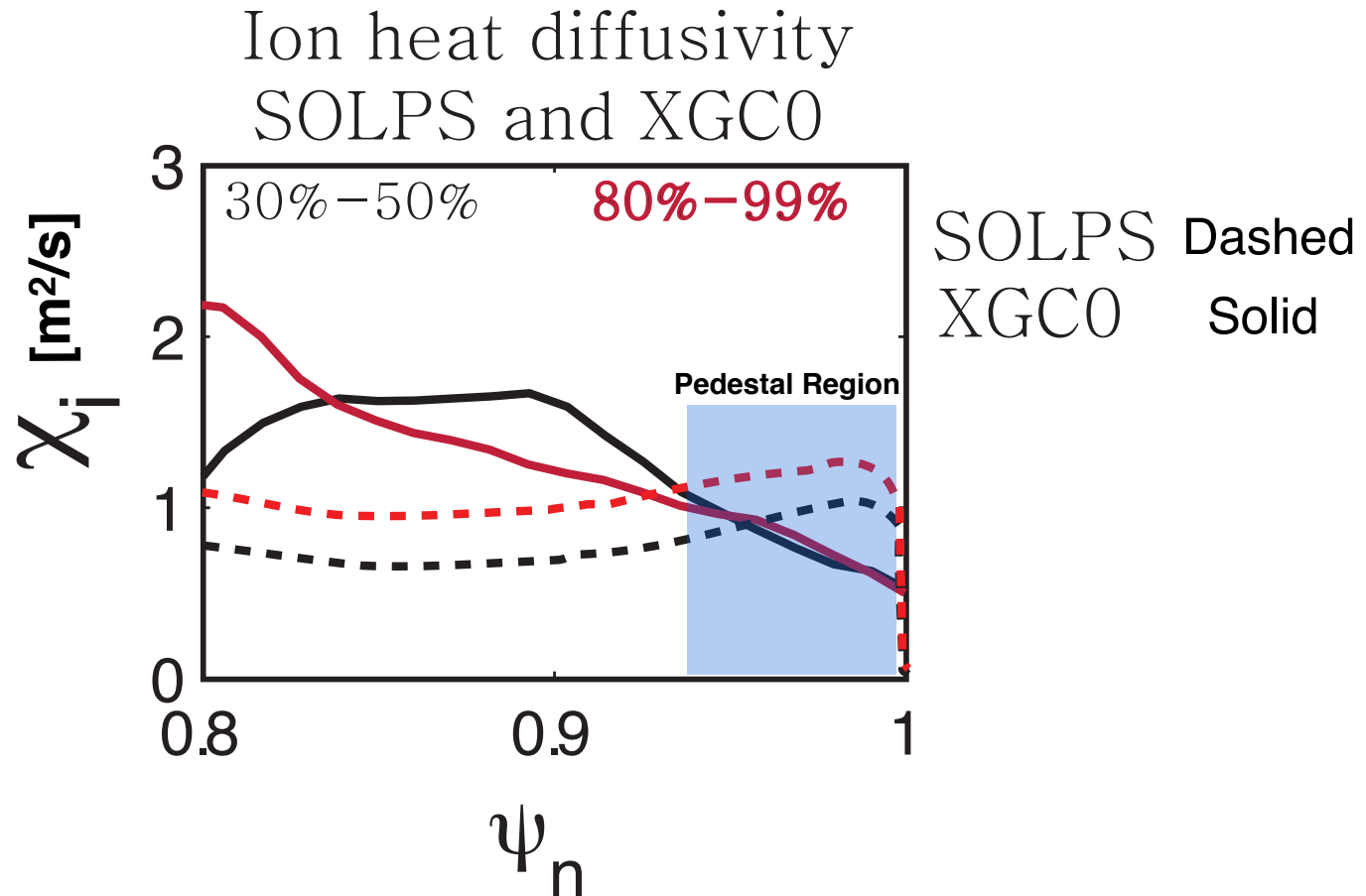
# 2D SOLPS\* modeling shows modest variations of the transport coefficients at the pedestal top



- Particle flux at the pedestal top is insignificant during the inter-ELM
- No clear trend for the electron heat flux
- Ion heat flux becomes larger later in the ELM cycle

\*R. Schneider et al, Contrib. Plasma Phys. 46 (2006)

# Ion heat diffusivity comparison: SOLPS and XGC0



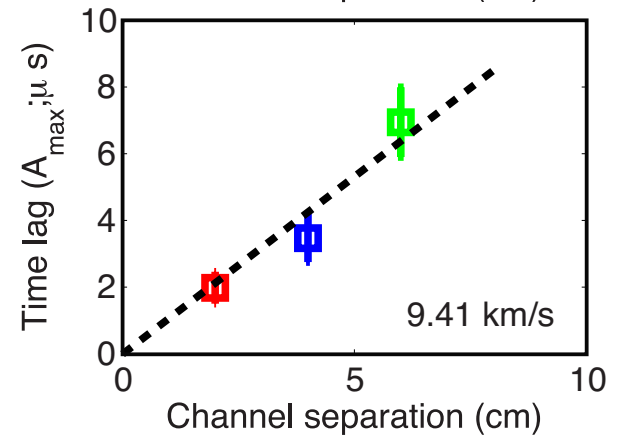
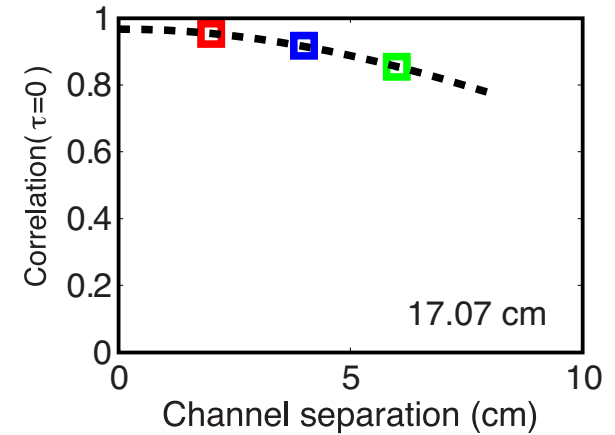
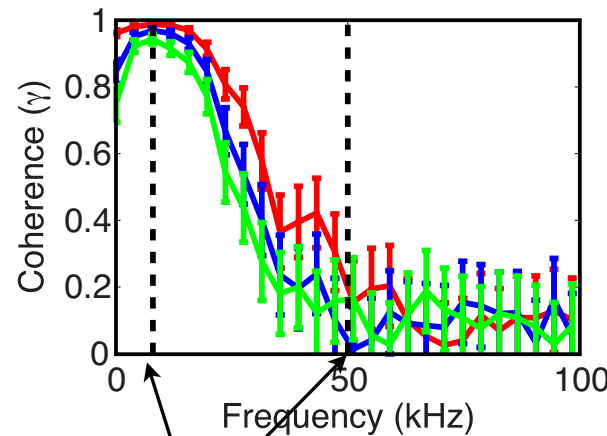
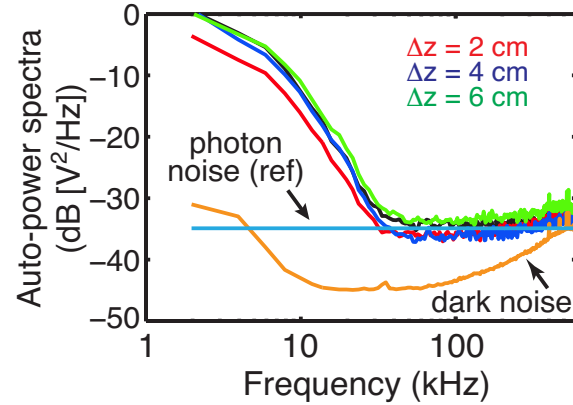
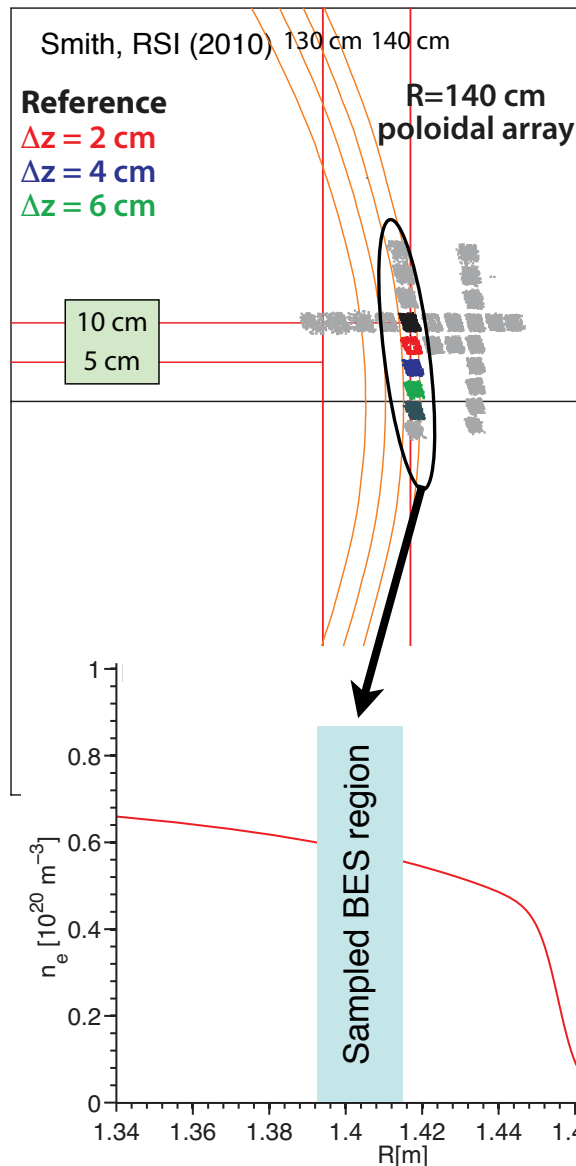
XGC0 computes the neoclassical transport parameters

- Neoclassical ion diffusivity remains unchanged during the inter-ELM phase in the pedestal region
- In the pedestal region SOLPS shows larger than neoclassical ion diffusivity

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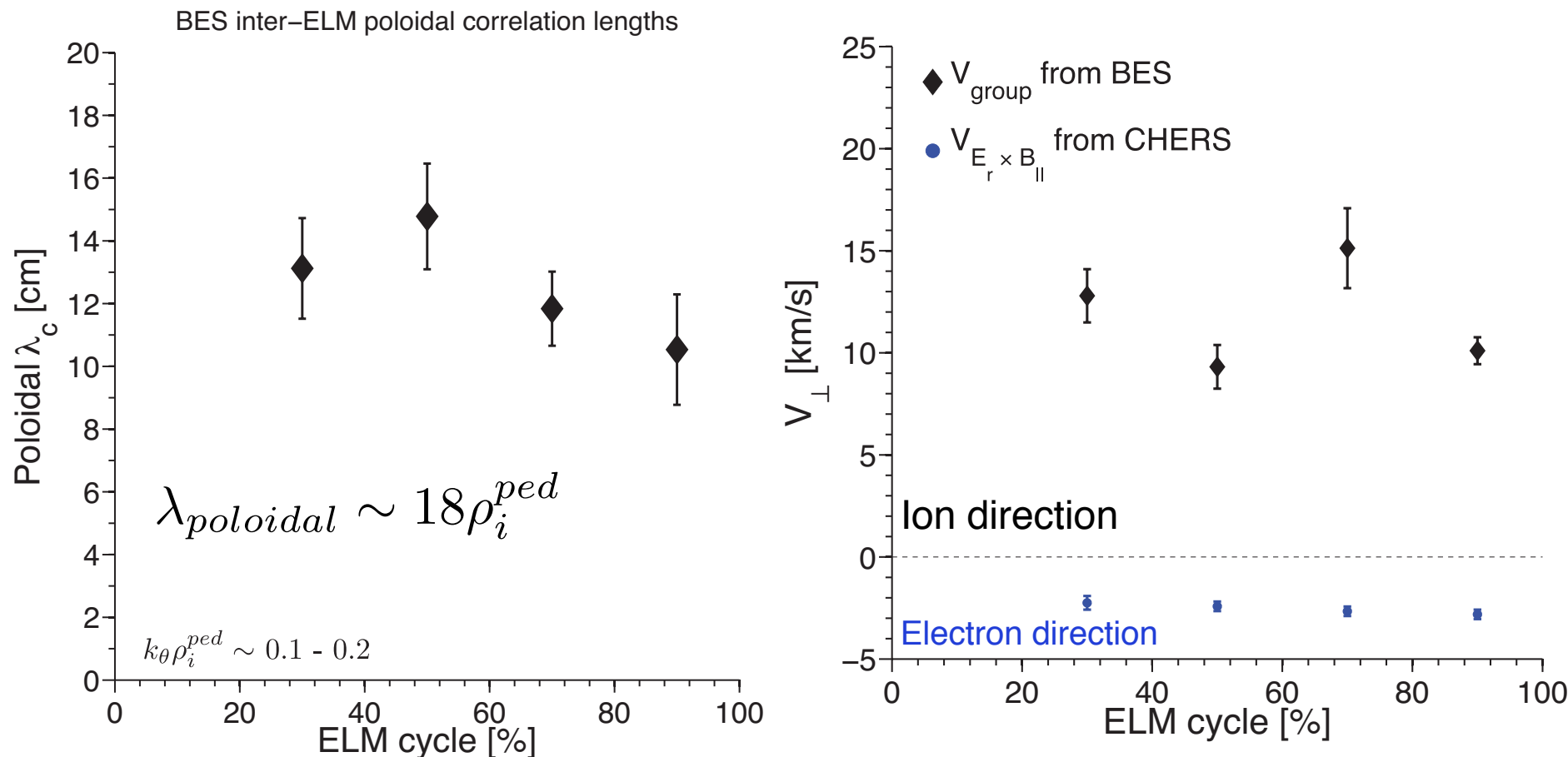
# BES yields measurements of the poloidal correlation at the density pedestal top



D. Smith: this meeting

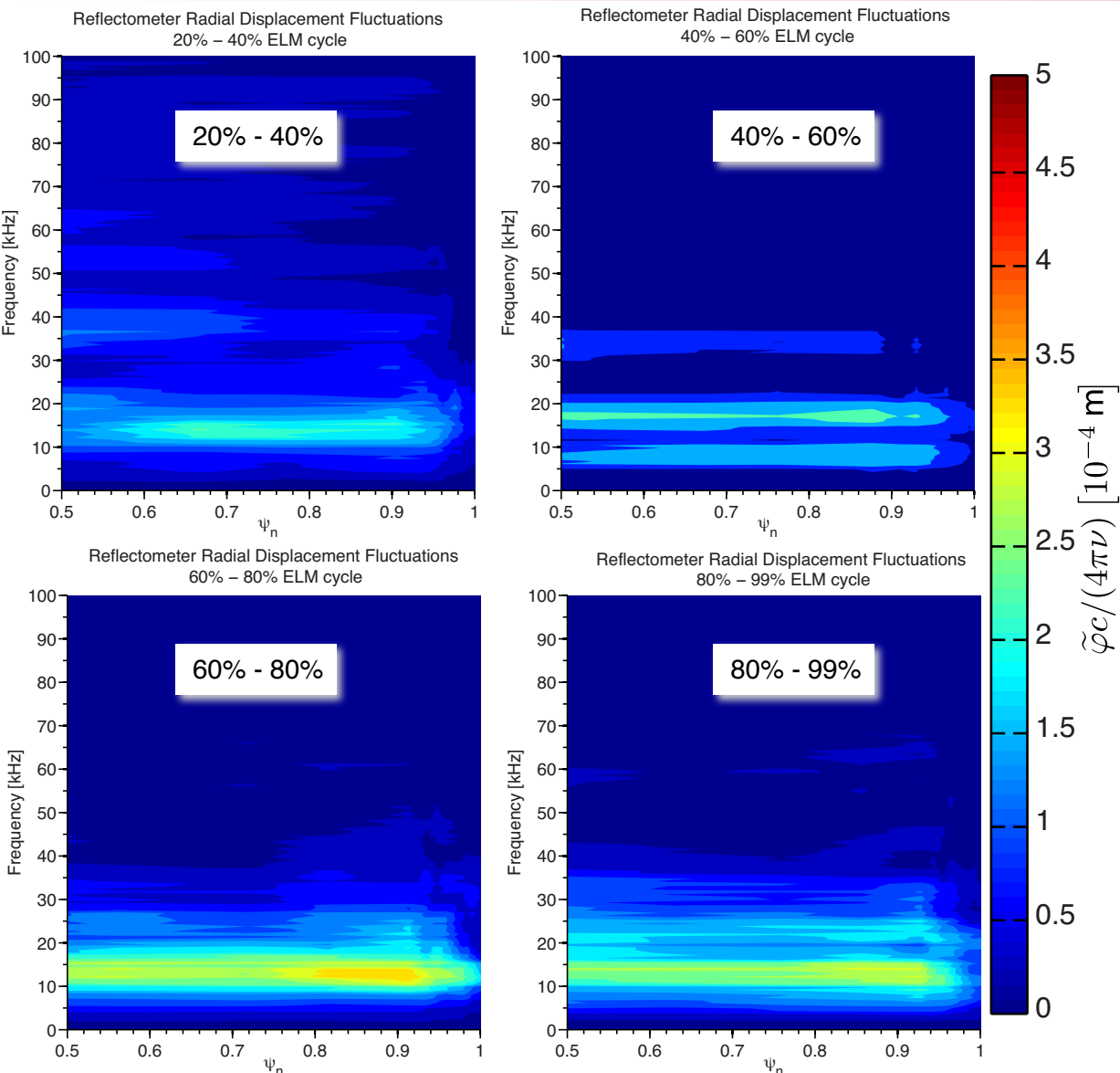
Shot 139047, Time 624-640 ms, Ref chan 3, Channels 11 12 13, freq = 8-50kHz

# BES provides measurements of the poloidal correlation length and poloidal velocity



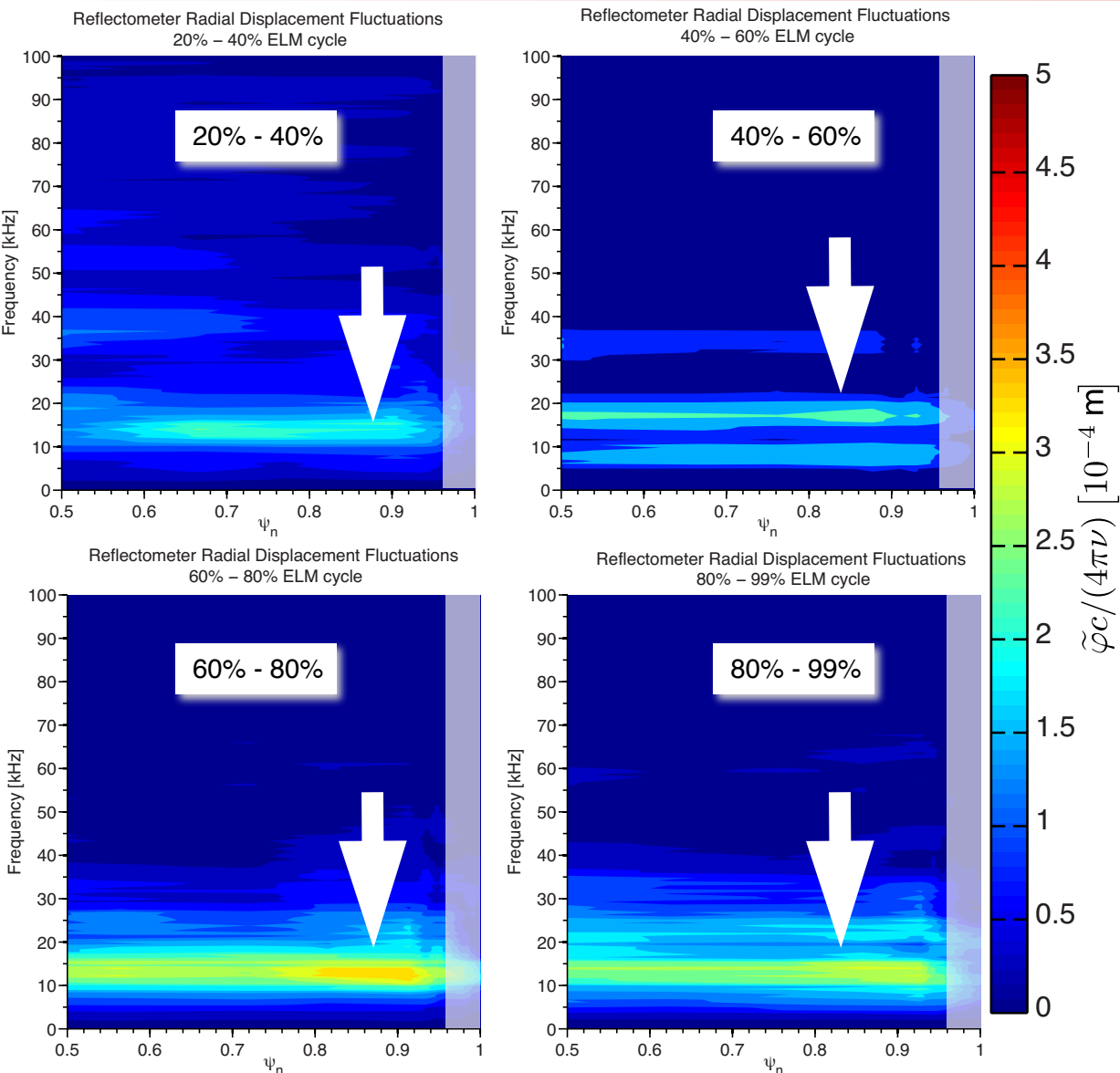
- Modest change in poloidal correlation length during the inter-ELM phase
  - Poloidal correlation length corresponds to toroidal mode number ( $rk_\theta/q$ )  $n = 2 - 3$
- Measurements show ion scale fluctuation in the pedestal top

# Evolution of the radial displacement power spectra indicates an increase of the fluctuation level during the last 40% of ELM cycle



- Radial displacement induced by fluctuations
  - Existence of broadband fluctuations centered around 12 kHz
- Increase of the overall mode amplitude late in ELM cycle
- Caution: Localization of the mode is difficult to assess as it could be due to the density scale length

# Evolution of the radial displacement power spectra indicates an increase of the fluctuation level during the last 40% of ELM cycle



● Radial displacement induced by fluctuations

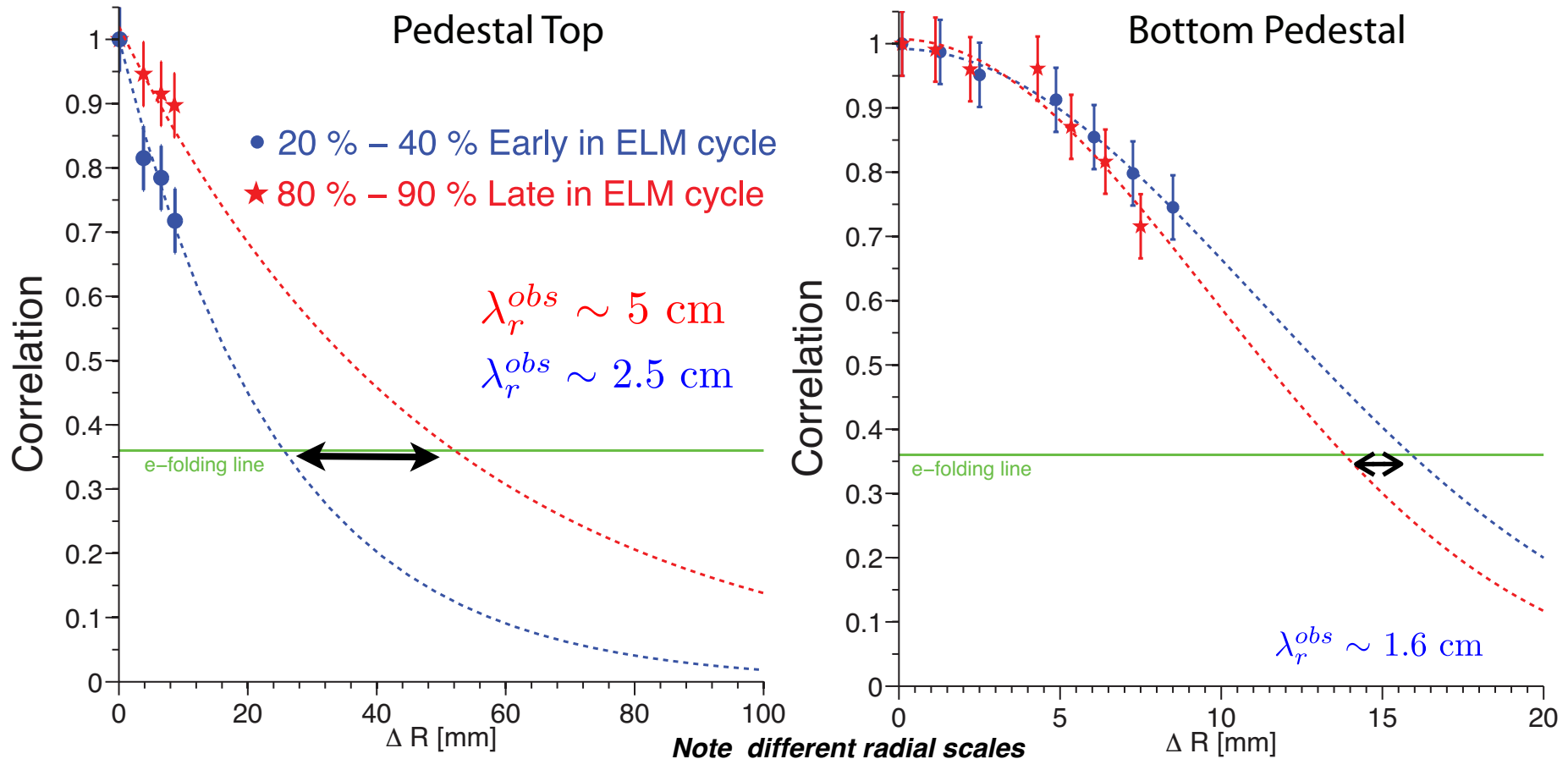
- Existence of broadband fluctuations centered around 12 kHz

● Increase of the overall mode amplitude late in ELM cycle

● Caution: Localization of the mode is difficult to assess as it could be due to the density scale length



# Radial correlation lengths at the pedestal top and steep gradient during the inter-ELM phase

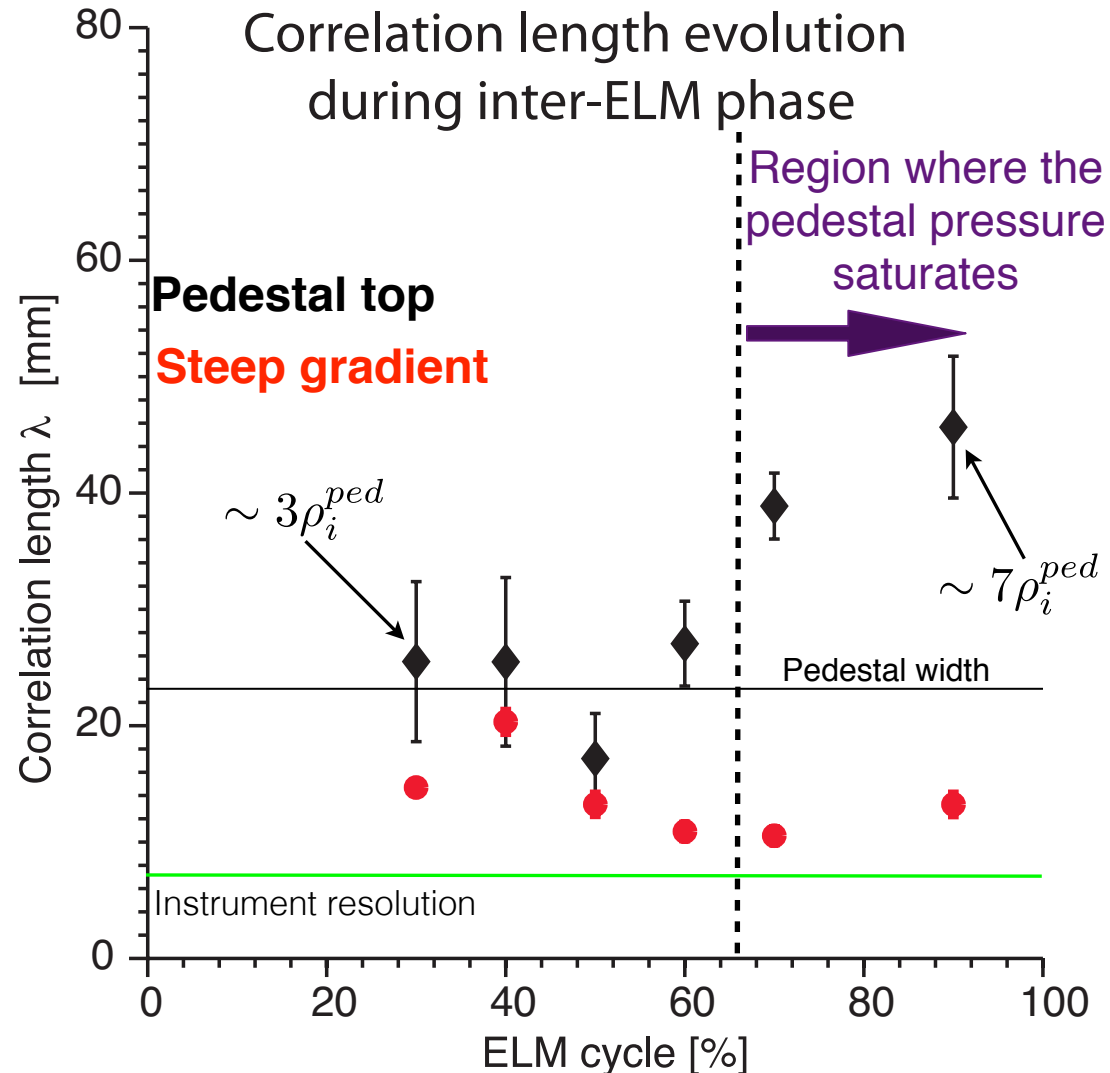


Radial correlation is valuable as it provides the spatial size of a turbulent perturbation (i.e. eddy).

Pedestal top correlation length is larger than that of the region of steep density gradient

# Radial correlation length evolution depends on location inside pedestal region

- Radial correlation length increases at the pedestal top
  - A factor of 2 increase during the last 50% of ELM cycle
  - Presumably due to the increase in radial displacement fluctuations
- Steep gradient correlation length is unchanged
- Quantify the geometric effects on the measured correlation?



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- Fluctuations characterizations during the inter-ELM phase
- Preliminary  $\delta f$  XGC1 calculations and comparison with observations**
- Summary/Discussion

# Preliminary simulations using XGC1 are performed for cases during the last part of the ELM cycle

## $\delta f$ mode in XGC1

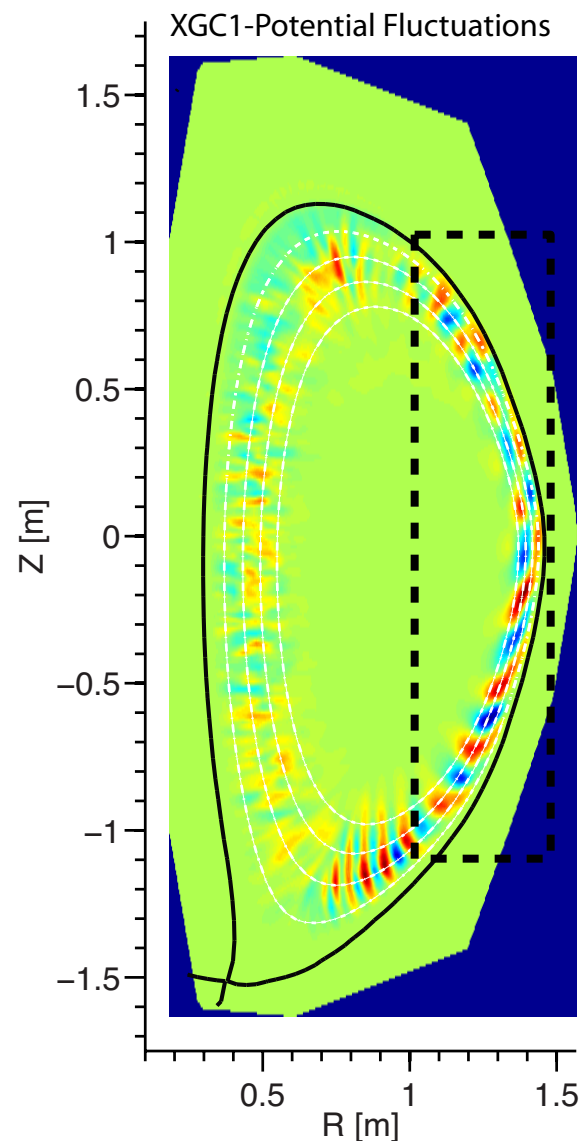
- 200 x 60 spatial grid
- simulation box up to  $\psi_n \sim 0.95$  to include the unstable region

Collisions and flows are not included in this simulation

Adiabatic electrons

Probing the fully nonlinear phase of the simulations

Full-f simulation in the whole edge with kinetic electrons, flows, and collisions will be performed.



# Preliminary simulations using XGC1 are performed for cases during the last part of the ELM cycle

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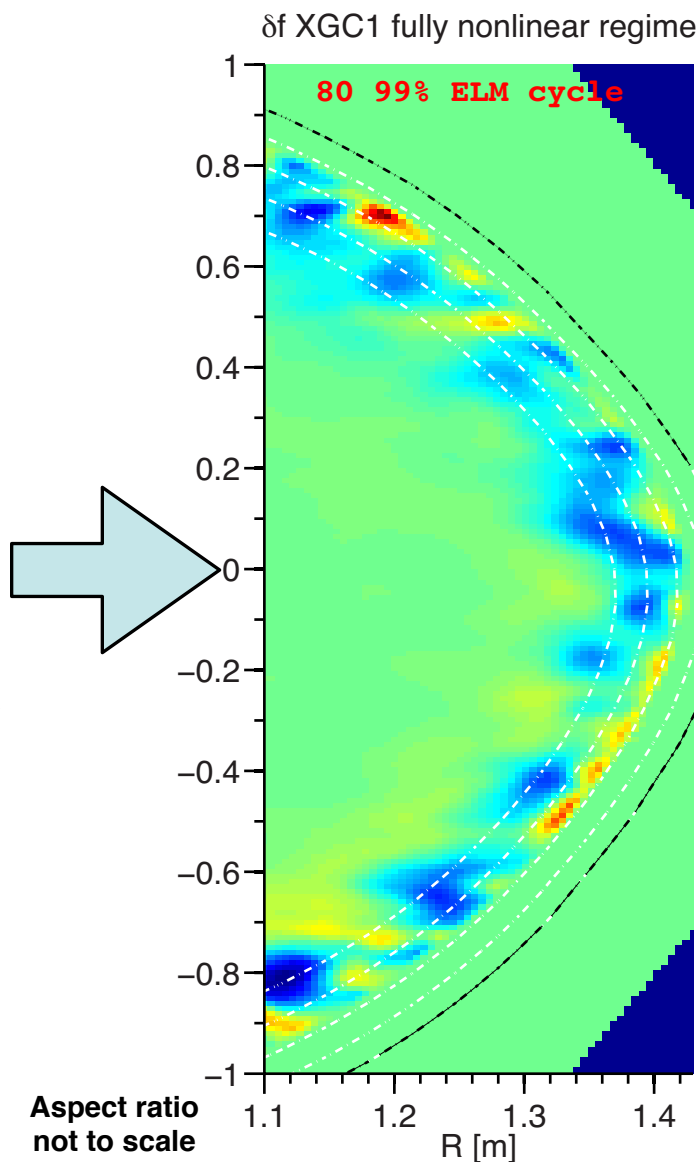
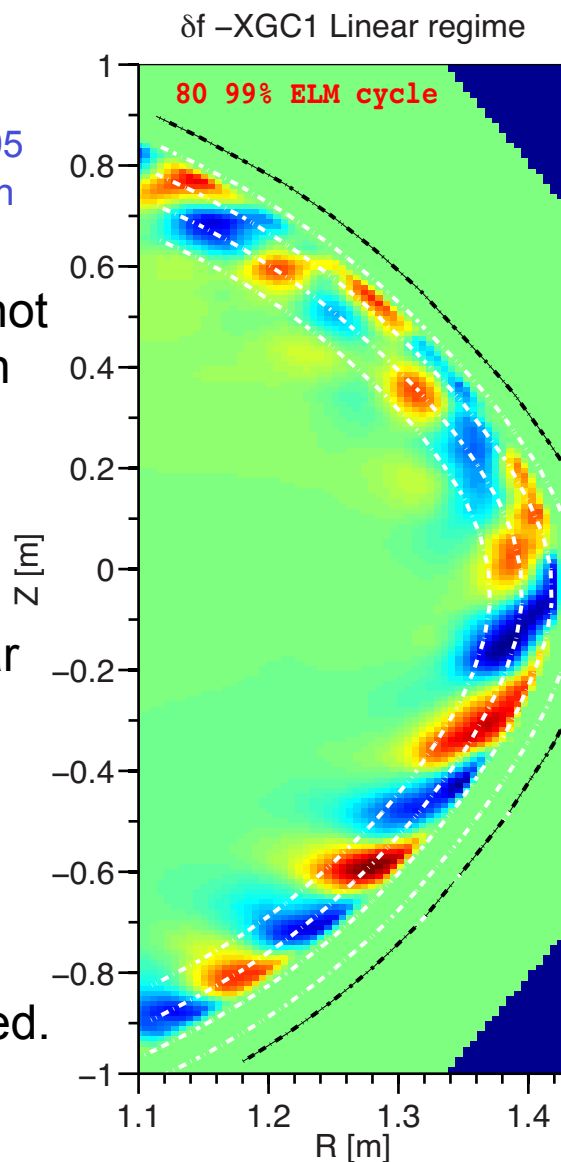
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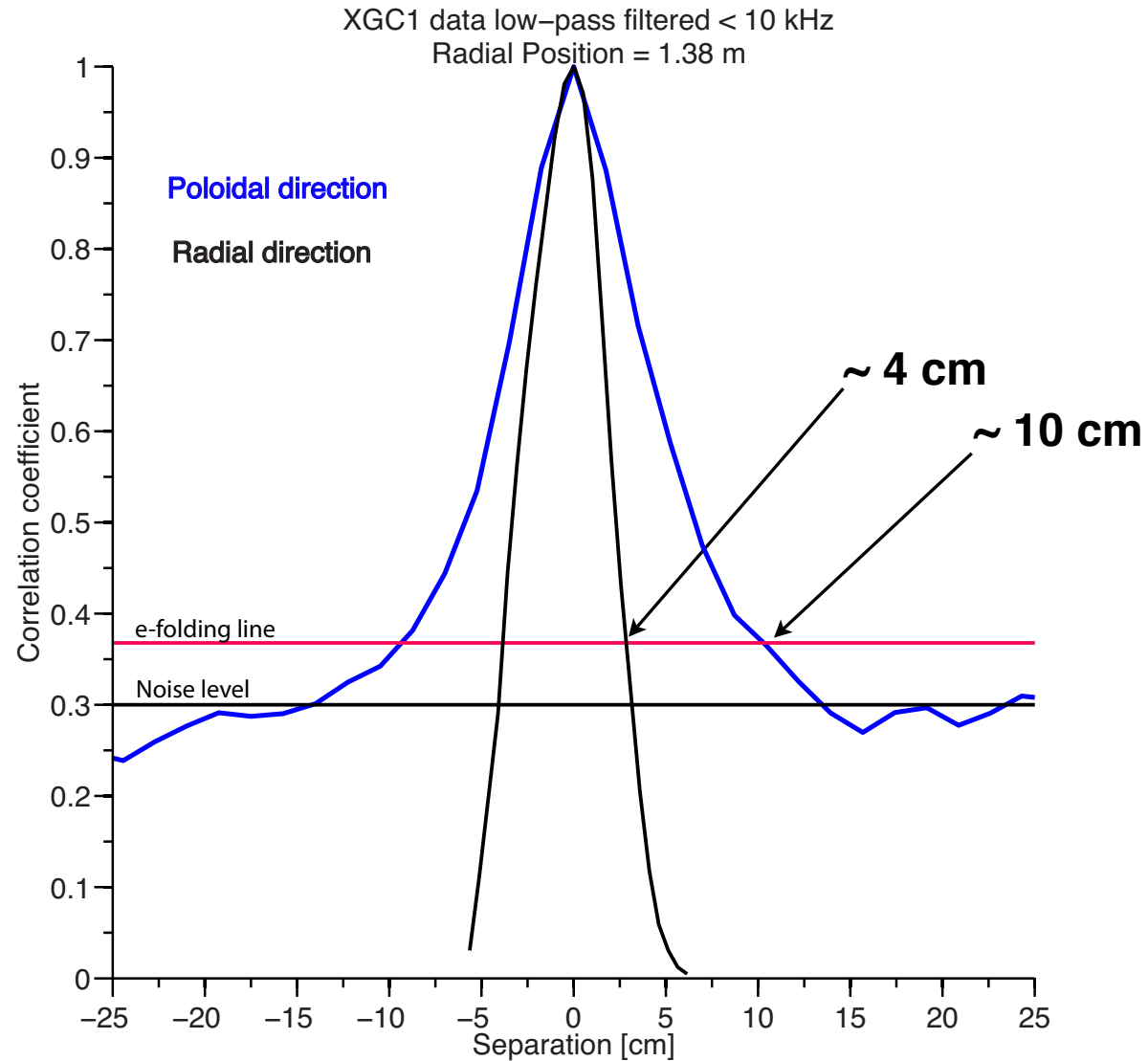
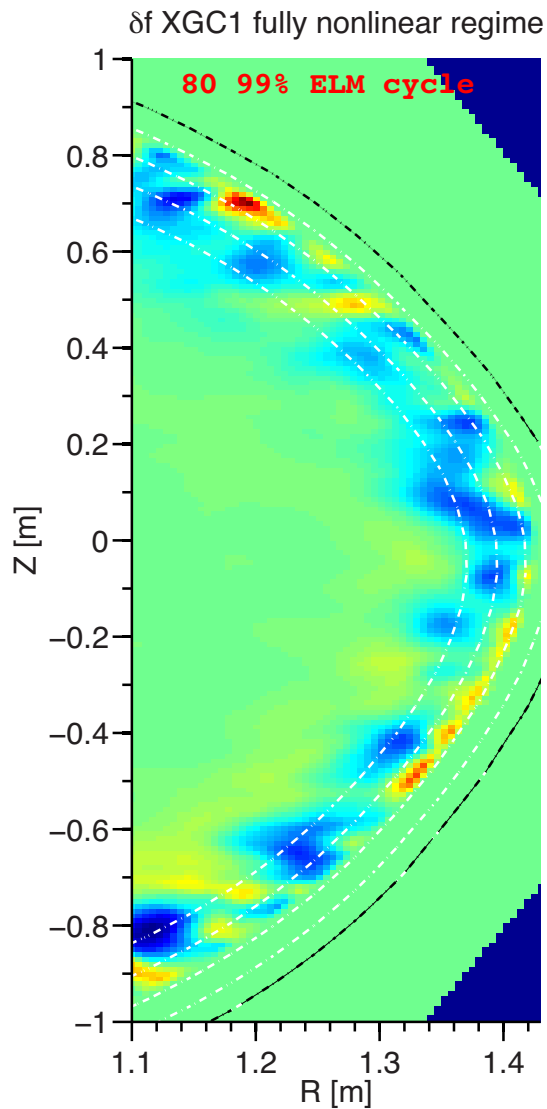
Adiabatic electrons

Probing the fully nonlinear phase of the simulations

Full-f simulation in the whole edge with kinetic electrons, flows, and collisions will be performed.



# Preliminary simulations from XGC1 show localized fluctuations with experimental level radial and poloidal correlation lengths



# Summary (I)

Using ELMy discharges, we observed during the ELM cycle:

- Continuous increase of the electron temperature and much less variation in electron density and pressure build up and at time saturation prior to the ELM onset
- Consistent with barrier expansion phenomenology observed in other tokamaks
- Pressure gradient, however, is clamped during most of the ELM cycle

Interpretive transport analysis using SOLPS, and XGC1

- No convincing correlation between the electron heat flux and ELM cycle
- Ion heat flux is larger in the pedestal region late in the ELM cycle
  - In addition, Ion heat flux larger than neoclassical estimates from XGC0
- Particle flux throughout remains difficult to assess

# Summary (II)

Characterization of the fluctuations during the inter-ELM phase

- BES and reflectometry confirm ion scale turbulence  $0.2 \leq k_{\perp} \rho_i \leq 0.7$
- Poloidal correlation is larger than radial correlation length

Preliminary simulations with XGC1 are performed

- Used in  $\delta f$  mode in XGC1
- Thus far, simulation results correlation lengths measurements agree with experimental observations
  - » Most unstable mode is ITG in simulation: identification of turbulence in experiment, however, will require full-f XGC1.
  - » Addition of electrons and flows will give a better sense of the heat transport responsible for clamping the pedestal gradient

More work: A continuation and expansion of these comparisons to simulations with the ultimate goal of identification of types of turbulence

- Further 3D analysis of reflectometry are underway to verify the 2D approximation initially performed
  - How well the diagnostic determine the true radial correlation length
- Extend the simulation to full-f mode using XGC1 and account for measured flows and add collision.



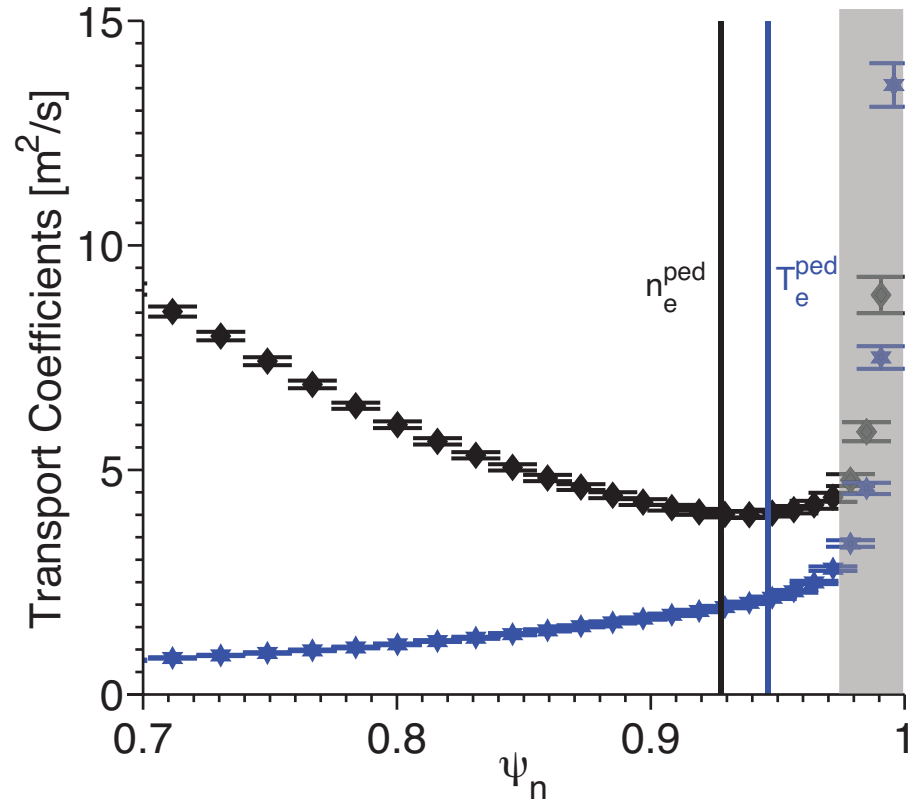
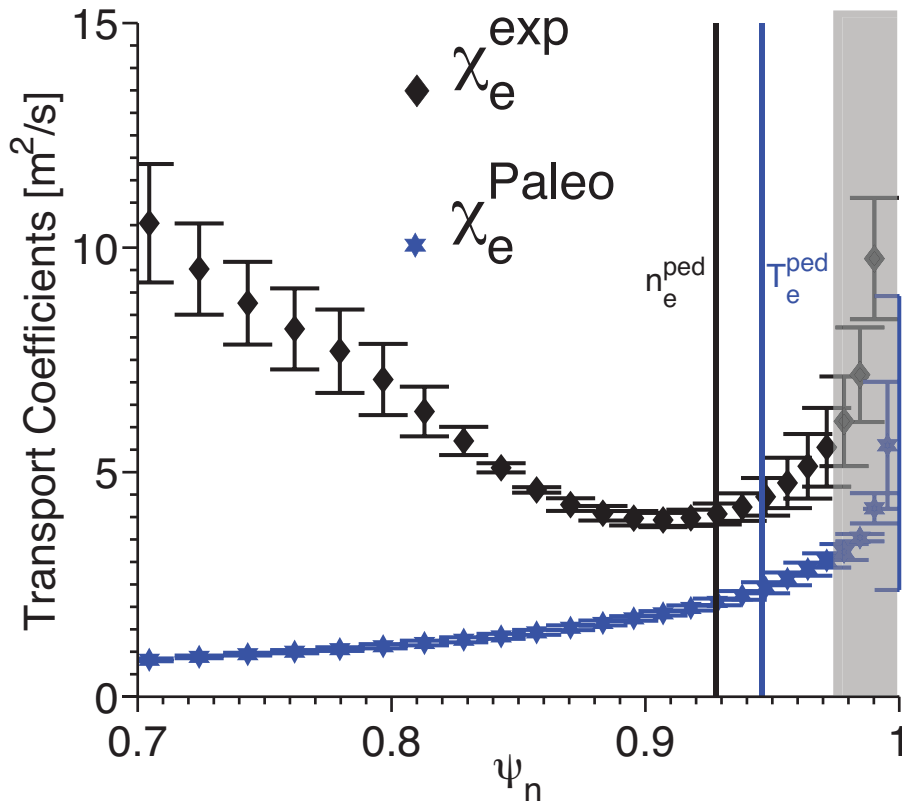
# Backup

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# Using TRANSP, no significant change in the electron heat diffusivity at the pedestal prior to the onset of ELM

Heat Diffusivities 30% – 50% ELM cycle

Heat Diffusivities 80% – 99% ELM cycle

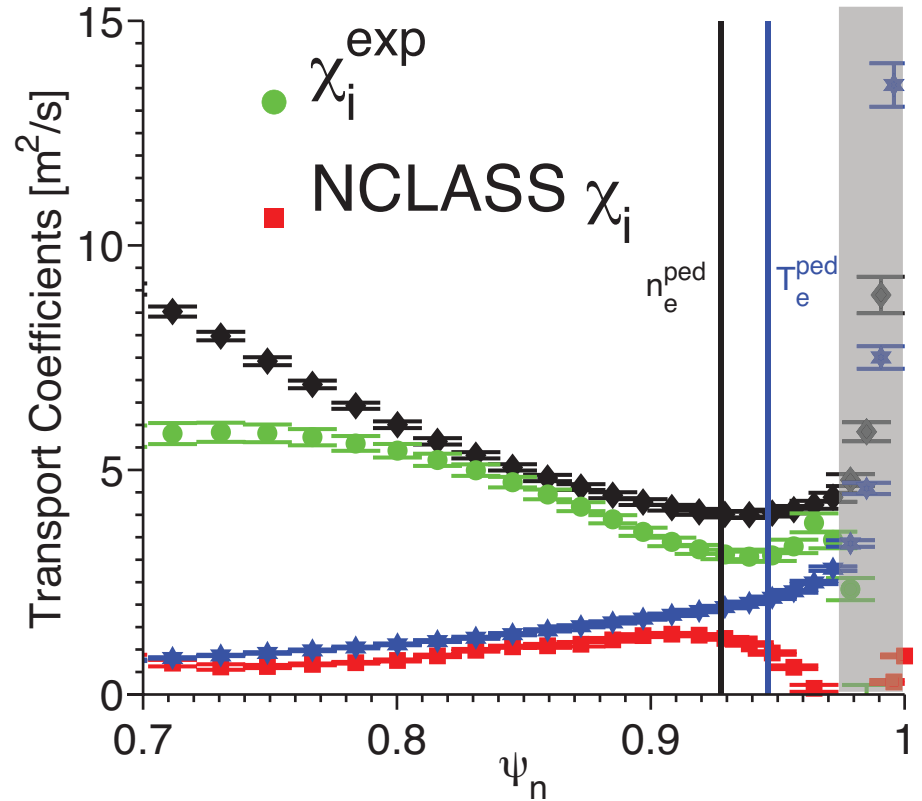
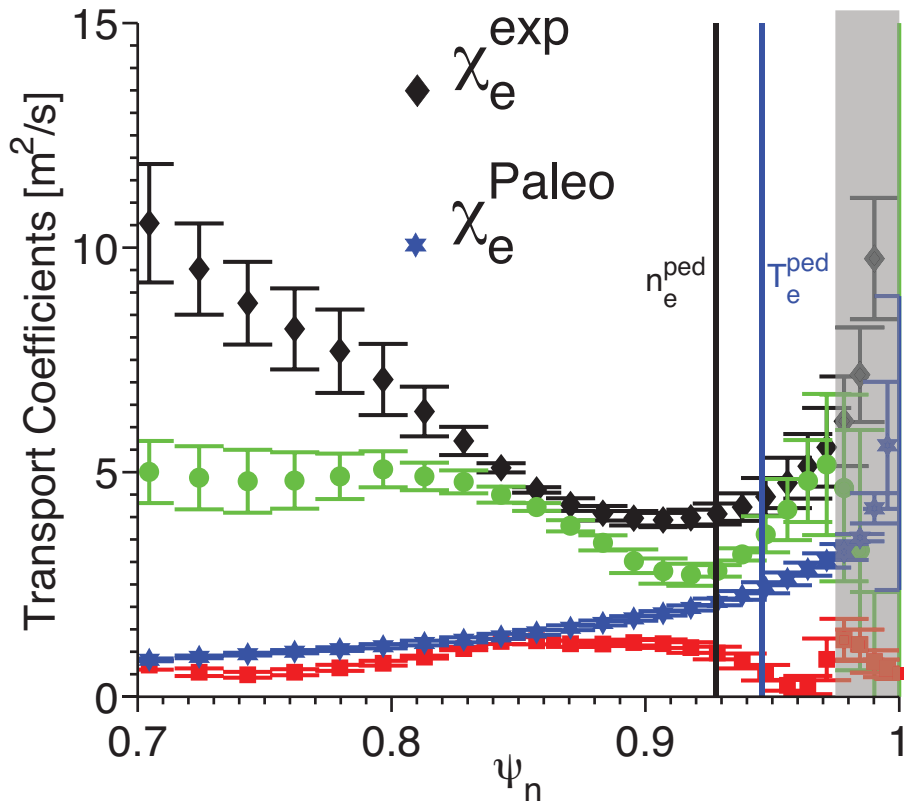


- Bottom half of the pedestal suggests that paleoclassical transport could be a major contributor to the electron heat transport, but not at the pedestal top.
- In the pedestal top, neoclassical heat diffusivities appear to overestimate the experimental ion  $\chi$ .

# Using TRANSP, no significant change in the electron heat diffusivity at the pedestal prior to the onset of ELM

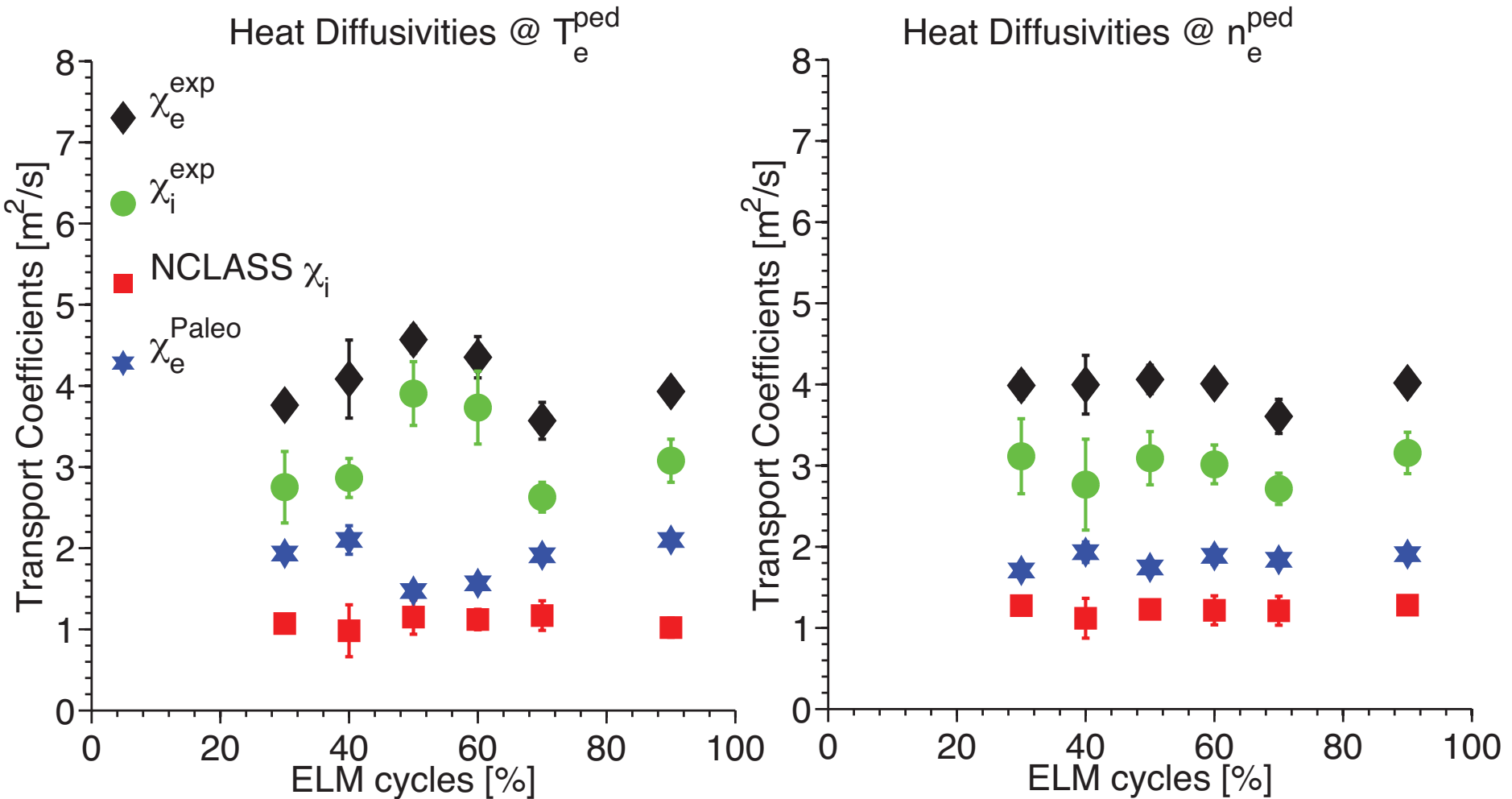
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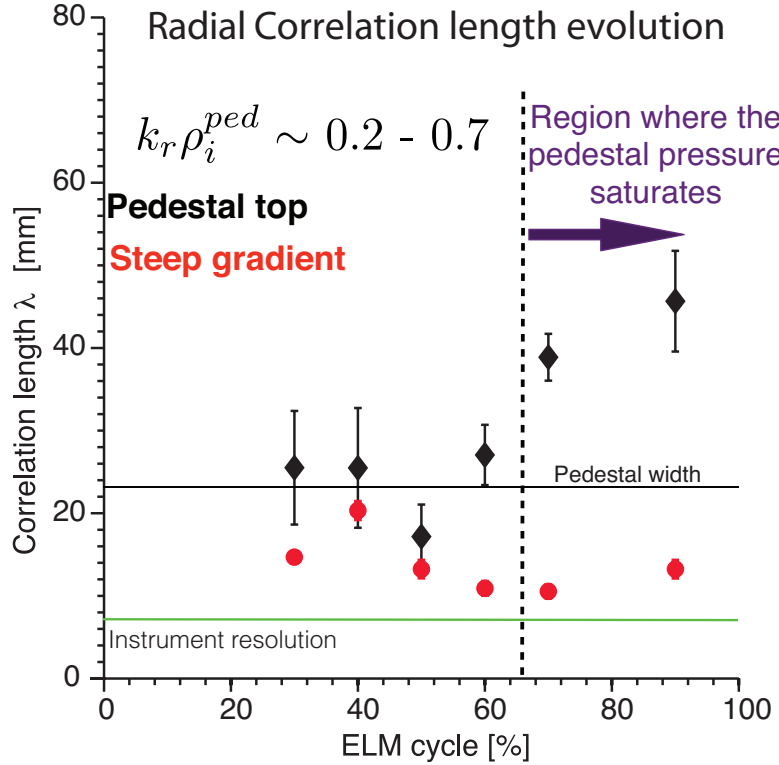
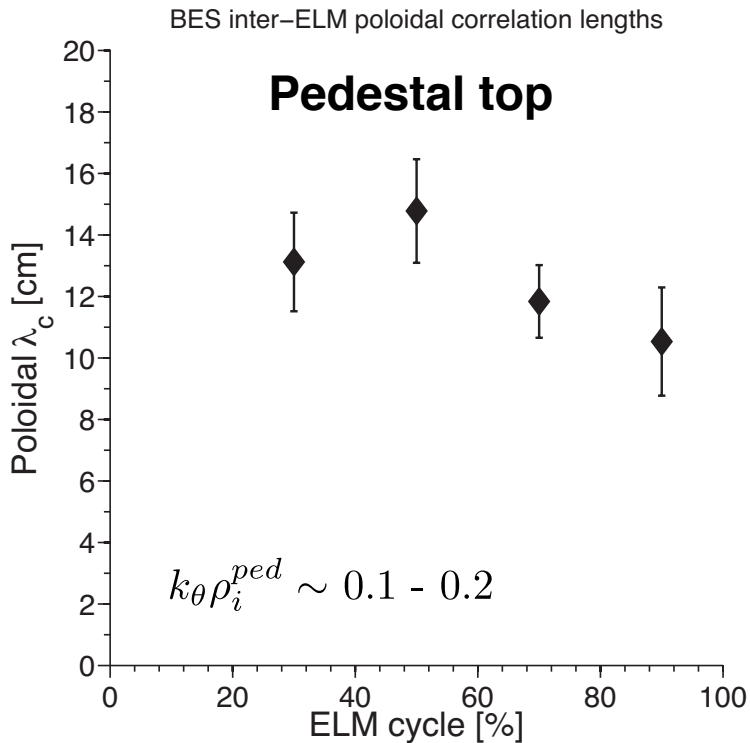
- Bottom half of the pedestal suggests that paleoclassical transport could be a major contributor to the electron heat transport, but not at the pedestal top.
- In the pedestal top, neoclassical heat diffusivities appear to overestimate the experimental ion  $\chi$ .

# The trends during the inter-ELM phase in heat diffusivities are difficult to unravel



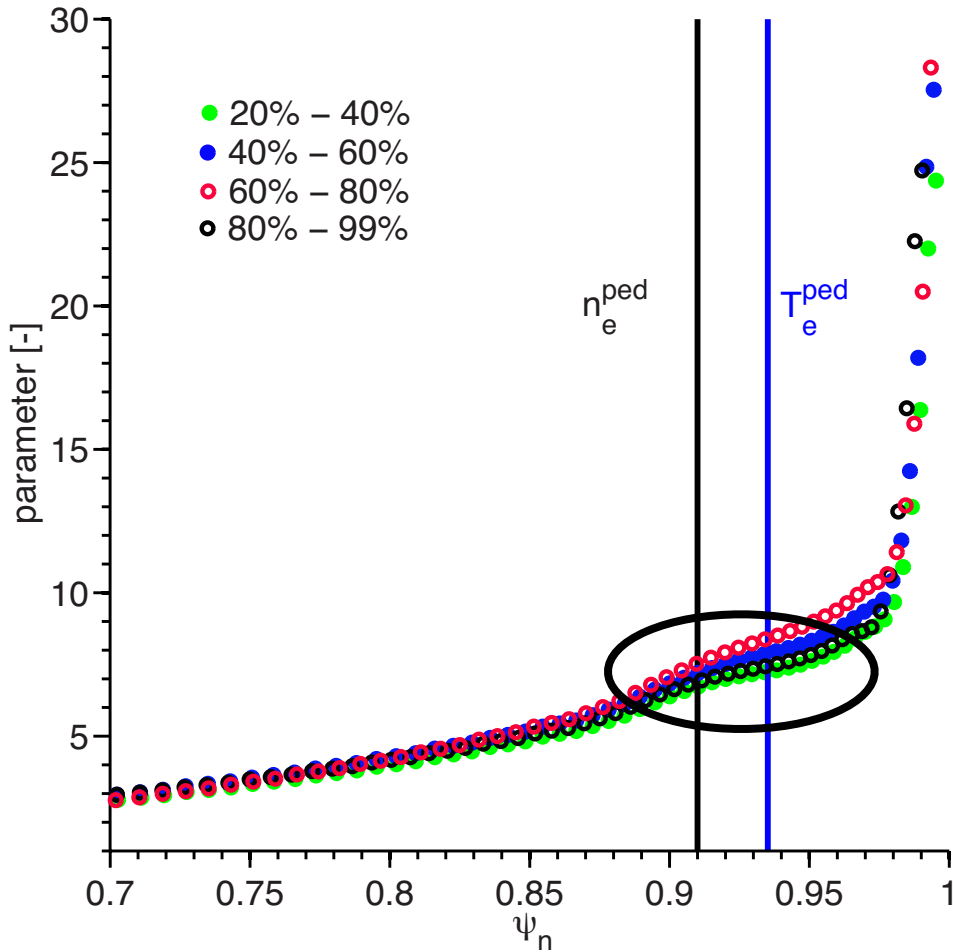
# Using both BES and correlation reflectometry, the inter-ELM spatial structure of fluctuations exhibit ion-scale microturbulence

- Strong anisotropy of the turbulence is observed during the inter-ELM phase
- Turbulence data suggests microturbulence with  $0.2 \leq k_{\perp} \rho_i \leq 0.7$  propagating in ion diamagnetic direction

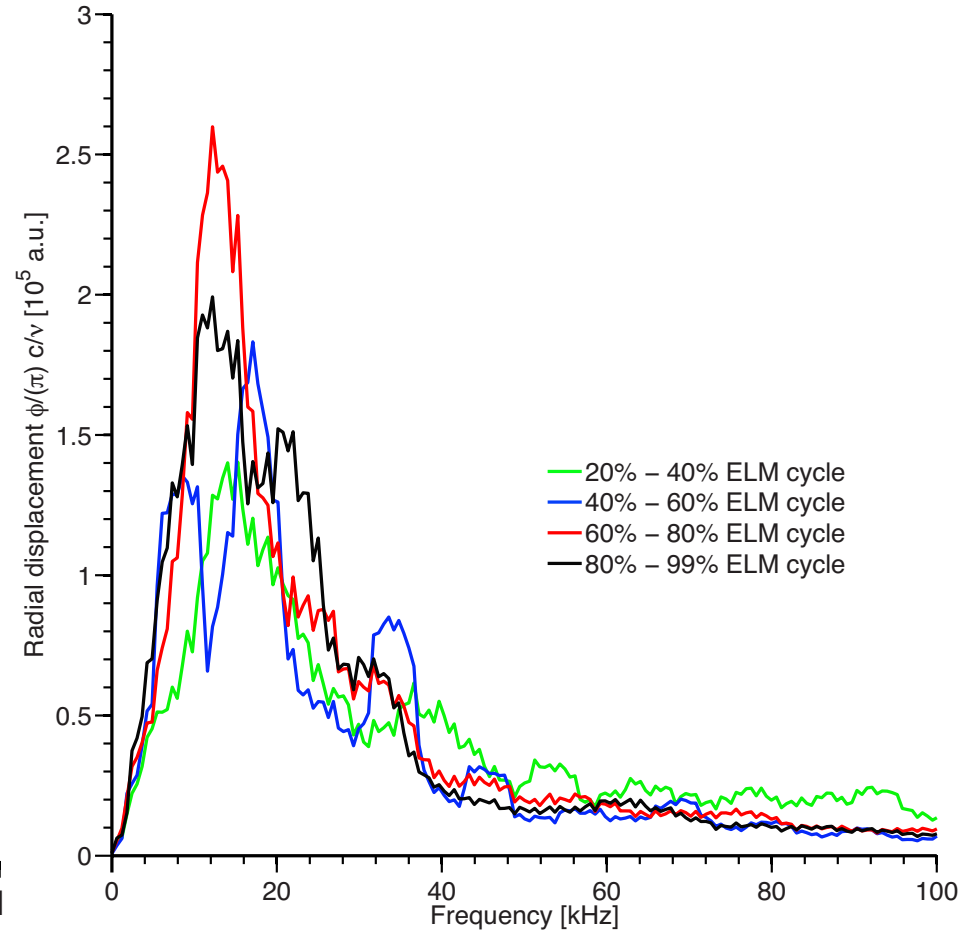


# Apparent correlation between the magnetic shear at the pedestal top and fluctuations amplitude increase during the ELM cycle

$S = q\partial_r q$  parameter during ELM cycle

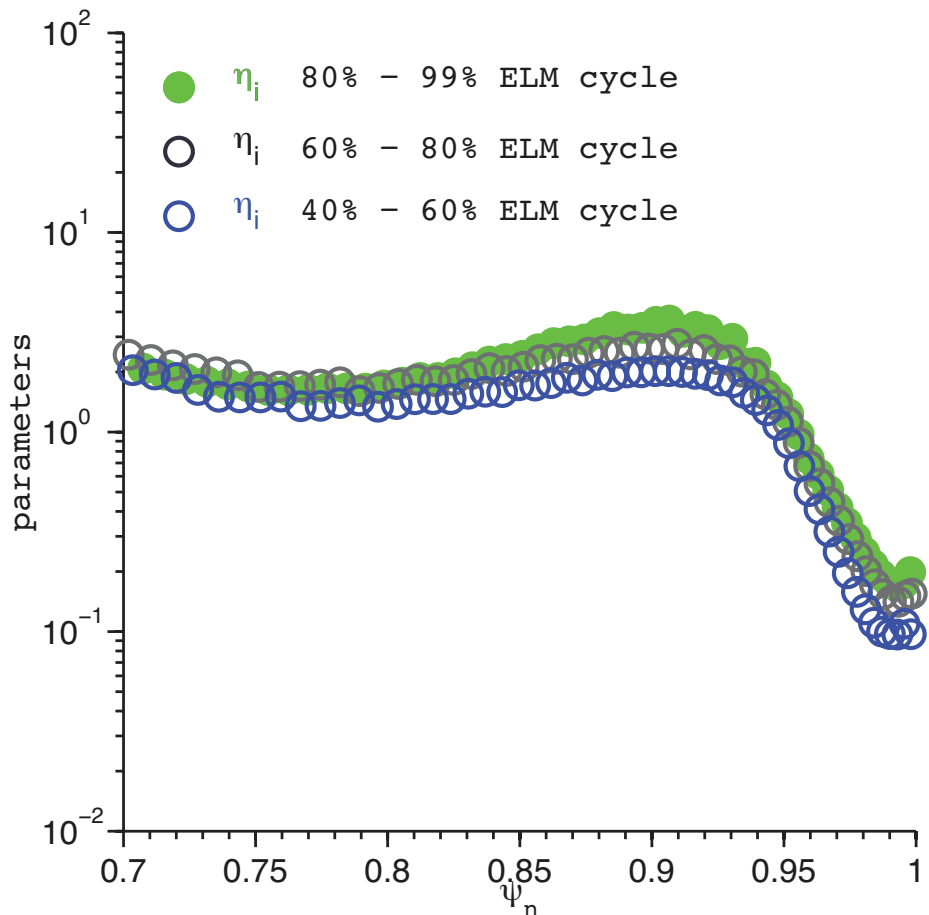


Radial averaging of displacement fluctuation spectra

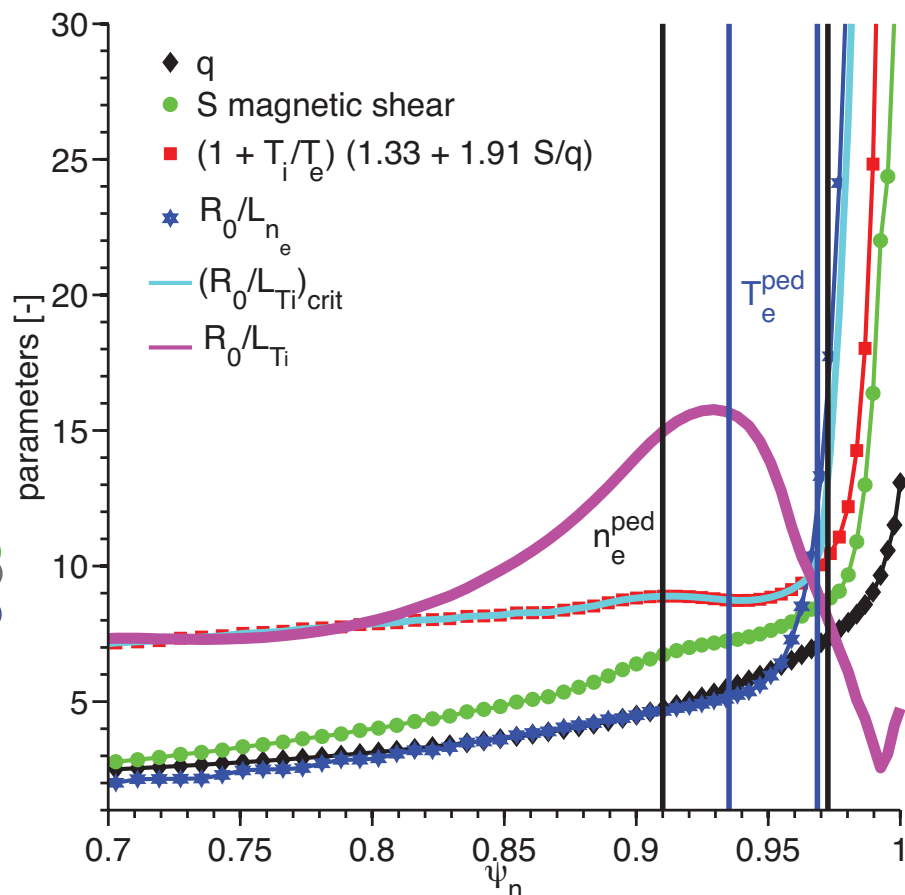


# ITG drives peaks at the pedestal top and $R/L_{Ti}$ is larger than $R/L_{Ti\_crit}$

## ITG drive during inter-ELM phase

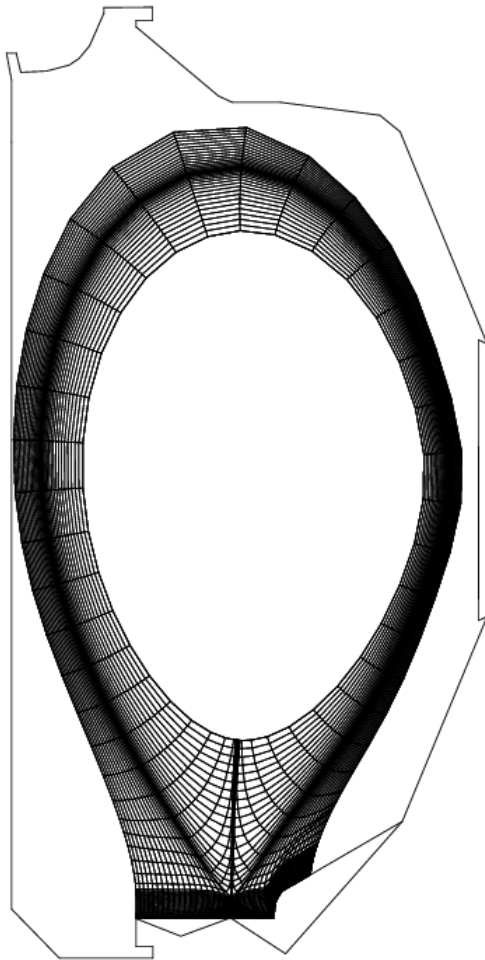


## Critical ITG Gradient



 Jenko's approximation suggests ITG could be unstable in the pedestal top.

# 2-D interpretive modeling with SOLPS\*

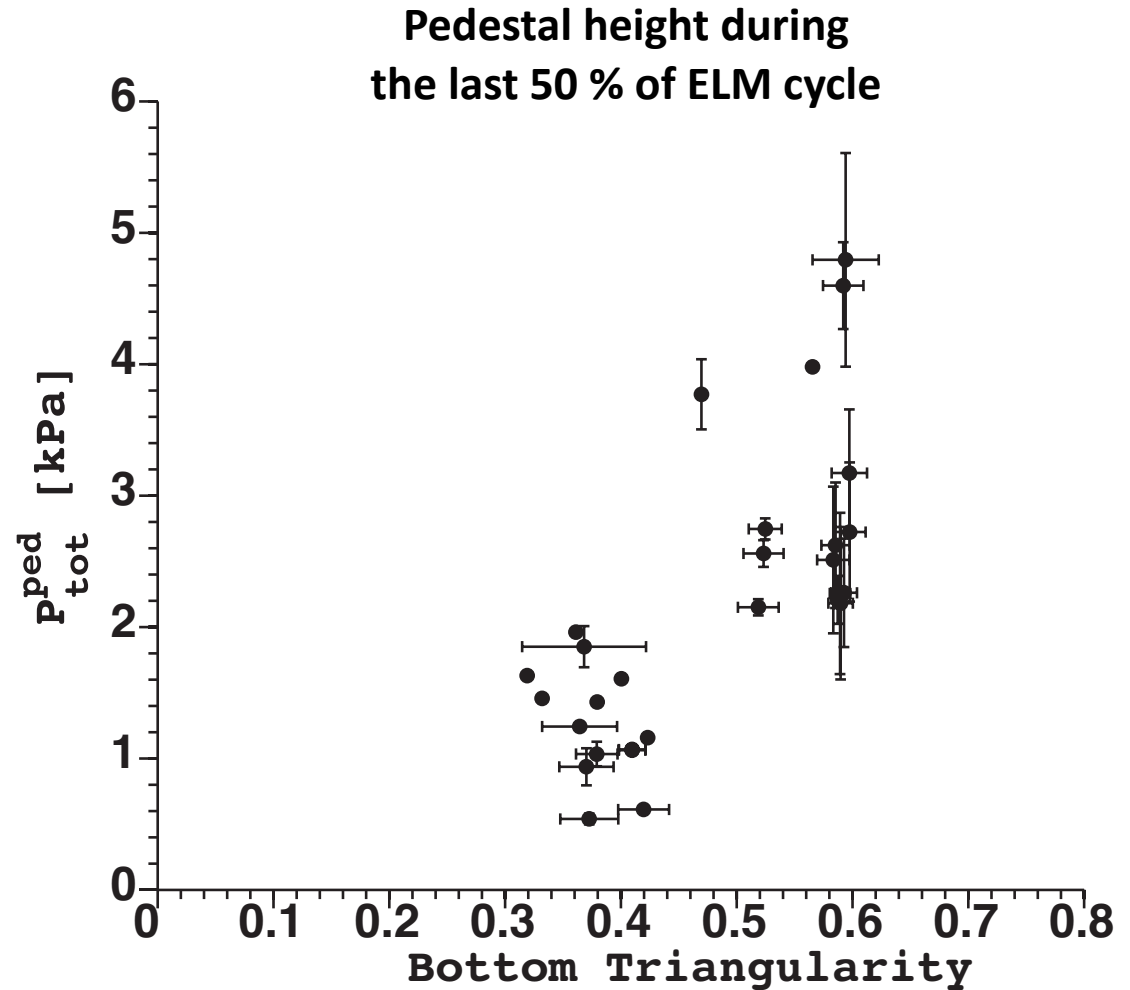
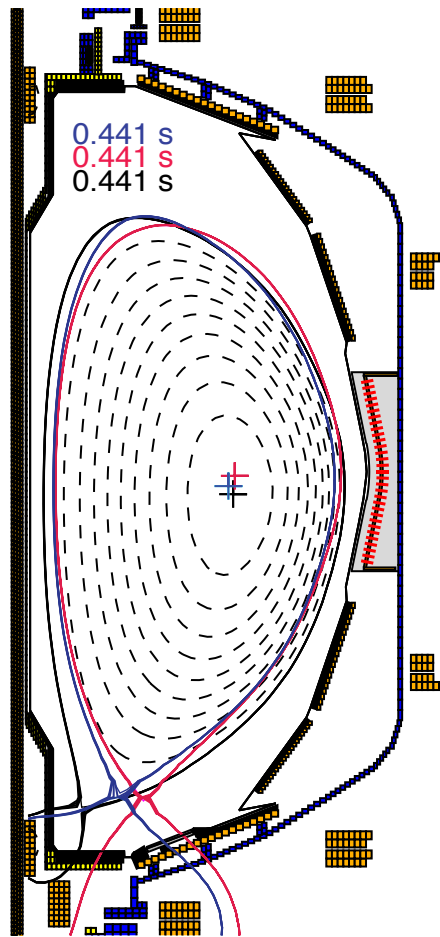


- Solves conservation equations for
  - Density, parallel momentum of each charge state, electron energy, ion energy, charge
- Includes models for plasma transport
  - Parallel: classical along field lines
    - with particle and heat fluxes limited to simulate kinetic effects
  - Radial:  $D$ ,  $\chi$  adjusted to fit measured plasma density and temperature profiles.
  - ExB and grad B drift effects not yet included.
  - Neutral transport: Kinetic, using the Eirene Monte Carlo code
- Yields edge transport coefficients, self-consistent neutral fueling profiles

\*R. Schneider et al, Contrib. Plasma Phys. 46 (2006) 3.

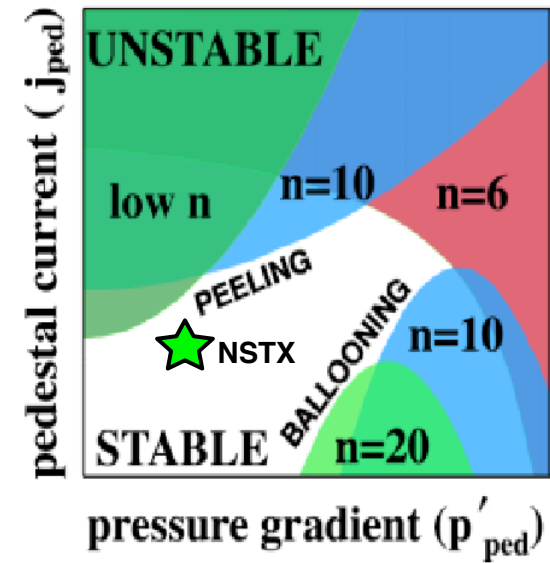


# Pedestal pressure height increases with shaping (triangularity $\delta$ )

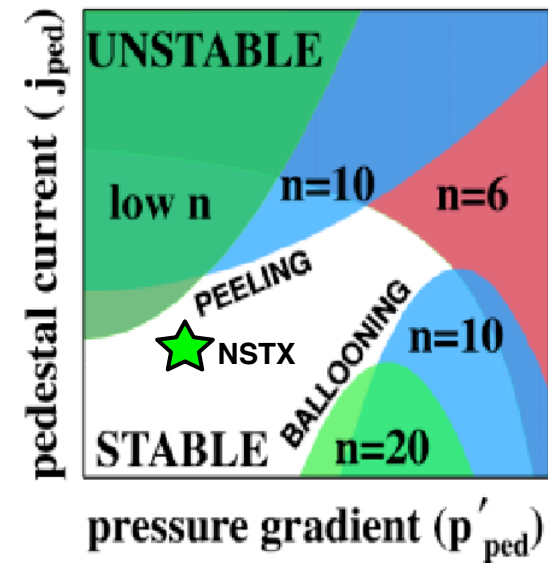


Increasing shaping leads to stability limits at higher  $P'_{ped}$

# Stability diagram with and without lithium: Lithium cases are farther away from the kink/peeling boundary



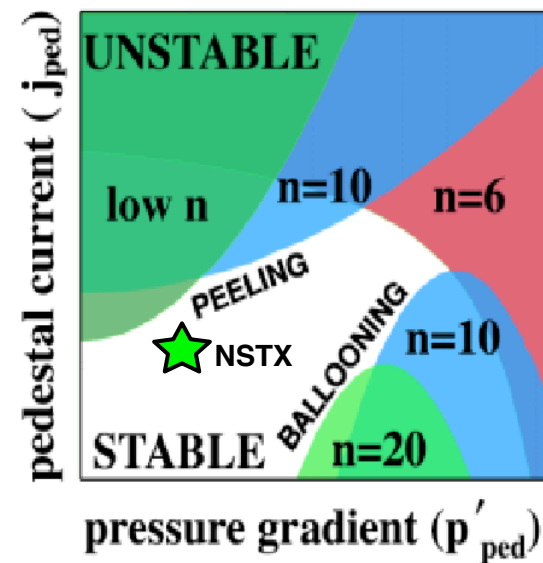
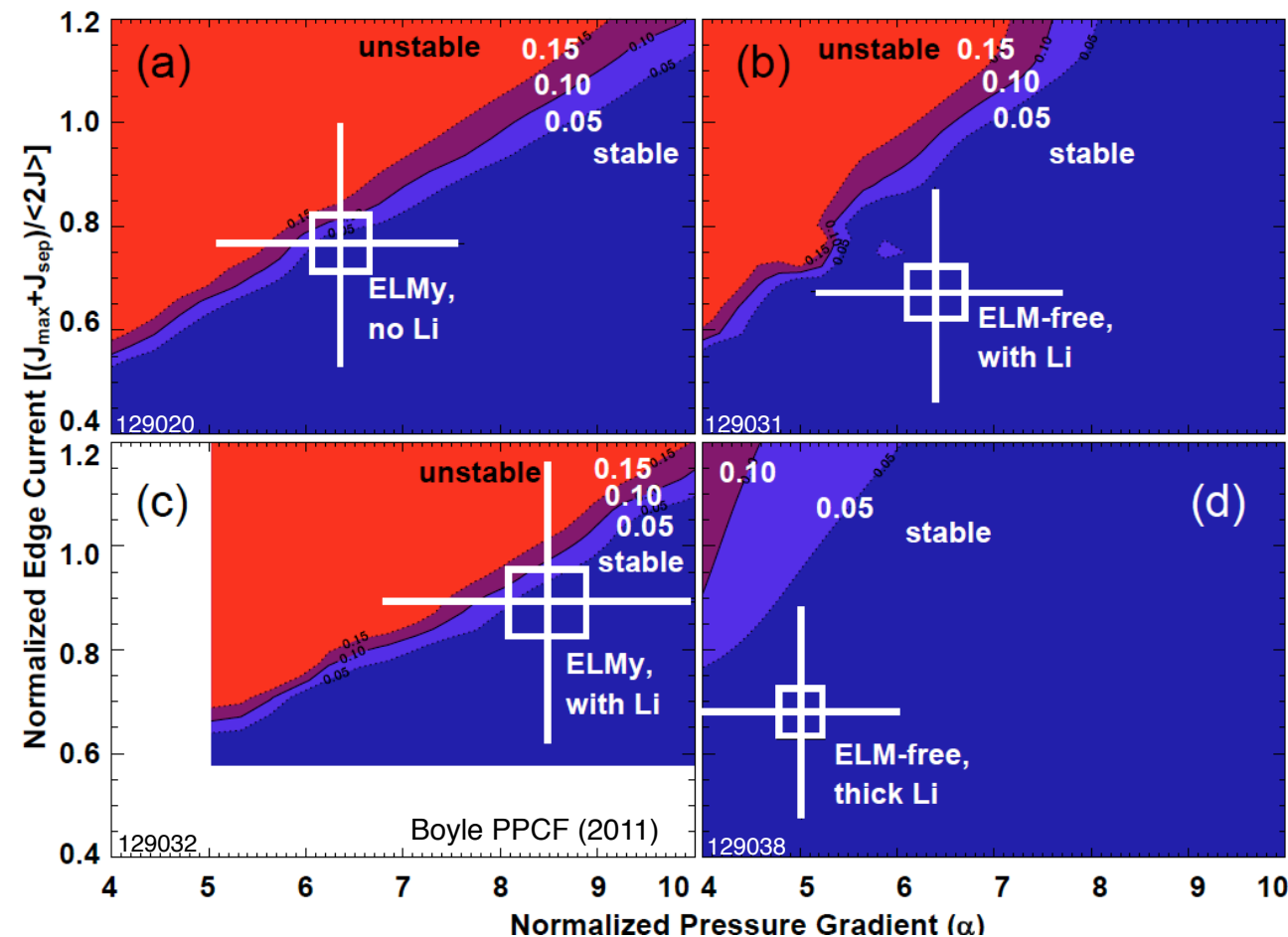
# Stability diagram with and without lithium: Lithium cases are farther away from the kink/peeling boundary



Consistent with NSTX close to the kink/peeling stability boundary

Lithium coatings are a useful tool for shifting peak pressure gradient inward and stabilizing kink/peeling modes.

# Stability diagram with and without lithium: Lithium cases are farther away from the kink/peeling boundary



Consistent with NSTX close to the kink/peeling stability boundary

Lithium coatings are a useful tool for shifting peak pressure gradient inward and stabilizing kink/peeling modes.

# Role of the edge density fluctuations on setting the pedestal structure during the inter-ELM phase

● Pedestal gradient has been predicted to be constrained by the onset of kinetic ballooning mode (KBM)\*

\*Snyder PoP 9 (2002)

– Recent DIII-D work has shown observations of modes localized in the pedestal region with features similar to KBM

Yan PRL 107 (2011)

– KBM characterized by:

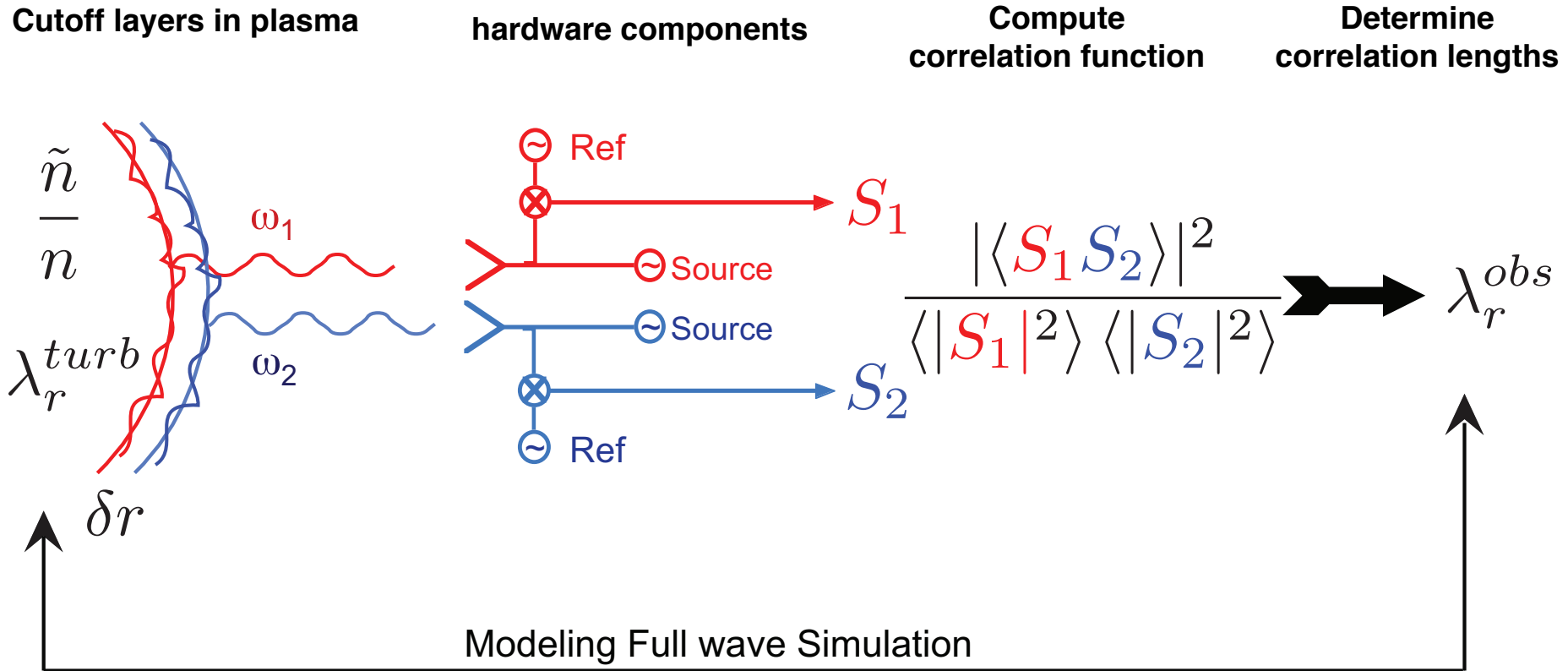
- $k_{\perp}\rho_i < 1$
- modes have radial scales of the order few cm in the pedestal region of NSTX
- fast rising growth rate increasing with electron  $\beta$
- propagation in the ion diamagnetic direction.

● NSTX: We look for evidence of pedestal-localized microinstabilities, and their correlation with the ELM cycle

– Use both reflectometry and BES

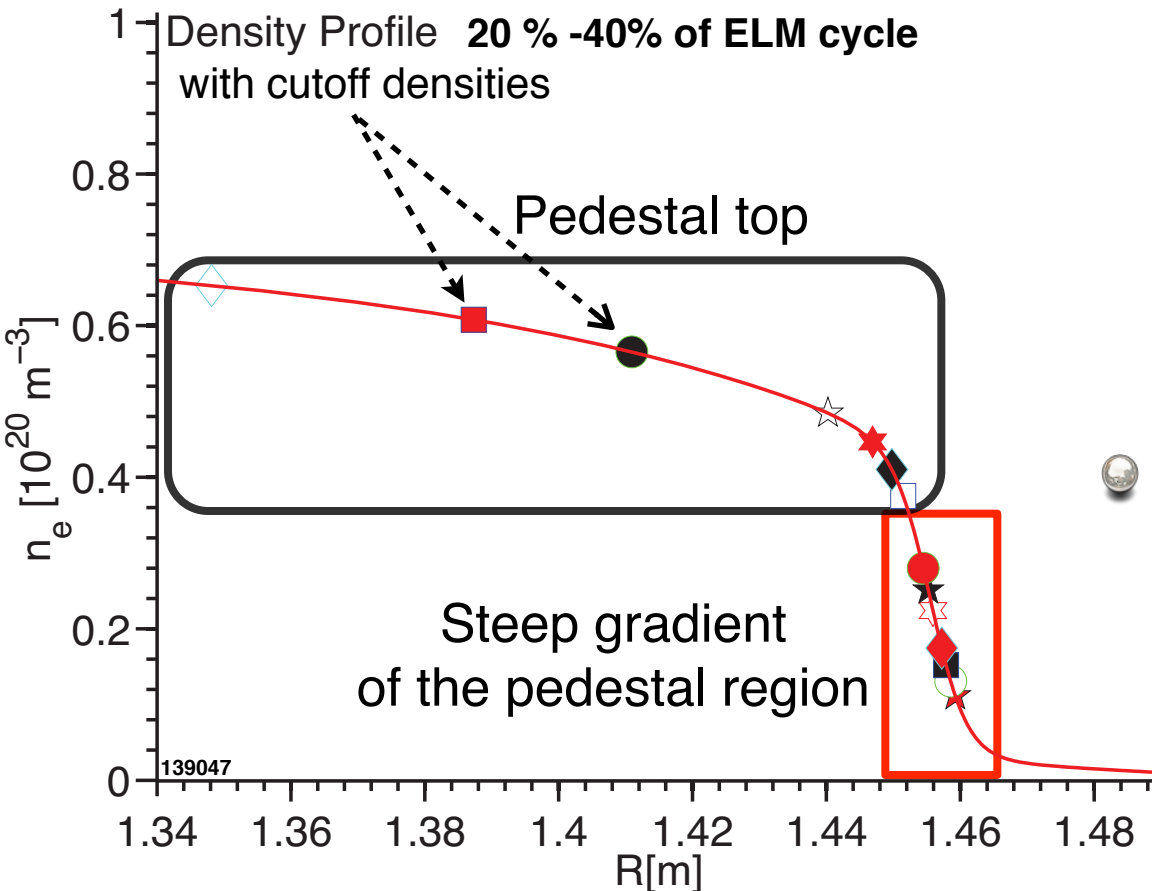
– Because it's hard to conclusively identify KBM, we characterize our instabilities in terms of radial scale, wave number, and propagation direction

# Turbulent fluctuations during the inter-ELM dynamics determined using the correlation reflectometry (UCLA)



- Compare the correlation length measurements with 2D full wave simulations to remove potential instrument function  $\lambda_r^{turb} \neq \lambda_r^{obs}$ 
  - density fluctuation level, equilibrium profiles, and turbulent correlation lengths.

# Radial density correlation lengths at the pedestal top and steep gradient region



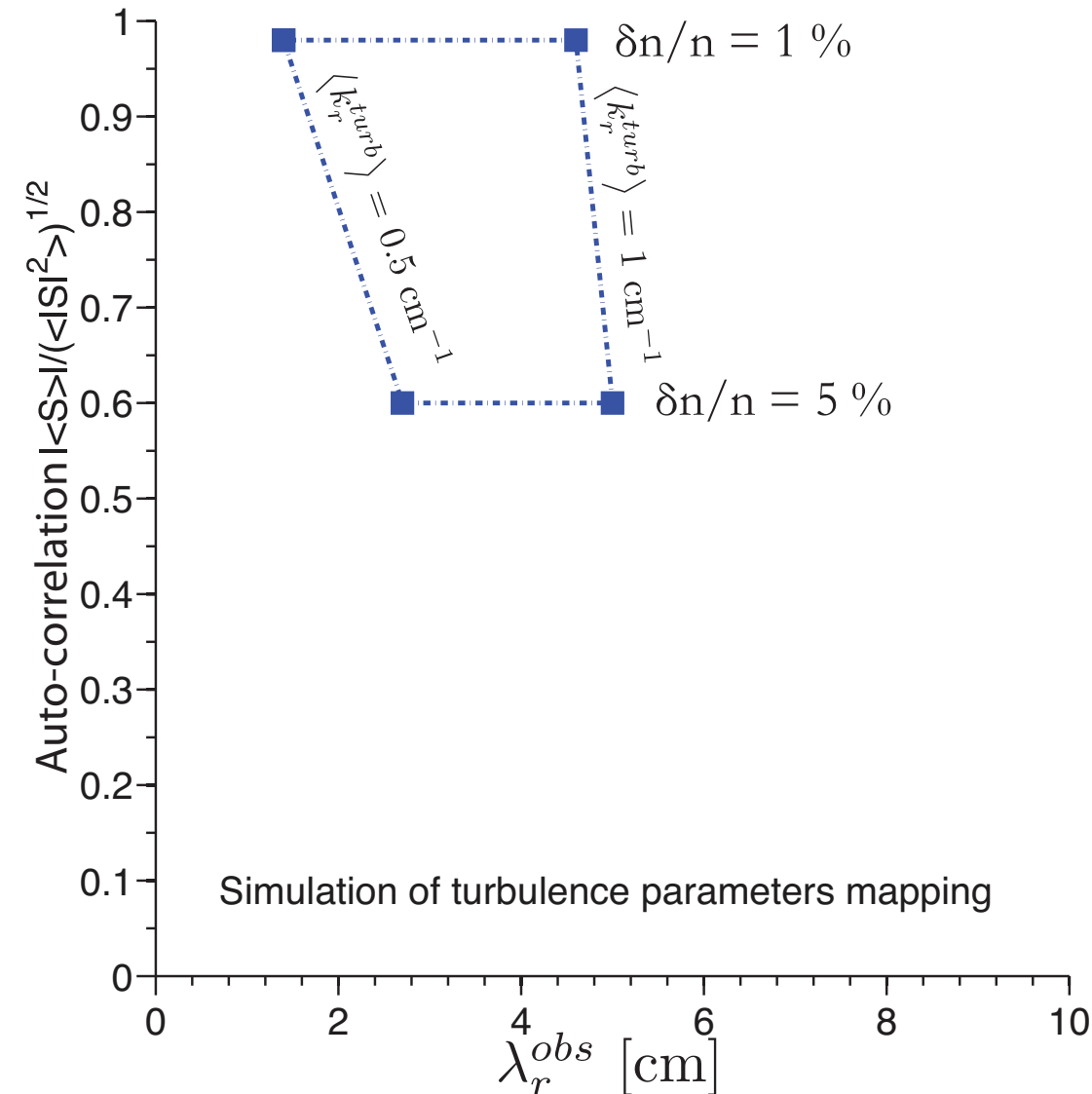
- The density fluctuations are measured using a 16-channel O-mode reflectometer

Crocker, PPCF (2011)

- Using two-point correlation the radial correlation function is determined.
  - tracks the equilibrium plasma motion

# 2D full wave modeling (FWR2D\*) provides correspondence between observed quantities and turbulent parameters

\*Valeo, PPCF (2002)



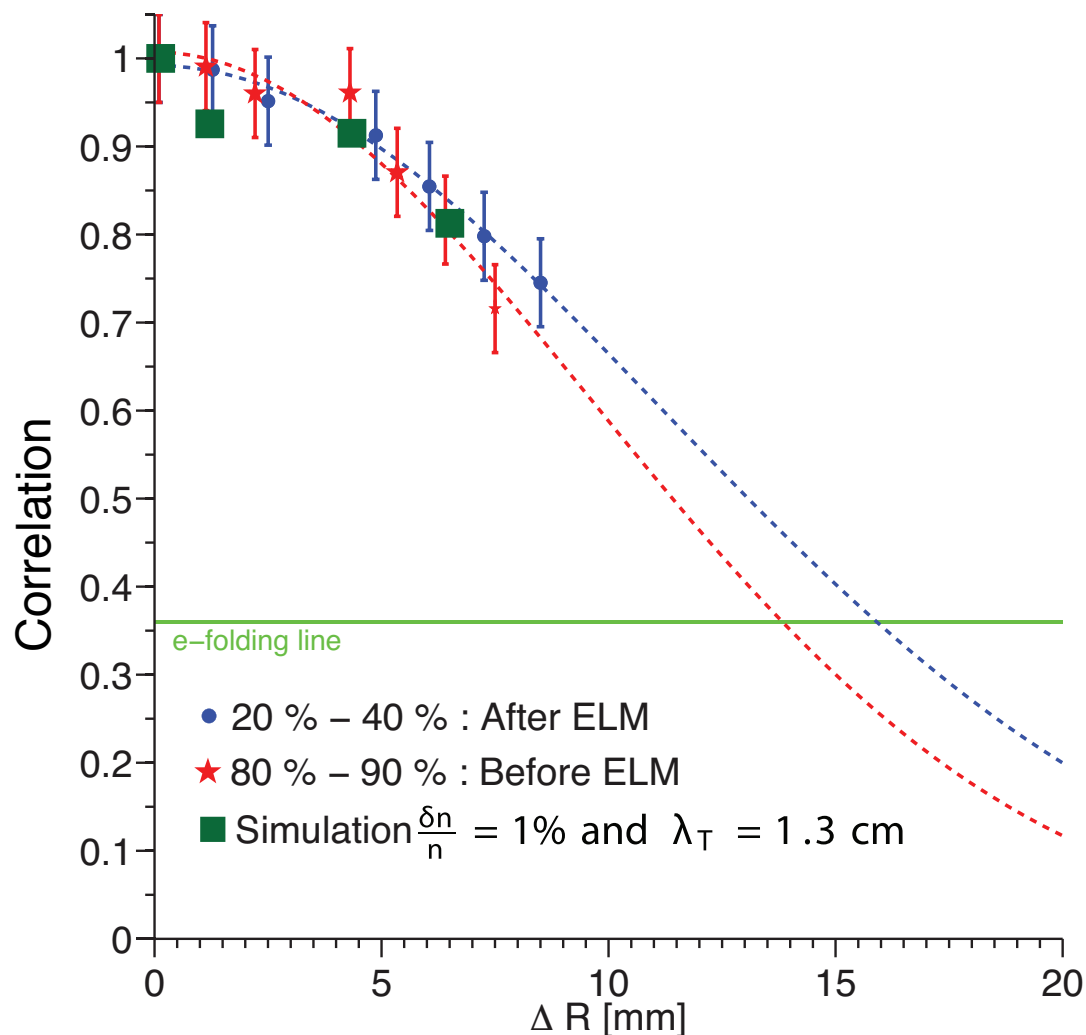
Density fluctuations are generated using a model based gaussian distribution

Kramer RSI 74 (2003)

$$\left(\frac{\delta n}{n}\right)^2 e^{-\delta R^2 / \lambda_{turb}^2}$$



# 2D full wave simulation of correlation function inside pedestal region reproduces measurements



Observed correlation length corresponds to an average eddy size of  $\sim 1.3$  cm with fluctuation level in the vicinity of 1% in the gradient region.