

# Progress Toward Stabilization of Low Internal Inductance Spherical Torus Plasmas in NSTX

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# NSTX is Addressing Global Stability Needs for Maintaining Low $I_i$ , High Beta Plasmas for Fusion Applications

## □ Motivation

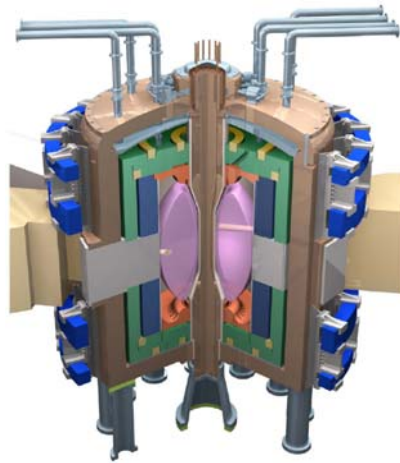
- Sustain high  $\beta_N$  with sufficient physics understanding to allow confident extrapolation to spherical torus applications (e.g. ST Component Test Facility, ST-Pilot plant, ST-DEMO)  
NP9.00011 Peng  
UP9.00006 Hawryluk
- Operate at low internal inductance ( $I_i$ ), consistent with desired high elongation and non-inductive (bootstrap, etc.) current fraction
- Demonstrating / understanding global mode stability at low  $I_i$  is key

## □ Physics Research Addressed in this Talk

- Global mode stability at low internal inductance
- Resistive wall mode (RWM) destabilization at high plasma rotation,  $\omega_\phi$
- RWM active control advances to maintain high  $\beta_N$  at low  $I_i$ , varied  $\omega_\phi$
- Multi-mode RWM spectrum in high  $\beta_N$  plasmas

# Future ST fusion applications will have high elongation, broad current profiles, high normalized beta

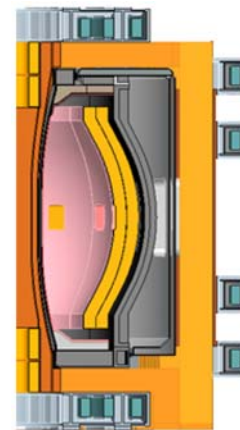
## FNSF / ST-CTF



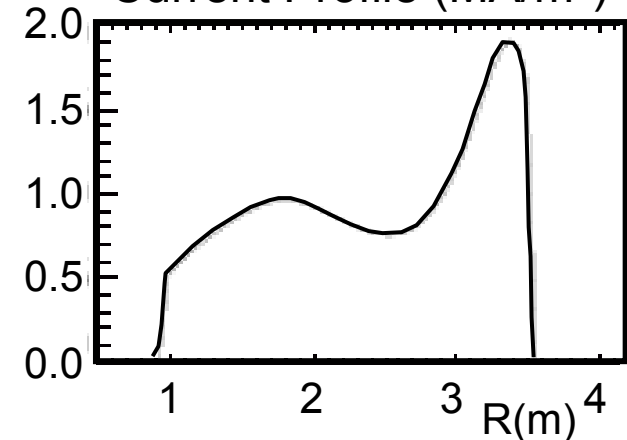
Y.K.M. Peng, et al., PPCF 47 (2005) B263

## ST-Pilot ( $Q_{eng} = 1$ )

J. Menard, et al., IAEA FEC 2010 Paper FTP/2-2



### Current Profile (MA/m<sup>2</sup>)



$$I_i = 0.47, \kappa = 3.2$$

$$R = 2.23 \text{ m}, A = 1.7$$

$$I_p = 16 \text{ MA}, B_t = 2.4 \text{ T}$$

$$\beta_N = 5.2, \beta_t = 30\%$$

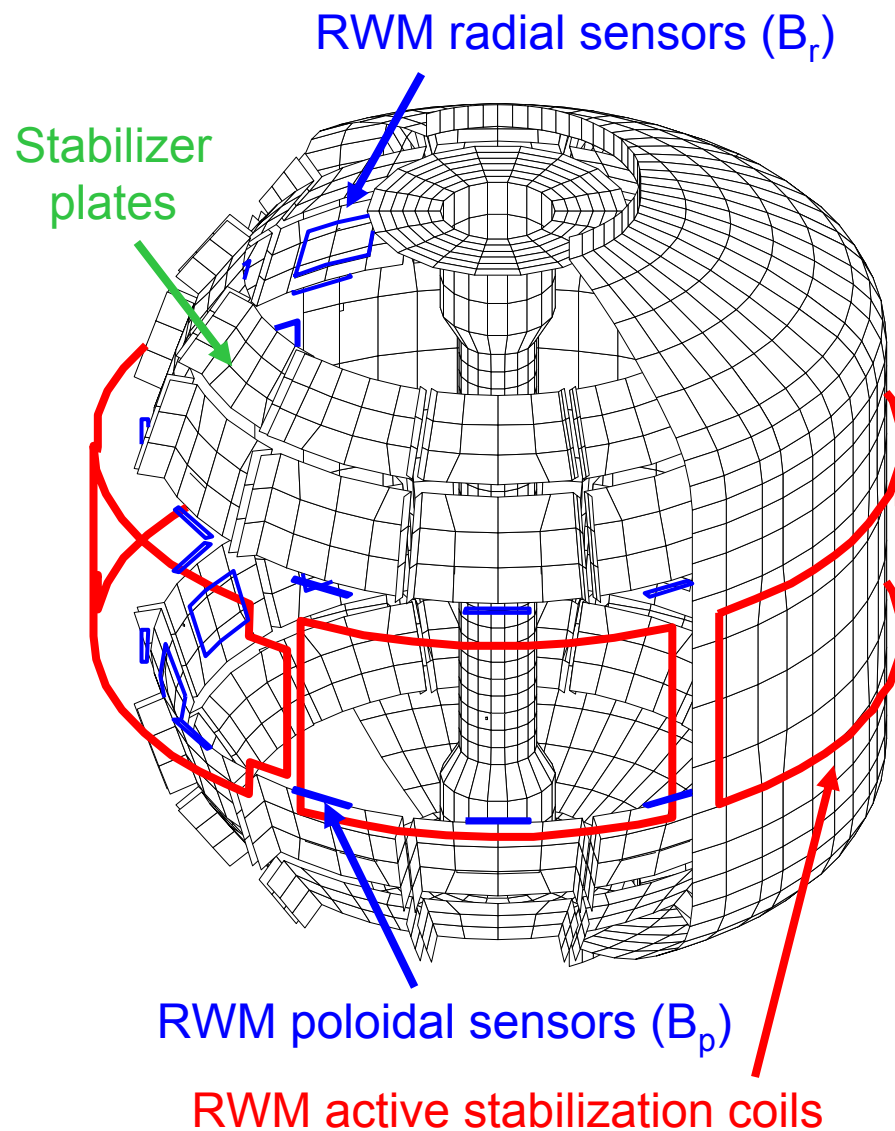
- ❑ Broad current profile  $\rightarrow$  low  $I_i = \langle B_p^2 \rangle / \langle B_p \rangle_\psi^2$ , has global mode stability implications
  - ❑ Improved vertical ( $n = 0$ ) and wall-stabilized, rotating kink ( $n = 1$ ) stability
  - ❑ Decreased RWM stability, influenced by  $\omega_\phi$
  - ❑ “Troyon limit” (ideal, static  $n = 1$  no-wall stability limit)

$$\beta_N = 10^8 \beta_t a B_t / I_p = \text{constant} \sim 3; \text{ variants: } \beta_N \propto I_i, \beta_N \propto 1/(\text{pressure peaking})$$

Operation at higher  $\beta_N$  possible by passive or active RWM stabilization ( $\beta_t = 2\mu_0 \langle p \rangle / B_t^2$ )

# NSTX is a spherical torus equipped to study passive and active global MHD control, rotation variation by 3D fields

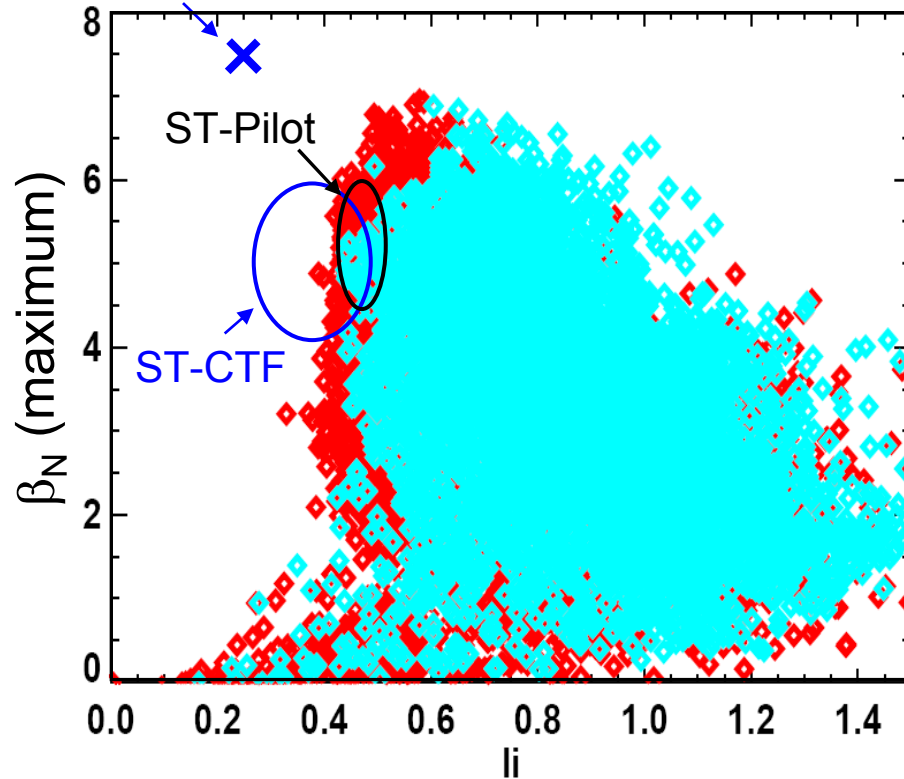
- High beta, low aspect ratio
  - $R = 0.86$  m,  $A > 1.27$
  - $I_p < 1.5$  MA,  $B_t = 5.5$  kG
  - $\beta_t < 40\%$ ,  $\beta_N > 7$
- Copper stabilizer plates for kink mode stabilization
- Midplane control coils
  - $n = 1 - 3$  field correction, magnetic braking of  $\omega_\phi$  by NTV
  - $n = 1$  RWM control
- Combined sensor sets now used for RWM feedback
  - 48 upper/lower  $B_p$ ,  $B_r$



# Operation has aimed to produce sustained low $I_i$ and high pulse-averaged $\beta_N$

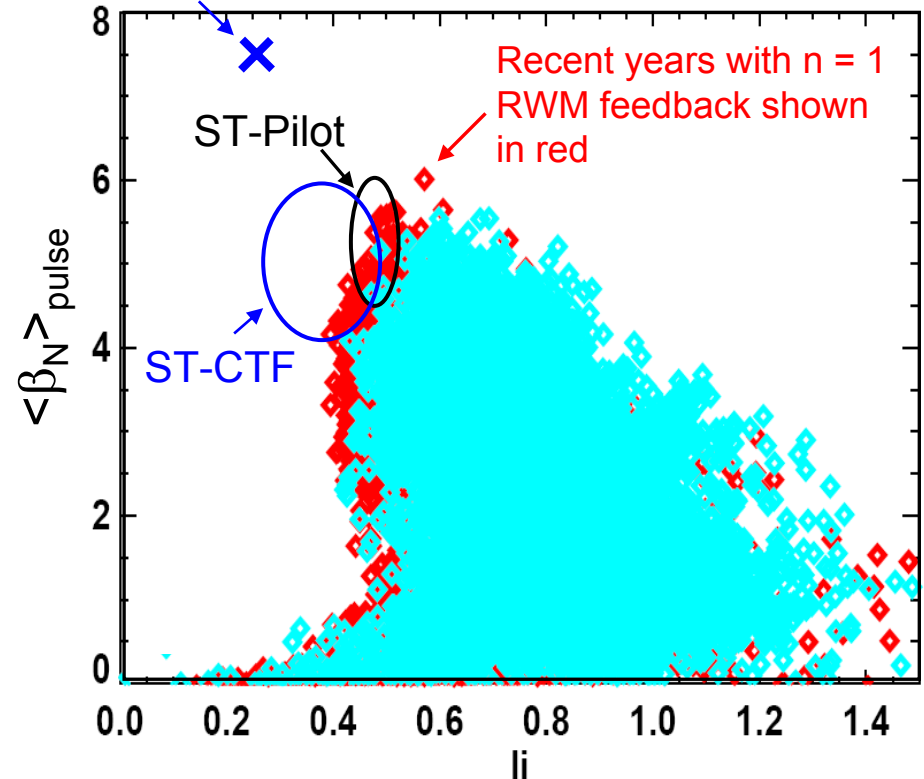
$\beta_N$  vs.  $I_i$  (maximum values)

ST-DEMO (ARIES-ST)



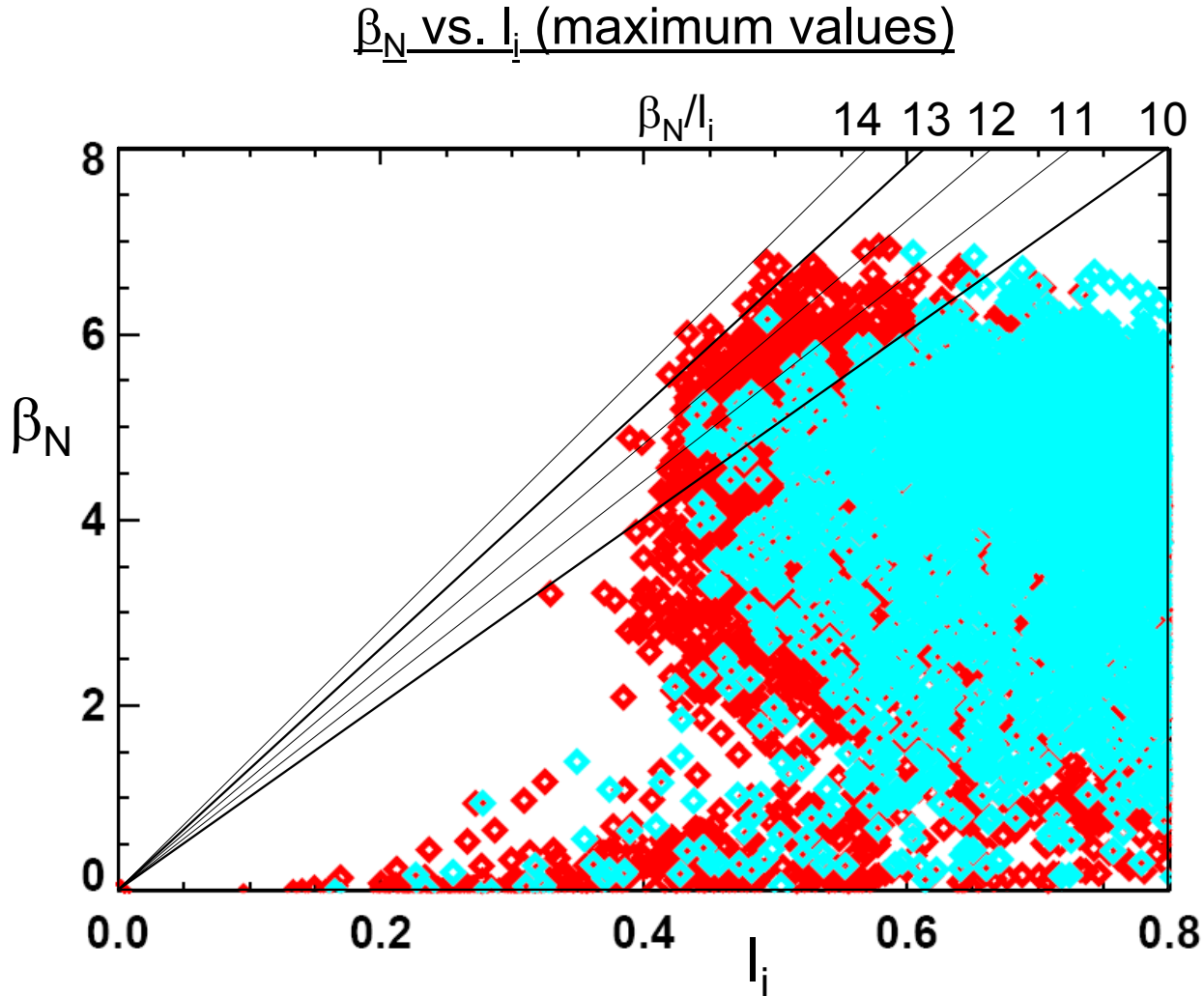
$\beta_N$  vs.  $I_i$  (pulse-averaged values)

ST-DEMO (ARIES-ST)



- Plasmas have begun to reach low  $I_i$  and high  $\langle \beta_N \rangle_{\text{pulse}}$  suitable for next-step ST fusion devices
  - Some parameters (e.g. elongation  $> 3$ ) still need to be reached self-consistently

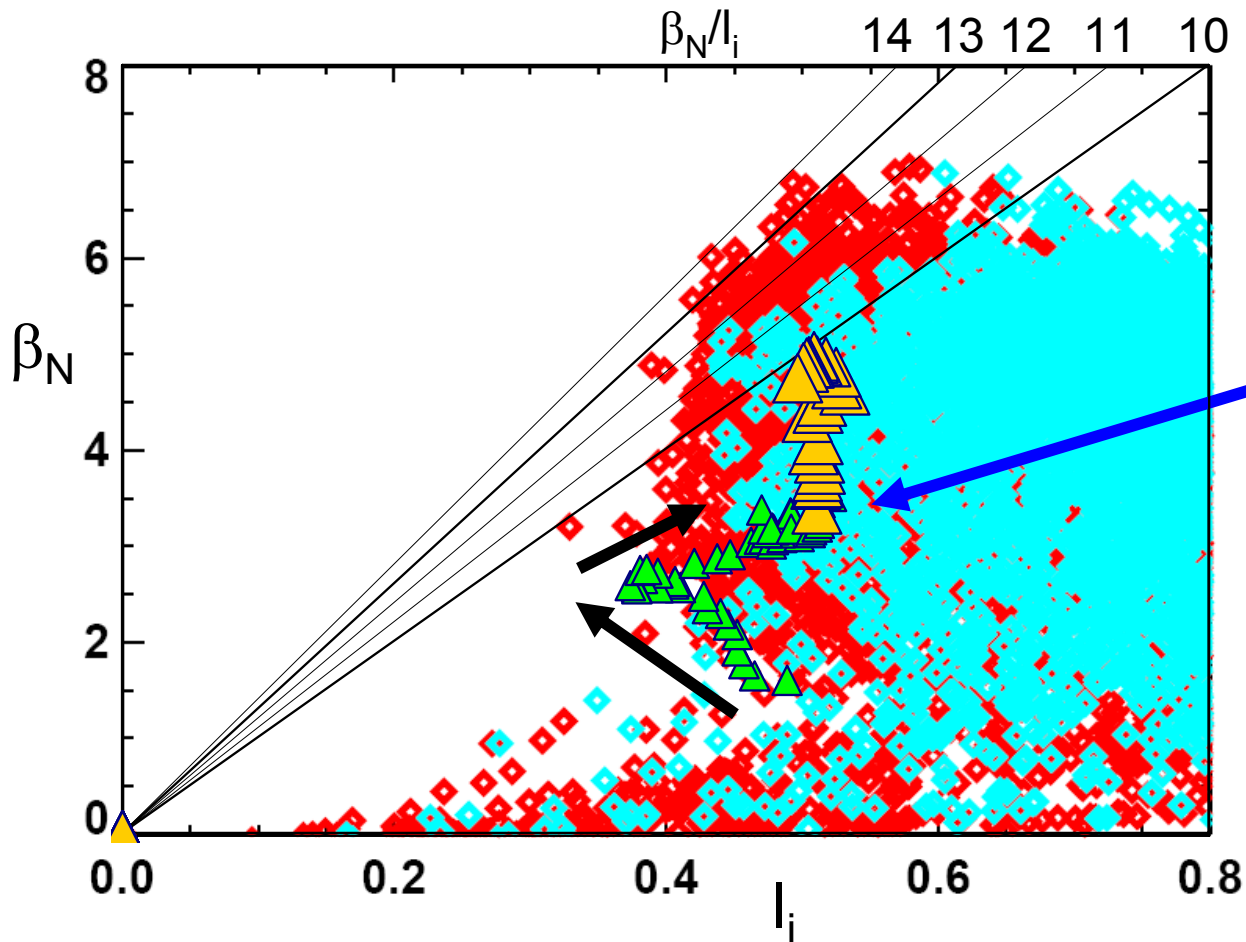
# Operational space is expanding to low $I_i$ and high $\beta_N$



- $\beta_N/I_i$  is a common parameter to evaluate global stability
  - Kink/ballooning and RWM stability
- Significant increase in maximum  $\beta_N/I_i$ 
  - Upper limit now between 13 - 14

# Ideal $n = 1$ no-wall stability examined for low $I_i$ plasmas

$\beta_N$  vs.  $I_i$  (maximum values)

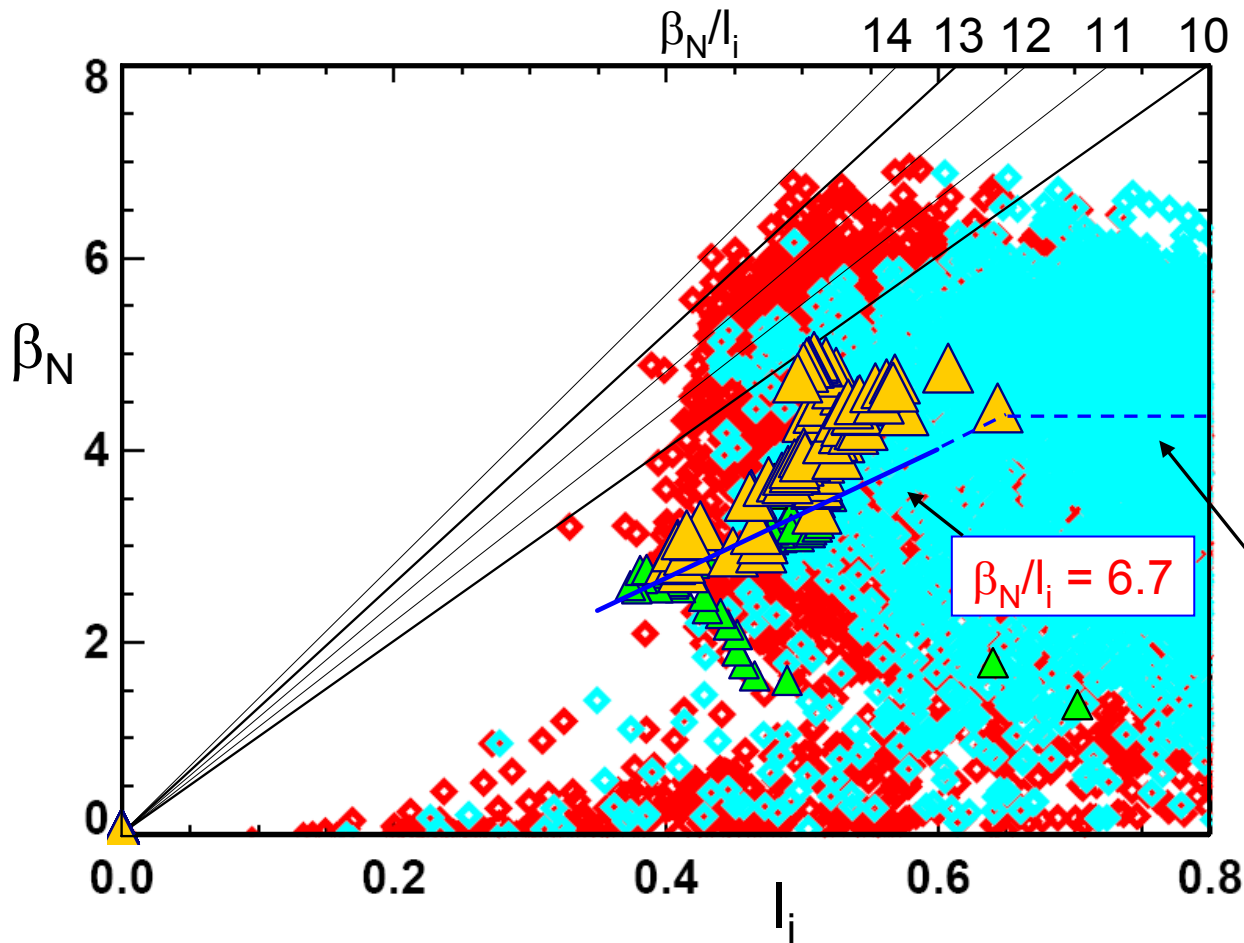


- Examine high plasma current,  $I_p \geq 1.0\text{MA}$ , high non-inductive fraction  $\sim 50\%$ ,  $I_i \sim 0.4 - 0.5$
- Ideal  $n = 1$  no-wall stability computed for discharge trajectory
  - Green: below limit
  - Gold: above limit
  - Exceeds no-wall limit at  $\beta_N = 3.4$ ,  $I_i = 0.51$

DCON (A. H. Glasser)

# Ideal $n = 1$ no-wall stability limit decreases for low $I_i$ plasmas

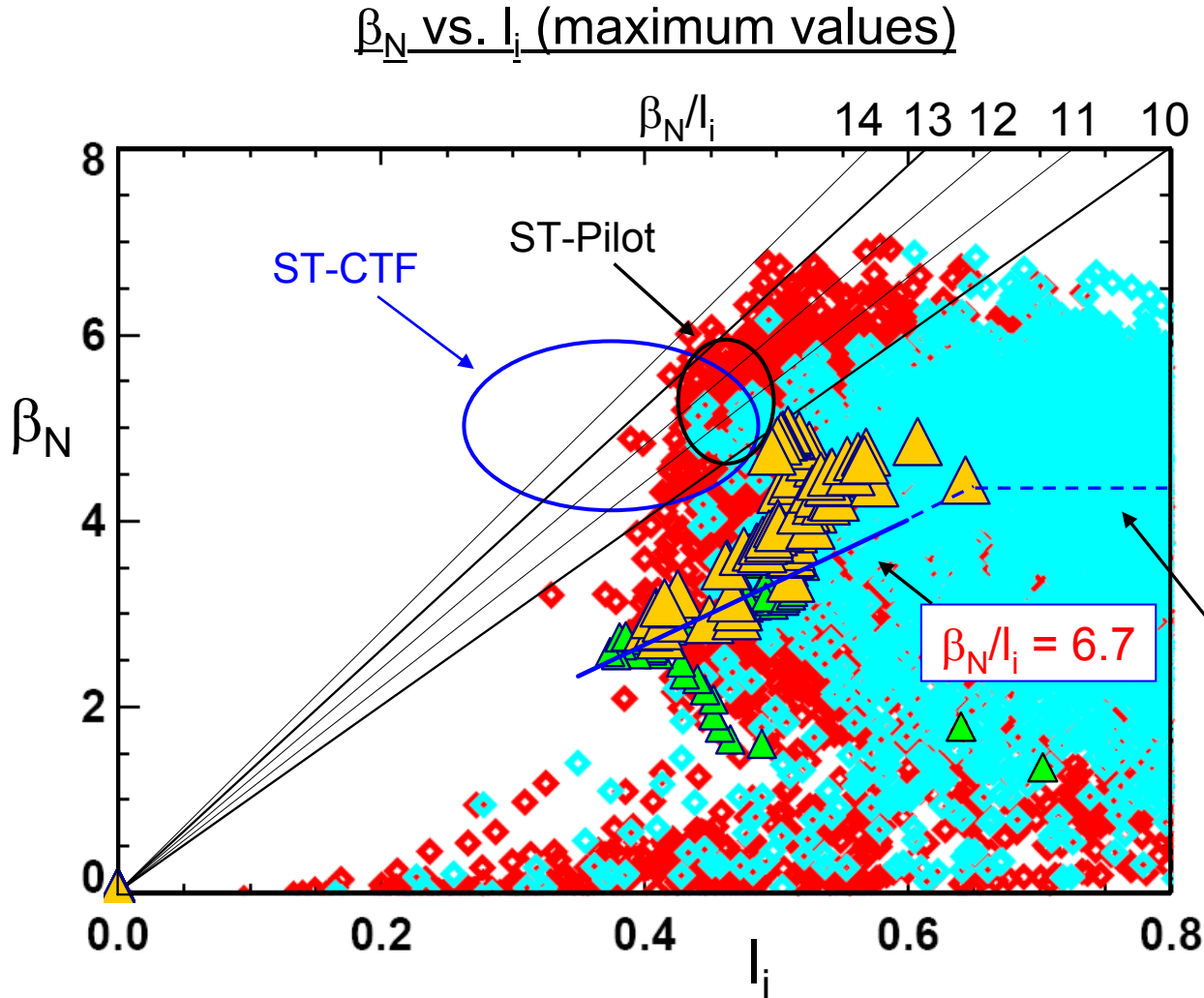
$\beta_N$  vs.  $I_i$  (maximum values)



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- Ideal  $n = 1$  no-wall stability computed for discharge trajectory
  - Adding trajectories yields  $\beta_N/I_i = 6.7$  for  $I_i = 0.38 - 0.6$
  - Significantly lower than no-wall limit at higher  $I_i$  ( $\beta_N = 4.3$ )

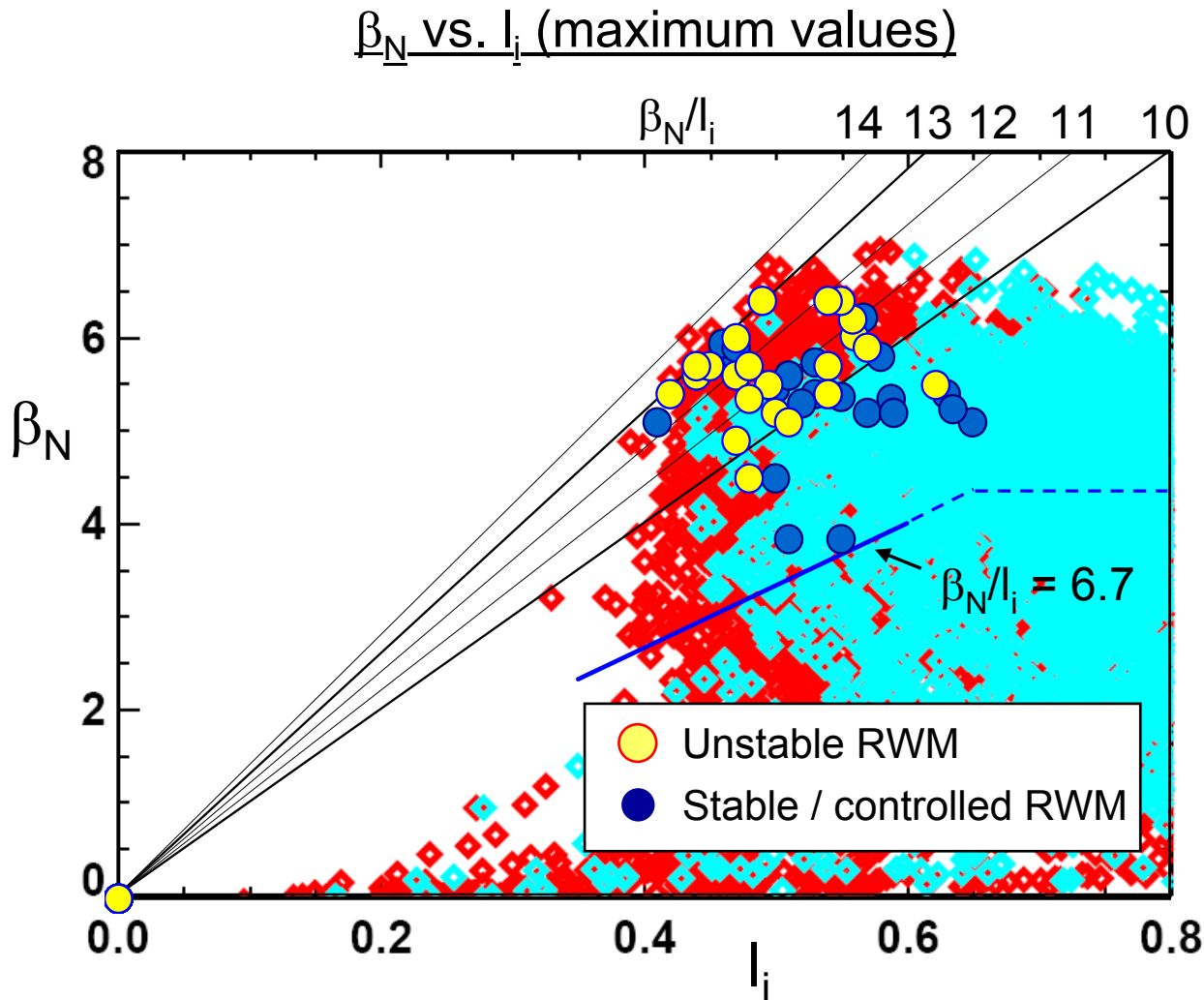


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- Examine high plasma current,  $I_p \geq 1.0\text{MA}$ , high non-inductive fraction  $\sim 50\%$ ,  $I_i \sim 0.4 - 0.5$
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  - Significantly lower than no-wall limit at higher  $I_i$  ( $\beta_N = 4.3$ )
  - RWM control will be important for future ST fusion devices

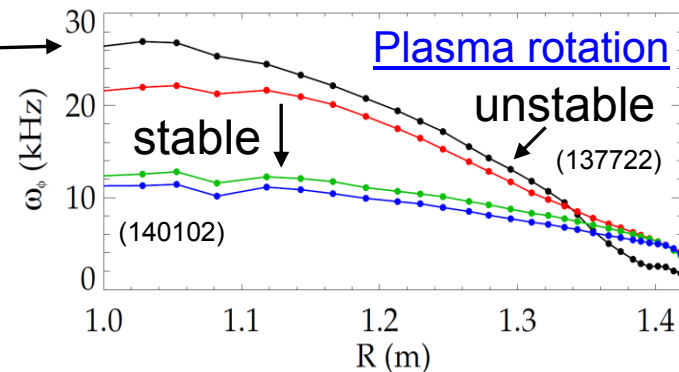
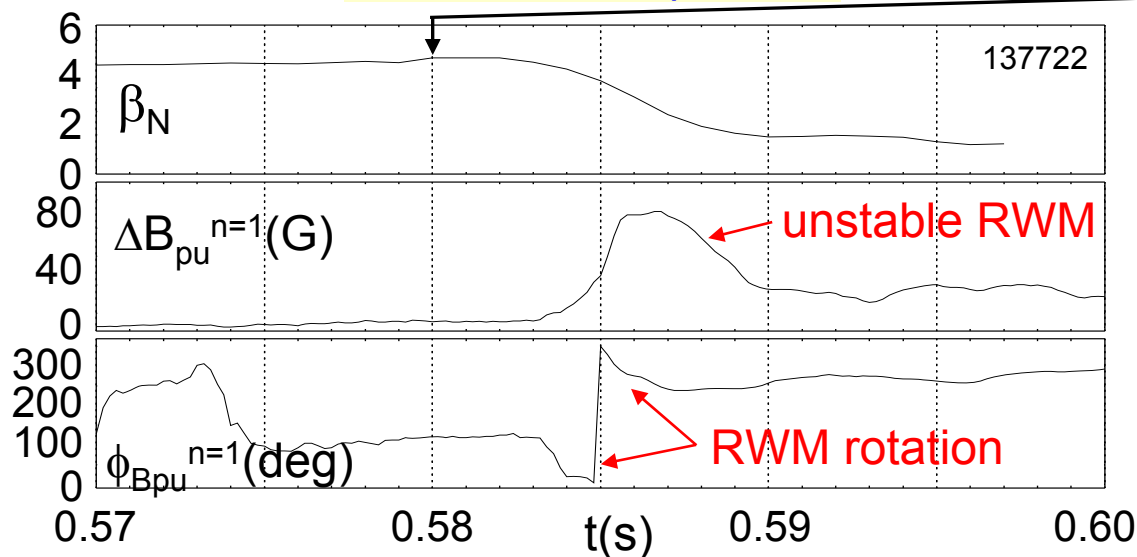
# Global stability examined for experiments aimed to produce sustained low $I_i$ and high $\beta_N$ at high plasma current



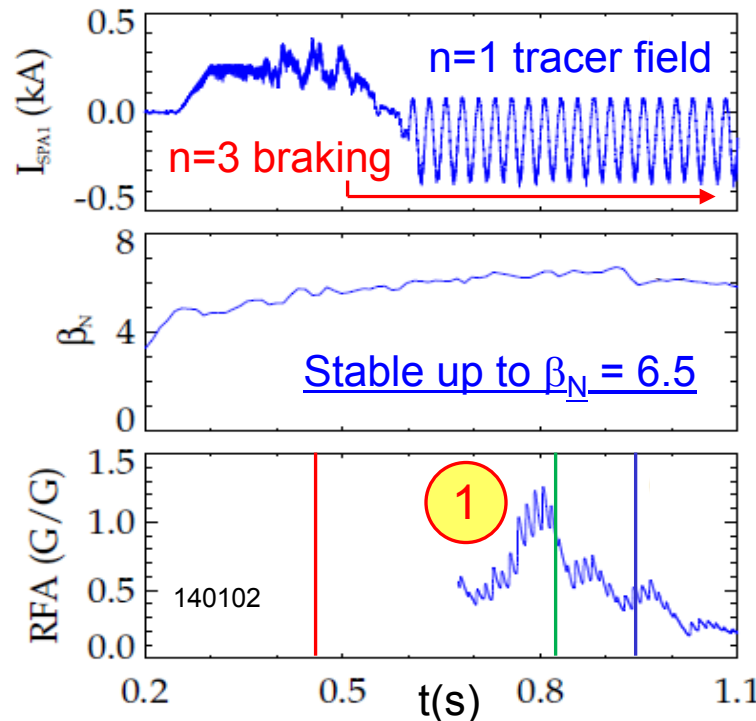
- High  $I_p \geq 1.0\text{MA}$ , high non-inductive fraction  $\sim 50\%$
- Initial experiments
  - Yielded low  $I_i$
  - Access high  $\beta_N/I_i$
  - High disruption probability
- Instabilities leading to disruption
  - Unstable RWM
    - 48% of cases run
  - Locked tearing modes

# Low plasma rotation level ( $\sim 1\% \omega_{\text{Alfvén}}$ ) is insufficient to ensure RWM stability, which depends on $\omega_\phi$ profile

## RWM unstable plasma



## MHD spectroscopy (stable plasma)



### RWM unstable plasma

- Instability occurs at relatively **high rotation** level, and **not** at highest  $\beta_N$  (4.7)

### RWM stable plasma

- MHD spectroscopy: increased resonant field amplification (RFA) indicates reduced stability **1**
- Plasma moves to more stable regime (lower RFA) at lower rotation ( $\beta_N$  up to 6.5)

# Modification of Ideal Stability by Kinetic theory (MISK code) investigated to explain experimental RWM stabilization

- Reason: simple critical  $\omega_\phi$  threshold stability models do not fully describe RWM marginal stability in NSTX Sontag, et al., Nucl. Fusion **47** (2007) 1005.

- Kinetic modification to ideal MHD growth rate

- Trapped / circulating ions, trapped electrons, etc.
- Energetic particle (EP) stabilization

$$\gamma\tau_w = -\frac{\delta W_\infty + \delta W_K}{\delta W_b + \delta W_K}$$

Hu and Betti, Phys. Rev. Lett **93** (2004) 105002.

- Stability depends on

- Integrated  $\omega_\phi$  profile: resonances in  $\delta W_K$  (e.g. ion precession drift)
- Particle collisionality, EP fraction  $\omega_\phi$  profile (enters through ExB frequency)

## Trapped ion component of $\delta W_K$ (plasma integral)

$$\delta W_K \propto \int \left[ \frac{\omega_{*N} + \left(\hat{\epsilon} - \frac{3}{2}\right)\omega_{*T} + \omega_E - \omega - i\gamma}{\langle \omega_D \rangle + l\omega_b - i\nu_{eff} + \omega_E - \omega - i\gamma} \right] \hat{\epsilon}^{\frac{5}{2}} e^{-\hat{\epsilon}} d\hat{\epsilon}$$

precession drift

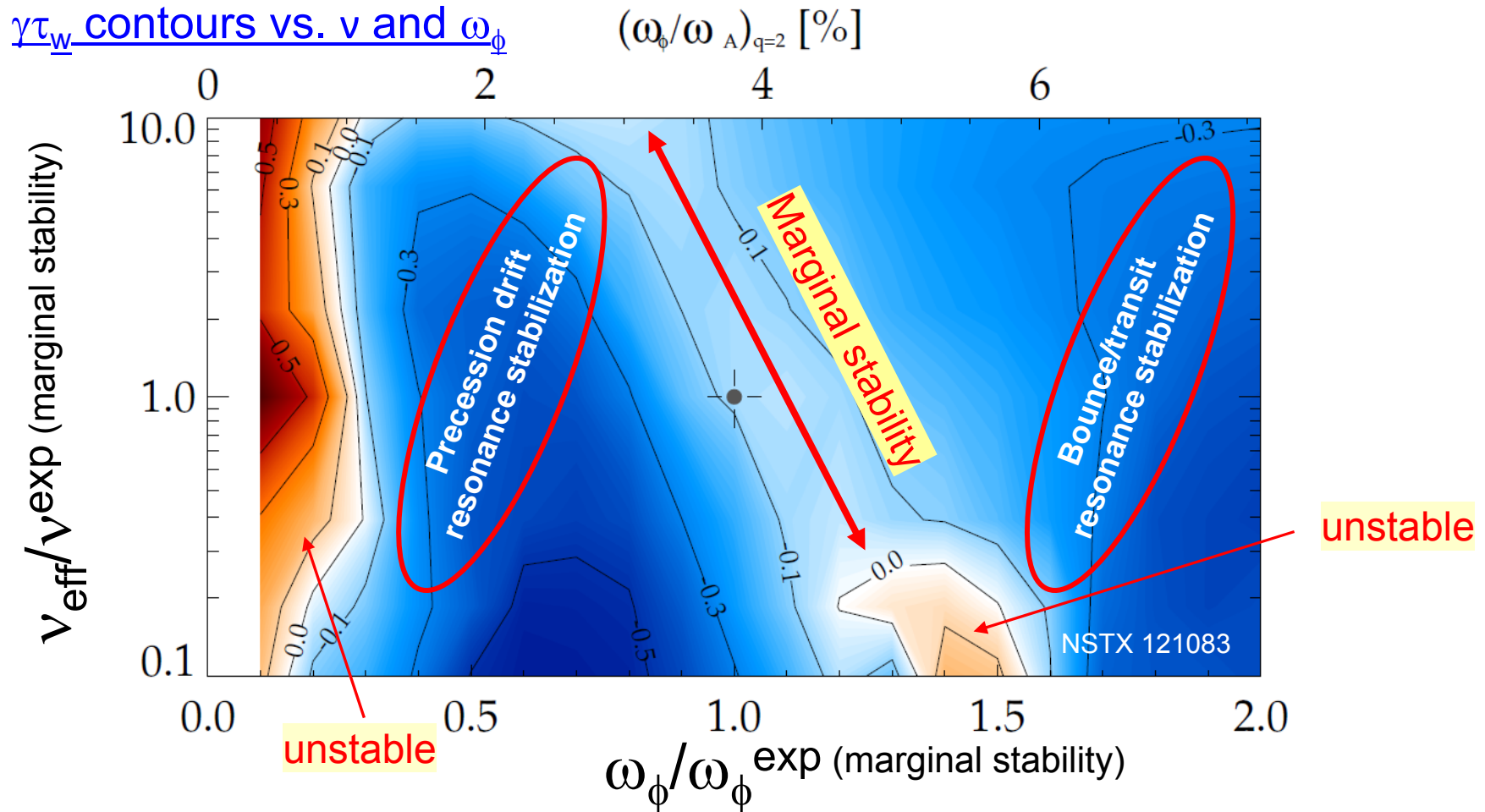
bounce

collisionality

← Energy integral

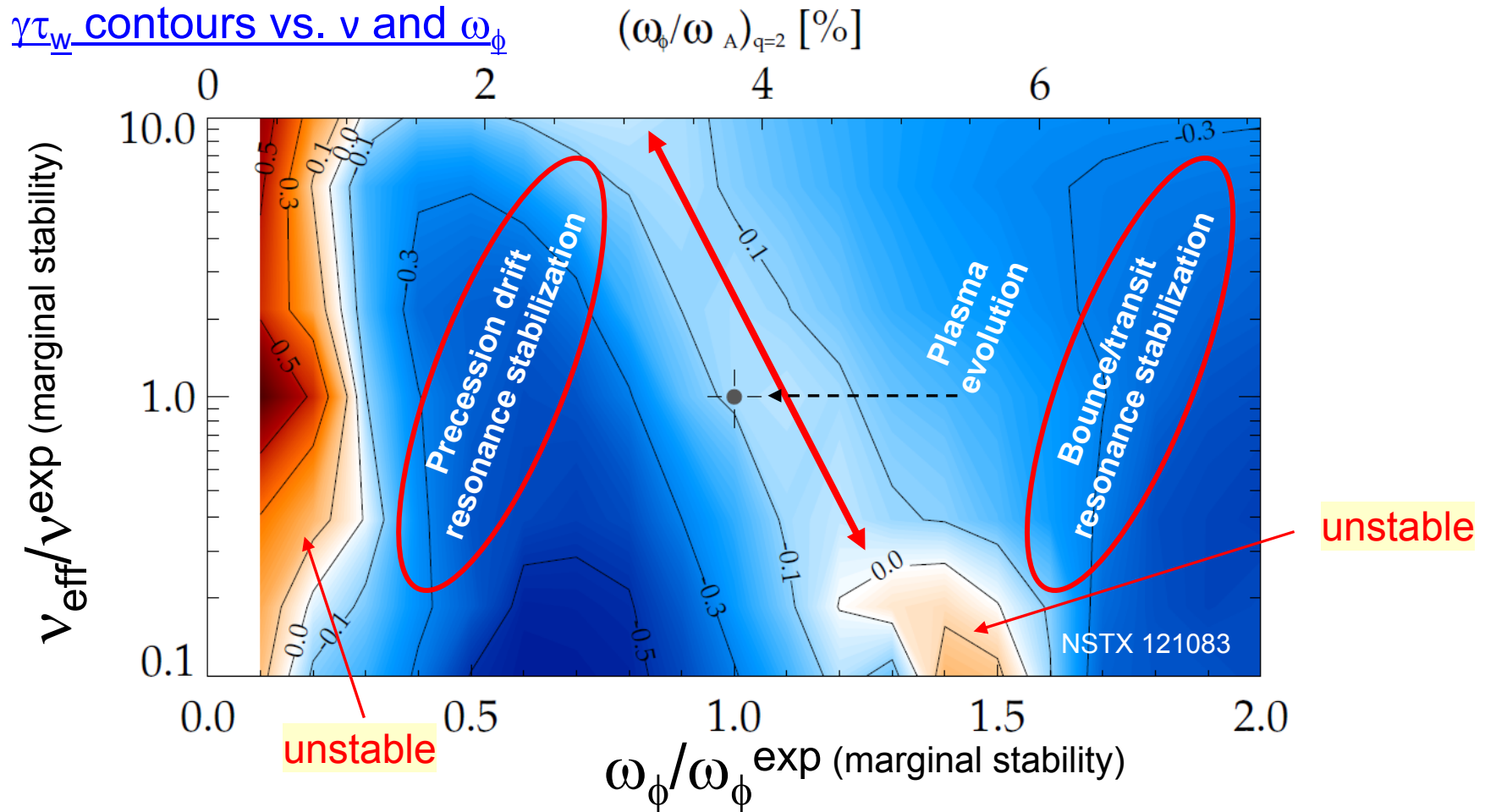
BP9.00057 J. Berkery, et al.

# MISK calculations consistent with RWM destabilization at intermediate plasma rotation; stability altered by collisionality



- Destabilization appears between precession drift resonance at low  $\omega_\phi$ , bounce/transit resonance at high  $\omega_\phi$
- J.W. Berkery, et al., PRL **104** (2010) 035003  
 S.A. Sabbagh, et al., NF **50** (2010) 025020

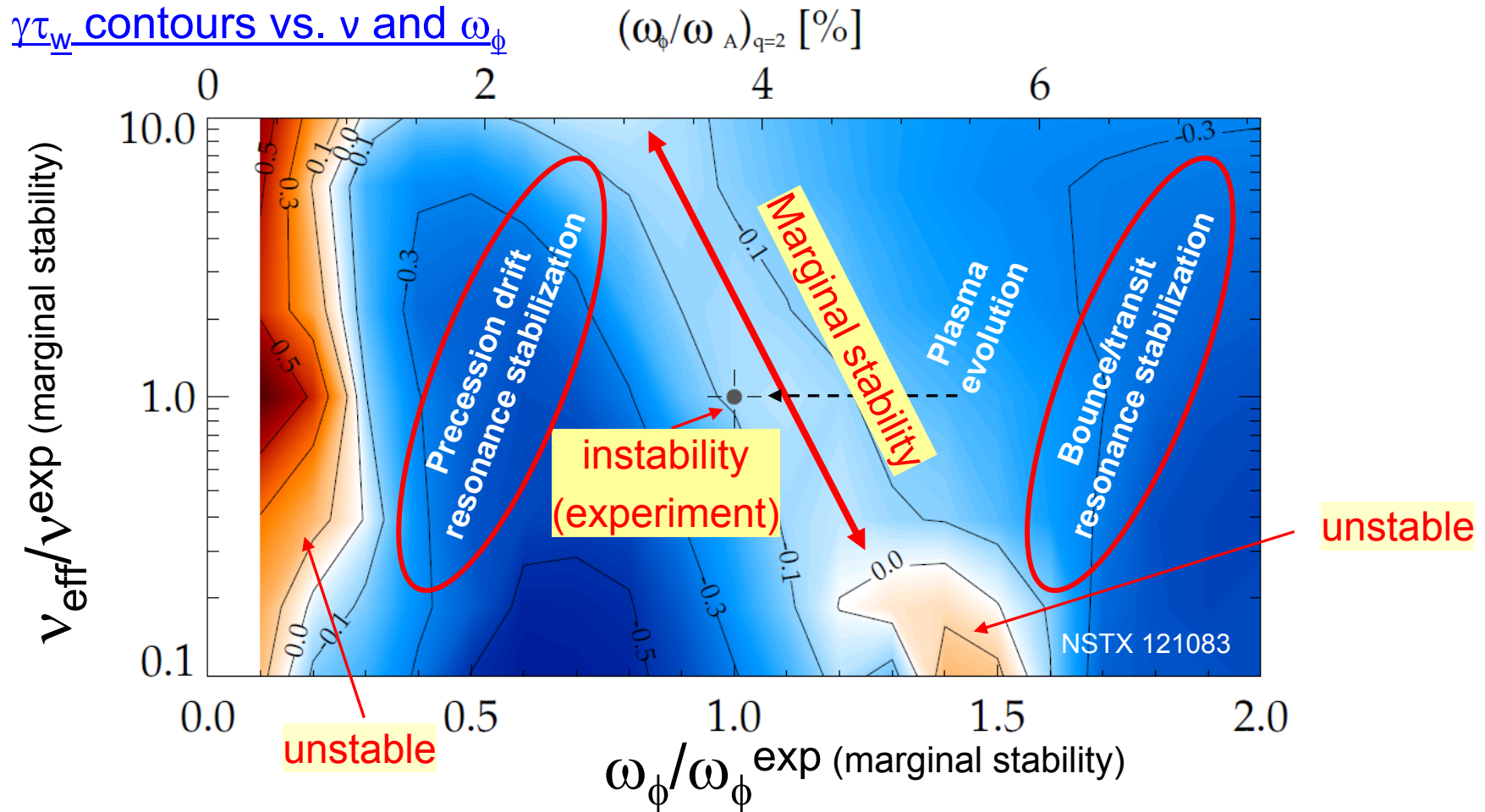
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J.W. Berkery, et al., PRL **104** (2010) 035003  
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# MISK calculations consistent with RWM destabilization at intermediate plasma rotation; stability altered by collisionality



- ❑ Destabilization appears between precession drift resonance at **low**  $\omega_\phi$ , bounce/transit resonance at **high**  $\omega_\phi$ 

J.W. Berkery, et al., PRL **104** (2010) 035003  
S.A. Sabbagh, et al., NF **50** (2010) 025020
- ❑ Destabilization moves to increased  $\omega_\phi$  as  $v$  decreases

# MISK calculations show reduced stability in low $I_i$ target plasma as $\omega_\phi$ is reduced, RWM instability is approached

## Stability evolves

- MISK computation shows plasma to be stable at time of minimum  $I_i$
- Region of reduced stability vs.  $\omega_\phi$  found before RWM becomes unstable ( $I_i = 0.49$ )

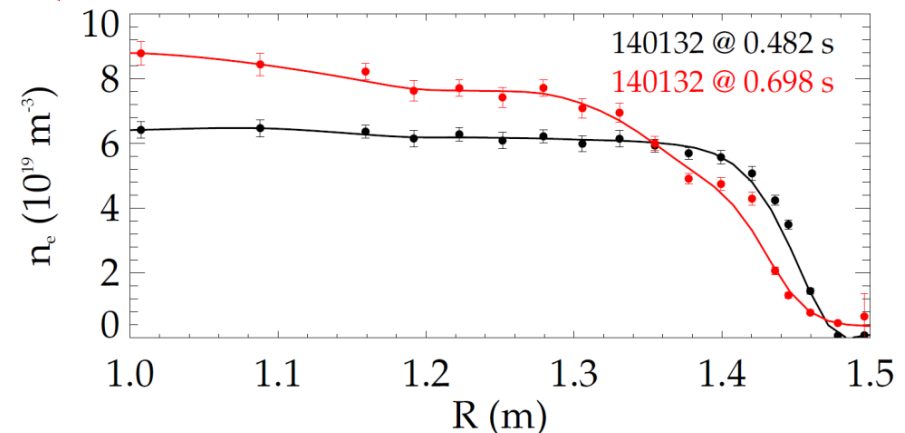
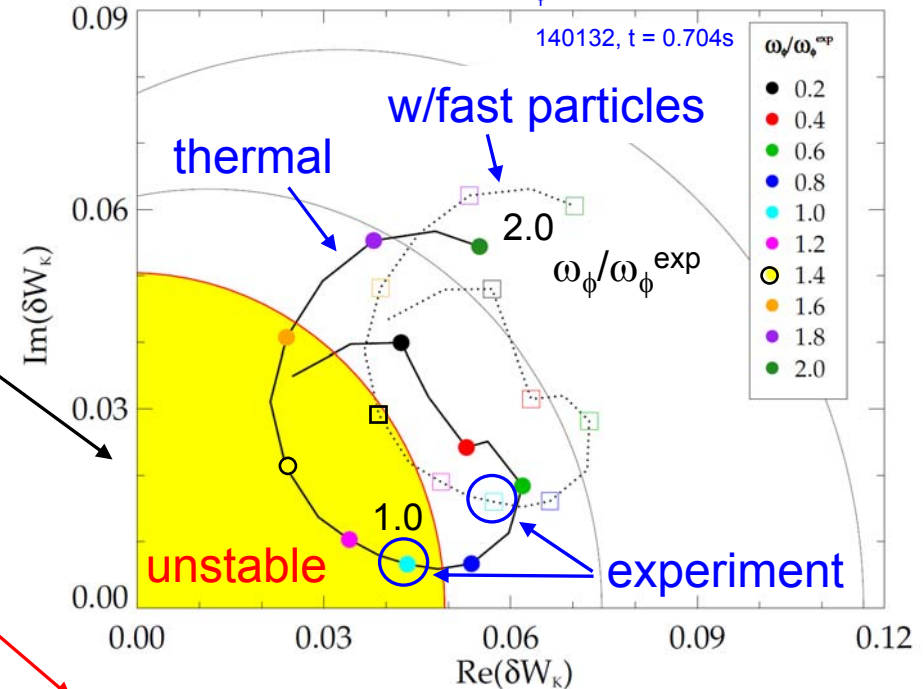
- Co-incident with a drop in edge density gradient – reduces kinetic stabilization

## MISK application to ITER (advanced scenario IV)

- RWM unstable at expected rotation
- Only marginally stabilized by alphas at  $\beta_N = 3$

- BP9.00057 J. Berkery, et al.
- Also, see poster for detail

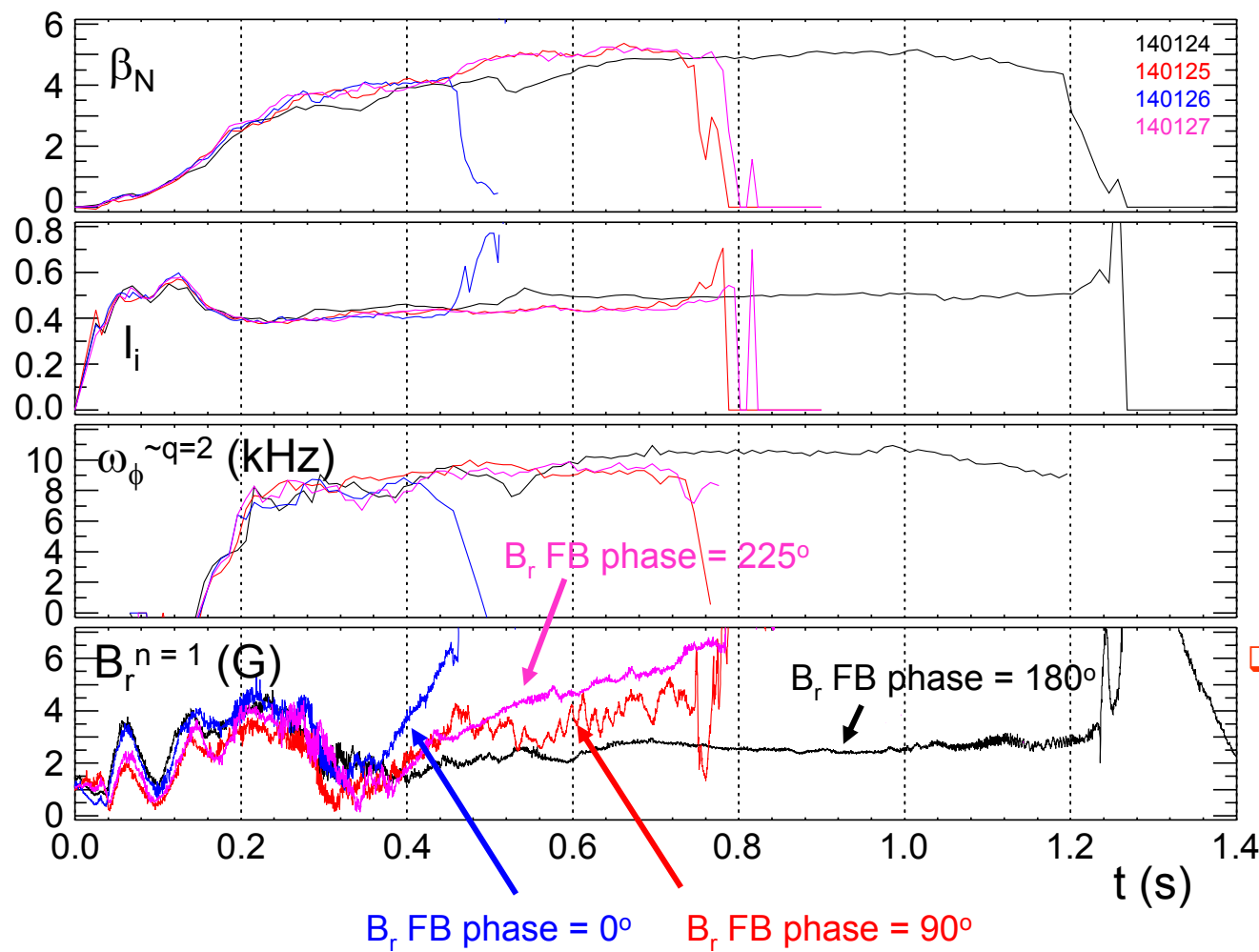
RWM stability vs.  $\omega_\phi$  (contours of  $\gamma\tau_{wv}$ )





# RWM $B_r$ sensor $n = 1$ feedback phase variation shows clear settings for improved feedback when combined with $B_p$ sensors

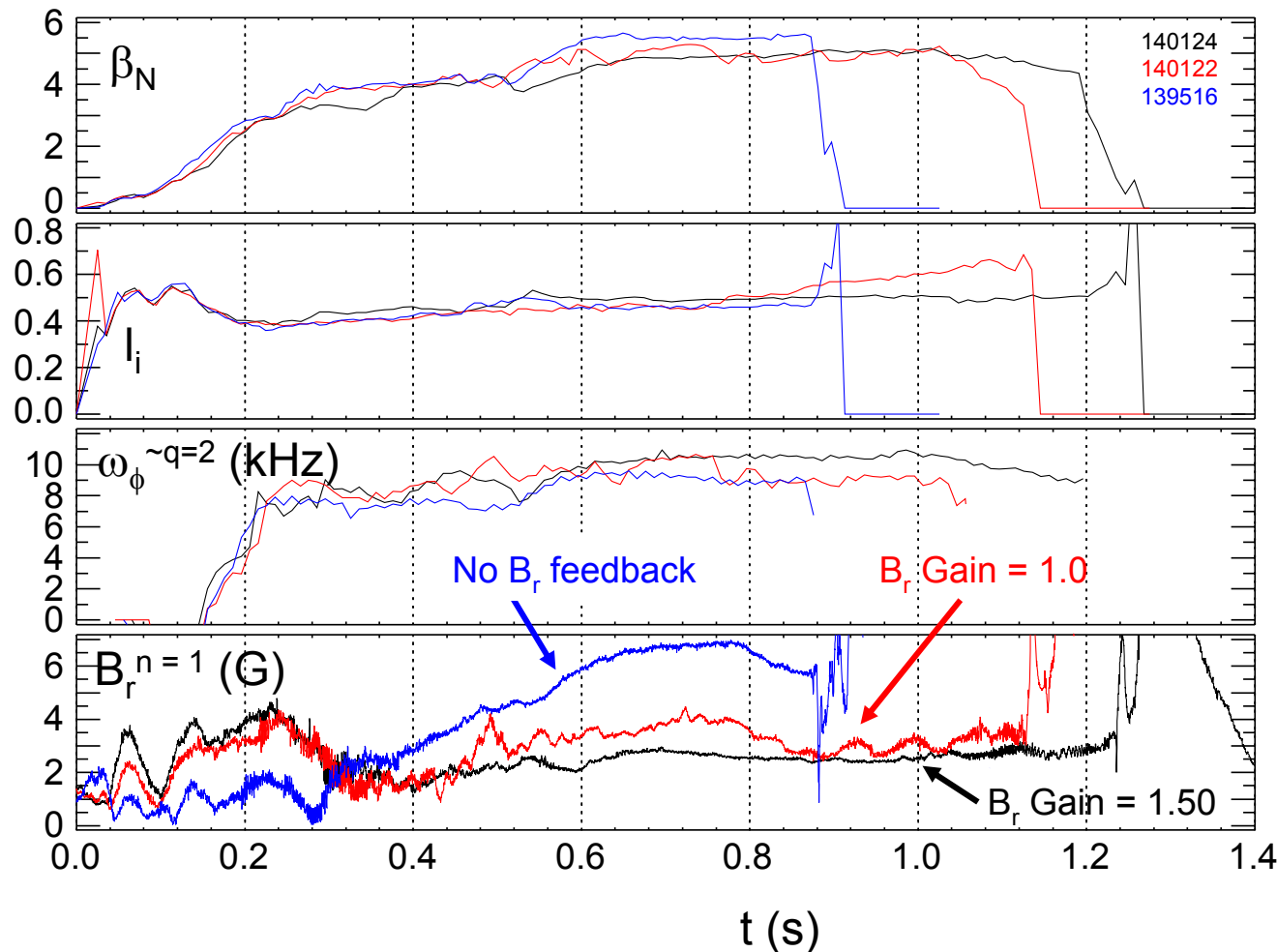
$n = 1 B_r + B_p$  feedback  
( $B_p$  gain = 1,  $B_r$  gain = 1.5)



- Recent corrections to  $B_r$  sensors improve measurement of plasma response
  - Removed significant direct pickup of time-dependent TF intrinsic error field
  - Positive/negative feedback produced at theoretically expected phase values
- Adjustment of  $B_p$  sensor feedback phase from past value further improved control performance

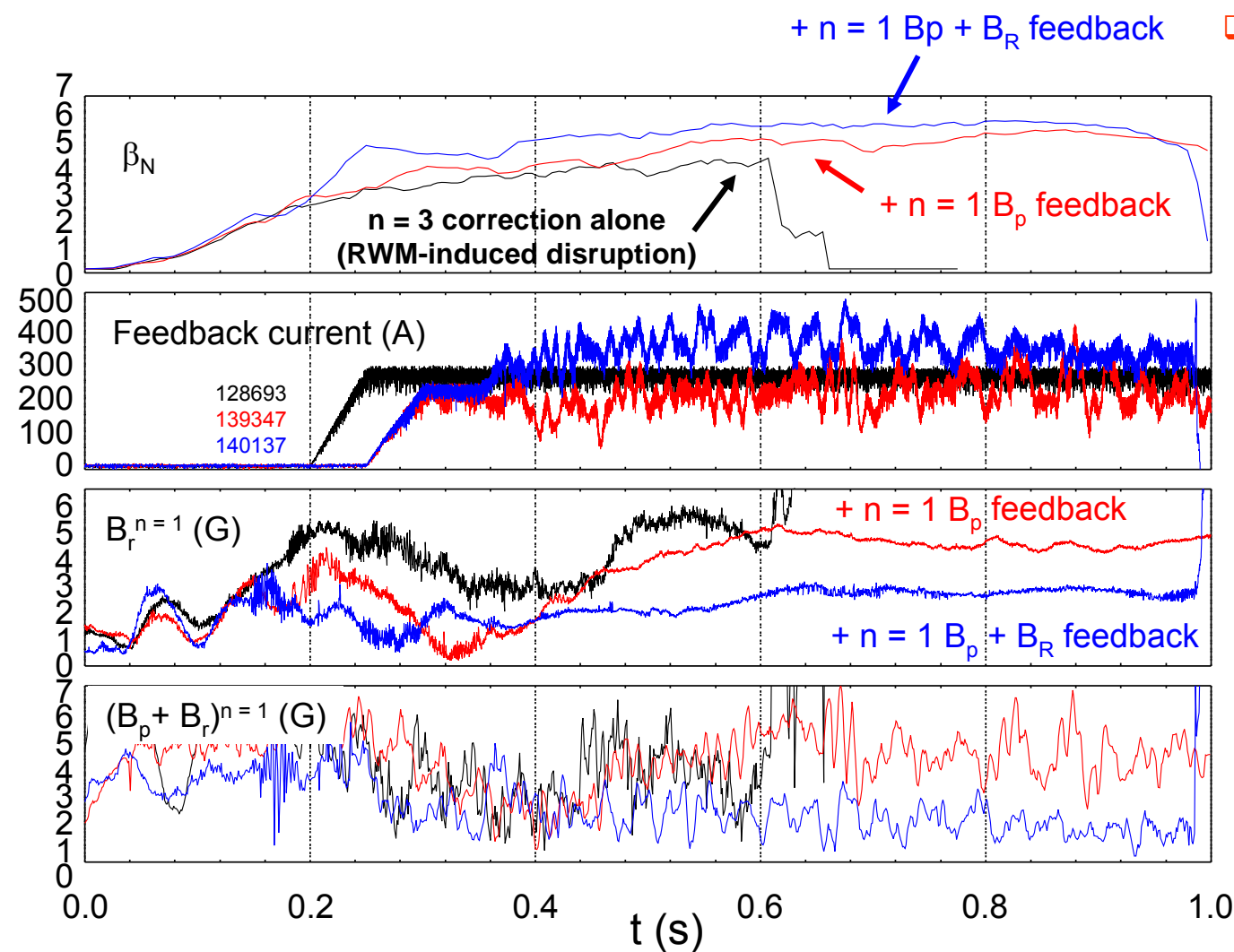
# RWM $B_R$ sensor feedback gain scan shows significantly reduced $n=1$ radial error field

$n=1 B_R + B_p$  feedback  
( $B_p$  gain = 1)



- New  $B_r$  sensor feedback gain scan on low  $I_i$  plasmas
  - Highest gain attempted (1.5) most favorable
  
- $B_r$  feedback constrains slow (10's of ms)  $n=1$  radial field growth
  - Addition of  $n=1$   $B_R$  sensors in feedback prevents disruptions when  $|\delta B_r^{n=1}| \sim 9\text{G}$ ; better sustains plasma rotation

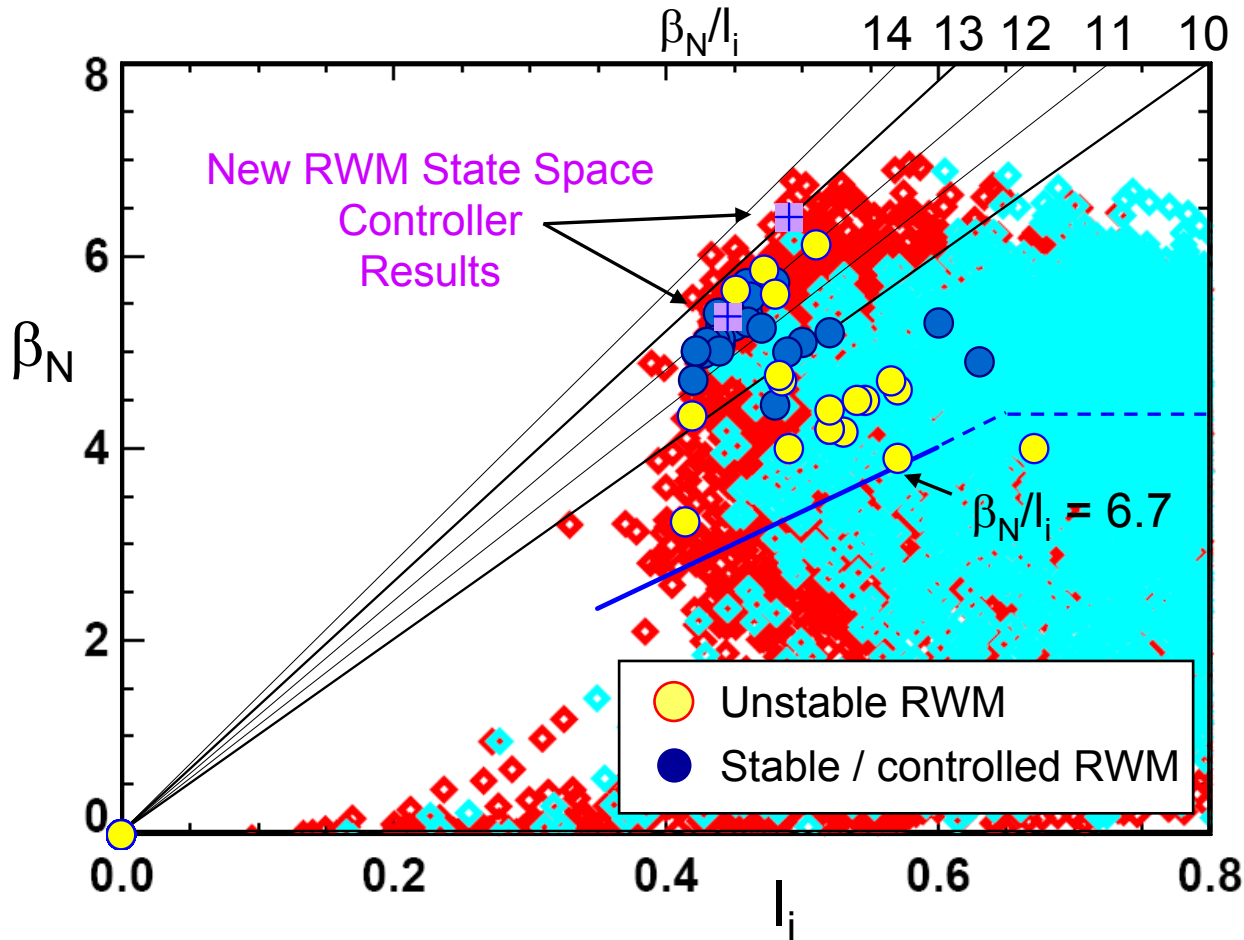
# Use of combined RWM sensor $n=1$ feedback yields best reduction of $n=1$ field amplitude / improved stability



□ Combination of DC error field correction,  $n=1$  feedback

- $n=3$  DC error field correction alone more subject to RWM instability
- $n=1$   $B_p$  sensor fast RWM feedback sustains plasma
- Addition of  $n=1$   $B_R$  sensors in feedback reduce the combined  $B_p + B_r$   $n=1$  field to low level (1–2 G)

# Improvements in stability control techniques significantly reduce unstable RWMs at low $I_i$ and high $\beta_N$



## Subset of discharges

- High  $I_p \geq 1.0\text{MA}$
- $n = 1$  control enhancements
- Mild  $\omega_\phi$  alteration

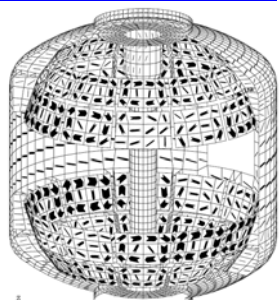
## Latest results

- Yielded low  $I_i$
- Access high  $\beta_N/I_i$
- Significantly reduced disruption probability due to unstable RWM

- 14% of cases with  $\beta_N/I_i > 11$
- Much higher probability of unstable RWMs at lower  $\beta_N$ ,  $\beta_N/I_i$

# New RWM state space controller implemented to sustain high $\beta_N$

Full 3-D model ~3000+ states



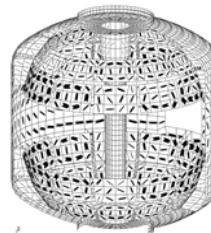
Balancing transformation

- device R, L, mutual inductances
- instability B field / plasma response
- modeled sensor response

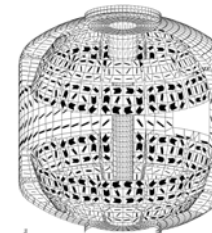
State reduction (< 20 states)

RWM eigenfunction (2 phases, 2 states)

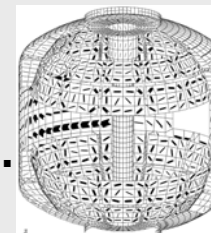
$(\hat{x}_1, \hat{x}_2)$



$\hat{x}_3$



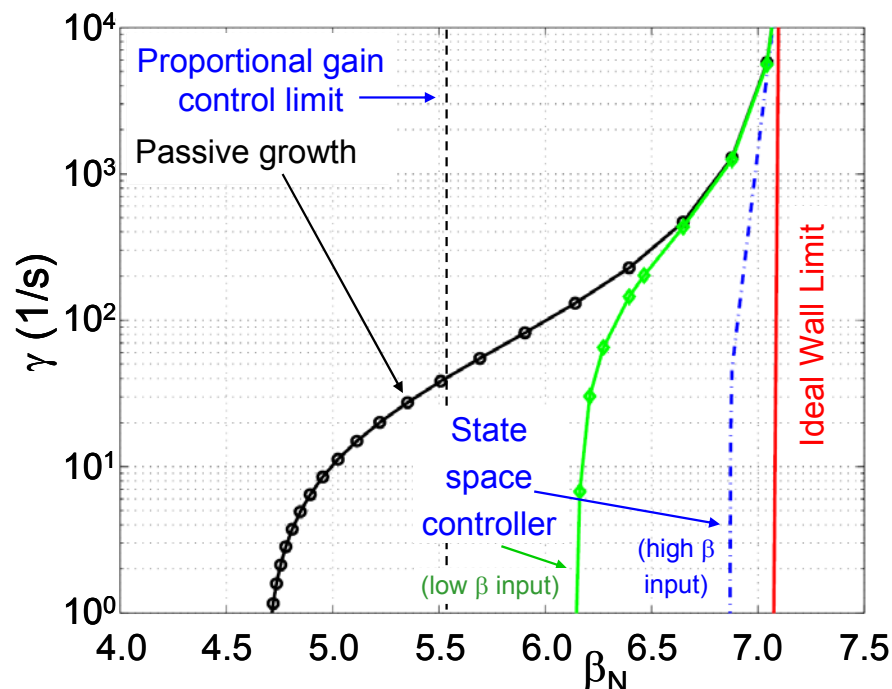
$\hat{x}_4$



$\hat{x}_N$   
truncate

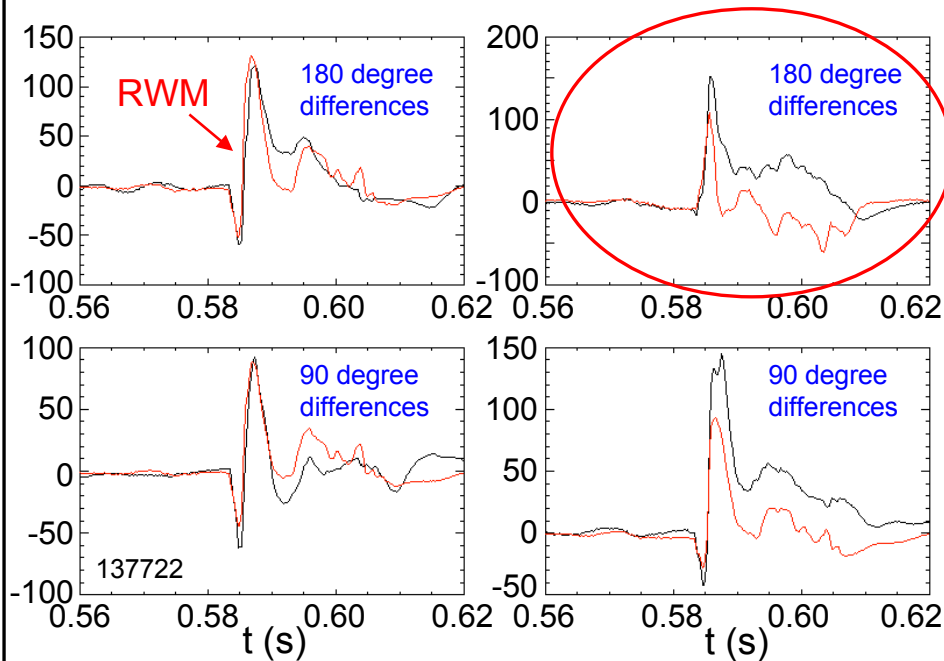
- Controller can compensate for wall currents
    - Including mode-induced current
  - Potential to allow more flexible control coil positioning
    - May allow control coils to be moved further from plasma, shielded
    - Examined for ITER
- Katsuro-Hopkins, et al., NF 47 (2007) 1157

Theoretical feedback performance ( $\omega_\phi = 0$ , 12 states)



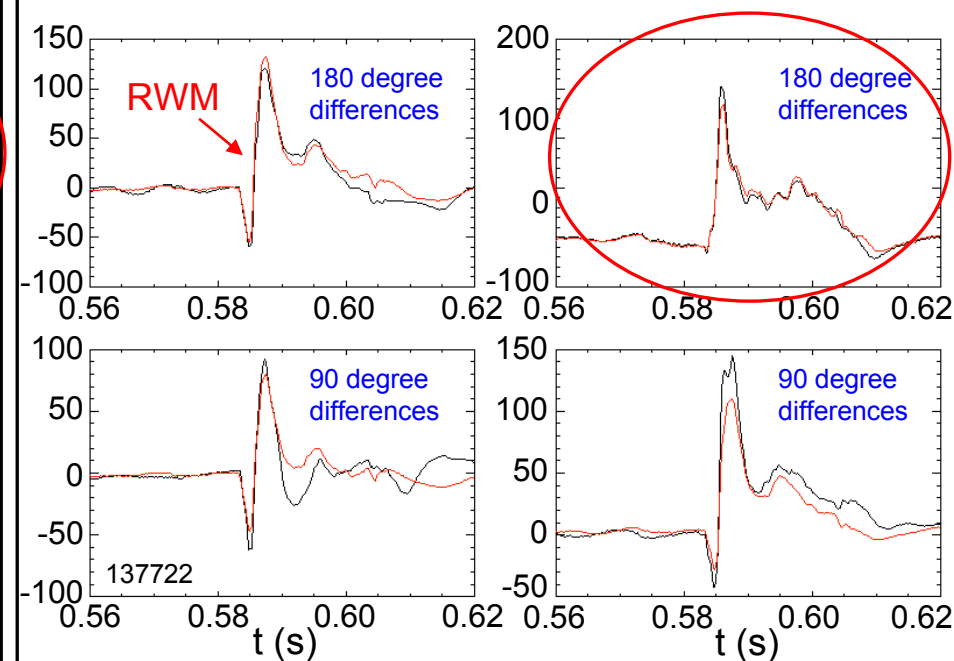
# Increased number of states in RWM state space controller improves match to sensors over entire mode evolution

RWM Upper  $B_p$  Sensor Differences (G) – 2 States



- Reasonable match to all  $B_p$  sensors during RWM onset, large differences later in evolution

RWM Upper  $B_p$  Sensor Differences (G) – 7 States

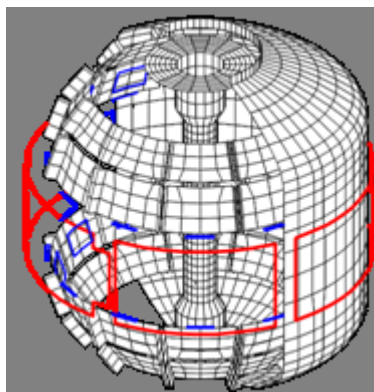


- Some 90 degree differences not as well matched
  - May indicate the need for an  $n = 2$  eigenfunction state

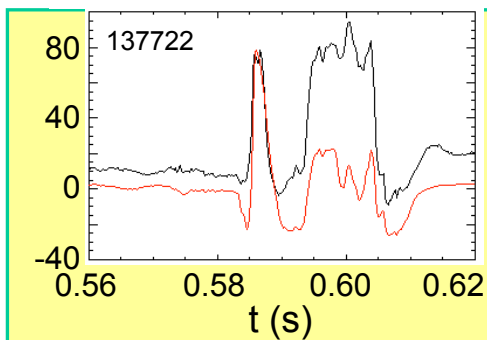
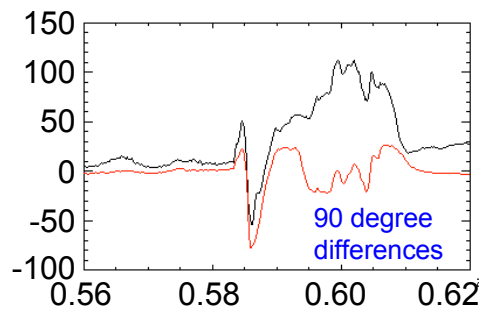
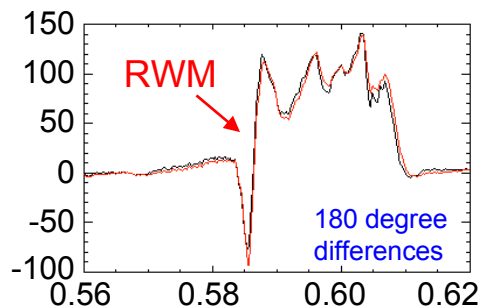
Black: experiment Red: offline RWM state space controller

# 3-D conducting structure detail can improve RWM state space controller match to sensors

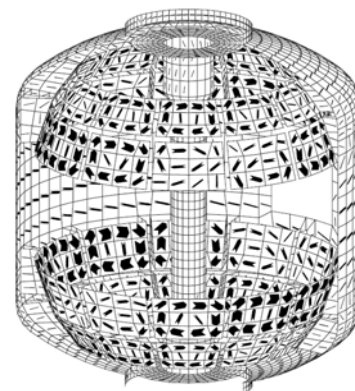
RWM Lower  $B_p$  Sensor Differences (G) – NO PORT



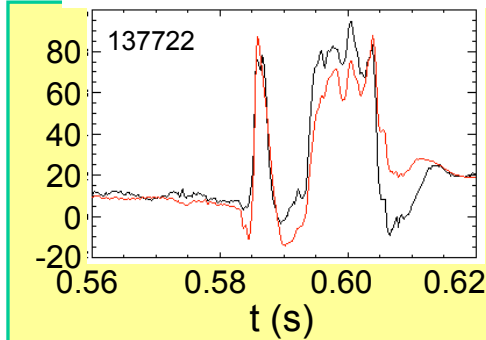
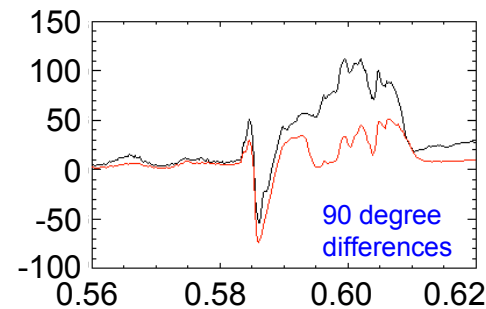
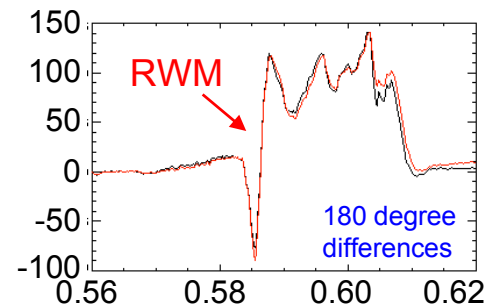
Some 90 degree differences not well matched



RWM Lower  $B_p$  Sensor Differences (G) – NBI PORT



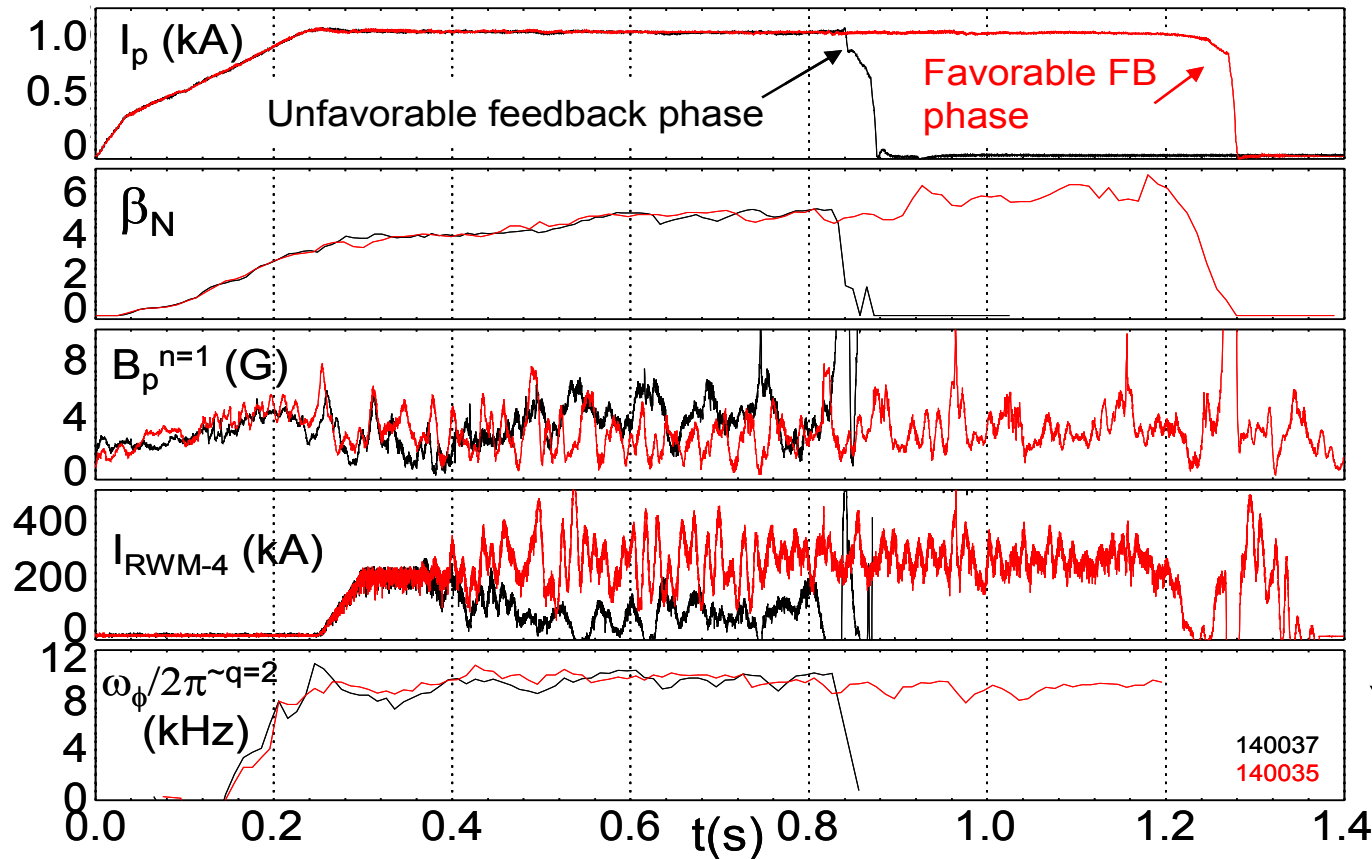
Adding NBI port leads to greater match on some sensors



Black: experiment Red: offline RWM state space controller

# New RWM state space controller sustains high $\beta_N$ , low $I_i$ plasma

## RWM state space feedback (12 states)



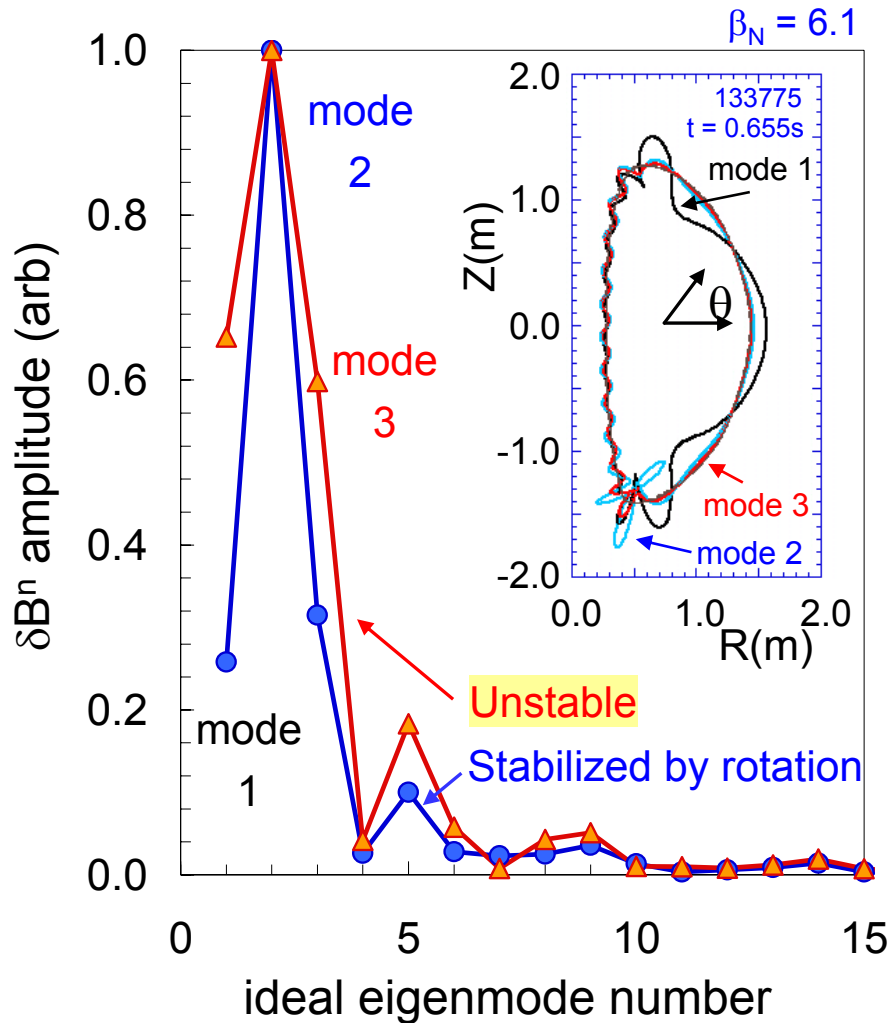
## Successful First Experiments

- $n = 1$  applied field suppression
  - Suppressed disruption due to  $n = 1$  field (See poster for detail)
- Feedback phase scan
  - Best feedback phase produced long pulse,  $\beta_N = 6.4$ ,  $\beta_N/I_i = 13$

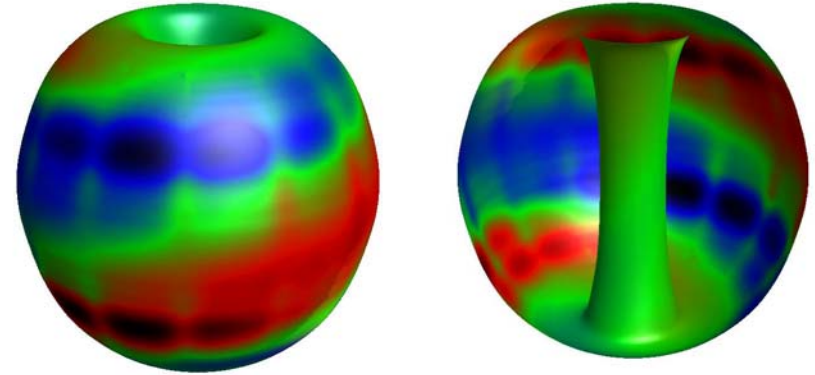


# Multi-mode RWM computation shows 2<sup>nd</sup> eigenmode component has dominant amplitude at high $\beta_N$ in NSTX stabilizing structure

## $\delta B^n$ RWM multi-mode composition



## $\delta B^n$ from wall, multi-mode response



- **NSTX RWM not stabilized by  $\omega_\phi$** 
  - Computed growth time consistent with experiment
  - 2<sup>nd</sup> eigenmode (“divertor”) has larger amplitude than ballooning eigenmode
- **NSTX RWM stabilized by  $\omega_\phi$** 
  - Ballooning eigenmode amplitude decreases relative to “divertor” mode
  - Computed RWM rotation  $\sim 41$  Hz, close to experimental value  $\sim 30$  Hz
- **ITER scenario IV multi-mode spectrum**
  - Significant spectrum for  $n = 1$  and 2

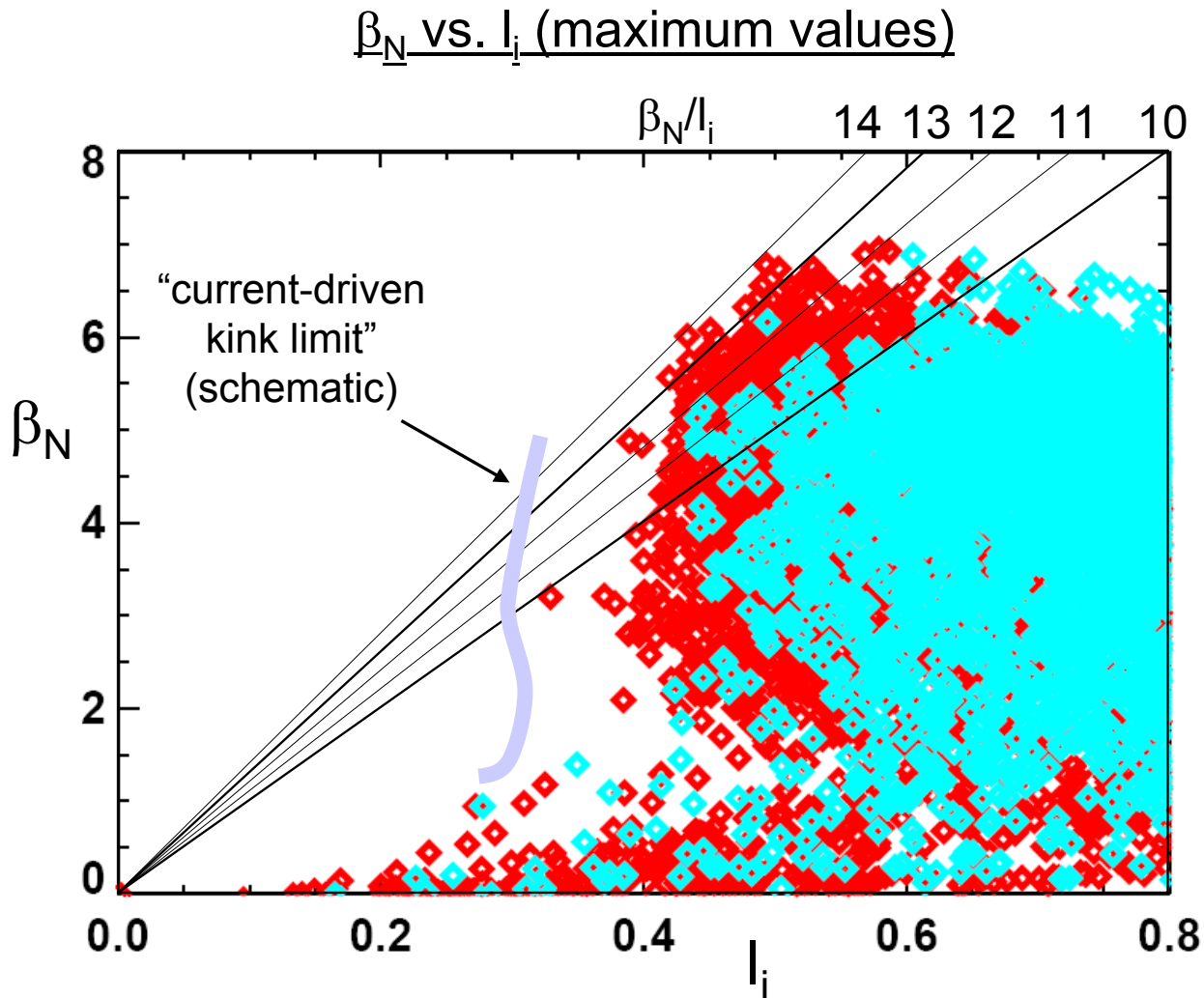
# NSTX is Addressing Global Stability Needs Furthering Steady Operation of High Performance ST Plasmas

- Success in producing and stabilizing high  $\beta_N$  plasmas with reduced  $I_i$ 
  - Approaching conditions needed for ST fusion applications
  - Incidence of RWM-induced disruption greatly reduced using control upgrades
- RWM instability observed at intermediate plasma rotation correlates with kinetic stability theory
  - Potential need for rotation control in future ST devices; evaluation of energetic particle (EP) stabilization
  - Initial analysis of low  $I_i$  plasmas indicates similar stability dependence on  $\omega_\phi$
- New RWM state space controller sustains low  $I_i$ , high  $\beta_N$  plasma
  - Potential for greater flexibility of RWM control coil placement in future burning plasma devices
- Computed multi-mode RWM spectrum at high  $\beta_N$  with 3D conducting structure shows significant amplitude of higher order ideal eigenmodes

# Backup and Poster Slides

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# Operational space is expanding to low $I_i$ and high $\beta_N$

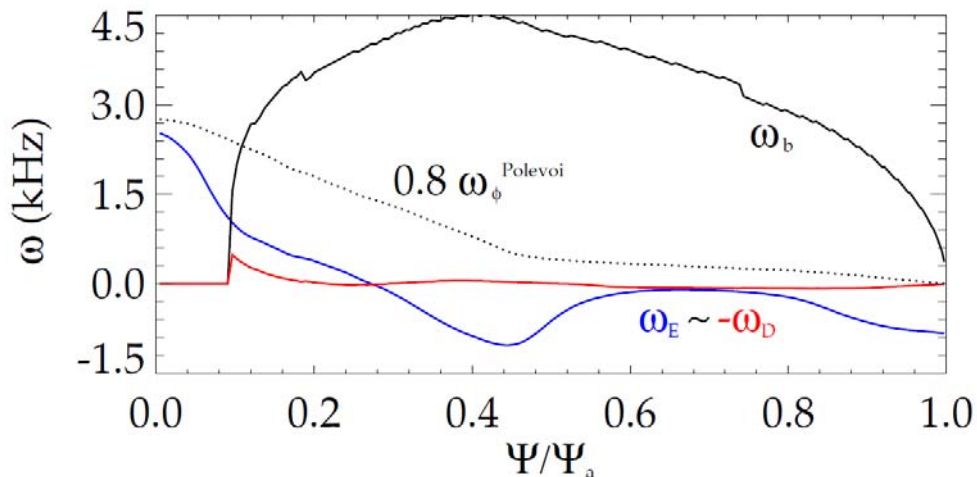


- $\beta_N/I_i$  is a common parameter to evaluate global stability
  - Kink/ballooning and RWM stability
- Significant increase in maximum  $\beta_N/I_i$ 
  - Upper limit now between 13 - 14
- At sufficiently low  $I_i$ , “current driven kink” limit exists
  - Plasma unstable at any  $\beta_N$  value without conducting wall, or feedback control

# ITER Advanced Scenario IV: RWM just reaches marginal stability by energetic particles with $\beta_N = 3$

## Equilibrium

- With  $\beta_N = 3$  (20% above  $n = 1$  no-wall limit)
- Plasma rotation profile linear in normalized poloidal flux

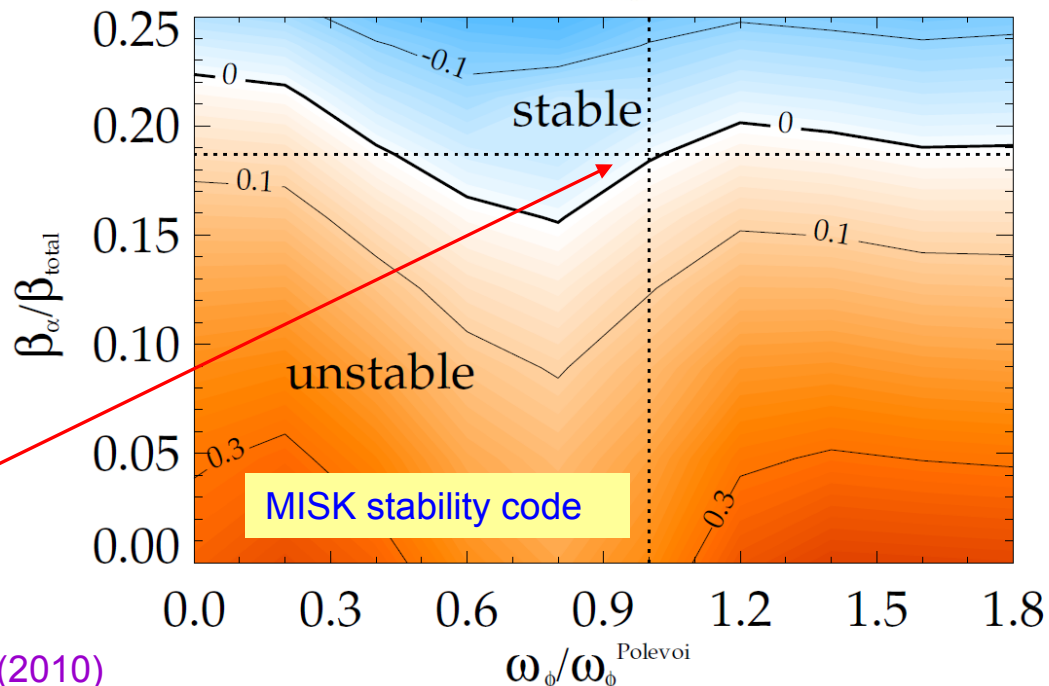


## Plasma rotation effect

- Stabilizing precession drift resonance weakly enhances stability near  $\omega_\phi = 0.8 \omega_\phi^{\text{Polevoi}}$

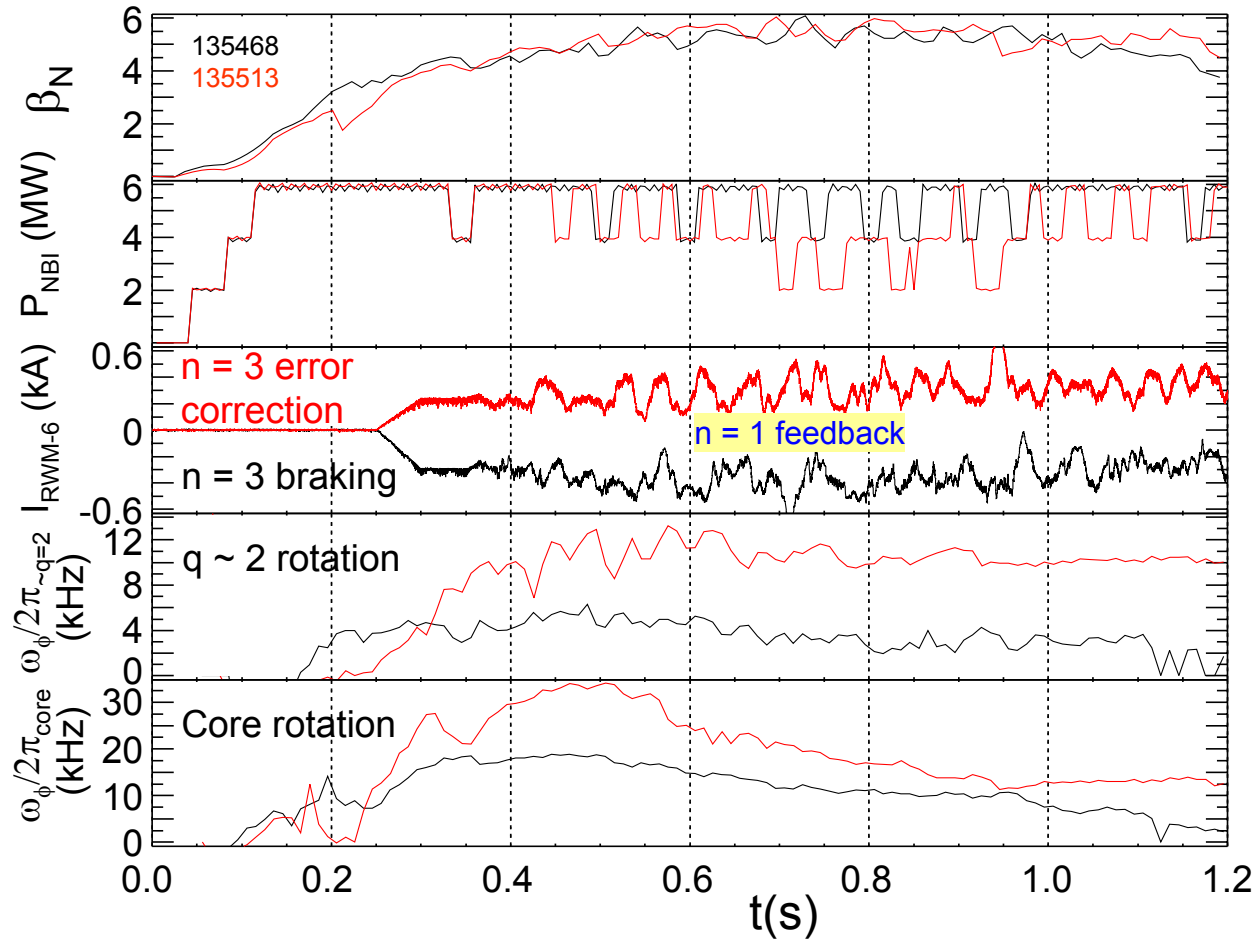
## Energetic particle (EP) effect

- Alpha particles are required for RWM stabilization at all  $\omega_\phi$
- Near RWM marginal stability at ITER expected  $\beta_\alpha/\beta_{\text{total}} = 0.19$  at  $\omega_\phi = \omega_\phi^{\text{Polevoi}}$



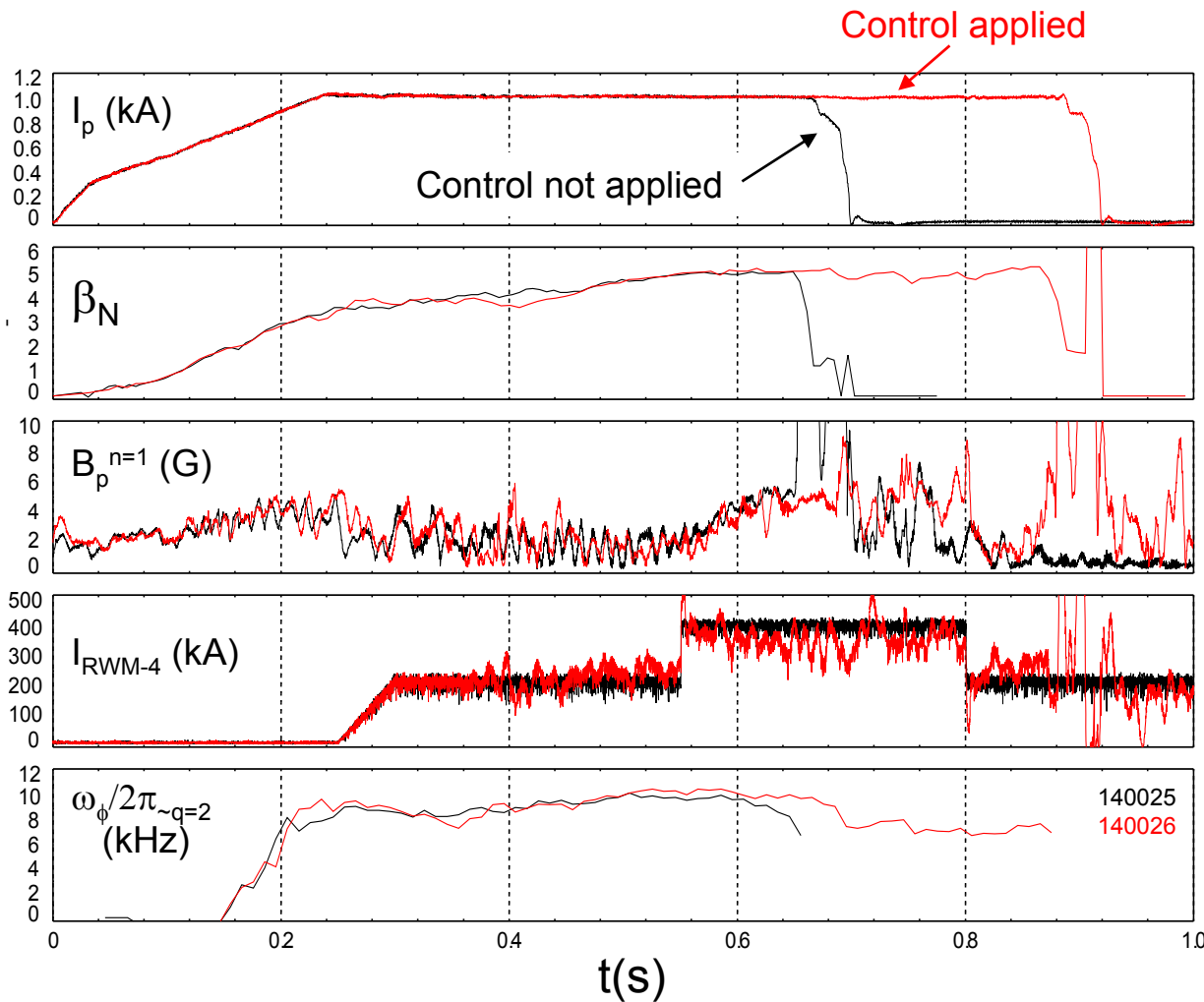
J.W. Berkery, et al., Phys. Plasmas **17**, 082504 (2010)

# $\beta_N$ feedback combined with $n = 1$ RWM control to reduce $\beta_N$ fluctuations at varied plasma rotation levels



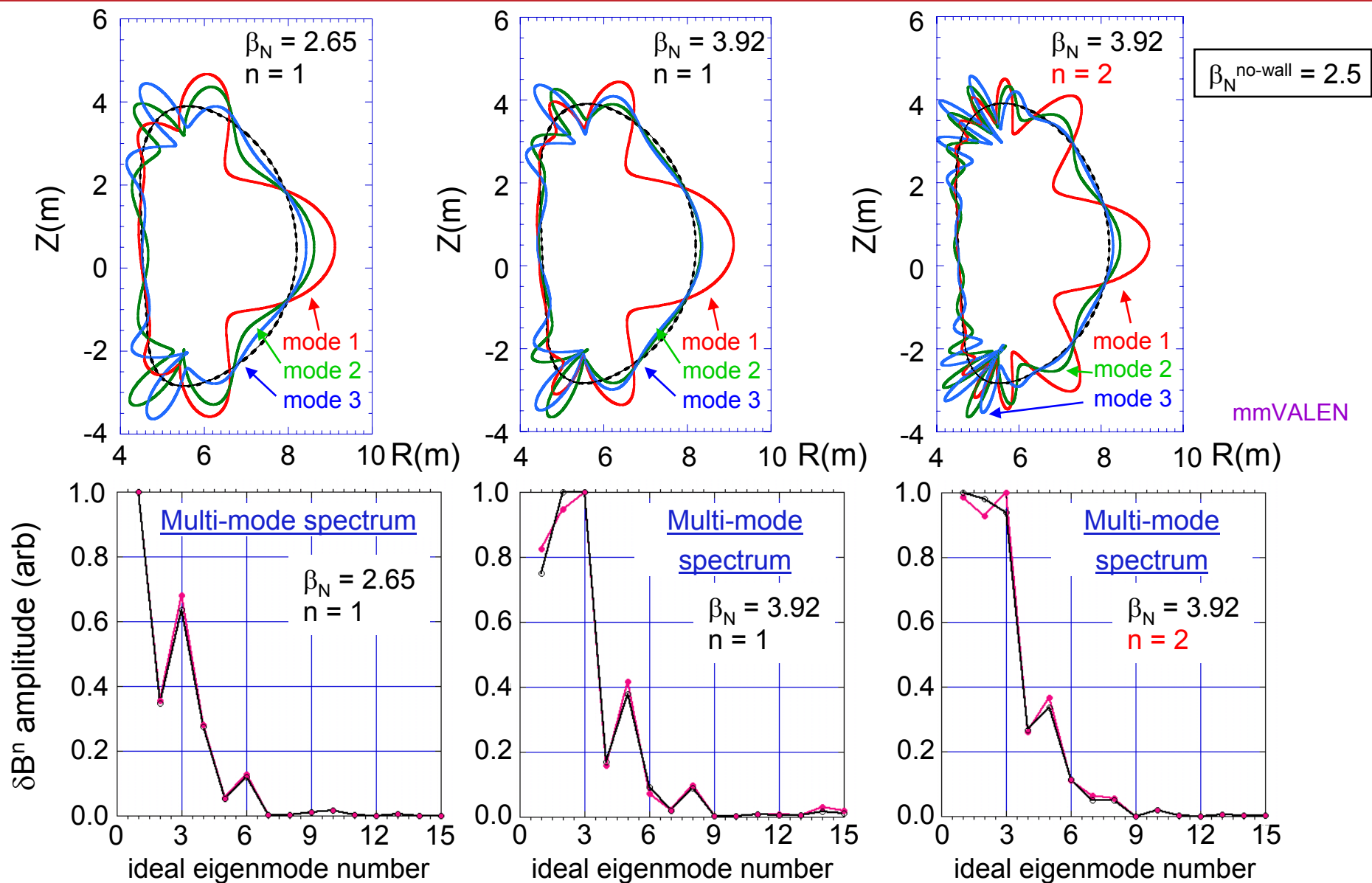
- Prelude to  $\omega_\phi$  control
  - Reduced  $\omega_\phi$  by  $n = 3$  braking is compatible with  $\beta_N$  FB control
- Steady  $\beta_N$  established over long pulse
  - independent of  $\omega_\phi$  over a large range
- Radial field sensors recently added to  $n = 1$  feedback

# RWM state space controller sustains otherwise disrupted plasma caused by DC n = 1 applied field



- n = 1 DC applied field
  - Simple method to generate resonant field amplification
  - Can lead to mode onset, disruption
- RWM state space controller sustains discharge
  - With control, plasma survives n = 1 pulse
  - n = 1 DC field reduced
  - Transients controlled and do not lead to disruption
  - **NOTE: initial run – gains NOT optimized**

# ITER Advanced Scenario IV: multi-mode RWM spectra computation shows significant ideal eigenfunction amplitude for several components





# ABSTRACT

Steady-state spherical torus plasmas for fusion applications, such as a component test facility or demonstration power plant, target operation with high non-inductive current fraction. These broad current profile targets have low values of plasma internal inductance,  $I_i$ , less than 0.4, near to the lower end of present NSTX operation. A key significance of this operation is that it approaches the purely current-driven ideal kink limit, which by definition exceeds the no-wall stability limit for all values of plasma normalized pressure (beta). In this regime, passive or active kink and resistive wall mode (RWM) stabilization is critical. Experiments on NSTX have recently approached this condition, evidenced by a significant reduction of the  $n = 1$  no-wall stability limit computed by DCON. This limit drops from normalized beta of 4.2 – 4.6 at  $I_i \sim 0.6$ , to 3.4 at  $I_i \sim 0.5$ , to below 2.8 for  $I_i \sim 0.4$ . Nevertheless, passive and active RWM control has produced high toroidal beta up to 28 percent, and normalized beta up to 6.5 (nearly double the no-wall limit), closely following a record normalized beta to  $I_i$  ratio of 13 between  $I_i = 0.4 - 0.5$ . Non-inductive current fraction reaches 0.5 in these high normalized current plasmas. However, the disruption probability of these plasmas increases significantly, with about half of the discharges suffering terminating instabilities. Alteration of  $n = 1$  RWM control system parameters, plasma rotation profile, and the role of beta feedback is examined to potentially improve mode stability. Ion precession drift and bounce frequency resonance stabilization is examined for these plasmas and compared to the identified stabilization reduction at intermediate plasma rotation and higher  $I_i$  [1,2].

[1] J.W. Berkery, et al., Phys. Rev. Lett. **104**, 035003 (2010); J.W. Berkery, et al., PoP **17**, 082504 (2010)

[2] S.A. Sabbagh, et al., Nucl. Fusion **50**, 025020 (2010)

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