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Non-linear dynamics of toroidicity-induced Alfvén eigenmodes on NSTX

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23rd IAEA Fusion Energy Conference

Daejeon, Rep. of Korea

11-16 October 2010

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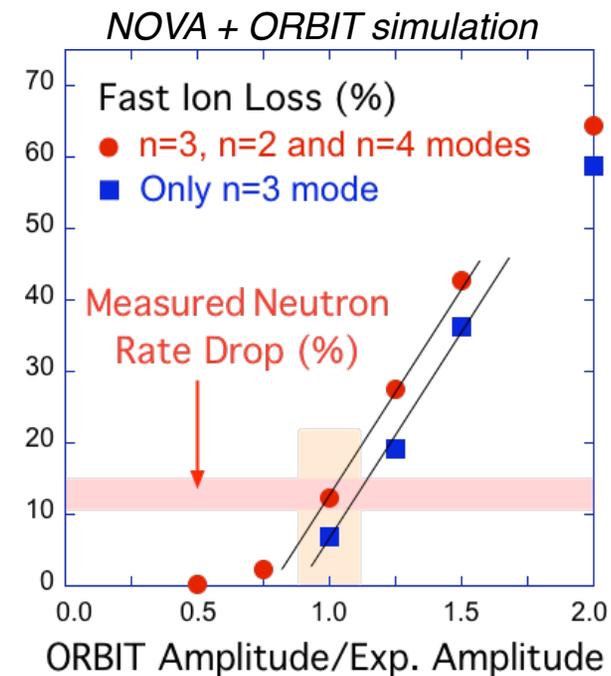
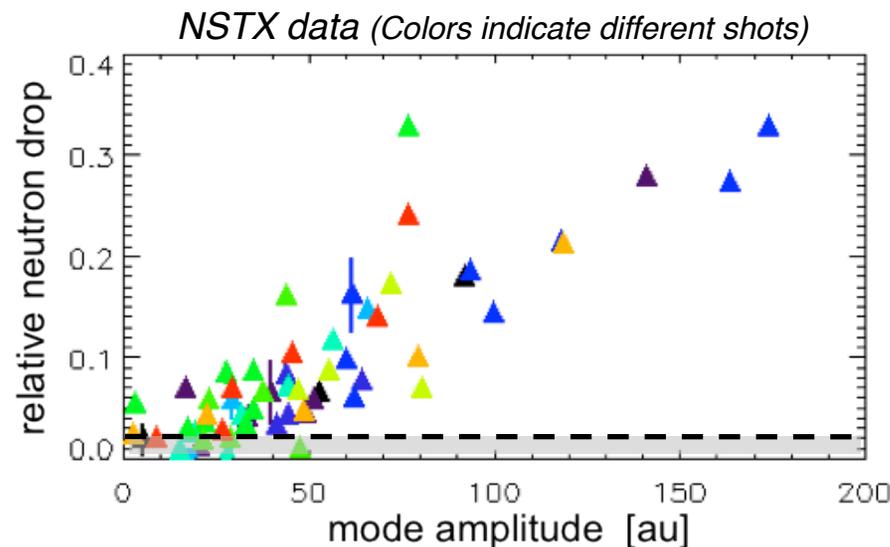
Abstract

The National Spherical Torus Experiment (NSTX) routinely operates with neutral beam injection as the primary system for heating and current drive. The resulting fast ion population is super-Alfvénic, with velocities $1 < v_{fast}/v_{Alfvén} < 5$ and normalized Larmor radius comparable to that of alphas in future reactors. Fast ions provides a strong drive for toroidicity-induced Alfvén eigenmodes (TAEs) with toroidal mode number $n = 2-8$ and frequency $60 < f < 250$ kHz. As the discharge evolves, the fast ion population builds up and TAEs exhibit increasing bursts in amplitude and down-chirps in frequency, which eventually lead to a so-called TAE *avalanche*. Avalanches cause large (up to ~30%) losses over < 1 ms, as inferred from the neutron rate and other fast ion diagnostics (FIDA, sFLIP, NPA). The increased fast ion losses correlate with a stronger activity in the TAE band. In addition, a $n = 1$ mode with frequency well below the TAE gap appears in the Fourier spectrum of magnetic fluctuations during avalanche events. The non-linear coupling between modes, which correlates with an enhanced fast ion transport during avalanches, is investigated.

Work supported by U.S. DOE contracts DE-AC02-09CH11466, DE-FG02-06ER54867 and DE-FG02-99ER54527

Bursting toroidicity-induced Alfvén eigenmodes (TAEs) can lead to enhanced fast-ion transport

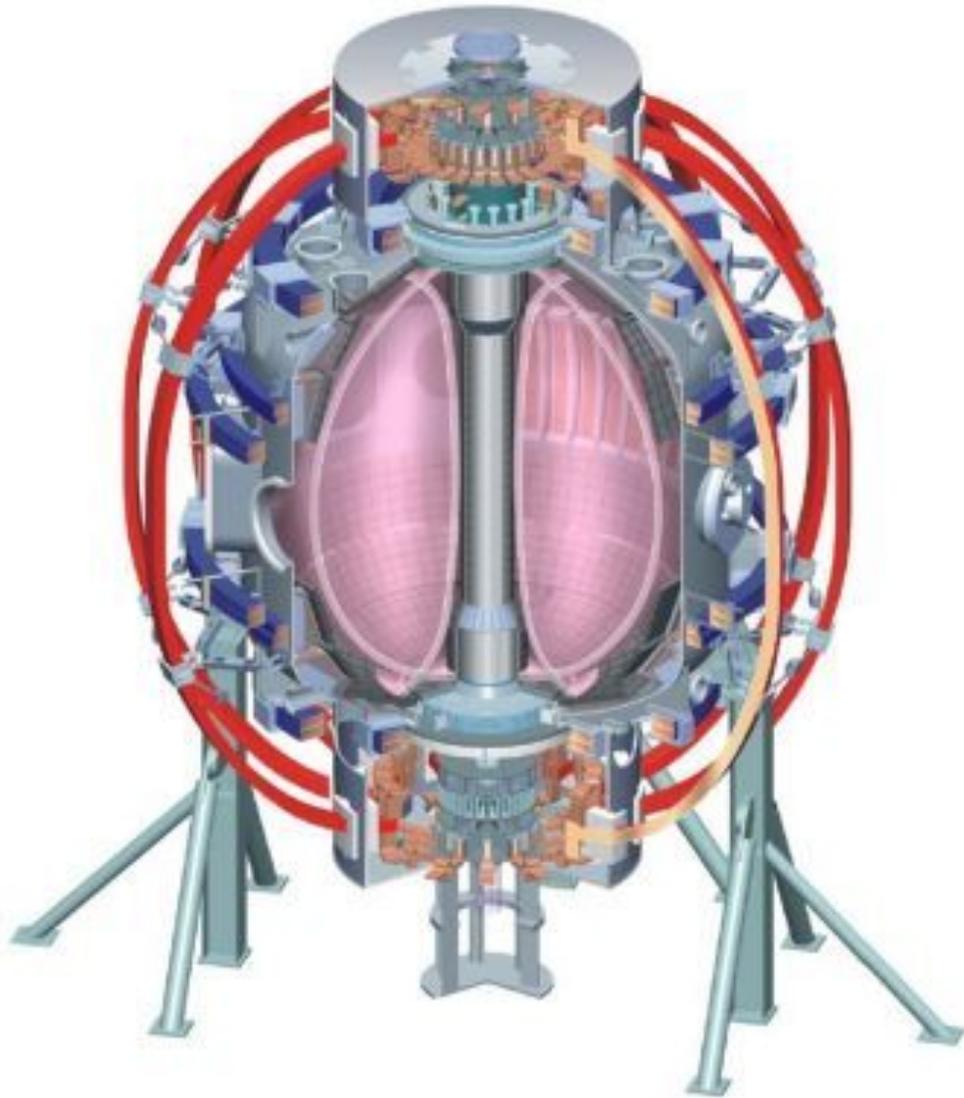
- Multiple TAEs can be simultaneously destabilized
 - Possible overlap of many resonances in phase space
 - Non-linear development into “TAE *avalanches*”



E. Fredrickson *et al.*,
Phys. Plasmas **16**, 122505 (2009)

- Must control/limit TAEs in future reactors (ITER, STs)
 - Need to understand the causes of bursting TAE behavior
 - Need to improve predictive capability

NSTX parameters



Major radius 0.85 m

Aspect ratio 1.3

Elongation 2.7

Triangularity 0.8

Plasma current ~ 1 MA

Toroidal field < 0.6 T

Pulse length < 2 s

3 Neutral Beam sources

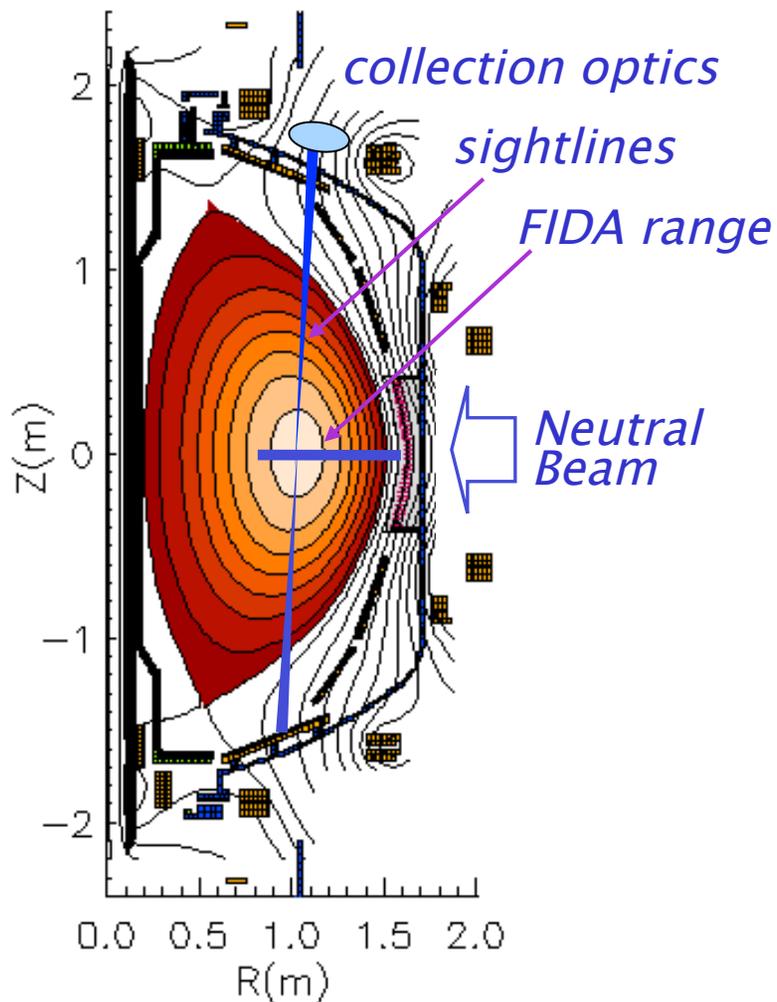
$P_{\text{NBI}} \leq 6$ MW, $E_{\text{injection}} \leq 95$ keV

$1 < v_{\text{fast}}/v_{\text{Alfvén}} < 5$

This work:
Focus on TAEs in L-mode plasma
Center-stack limited
Deuterium plasma
 $B_{\text{tor}} = 0.55$ T, $I_p = 0.7-0.9$ MA

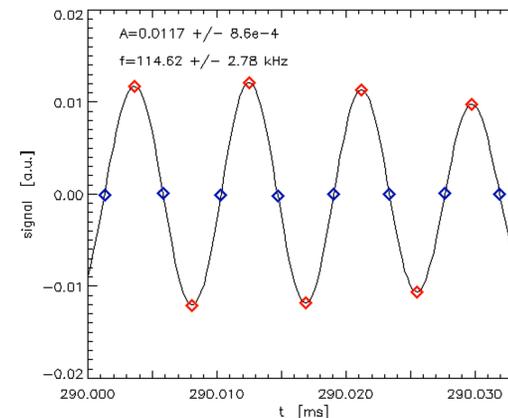
Mode activity and fast ion diagnostics on NSTX

shot#135404, t=320 ms



- Mirnov coils
 - Magnetic fluctuations up to 2.5 MHz
- Multi-channel reflectometer
 - Mode structure (L-mode)

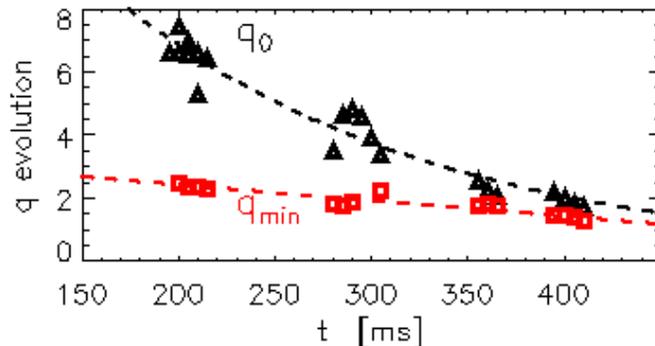
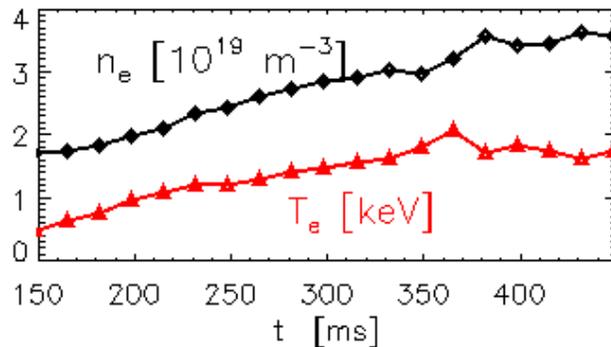
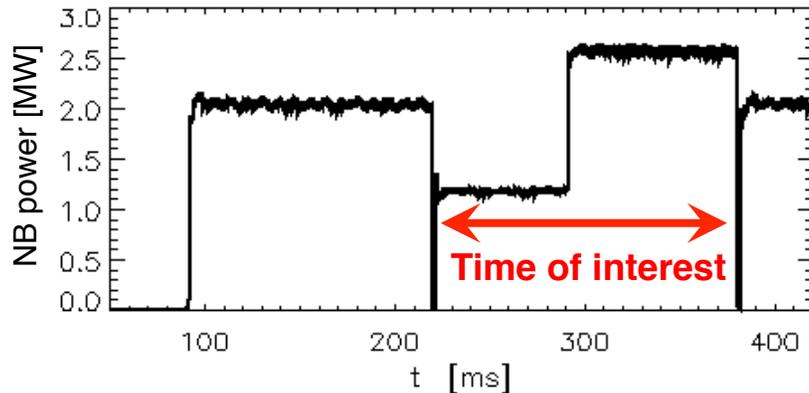
FFT analysis complemented by analysis in time domain to study mode dynamics over short time scale



- Fast Ion D-Alpha (FIDA) system
 - Fast ion profile and spectrum through active charge-exchange recombination spectroscopy
 - Weighted toward small pitch (perp. component)
- Neutron rate, NPA, sFLIP

Experimental scenario :

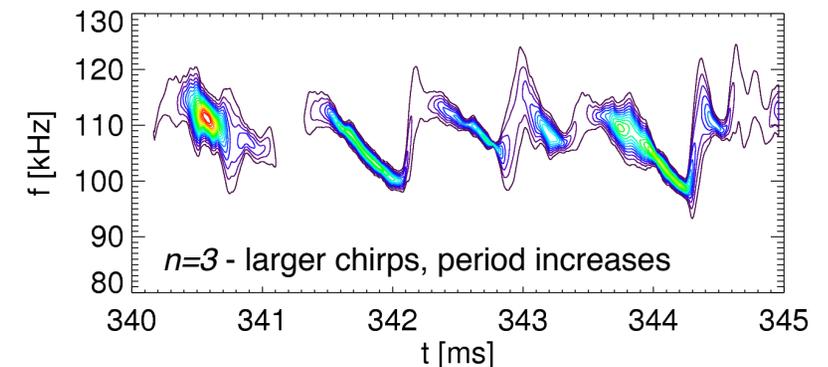
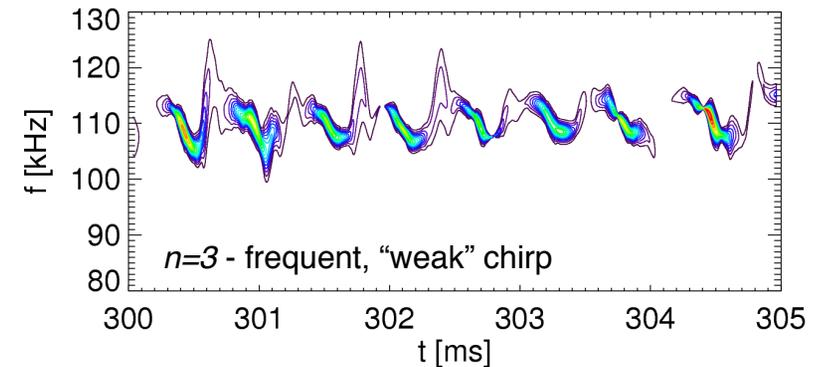
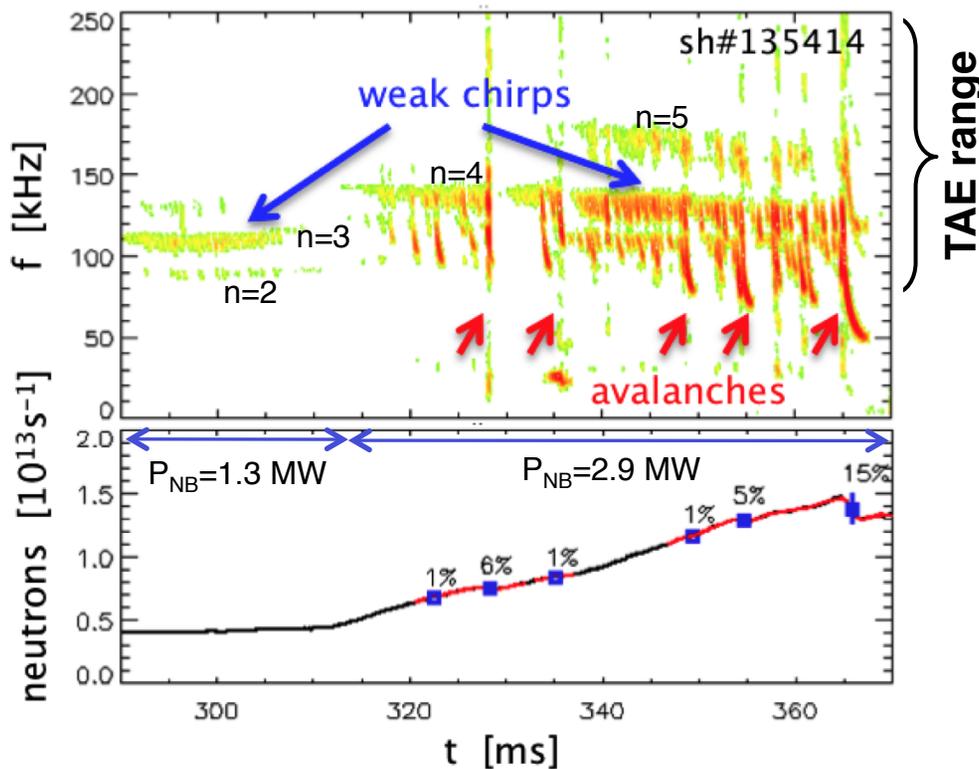
$P_{NB} < 3\text{MW}$, $n_e \sim 3 \times 10^{19} \text{m}^{-3}$, $T_i \sim T_e = 1 - 1.5 \text{keV}$



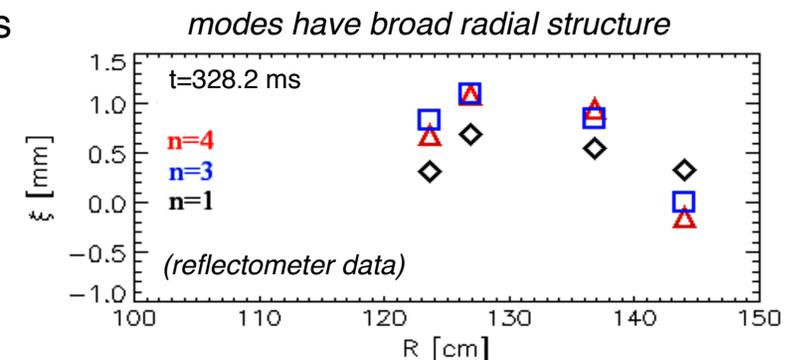
- NB-heated, L-mode plasmas

- Plasma limited on center-stack
- NB power and timing varied to affect mode stability
- Plasma profiles evolving in time
- Reversed-shear q profile
- Safety factor evolution reconstructed from four similar discharges through LRDFIT code constrained by MSE data

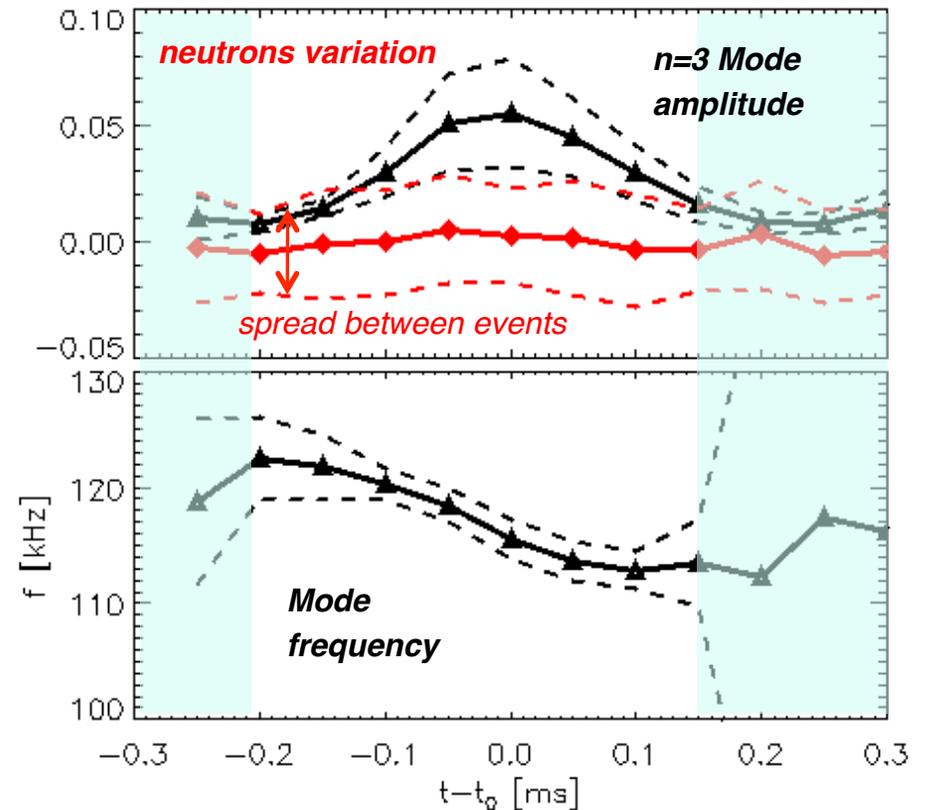
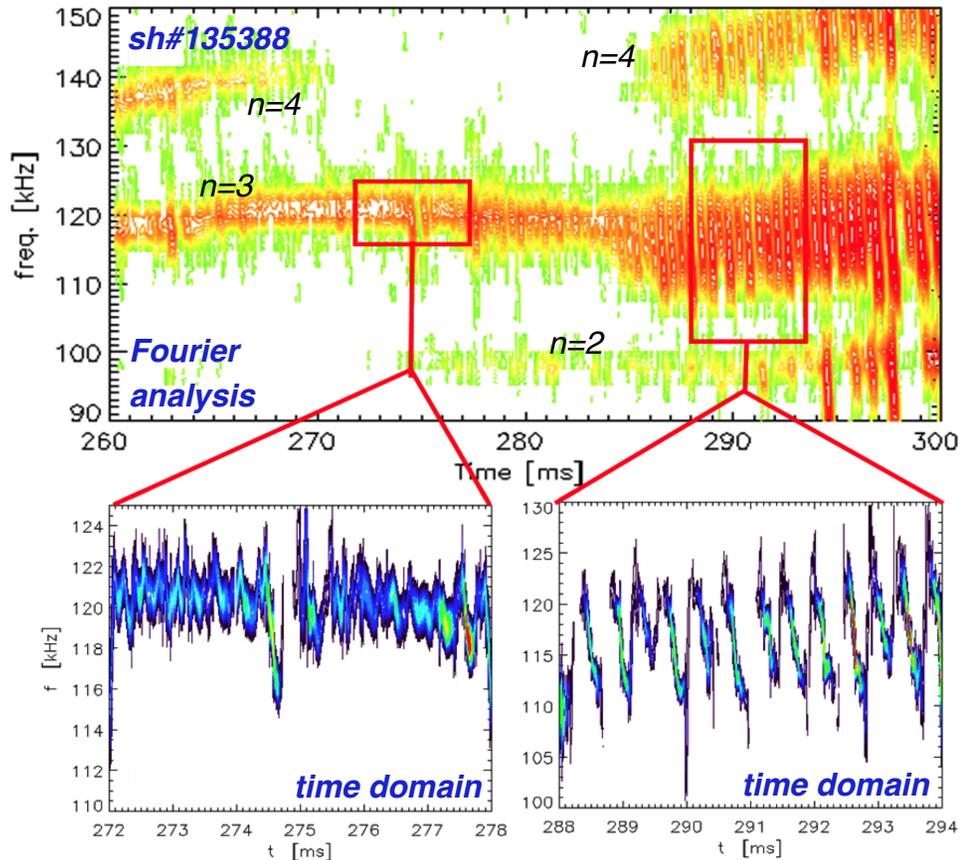
TAEs with low toroidal mode number ($n=2 \rightarrow 7$) are observed, with dominant $n=2-4$ modes



- Modes show more bursting character as discharge evolves
 - NB power increases, fast ion population builds up
- Usually, each mode chirps independently of the others...
- ... but, eventually, *avalanches* occur:
 - Modes lock on similar dynamic, multiple TAEs involved
 - Drop in neutron rate, FIDA

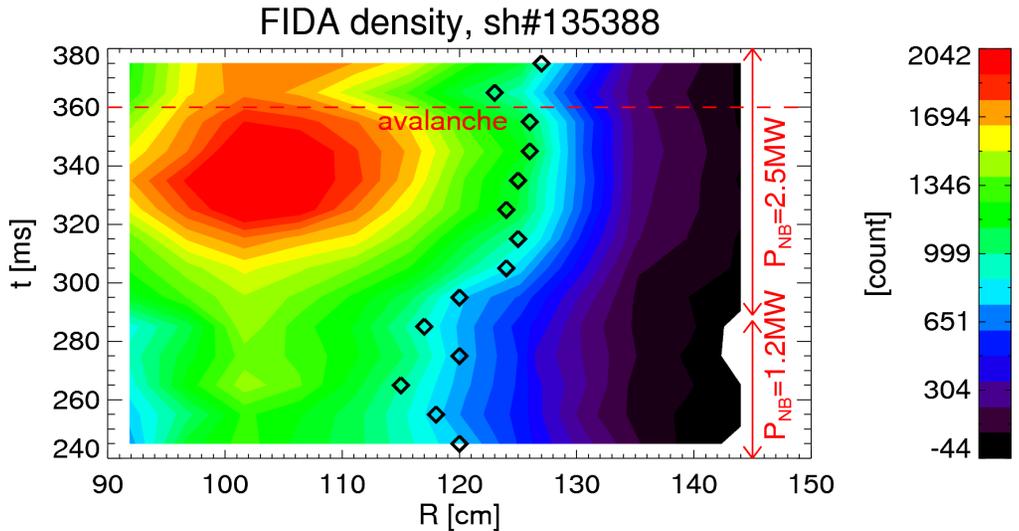


No detectable fast ion losses are observed during weakly bursting/chirping phase

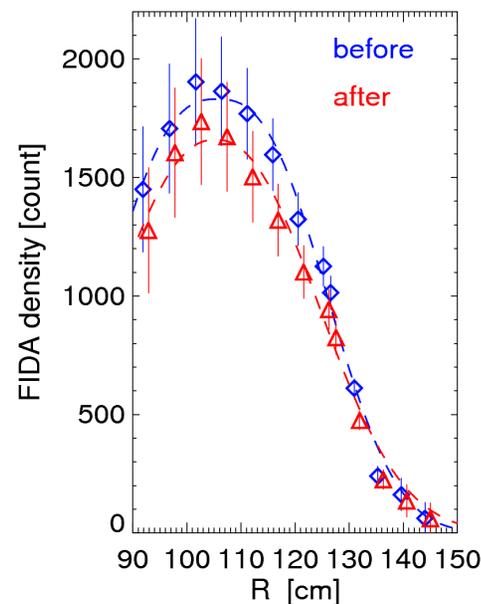
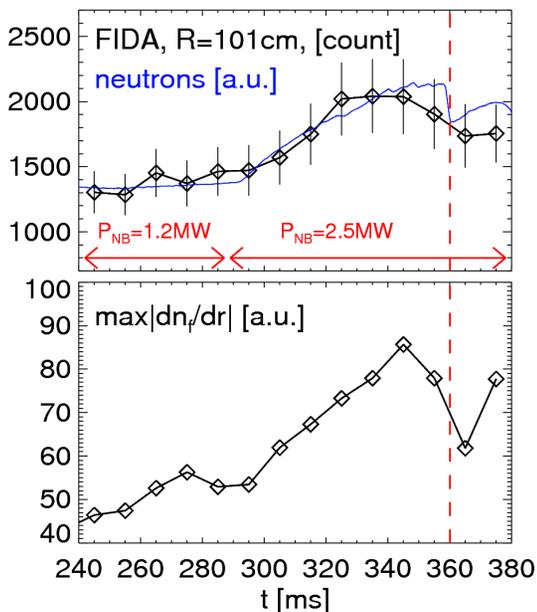


- Statistical average over ~20 events (~10ms)
- No clear evidence of losses from neutrons, FIDA
 - Does not exclude “continuous” (non-bursting) losses

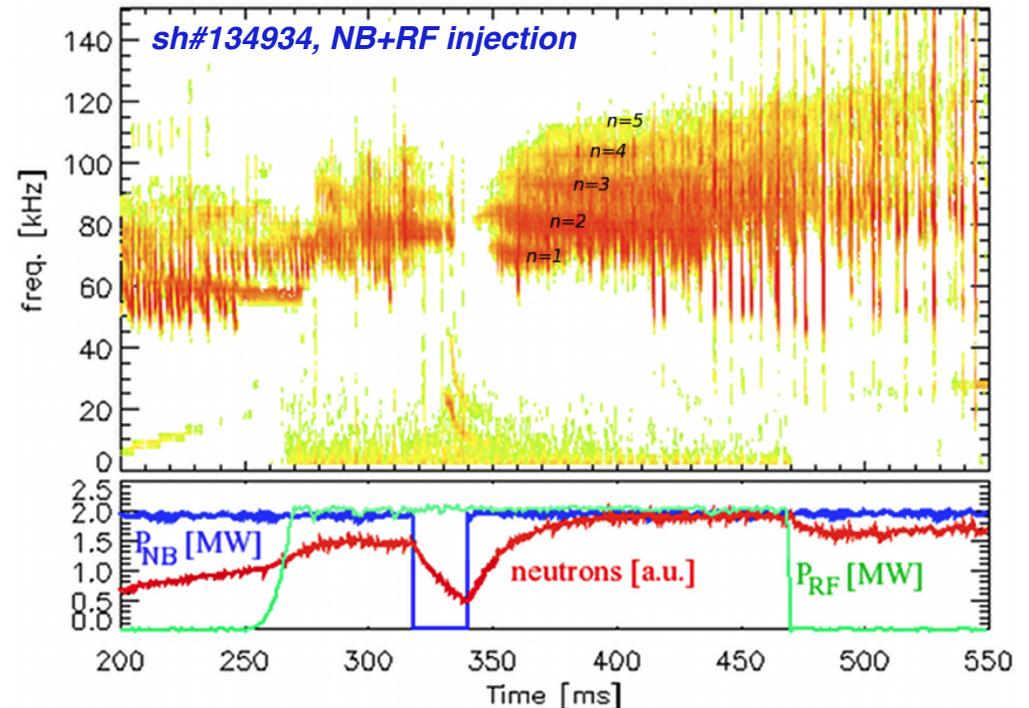
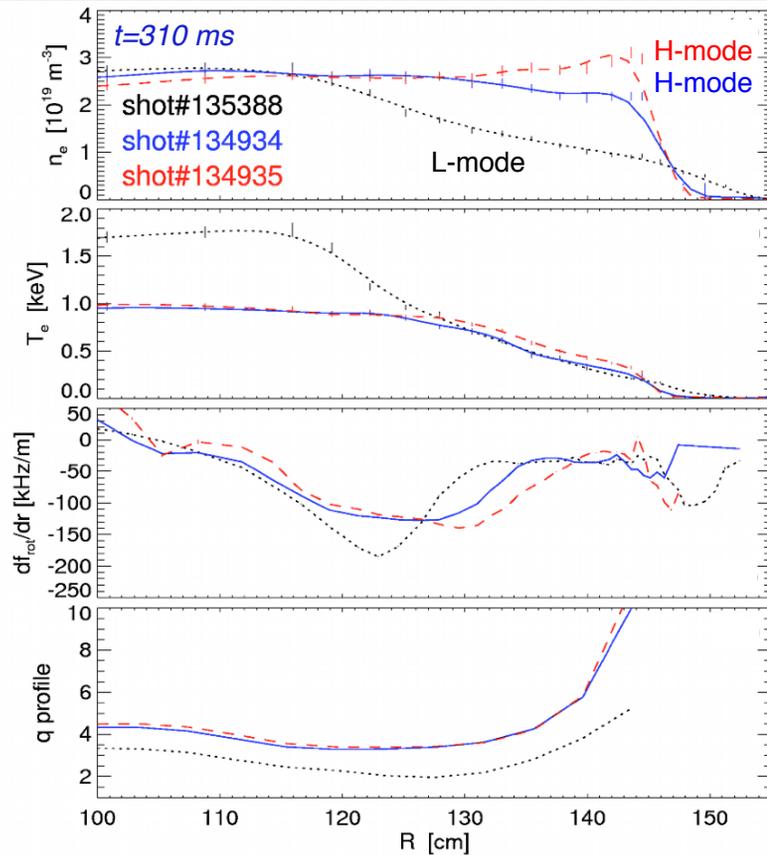
Up to ~30% of fast ions can be lost during a single TAE avalanche



- Fast ion density (FIDA) drops over most of minor radius
- Loss results in a relaxation of the radial gradient \rightarrow drive for TAEs is reduced
- Comparable losses estimated from FIDA and neutron rate
- Losses increase with (total) mode amplitude



Similar features are observed in L- and H-mode plasmas and during combined NB+RF : robust dynamics



- Example: H-mode discharges with NB and NB+RF heating
 - Different profiles with respect to L-mode
 - Higher safety factor than for L-mode discharges
 - Reversed shear in both L- and H-mode

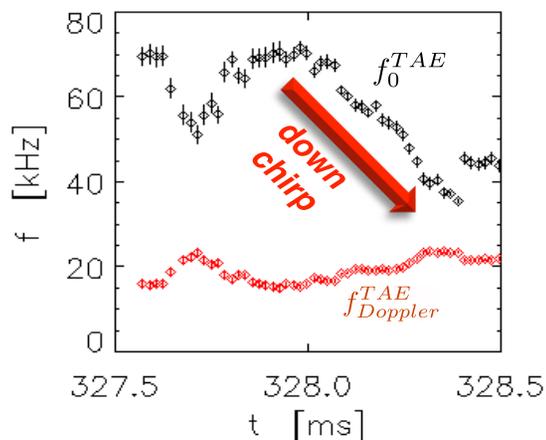
On average, TAE frequencies are consistent with a common frequency *in the plasma frame*

- $f_{n=2-6}$ consistent with:

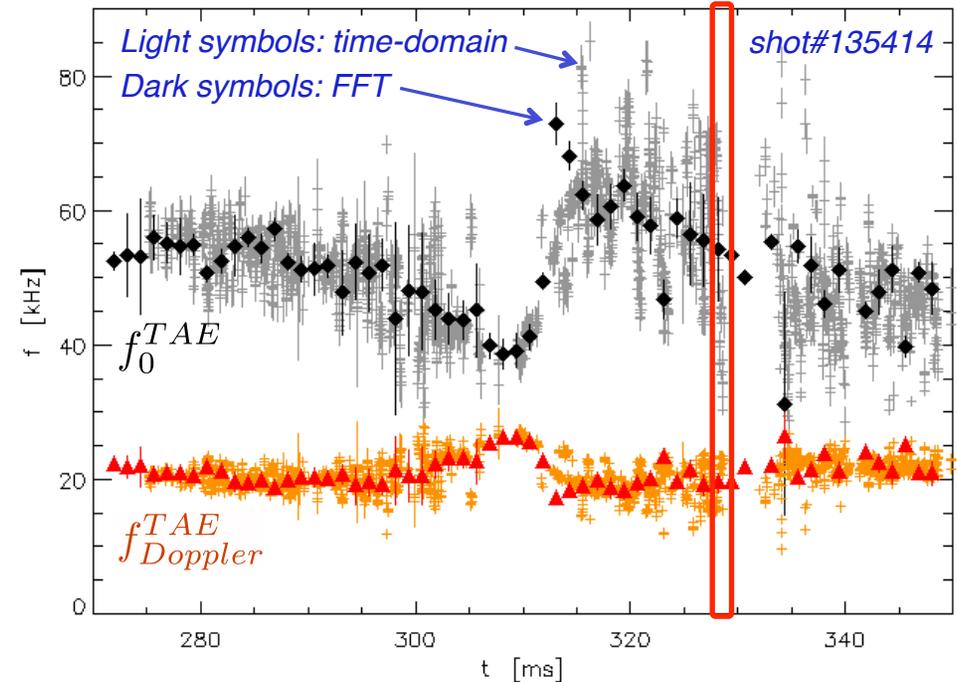
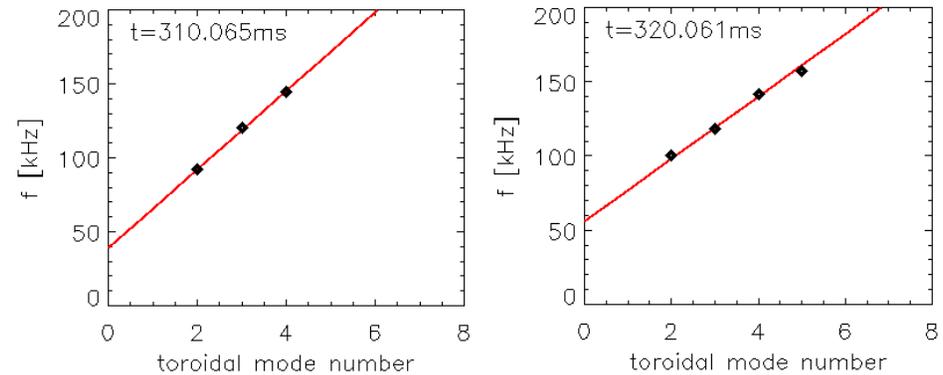
$$f_{lab,n}^{TAE} = f_0^{TAE} + n f_{Doppler}^{TAE}$$

\downarrow lab frame \downarrow plasma frame \downarrow shift from plasma rotation

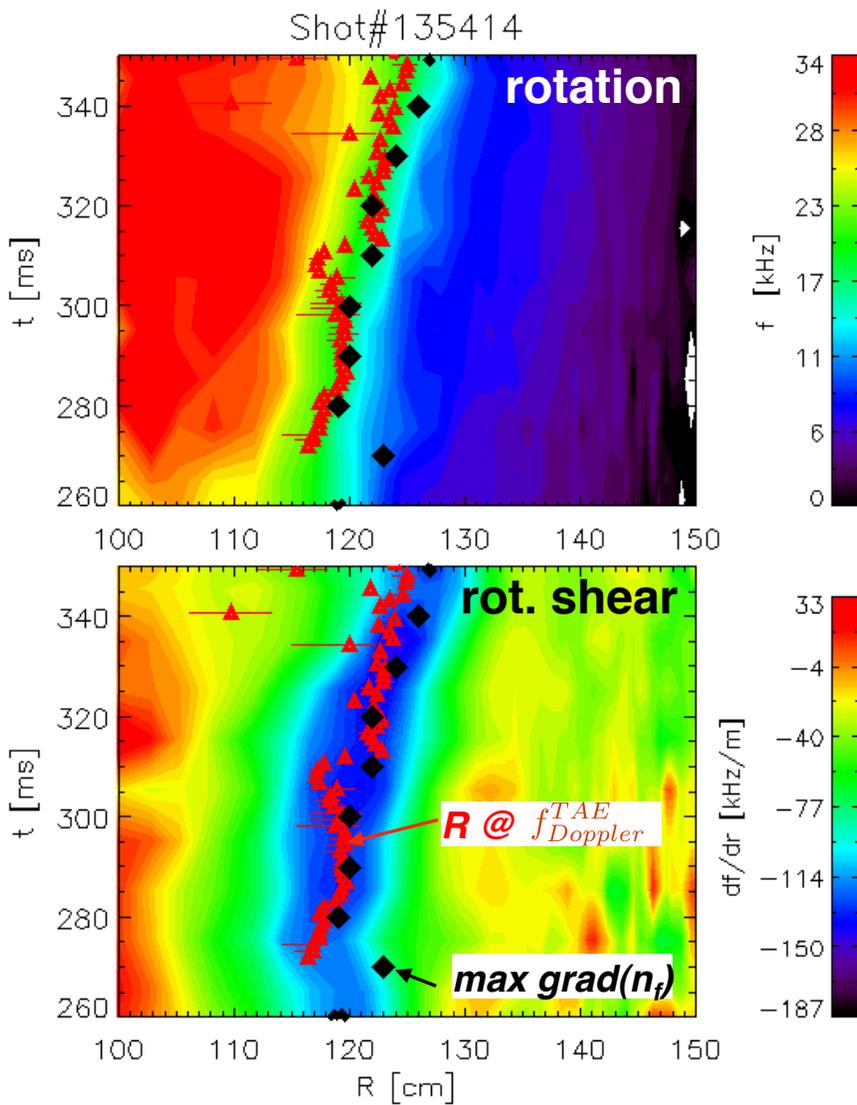
- Valid for time scales >1 ms
- In general, each mode show a different sub-millisecond dynamic...
- ...except during large bursts:
 - Doppler shift only slightly changed
 - Chirp mainly due to decrease in f_0^{TAE}



FFT window 1.3ms



Understanding TAE dynamic requires detailed knowledge of fast ion drive



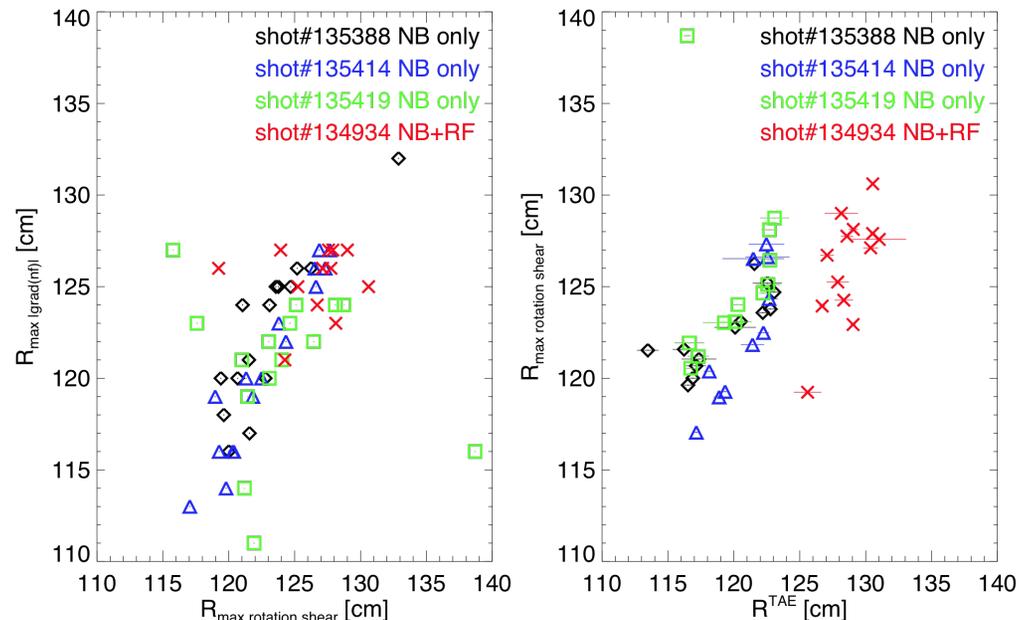
- Modes' location, R^{TAE} , obtained by matching with measured rotation profile:

$$f_{rot}(R^{TAE}) = f_{Doppler}^{TAE}$$

- Correlation between

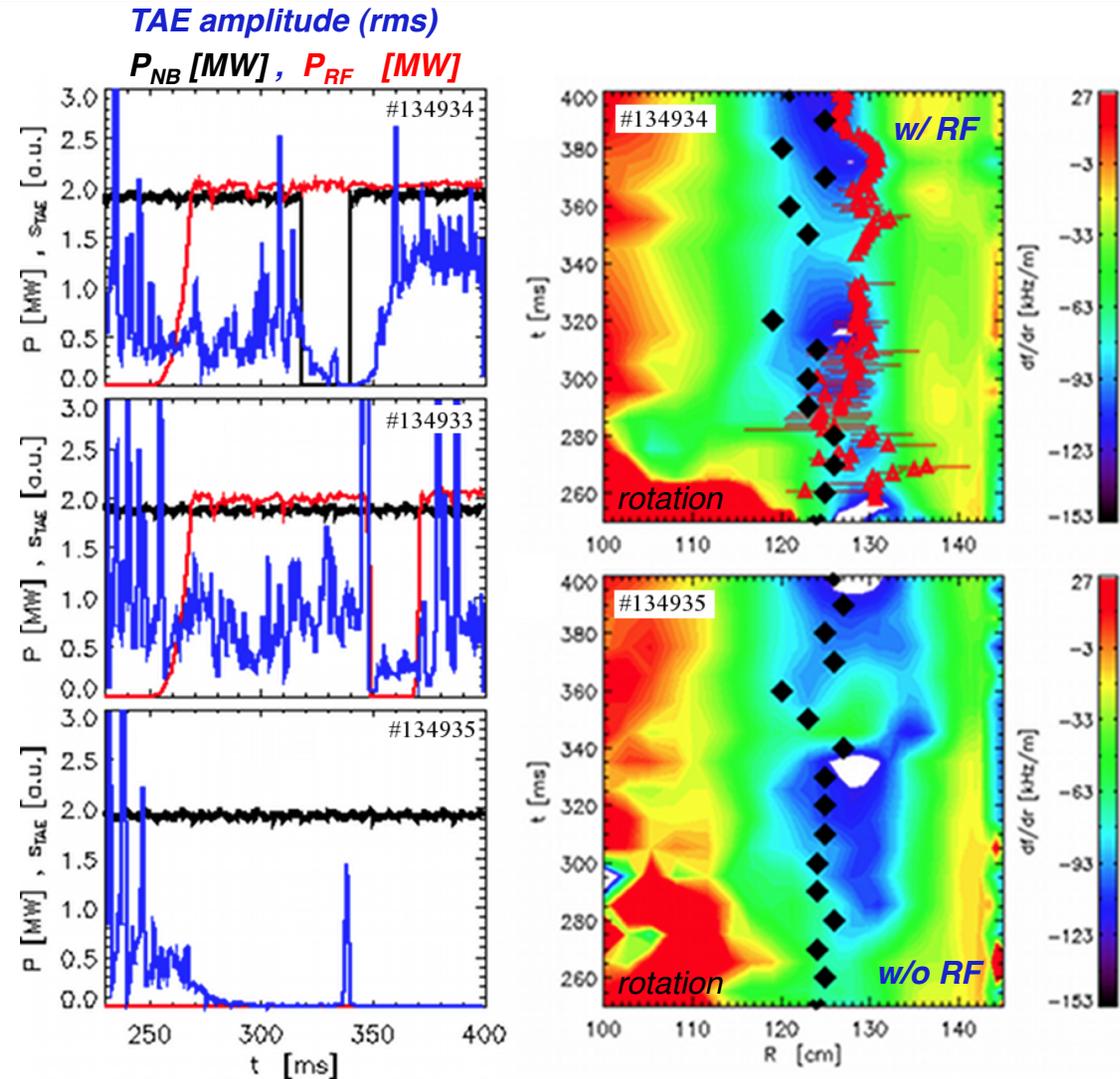
- Mode location
- Max rotation shear
- Steepest fast ion gradient

Coupling through common "source term", i.e. NB injection



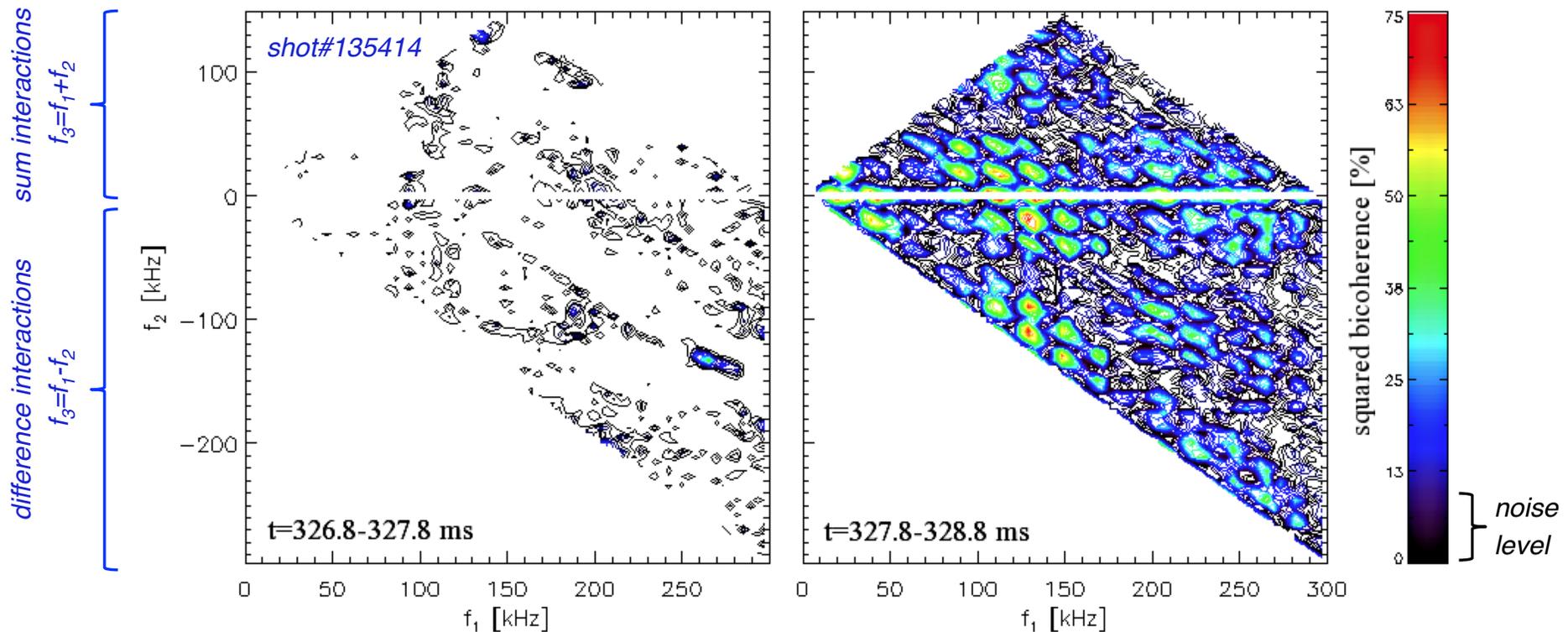
Effects of TAE drive are key factor in determining the observed bursting dynamics

- Bursting dynamics is preserved when drive, $f_{Doppler}^{TAE}$ and shear locations separate
- TAEs respond quickly to notches in NB, RF power
- NB alone is not enough here to drive TAEs unstable



Bicoherence suggests stronger coupling at play during large bursts

- High bicoherence $>70\%$ measured during burst
 - Average over 11 Mirnov coils distributed toroidally over 360°
 - Indicative of sum/difference interactions between modes
 - Both TAEs and low-frequency modes participate



Simple model based on quadratic interactions can be used to investigate coupling between TAEs

$$\dot{s}_{n_3} = \langle c(n_1, n_2) s_{n_1} s_{n_2} \rangle_{f_{n_3}}$$

$s_{n_2} \rightarrow s_{n_2}^*$ (complex conjugate) for difference interaction

Right-hand side filtered around frequency f_{n_3}

Modes must satisfy matching conditions $\begin{cases} n_3 = n_1 \pm n_2 \\ f_{n_3} = f_{n_1} \pm f_{n_2} \end{cases}$

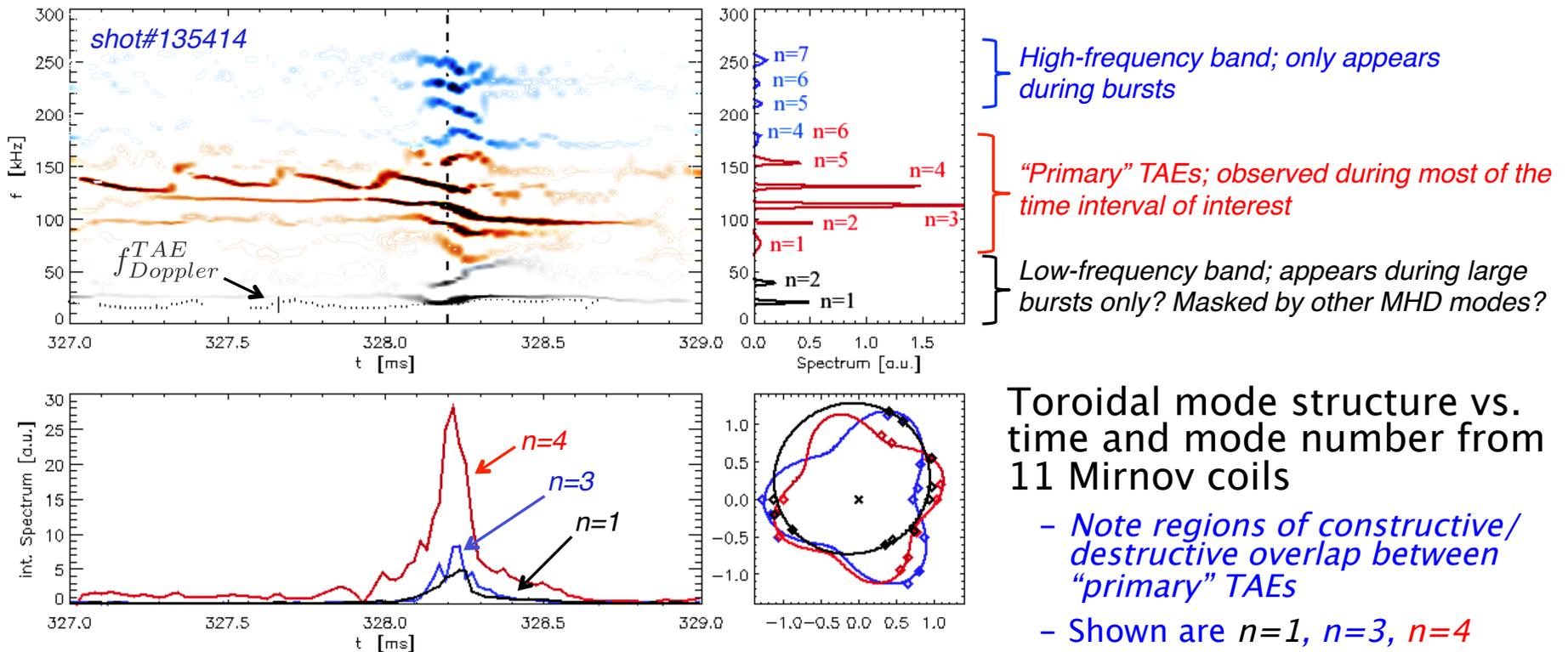
$c(n_1, n_2)$ is the coupling coefficient

In practice:

- Real signals s_{n_1} , s_{n_2} , s_{n_3} measured for each possible triplet, e.g. from Mirnov coils
- “Reconstruct” $\dot{s}_{n_3} \rightarrow \dot{s}_{n_3, rec}$ from measured s_{n_1} , s_{n_2}
- Compare measured and reconstructed \dot{s}_{n_3}
- Frequency match must be verified in the plasma frame:
 - Rotation profile and location of each mode must be accurately known

New modes appear in the spectrum above/below TAE range during large bursts

- Modes can be classified into three groups
 - Discriminants: frequency, temporal evolution*

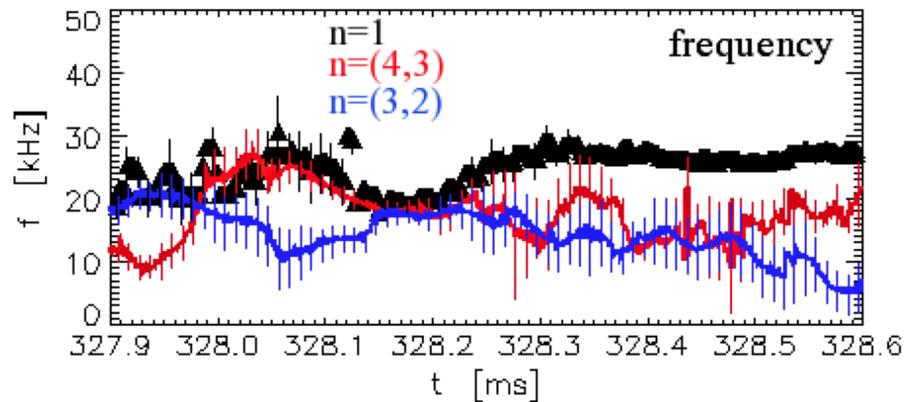
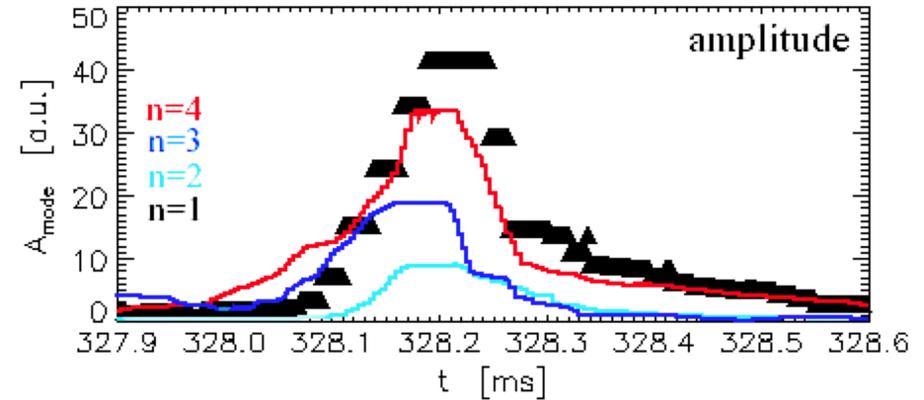
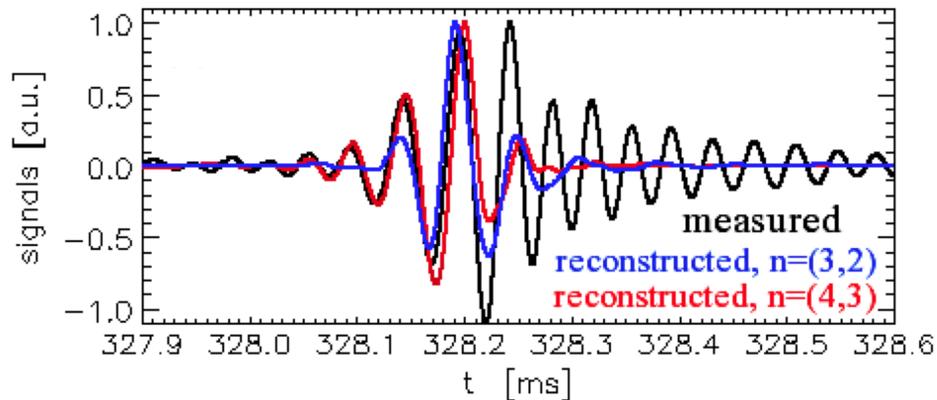


Toroidal mode structure vs. time and mode number from 11 Mirnov coils

- Note regions of constructive/destructive overlap between "primary" TAEs
- Shown are $n=1$, $n=3$, $n=4$

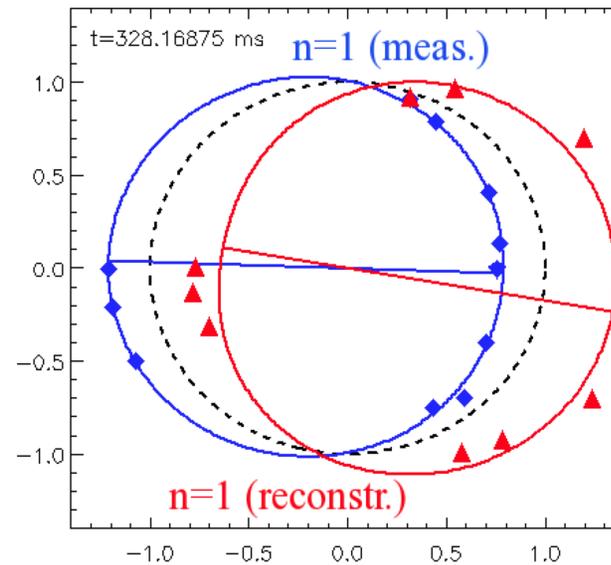
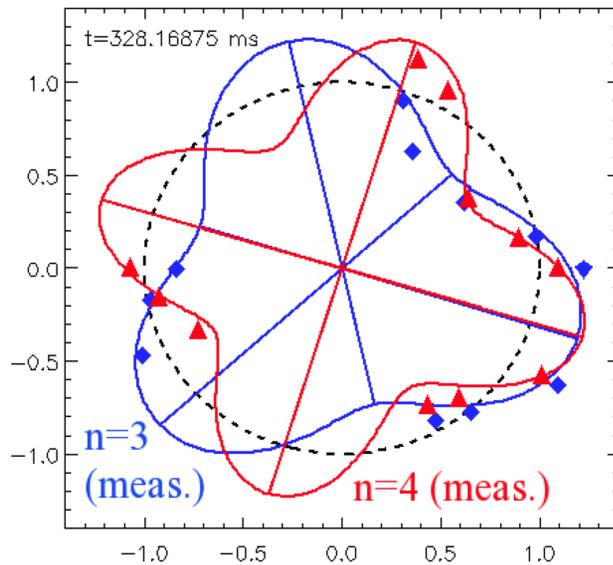
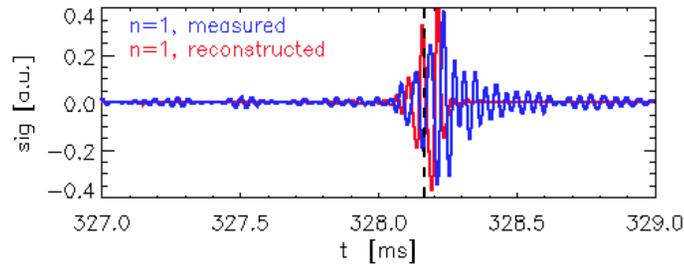
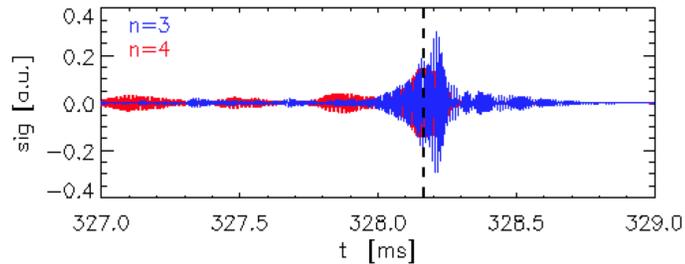
- Picture consistent with primary TAEs
 - coupling to each other
 - generating *secondary* modes through sum/difference with $\Delta n=1$

Good agreement with quadratic interactions' model: amplitude evolution and frequency matching



- “Reconstructed” $n=1$ mode agrees with measured one
 - $n=1$ mode fades away when either amplitude or frequency matching vanishes

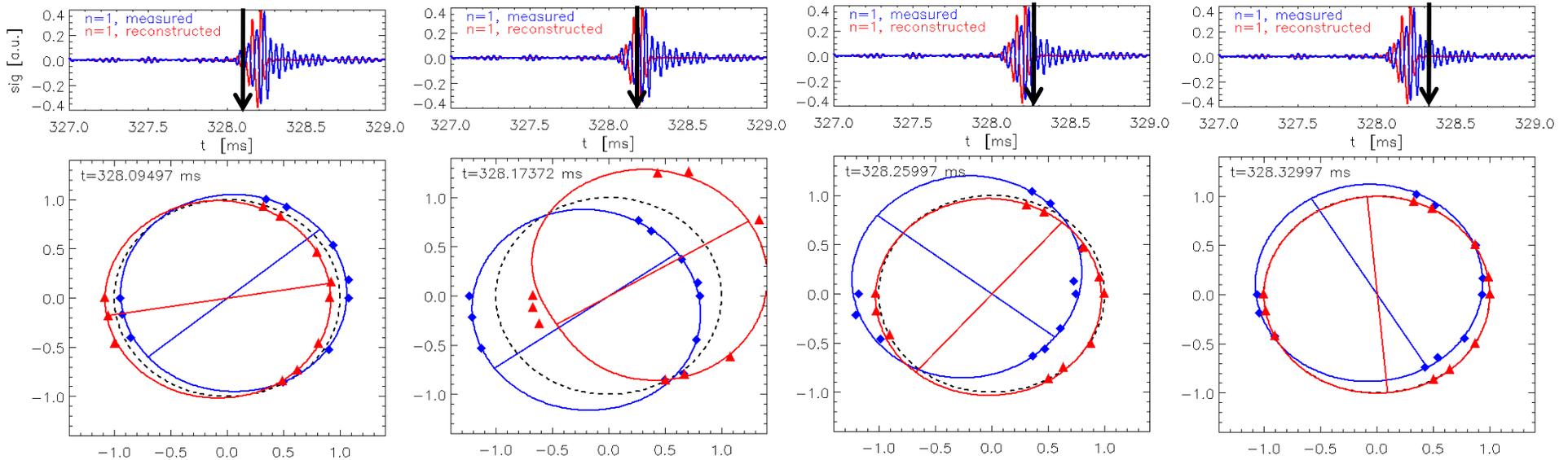
Mode number matching condition



- Symbols: rms mode amplitude data from 11 Mirnov coils
 - Solid lines: fit for a given n ($n=1$ here)
 - Dashed line: unit circle (zero-amplitude reference)

- “Reconstructed” toroidal structure of $n=1$ mode also agrees with measured one
 - Phase shift of 180 degrees, as expected for “difference” interaction (complex conjugate term)

Phase matching condition is transiently verified during large bursts



- Phase resulting from quadratic interaction is important!
 - $n=1$ mode fades away \Leftrightarrow phase deviates from 180 degrees
 - “Single mode” dynamic, with each mode following its own chirp/burst cycle, is effective in reducing efficiency of quadratic interactions
 - The result is a “semi-cahotic” scenario, with small bursts (single mode) and occasional large bursts (multi-mode avalanches)

Summary

- TAE bursts can cause large, intermittent fast ion transport
- Bursting TAE regime is “robust” against small variations of plasma parameters
 - L-mode vs. H-mode, NB only, NB+RF, ... : all show similar features
- Non-linearities occur in both single-mode and multi-mode (avalanching) TAE dynamic
 - Only avalanches seem to cause significant fast ion losses
- More experiments planned for near term
 - Systematic study of TAEs (and *avalanches*) in H-mode
 - Comparison with M3D-K code planned; plasma rotation included
 - Improve “linear” analysis (NOVA-K + ORBIT)