Characterization of Fast Ion Power Absorption of HHFW in NSTX

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Introduction and Motivation

- Ion absorption critically important to assessing viability of HHFW to heat and drive current in STs
- Experimental evidence of HHFW interaction with NBI
 - Neutral Particle Analyzer (NPA) scannable at midplane
 - neutron rates, Fast Lost Ion probes
- Thompson scattering measures T_e, n_e profiles
 X-ray Crystal Spectroscopy measures peak T_i
- Computational evidence
 - HPRT, TRANSP, CURRAY, AORSA, METS

NSTX Utilizes the TFTR ICRF system

- 30 MHz Frequency corresponds to $\omega/\Omega_D = 9-13$
- 6 MW from 6 Transmitters for up to 5 s

• 12 Element antenna with active phase control allows wide variety of wave spectra

$$-k_{\parallel} = \pm 3-14 \text{ m}^{-1}$$

HHFW 12 element antenna array



- Antenna takes up almost 90° toroidally
- Provides high power capability with good spectral selectivity

Phase Feedback Control Configuration



- Digital based phase feedback control is used to set the phase between the voltages of antenna elements 1 through 6
- Decouplers compensate for large mutual coupling between elements and facilitate phase control

HHFW can generate a fast ion tail with NBI

• Typical shot

• D⁺ tail extends to 130 keV

• Tail saturates in time during HHFW

VSTX



• Tail decays on collisional time scale

HHFW enhances neutron rate



- After RF turnoff, rate decays close to measured and predicted no RF value
- TRANSP neutron rate predictions without RF input fall shorter than measured rate for RF shot



- Larger β_t promotes greater off-axis electron absorption
- Reduces fraction of power available to core fast ion population

Ion loss with lower B-field can't account for reduction in tail

- Two codes used to check ion loss at B₀=4.5 kG vs. 3.5 kG for 80-120 keV ions
- CONBEAM Egedal (MIT-PSFC)

- Loss fraction at 120 keV $B_0 = 4.5 \text{ kG}: 21\%$

 $B_0 = 3.5 \text{ kG}: 25\%$

- EIGOL Darrow (PPPL) – Loss fraction at 120 keV $B_0 = 4.5 \text{ kG: } 17\%$ $B_0 = 3.5 \text{ kG: } 23\%$
- Small change in loss fraction insignificant compared to major tail reduction
 - More likely an RF effect

k_{\parallel} has little observed effect on fast ions

DNSTX



• Greater ion absorption predicted with lower k_{\parallel} , but surprisingly little variation in tail, small neutron enhancement with higher k_{\parallel}

Ray tracing predicts fast ion absorption competitive with electrons



- HPRT computes hot plasma absorption over cold ion/hot electron ray path
- 25-50 rays used
- TRANSP output used as input for fast ion temp and density distribution
- Fast ions dominate central absorption, electrons further off-axis
- $T_{i,th} = 2 T_e$ (XCS), no thermal ion absorption

Effective Maxwellian a good approximation for fast ions



- One effective Maxellian, exactly matching TRANSP energy density and temp., fits f(E) well, though neutron rate off +20%
- More Maxwellians (including negative) can match all moments

Observation of less fast ion absorption at higher β_t consistent with theory



• Lower on-axis absorption for lower B, higher β_t predicted

• Energy moment of Vlasov eq. + vector identities leads to [Menard, RF 1999]:

$$\frac{\partial}{\partial t} \int \mathbf{d}^3 \mathbf{v} \frac{\mathbf{m}_s v^2}{2} f_s + \nabla \cdot \int \mathbf{d}^3 \mathbf{v} \frac{\mathbf{m}_s v^2}{2} \mathbf{v} f_s = \mathbf{E} \cdot q_s \int \mathbf{d}^3 \mathbf{v} \mathbf{v} f_s$$

• The first non-vanishing terms of this are:

$$\frac{\partial W_{ps}}{\partial t} + \nabla \cdot \mathbf{T}_{s} = \mathbf{E} \cdot \mathbf{j}_{s}$$

• Kinetic flux term is necessary if one uses the full hot, complex dielectric tensor

Clarifying Kinetic Flux

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- In Stix (4-15) $\frac{\partial W}{\partial t} = \frac{1}{8\pi} \left[\omega \mathbf{E}^* \cdot \left(\mathbf{k}_i \cdot \frac{\partial}{\partial \mathbf{k}} \mathbf{\epsilon}_h \right) \cdot \mathbf{E} + \dots \right]$
- In Stix (4-17,19) $\frac{\partial W}{\partial t} = -2\mathbf{k}_i \cdot \mathbf{T} + \dots \quad \mathbf{T} = -\frac{\omega}{16\pi} \mathbf{E}^* \cdot \frac{\partial}{\partial \mathbf{k}} \mathbf{\varepsilon}_h \cdot \mathbf{E}$
- 1st important note Stix, p. 75: "In expression for **T**, the dot products are between \mathbf{E}^* and ε_h , and ε_h and $\mathbf{E}_.$ "

• So, unambiguous **T** is:
$$\mathbf{T} = -\frac{\omega}{16\pi} \frac{\partial}{\partial \mathbf{k}} \left(\mathbf{E}^* \cdot \boldsymbol{\varepsilon}_h \cdot \mathbf{E} \right)$$

• Does (4-15) = (4-17)? Yes! However...

Only $\nabla \cdot \mathbf{T}$ is unique

$$\mathbf{E}^* \cdot \left(\mathbf{k}_i \cdot \frac{\partial}{\partial \mathbf{k}} \mathbf{\varepsilon}_h \right) \cdot \mathbf{E} \quad \text{also equals:} \quad \mathbf{k}_i \cdot \left[\frac{\partial}{\partial \mathbf{k}} \left(\mathbf{E}^* \cdot \mathbf{\varepsilon}_h \cdot \mathbf{E} \right) + \mathbf{k}_i \times \text{anything} \right]$$

- So expression for **T** is *not unique*; only $\nabla \cdot \mathbf{T}$ well-defined
- Common loss-free plasma approximation of group velocity: $v_g = \frac{P+T}{W}$ is invalid, as is any expression which uses **T** in this form.
- 2nd important note: $\mathbf{E}^* \cdot \left(\mathbf{k}_i \cdot \frac{\partial}{\partial \mathbf{k}} \mathbf{\varepsilon}_h\right) \cdot \mathbf{E} \neq \mathbf{k}_i \cdot \left(\mathbf{E}^* \cdot \frac{\partial}{\partial \mathbf{k}} \mathbf{\varepsilon}_h\right) \cdot \mathbf{E}$
- RHS gets wrong cross derivatives wrt k

NPA scan indicates induced tail well off-axis



- Depletion in particle flux with NPA R_{tan} further off-axis
- Tail extends to same energy range
- Future: scan over wider range of r/a

Summary



- Clear RF-induced fast ion tail observed with NBI
- Neutron rate and modeling support interaction
 Good agreement between HPRT, AORSA, CURRAY
- Tail formation suppressed with higher β_t
- Little effect with k_{\parallel} observed
- Effective Maxwellians can represent fast ion f(E)
- Kinetic flux clarified