

Modeling of Low-frequency MHD-induced Beam-ion Transport In NSTX

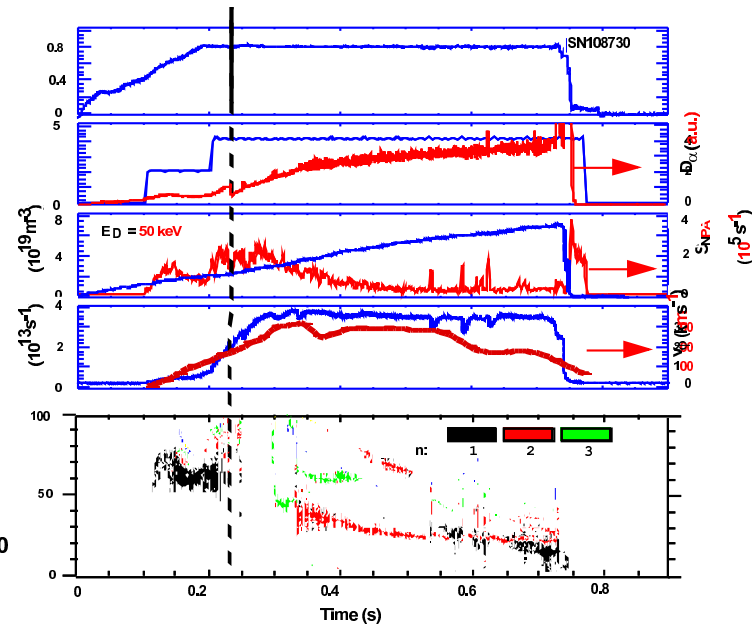
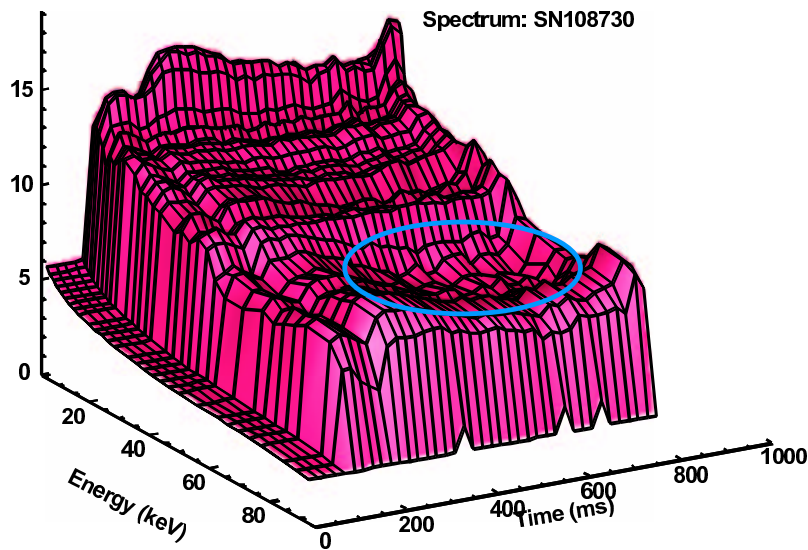
N. N. Gorelenkov and S.S.Medley

Princeton Plasma Physics Laboratory, Princeton

NSTX seminar, PPPL, 11 October, 2004



NPA measures beam ion signal depletion at 40 – 80keV



S.S.Medley et.al. NF'04 submitted

After H-mode transition $m = 4/n = 2$ mode is observed

Why NPA signal is depleted at those energies?

Motivation



NBI ion possible loss/redistribution raise question on

- first wall heat flux,
- heating efficiency.
- What about current drive?
 - ASDEX shows that at high P_{NBI} off-axis injected beam ions are flattened with the diffusivity of thermal plasma (Günter, EPS2004)

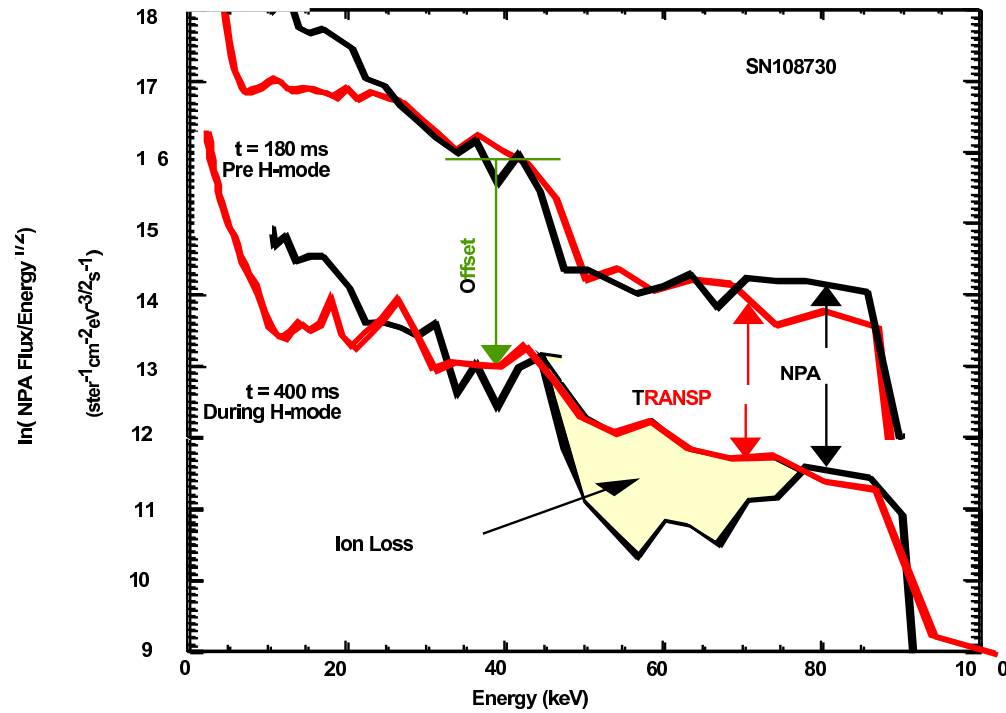
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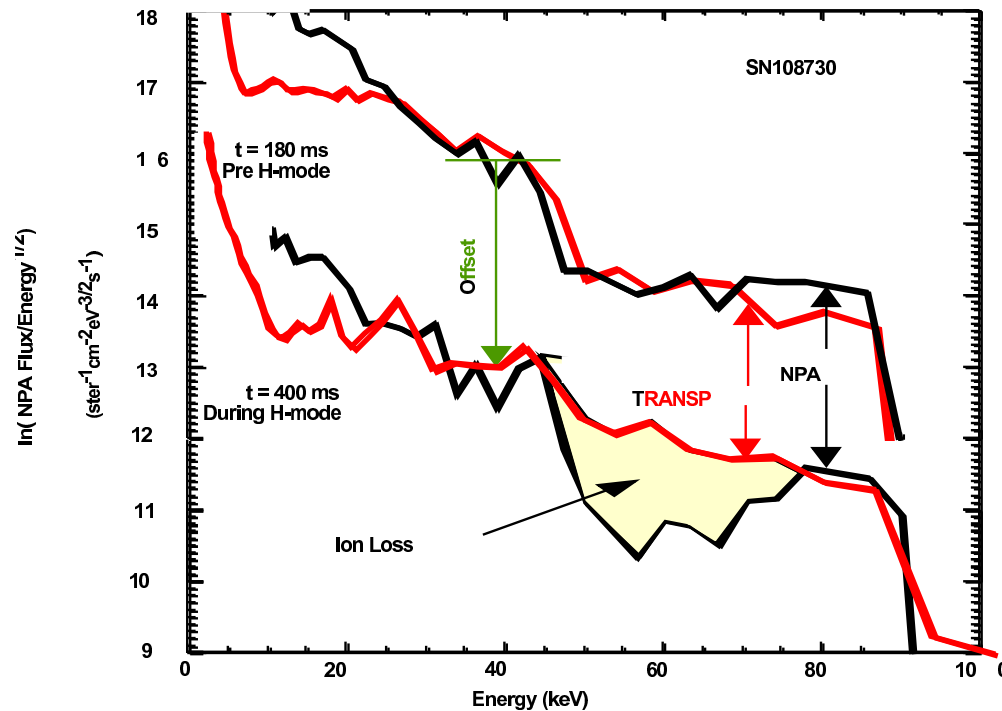
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 - ASDEX shows that at high P_{NBI} off-axis injected beam ions are flattened with the diffusivity of thermal plasma (Günter, EPS2004)
 - **Can ITER have steady state current drive? What can affect it?**

TRANSP slowing down beam ion distribution vs NPA signal



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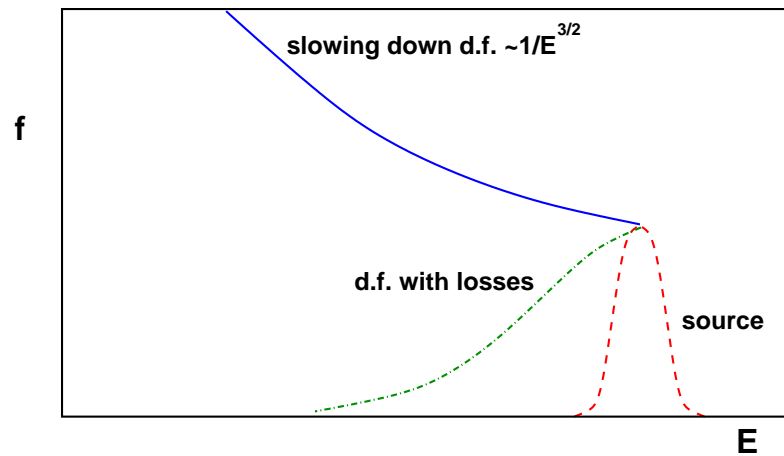


Why beam ions do not fill the gap and where do they go?

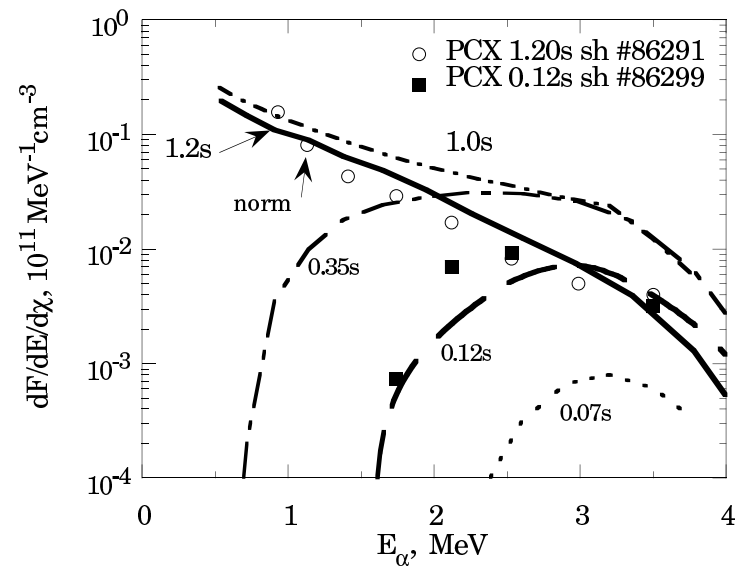
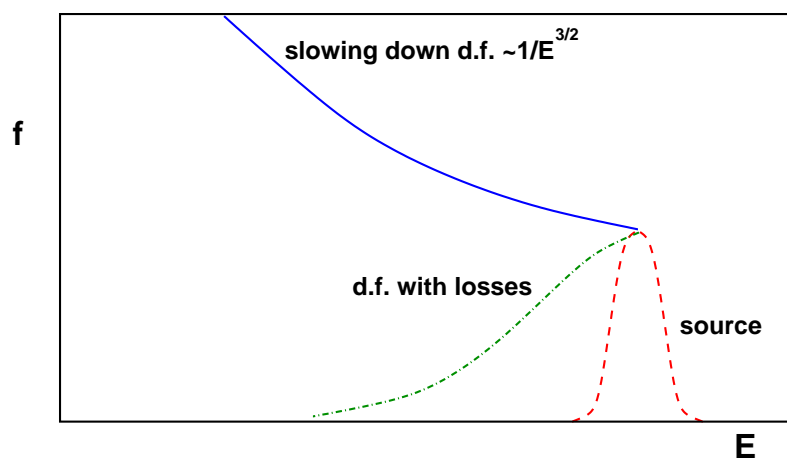
In this talk:

- (1) what is the confinement time of beam ions
- (2) can $m = 4/n = 2$ perturbation explain fast ion losses?

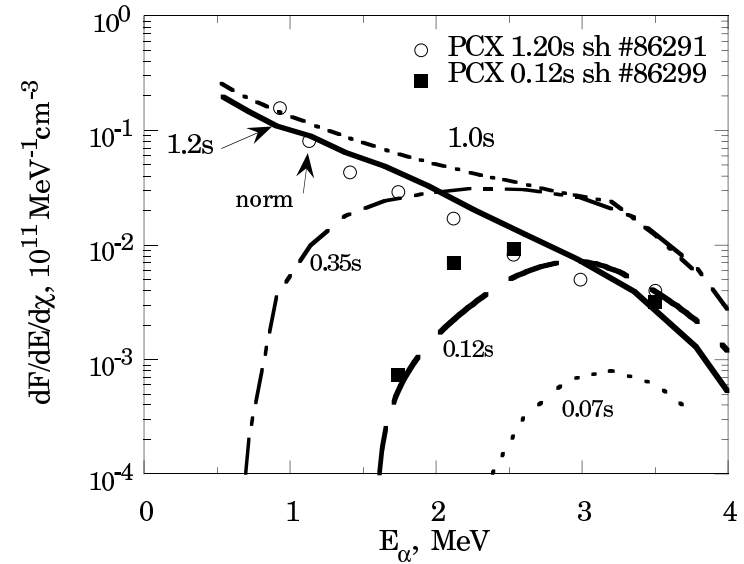
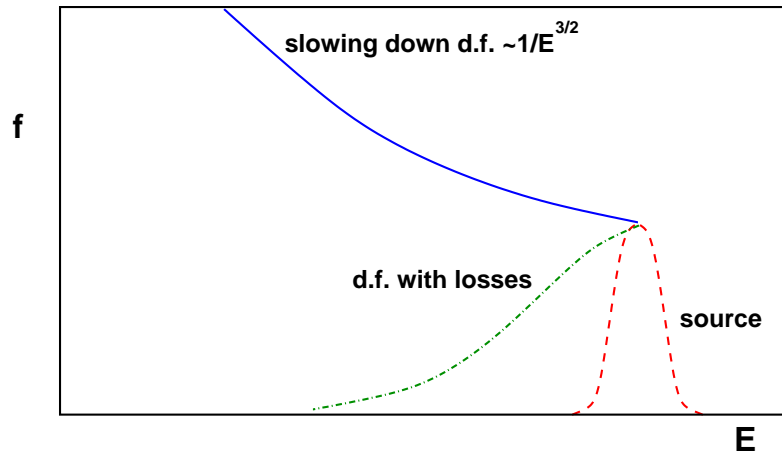
(I) Losses effect fast ion distribution function



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Kinetic equation in steady state (Cordey, Goldston, Mikkelsen, '81):

$$\frac{1}{\tau_{se} v^2} \frac{\partial}{\partial v} (v^3 + v_*^3) f - \frac{f}{\tau_{loss}} + S \delta(v - v_0) = 0 \quad (1)$$

Solution depends on the loss to drag time ratio



At finite τ_{loss} we obtain

$$f = \frac{Cn_b}{v^3 + v_*^3} \left(\frac{v^3 + v_*^3}{v_{b0}^3 + v_*^3} \right)^{\tau_{se}/3\tau_{loss}} \quad (2)$$

and $f \sim 1/(v^3 + v_*^3)$ if $\tau_{loss} \rightarrow \infty$.

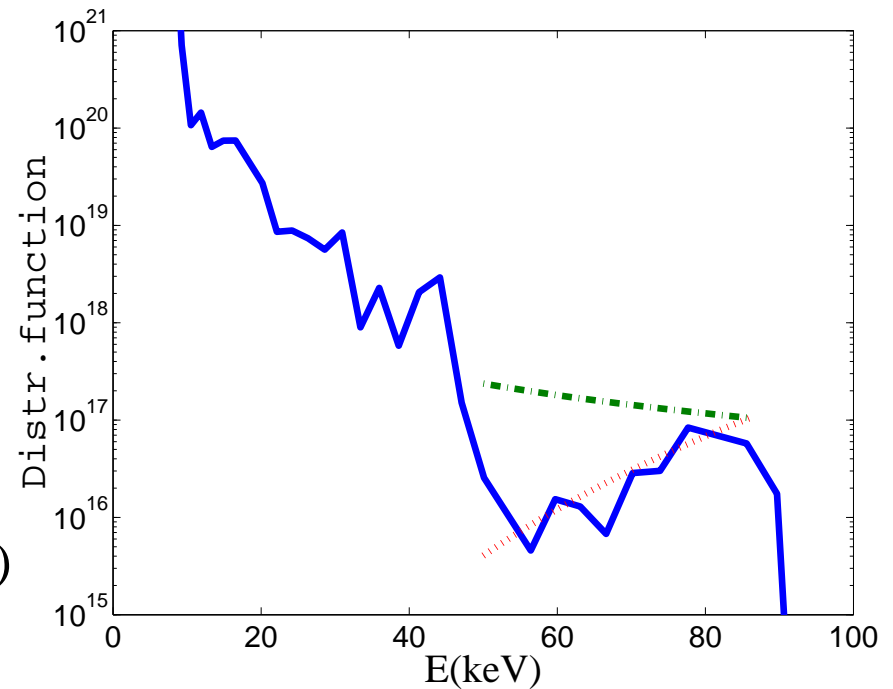
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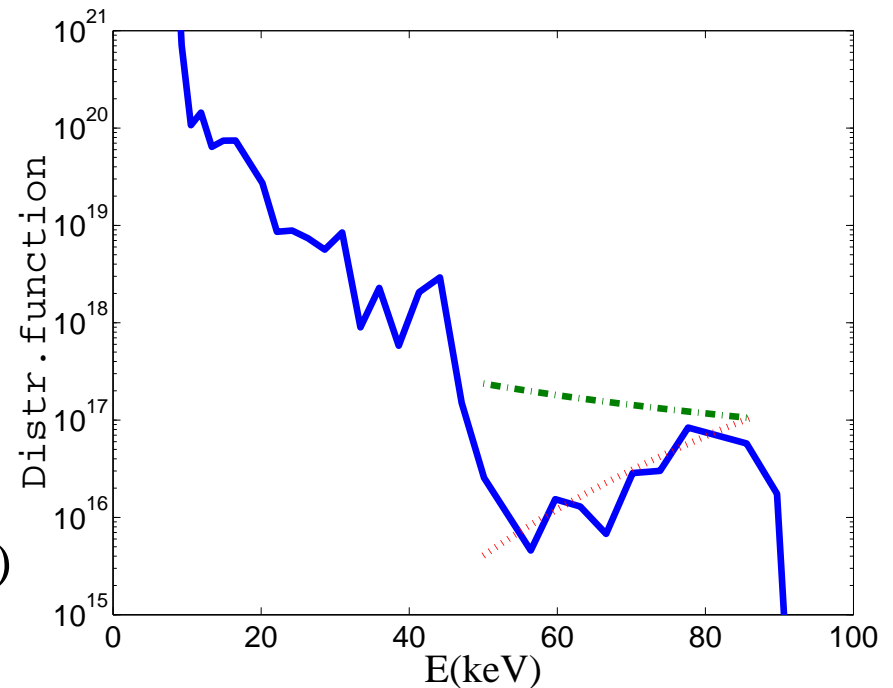
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Implies that $\tau_{loss} = \tau_{se}/15$, i.e. $\tau_{loss} = 4msec$.

(II) What is the mechanism for “losses”/redistribution



Numerical study includes

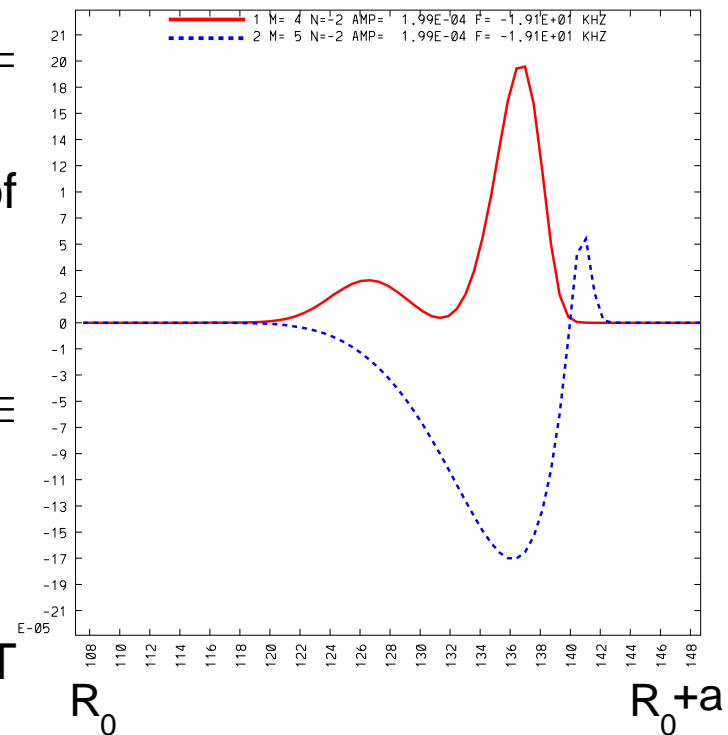
1. plasma zero frequency $m = 4/n = 2$ perturbation,
2. amplitude on the order of $\delta B/B \sim 10^{-4}$,
3. strong toroidal sheared rotation,
4. pitch angle for NPA sight line $\chi \equiv v_{\parallel}/v = 1.5125 - 0.629R_{cx}[m]$,
 $\chi = 0.9 \Rightarrow R_{cx} = 0.97m$,
 $\chi = 0.6 \Rightarrow R_{cx} = 1.45m$
5. realistic equilibrium and ORBIT code

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mode structure consistent with ideal MHD $\delta \mathbf{B} = \nabla \times \alpha \mathbf{B}$

$$\alpha \sim (1 - nq/m) (r/r_s)^m \sin(n\phi - m\theta), \text{ if } r < r_s$$

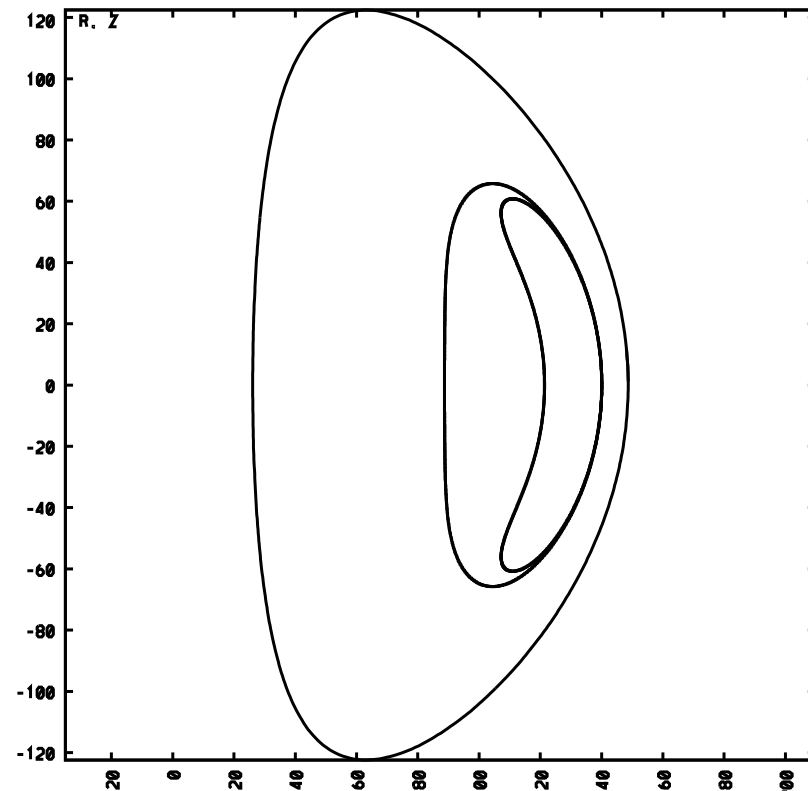
Beam ion orbits without perturbations



Example trapped ion orbit at $E = 70keV$ and $\chi = 0.55$.

Electric field in NSTX -
central potential $3.8keV$, cen-
tral rotation $\dot{\phi} = 8 \cdot 10^5 sec^{-1}$ -

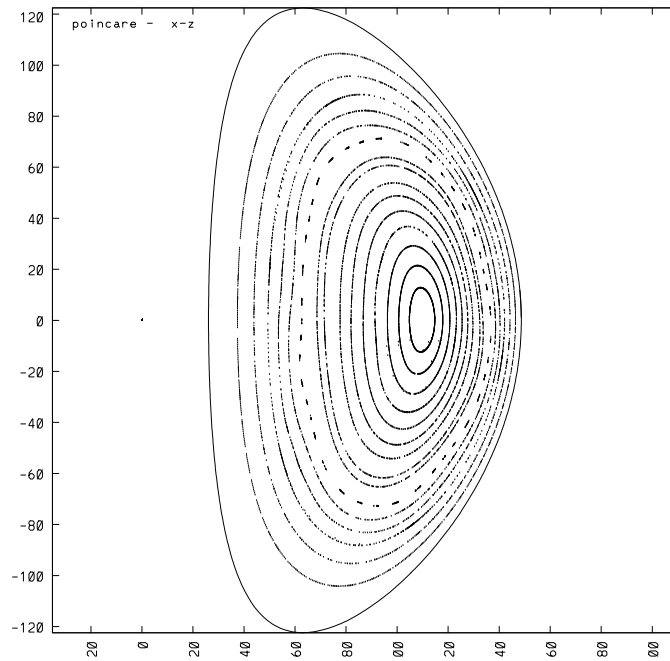
1. changes particle orbits
2. effects precession frequency
3. shifts mode frequency



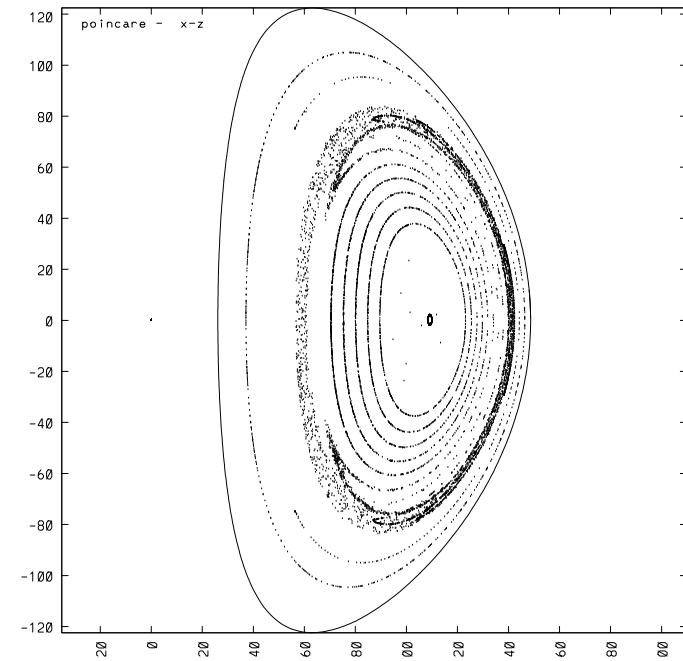
Islands in the real space (R-Z) with perturbations



w/out electric field



with electric field

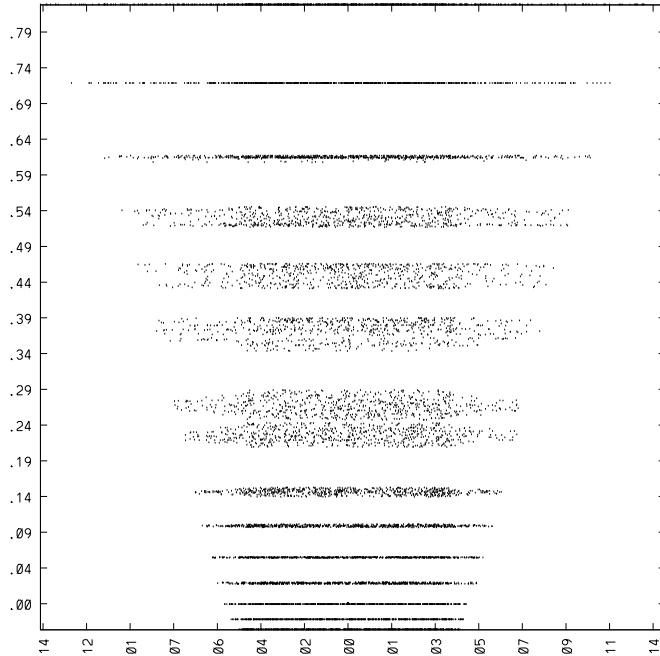


$$E = 0.1keV \text{ and } \chi = 1$$

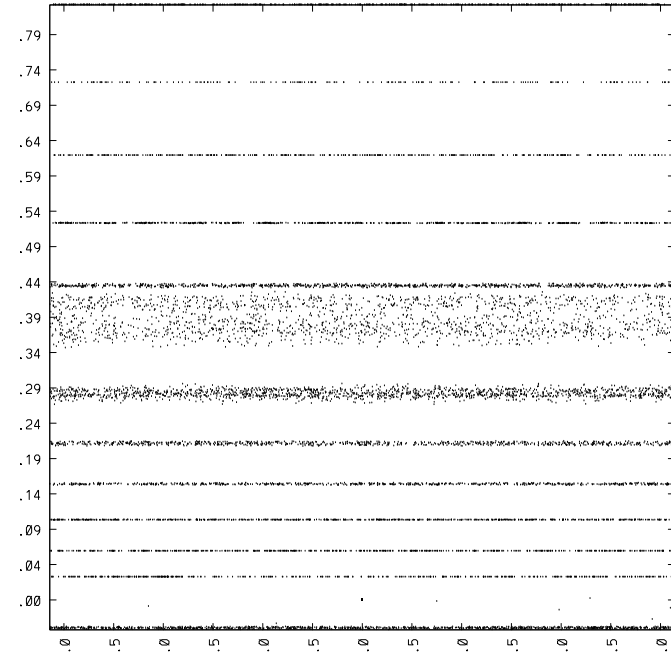
Islands in $P_\phi - (\omega - \phi t) \left(\sim (r/a)^2 - (\omega - \phi t) \right)$ space



w/out electric field



with electric field



Wave-particle approximate resonance condition



$$\omega - \omega_{E \times B} - (k_{\parallel} + l/qR) v_{\parallel} = 0, l = \pm 1, 2, \dots \quad (3)$$

Frequency effect

- If $\omega = 0$ and there is no electric field, resonance is $k_{\parallel} + l/qR = 0$ - in real space
- If $\omega \neq 0$ and/or $\omega_{E \times B} \neq 0$ resonance involves phase space.

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Orbit width effect

1. In zero orbit width case, $l = \pm 1$ due to toroidal drift velocity $\cos \theta$ -like modulation.
2. At large orbit width, only parts of particle orbit interact with the mode, $\Rightarrow |l| > 1$ appear.

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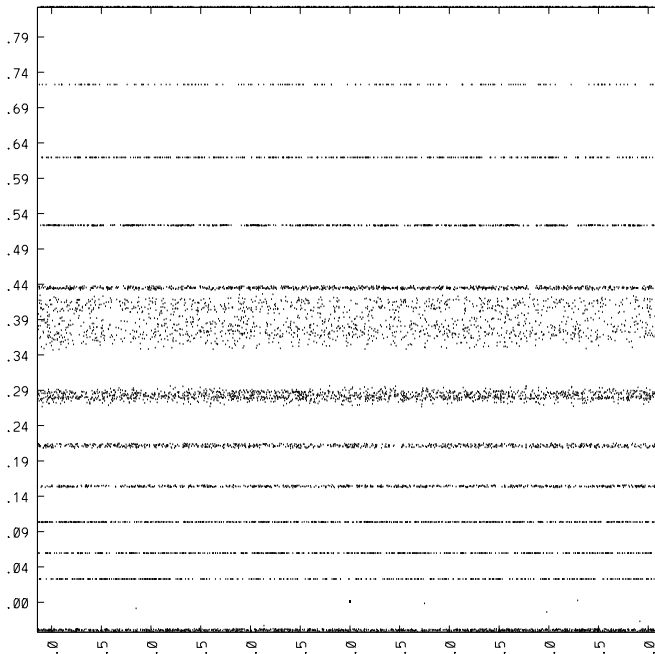
Since $|\omega - \omega_{E \times B}| \ll |v_{\parallel}|/qR$ the resonance is possible if $|k_{\parallel}qR + l| \ll 1$ at given magnetic surface.

Thus the resonance is selective (narrow in l) for low energies and broad for high energies.

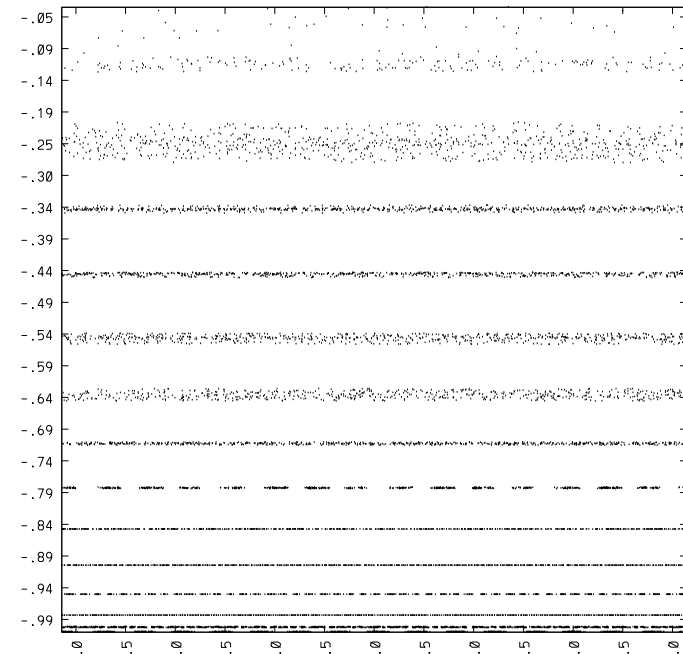
Islands in $P_\phi - (\omega - \phi t)$ for different energies



$E = 0.1keV$



$E = 70keV$



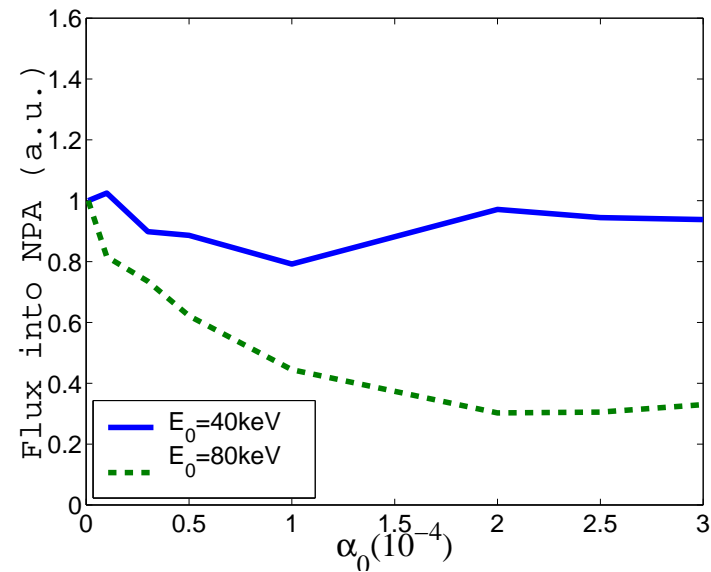
and $\chi = 1$, $\alpha_0 = 2 \times 10^{-4}$

Wide range of P_ϕ or r/a is affected.

Numerical results for injected ions at $E_0 = 40$ and 80keV



Allow for ion thermalization until $E = E_0/2$:



Fluxes vs.

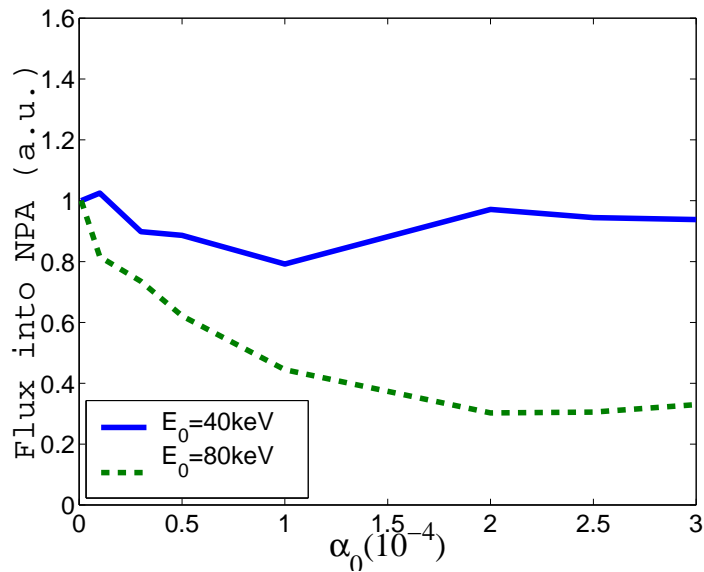
perturbation amplitude

Particles are effected above 40keV .

Numerical results for injected ions at $E_0 = 40$ and 80keV

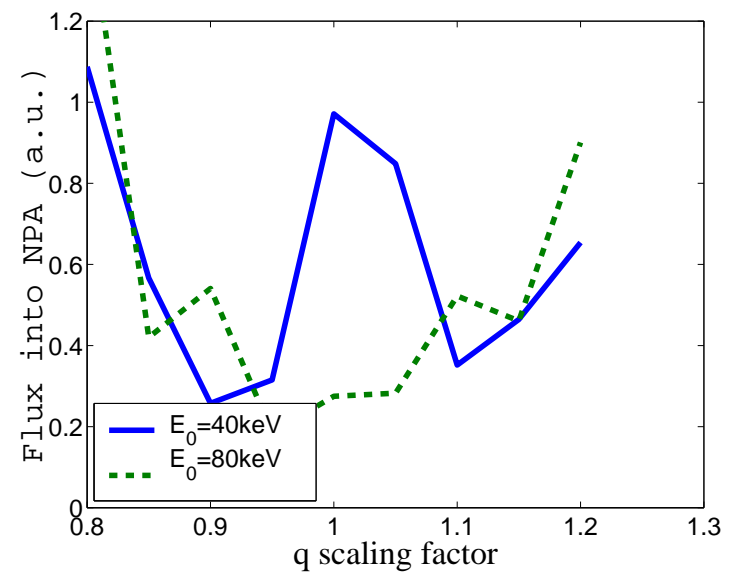


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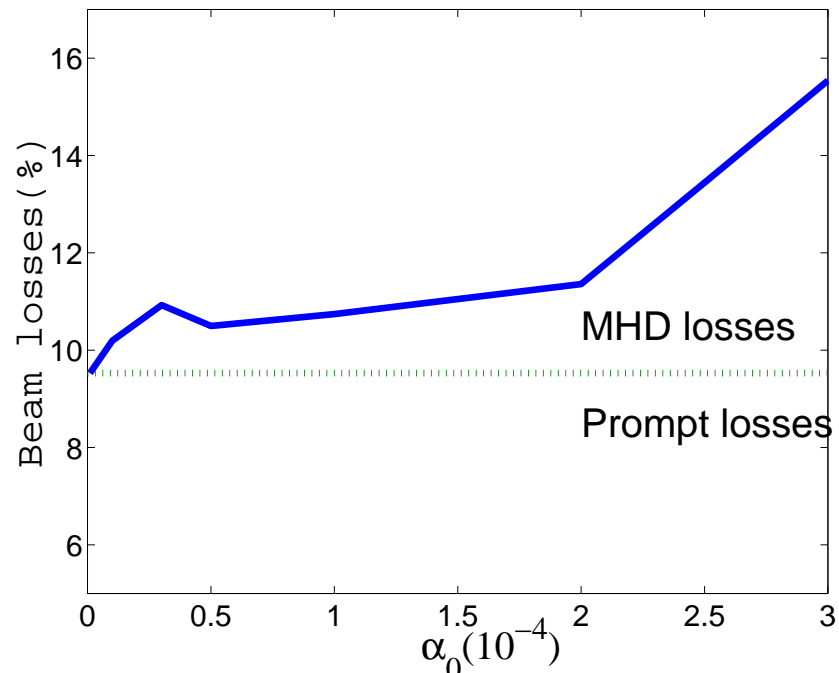
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Fluxes vs. q-factor
($q_{new} = q * (q - factor)$)

Shows sensitivity to resonant $k_{||}$

Are there any losses due to MHD



At expected amplitudes $\alpha_0 = 2 - 3 \times 10^{-4}$ ($\delta B/B$) $\sim 10^{-3}$, $m = 4/n = 2$ mode can induce losses comparable to prompt losses.

Summary and conclusions



MHD activity observed in NSTX H-mode plasma is shown to be responsible for the NPA signal loss.

- Beam ion redistribution is energy selective affecting ions at $E = 50 - 80keV$.
- Characteristic loss/redistribution time is $\tau_{loss} \simeq 4msec$.