

# Predictions and observations of global beta-induced Alfvén-acoustic modes in JET and NSTX

N.N. Gorelenkov, E.D. Fredrickson, S. Kaye, H. Park

*Princeton Plasma Physics Laboratory, Princeton*

H. L. Berk

*Institute for Fusion Studies, Austin, Texas*

S. E. Sharapov

*Euroatom/UKAEA Fusion Assoc., Culham Science Centre, Abingdon, Oxfordshire*

S. A. Sabbagh

*Columbia University, New York*

K. Tritz

*Johns Hopkins University, Baltimore, Maryland*

F. M. Levinton

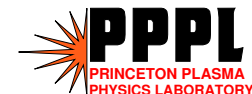
*Nova Photonics, Princeton, New Jersey*

and JET EFDA Contributors

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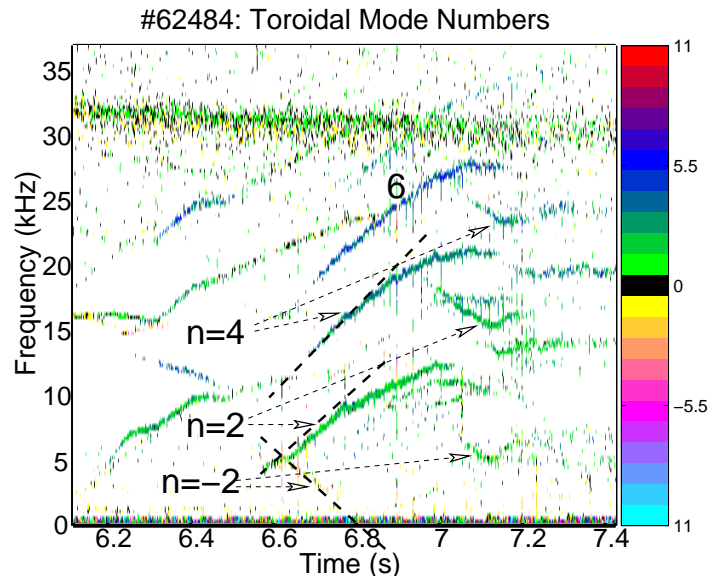
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## New experimental observations on JET and NSTX motivate low frequency mode study

JET

PPPL

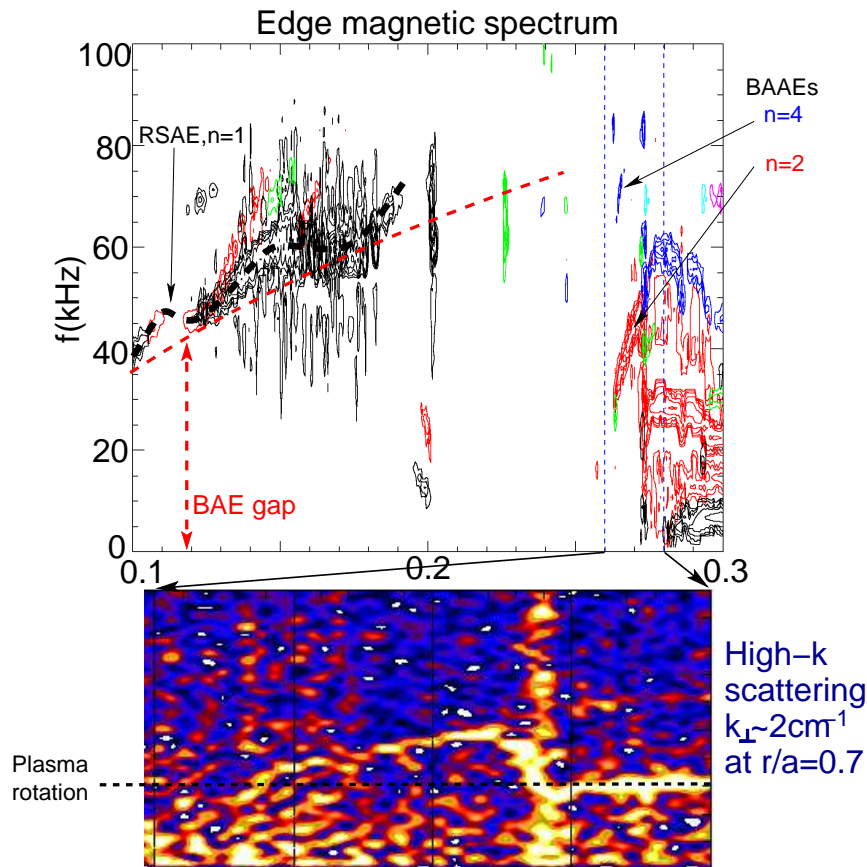


- JET:  
2MW ICRH in low plasma density  $\Rightarrow$  high H-minority beta:  
 $\beta_H \sim \tau_{se} \sim n_e^{-1}$ ;  
 $\beta_H \sim \beta_{plasma} (\sim 2\%)$ .
- New chirping frequency activity  
 $f_{pl} = 0 - 20\text{kHz}$   
 (for RSAE, Alfvénic cascades  
 $f_{pl} \neq 0$ )
- Structure at upper limit (another mode, same  $n$ ).
- Only even  $n$ 's were observed.

Frequency is much lower than RSAE/TAE frequency,  $v_A/qR$ :  
 $\omega_* = 0.5\text{kHz}$  is small

# Variety of MHD instabilities are routinely observed in NSTX

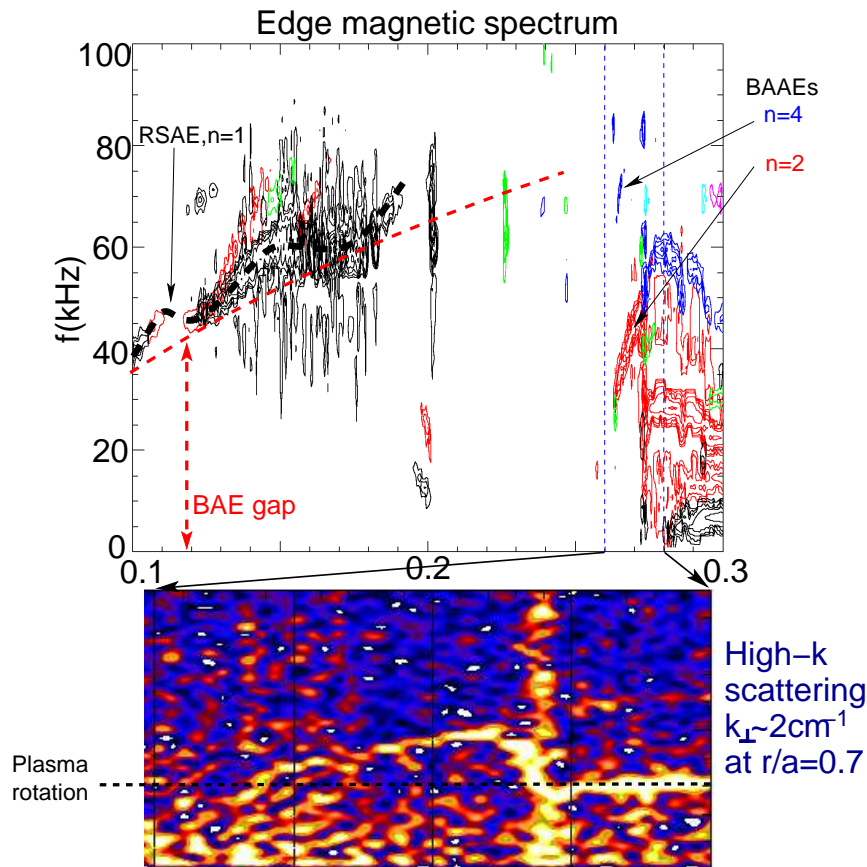
JET PPPL



- Edge Magnetic spectrum and internal high-k scattering see the same low- $f$  activity as in JET,  $t = 0.27s$ .
- only even numbers, sweeping frequency,  $n = 2, 4$
- reversed shear  $q$ -profile from MSE
- $f_{TAE} = v_A/2qR \simeq 90kHz$
- $t = 0.262sec$ ,  $n = 2$  mode frequency  $f_{lab} \simeq 30 - 50kHz$ ,  $f_{pl} \simeq f_{lab} - n \times 15 = 0 - 20kHz$ ,  $\Rightarrow$  **too low for TAE for observed  $n = 2$  activity.**

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What are these modes: EPs, fishbones, KBMs, TAEs?

## *What is the importance of low- $f$ instabilities?*

JET

PPPL

- New class of instabilities called here **Beta-induced Alfvén Acoustic Eigenmode (BAAE)** helps to study two fundamental MHD waves: Alfvén and acoustic.

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- As opposite to high- $f$  instabilities, low- $f$  MHD mostly result in radial particle transport:
  - On NSTX bursting low- $f$  modes lead to wide range of **losses up to 40%** of injected beam ions (Fredrickson'06).

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- **MHD spectroscopy** application for  $q$ -profile diagnostic:
  - BAAE can complement MHD spectroscopy in low-, medium- $\beta$  plasma
  - can be the only alternative in **high- $\beta$  plasma**, such as in STs.

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- Due to coupling to acoustic branch strong interaction with thermal ions is expected:
  - $\Rightarrow$  strong drive due to fast ions and strong damping due to thermal ions,
  - $\Rightarrow$  potential for **energy channeling** from beam ions directly to thermal ions ( **$\alpha$ -channeling**).



## TALK OUTLINE

JET

PPPL

1. Theory of Alfvén - acoustic continuum in ideal MHD
2. JET analysis and data comparison
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4. Discussion and Summary

*Theory of Alfvén/acoustic continuum*

JET PPPL

Simplified shear Alfvén and acoustic coupled equations in low- $\beta$ , large aspect ratio plasma, low  $\omega_*$ , (Cheng, Chance '86):

$$\Omega^2 y + \partial_{\parallel}^2 y + \gamma\beta \sin\theta z = 0 \text{ (Alfvénic)} \quad (1)$$

$$\Omega^2 \left(1 + \frac{\gamma\beta}{2}\right) z + \frac{\gamma\beta}{2} \partial_{\parallel}^2 z + 2\Omega^2 \sin\theta y = 0 \text{ (acoustic)}, \quad (2)$$

where  $\Omega \equiv \omega R_0 / v_A$ ,  $y \equiv \xi_s \epsilon / q$ ,  $\xi_s \equiv \vec{\xi} \cdot \frac{[\mathbf{B} \times \nabla \Psi]}{|\nabla \Psi|^2}$  and  $z \equiv \nabla \cdot \vec{\xi}$ ,  $\hat{k}_{\parallel} \equiv i\partial_{\parallel}$ .

Coupling is due to geodesic curvature:  $m$  Alfvénic and  $m \pm 1$  acoustic harmonics.

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**Coupling is due to geodesic curvature:  $m$  Alfvénic and  $m \pm 1$  acoustic harmonics.**

**Various solutions follows (Winsor'68, Goedbloed'75, [Mikhailovski'75,'98](#), [Chu'92](#), [Turnbull '92](#), [Zonca'96](#), [van der Holst'00](#), [Breizman'05](#), [Berk'06](#)):**

- Pure acoustic modes (AMs)  $\Omega^2 = \frac{1}{2} \gamma\beta k_{\parallel}^2$ .
- Pure Alfvénic branch  $\Omega^2 = k_{\parallel}^2 + \gamma\beta (1 + 1/2q^2)$ .
- GAMs:  $\Omega^2 = \gamma\beta (1 + 1/2q^2)$  in the assumption of  $\Omega^2 \geq \gamma\beta$ .
- **Modified shear Alfvén** branch  $\Omega^2 = k_0^2 / (1 + 2q^2)$  exists for  $\Omega^2 \ll \gamma\beta$ .

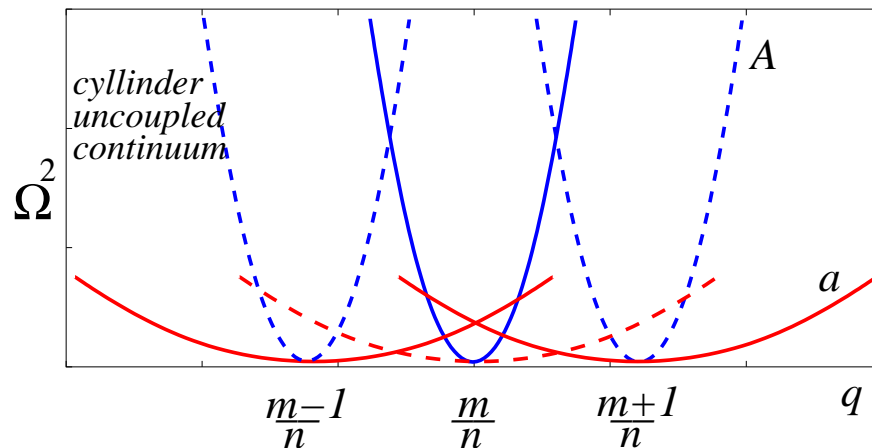
## Alfvén/acoustic coupling in toroidal equilibrium (schematic)

JET 

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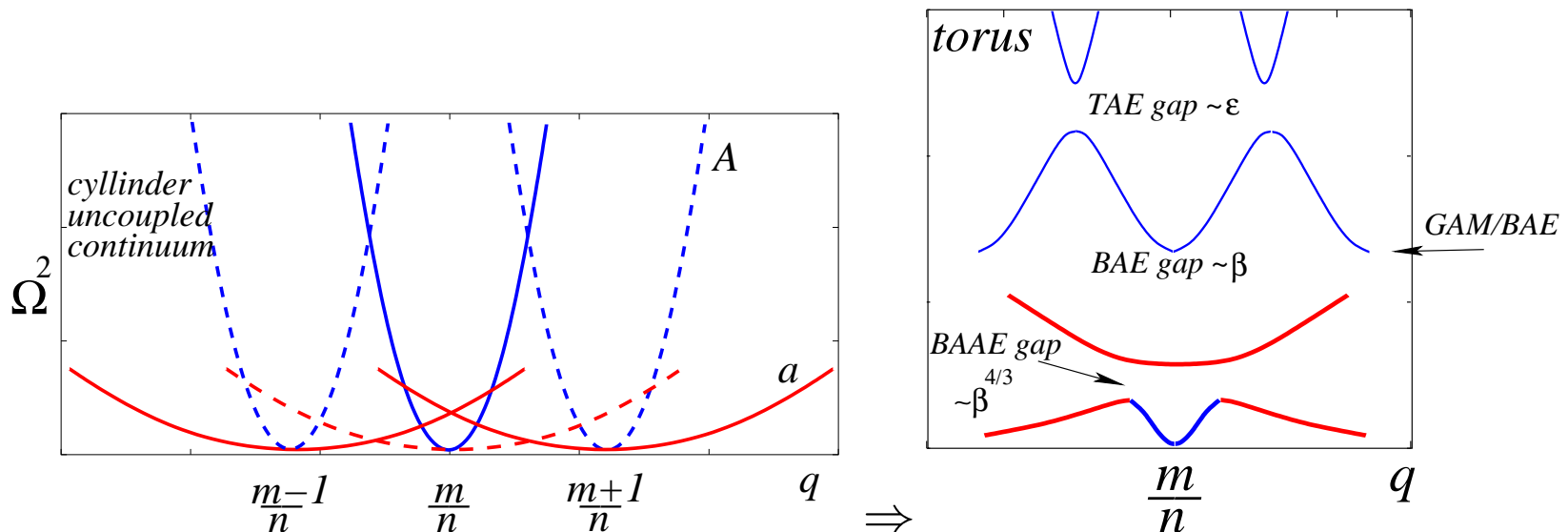
- Alfvén (A) continuum at low frequency:  $\Omega^2 = k_{0,\pm 1}^2$
- Acoustic (a) branch  $\Omega^2 = \gamma\beta k_{0,\pm 1}^2/2(1 + \delta)$



## Alfvén/acoustic coupling in toroidal equilibrium (schematic)

JET  PPPL

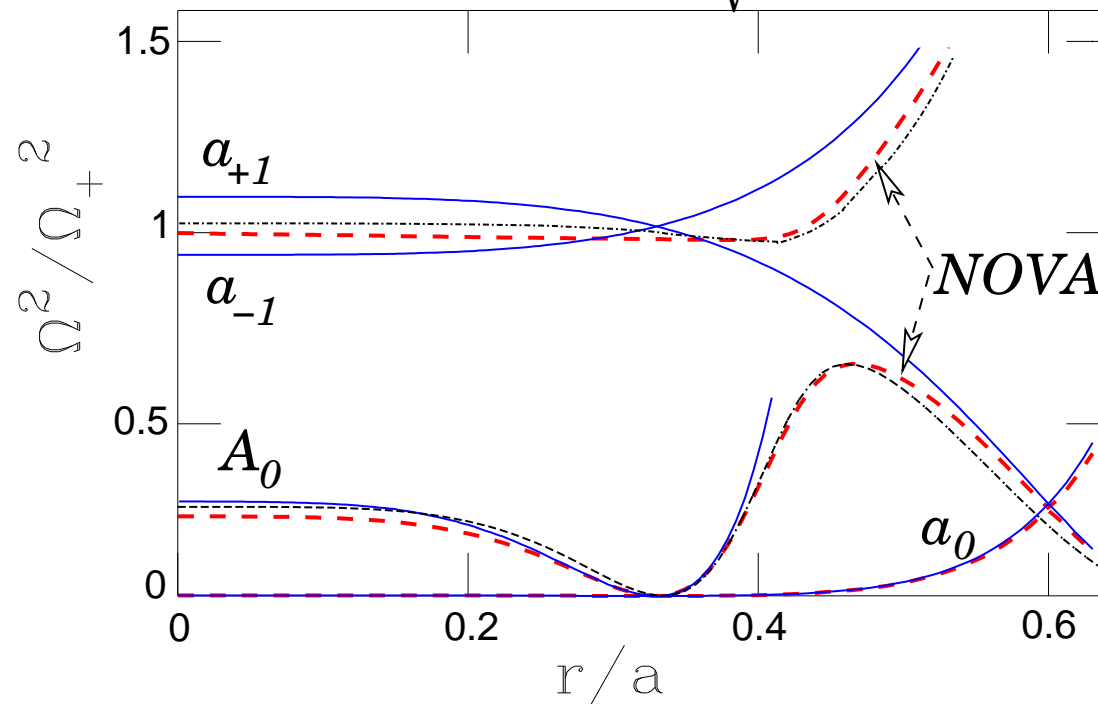
- Alfvén (A) continuum at low frequency:  $\Omega^2 = k_{0,\pm 1}^2 / (1 + 2q^2)$  (modified)
- Acoustic (a) branch  $\Omega^2 = \gamma\beta k_{0,\pm 1}^2 / 2(1 + \delta)$  is coupled via  $m \pm 1$  sidebands with modified Alfvén continuum ( $m$  harmonic) due to geodesic curvature and pressure.



## Analytic dispersion for Alfvén/acoustic continuum gap is derived

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- Consider JET, monotonic q-profile, ten times higher aspect ratio.
- $\Omega_+ = \sqrt{\gamma\beta/2}/q_r$  (compare with GAM  $\Omega = \sqrt{\gamma\beta(1 + 1/2q^2)}$ ).



Analytic dispersion - black, dashed lines.

NOVA continuum is in good agreement with theory.

## TALK OUTLINE

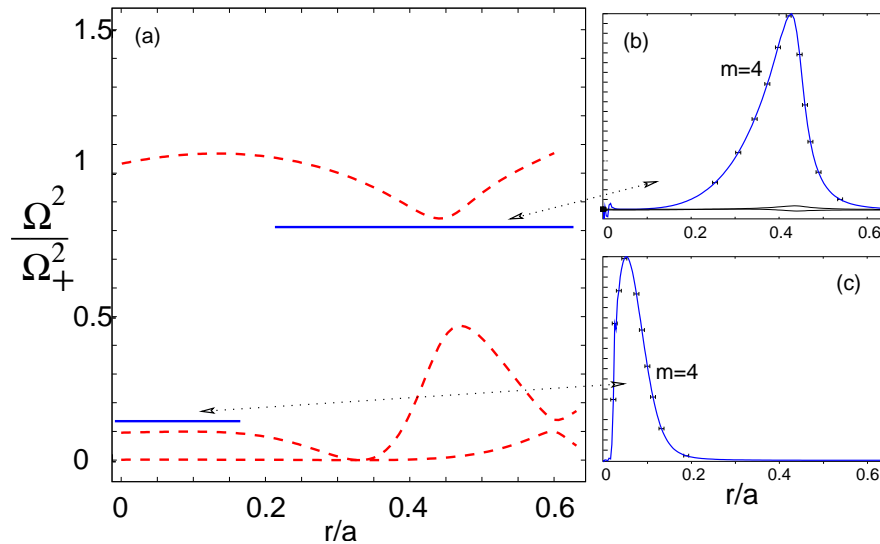
JET

PPPL

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3. NSTX analysis and data comparison
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# JET plasma analysis: two global BAAE modes are found numerically

JET PPPL

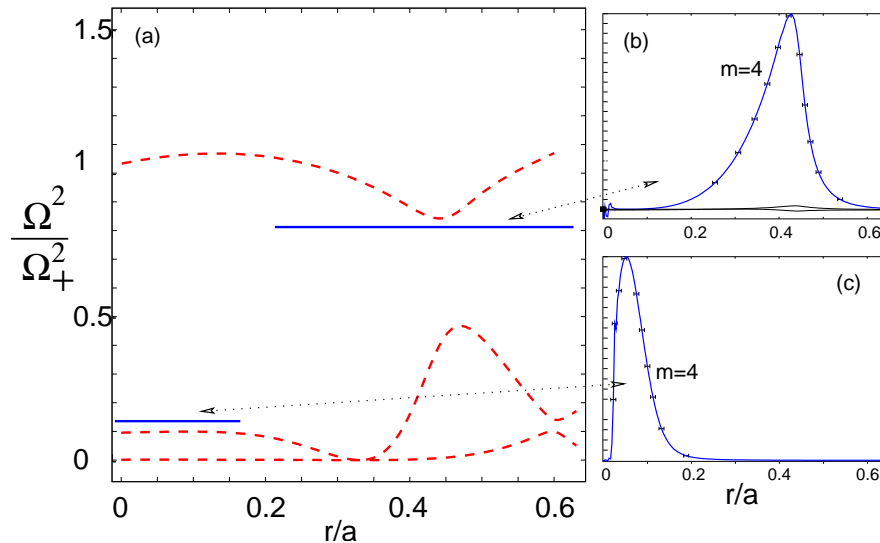


- Monotonic  $q$ -profile (EFIT),  $q_0 \simeq 1$ ,  $q_a = 4$ .
- Core localized and gap BAAEs are found with one dominant poloidal harmonic.
- Core localized BAAE has  $\omega = v_A k_{\parallel} / \sqrt{1 + 2q_{min}^2}|_{r=0}$ .
- Gap BAAE has  $\Omega_+ \simeq v_A \sqrt{\gamma\beta/2} / q_{min} R$ .
- $n = 4$ ,  $r\xi_n$  is shown.
- $\nabla\xi$ ,  $m \pm 1$  sidebands are present ( $\sim \xi_{\theta}/a$ ).



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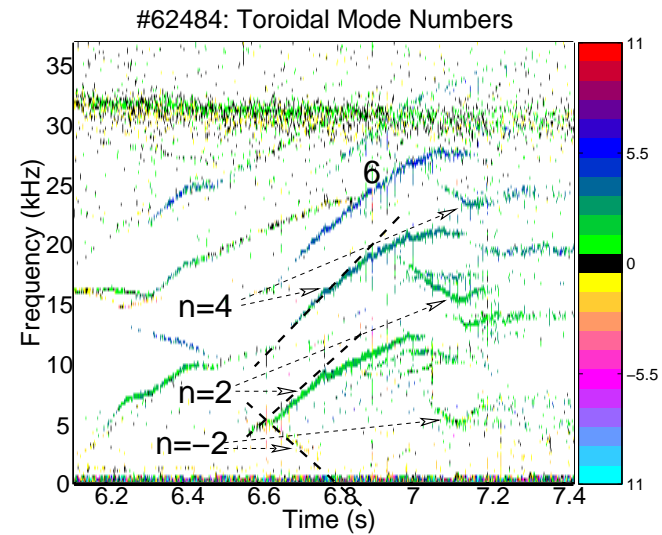
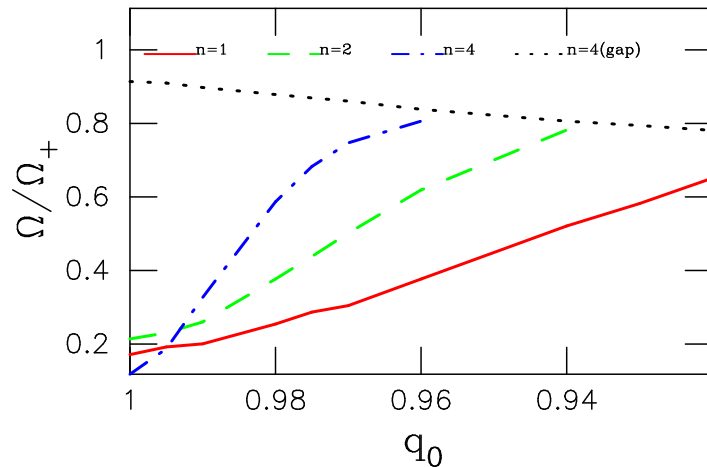
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**BAAE frequency is related to  $q_{min}$  value  $\Rightarrow$  useful for diagnostic**

## Relaxing $q$ -profile results in BAAE frequency up-sweep

JET PPPL

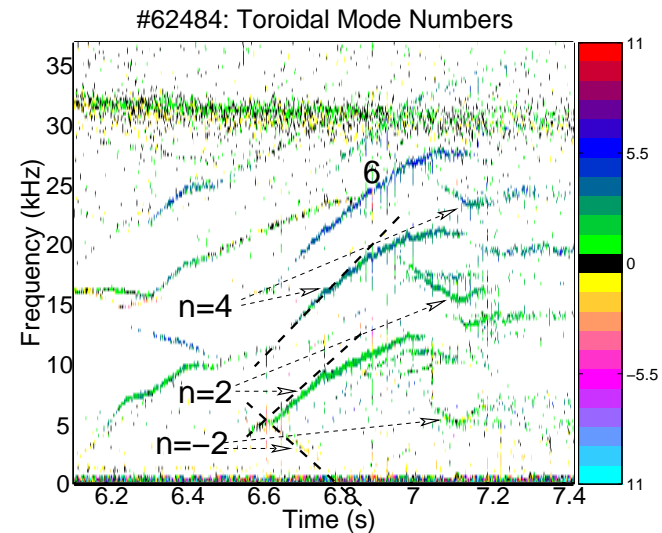
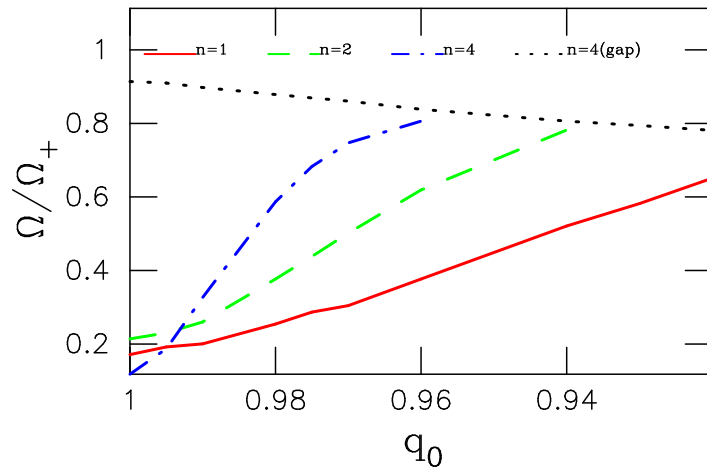
- Core BAAE activity is predicted to have sweeping frequency
- Up-chirp is limited by the gap,  $\Omega_+ \simeq v_A \sqrt{\gamma\beta/2}/qR$ .  
 $\Omega_+ = 31\text{kHz}$  for pure electron plasma.
- Rotation is significant  $f_{rot} = 2.5\text{kHz}$ .



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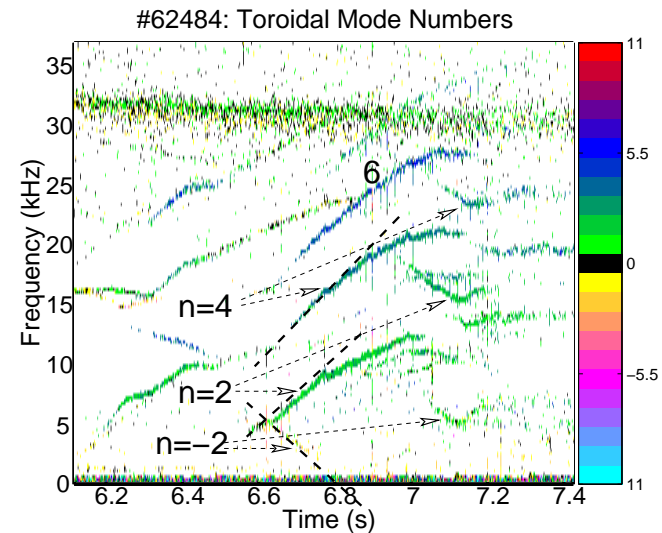
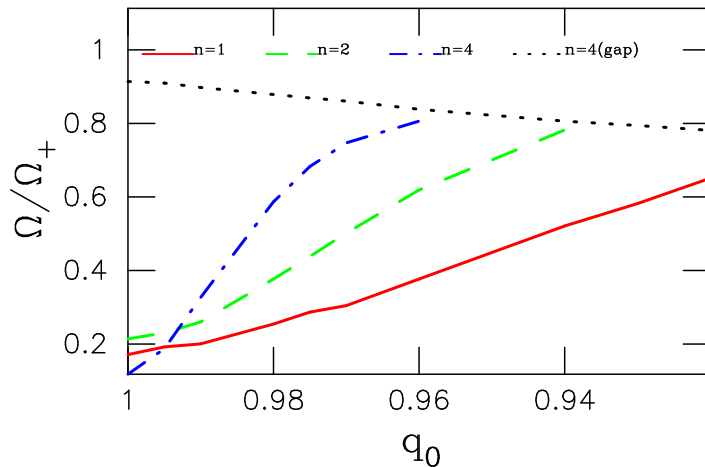


**NOVA predicts  $f_{BAAE} = 24.8\text{kHz}$  against observed  $14\text{kHz}$ , all  $n$ 's exist ( $q_0 = 1$ ).**  
**Possible way to resolve this is to assume local negative shear with  $q_{min} = 3/2$ :**  
**1) frequency  $\sim q^{-1}$ , goes down to  $\sim 16\text{kHz}$ ,**  
**2) only even  $m$ 's are expected:  $m = nq_{min}$  is integer.**

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MSE was available only later in the shot.

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PPPL

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## *NSTX experiments with MSE address frequency mismatch*

*JET* 

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*PPPL*

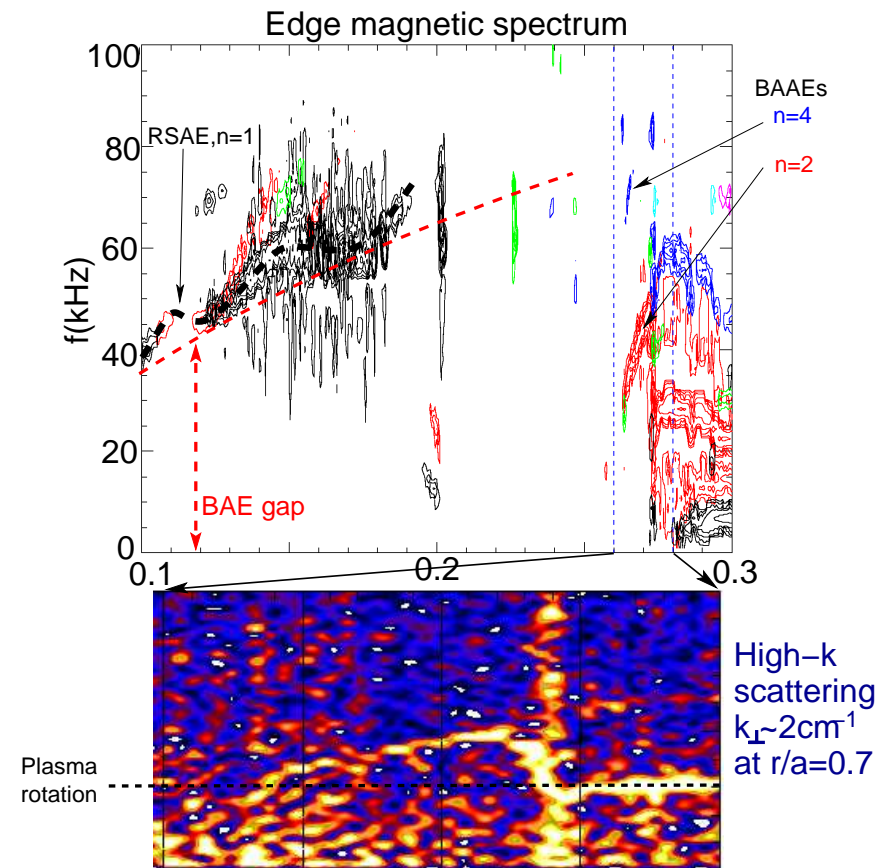
- low density  $n_e \simeq 3 \times 10^{19} m^{-3}$ ,  
 $P_{NBI} = 2MW$ ,  $E_{NBI} = 90keV$ .
- 16? channel MSE measures  $q$   
profile (reversed shear).
- Goal is to validate theory.

# NSTX experiments with MSE address frequency mismatch

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- 16? channel MSE measures  $q$  profile (reversed shear).
- Goal is to validate theory.
- Low frequency oscillations (BAAEs) are seen unstable:
  - Characteristic upshift frequency evolution from zero (plasma frame).
  - BAAEs reside in wider BAE gap  $f \sim \sqrt{\beta_{pl}}$ .
  - High-k component of BAAE at  $r/a = 0.7 \Rightarrow$  conversion to KAW (see H.Park poster).

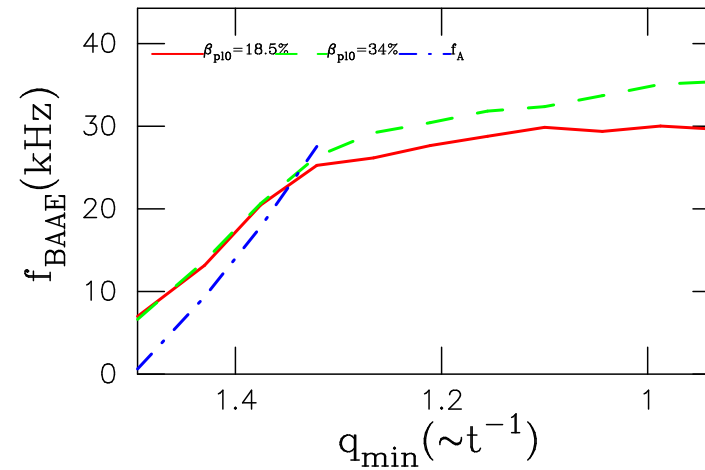
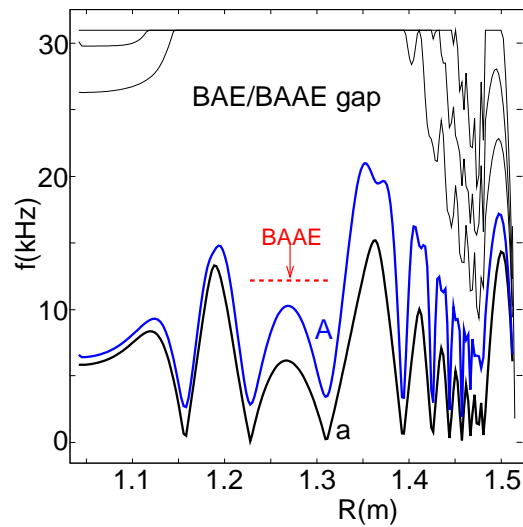


**TAE/RSAEs are suppressed (see E. Fredrickson poster) and BAAEs are excited by beams in high- $\beta$  NSTX plasmas (typically  $\beta_{pl} > \sim 15\%$ ).**

## Global BAAE modes are found at $q_{min}$ surface in NSTX

JET

PPPL



MSE measured inversed  $q$ -profile is used in NOVA modeling.

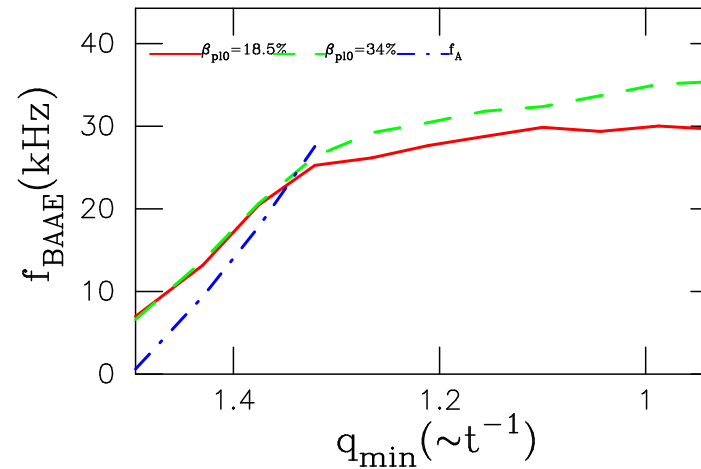
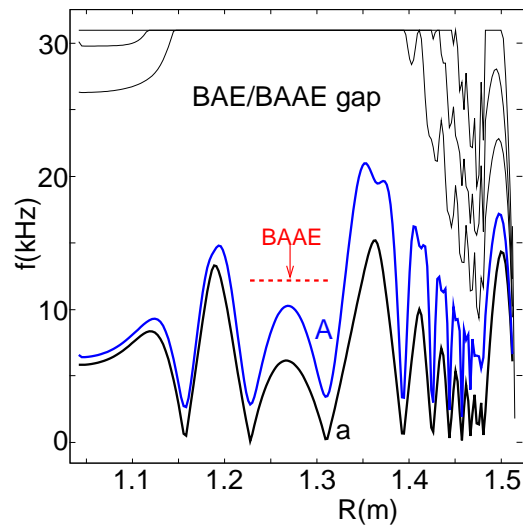
- At high- $\beta_0 = 34\%$ , BAE gap opens up to TAE gap.



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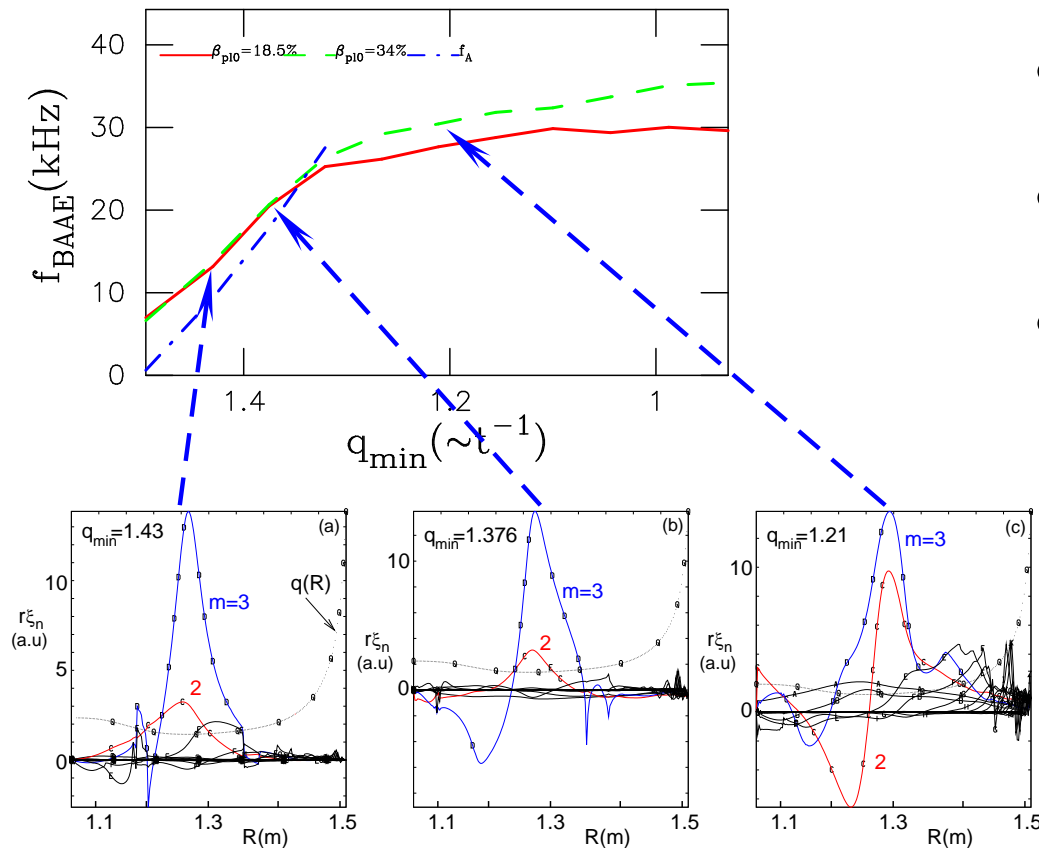
MSE measured inversed  $q$ -profile is used in NOVA modeling.

- At high- $\beta_0 = 34\%$ , BAE gap opens up to TAE gap.
- Two Alfvén/acoustic (A/a) continuum branches are found with  $\Omega^2 < \gamma\beta$ ,  $n = 2$ .
- Low shear BAAE frequency does not depend on  $\beta$  for  $q$  close to rational
- $f_{BAAE}$  is close to theory dispersion  $f_A = v_A k_{||} / \sqrt{1 + 2q_{min}^2}|_{r=0}$ .

## NOVA: BAAE broadens radially as $q_{min}$ decreases

JET

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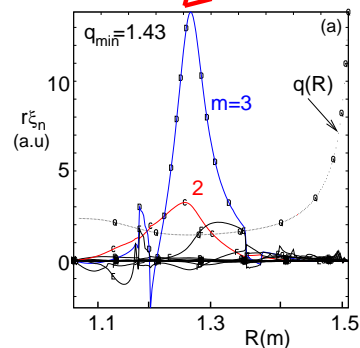
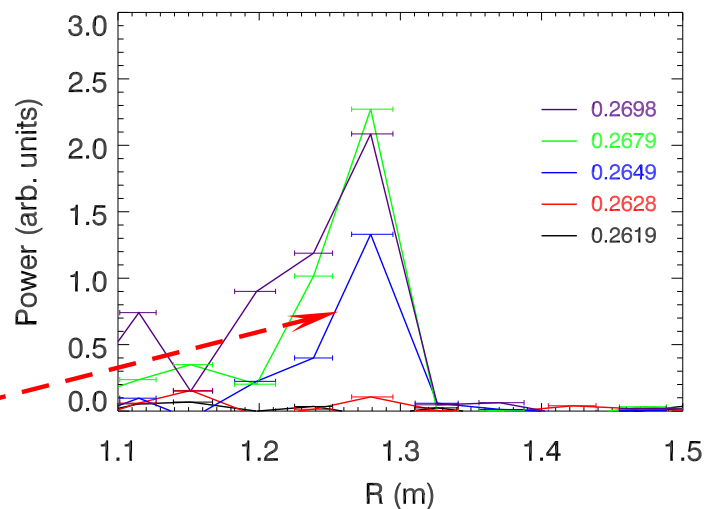
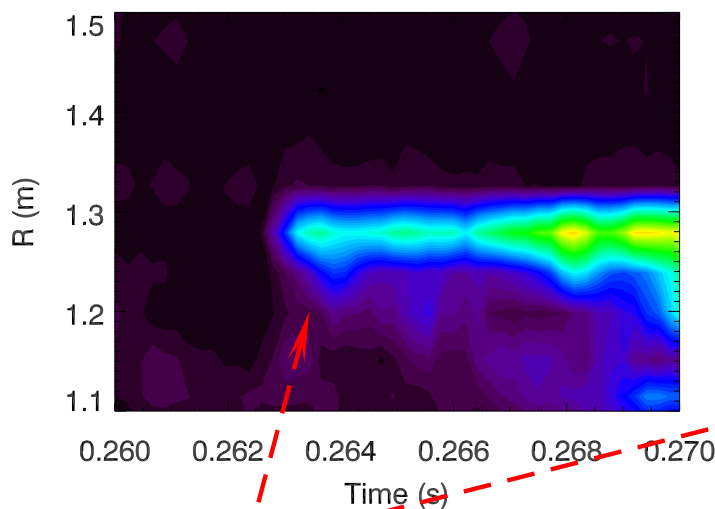
- BAAE frequency sweeps as  $q$ -profile relaxes.
- One dominant harmonic  $m = nq_{\text{min}} = 3$ .
- BAAEs interact with the continuum.

# Ultra SXR measures the same radial structure broadening

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Raw USXR signal ( $\sim$ BAAE structure)

Radial profile evolution



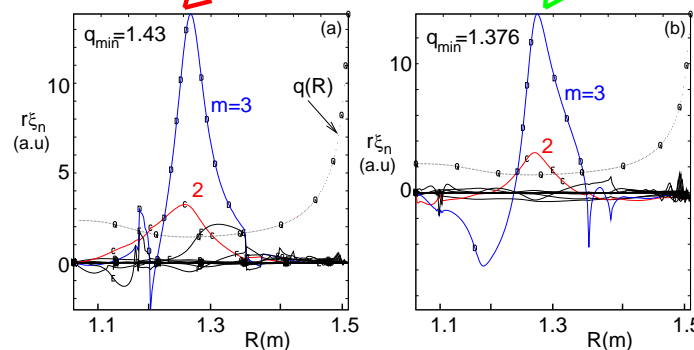
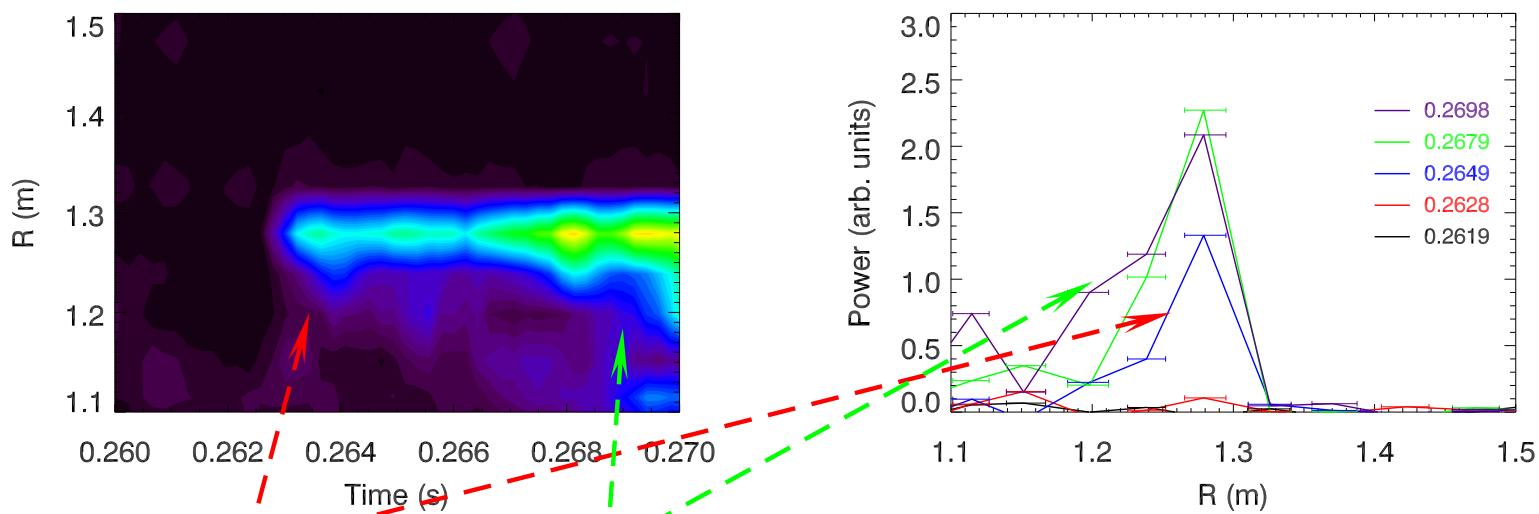
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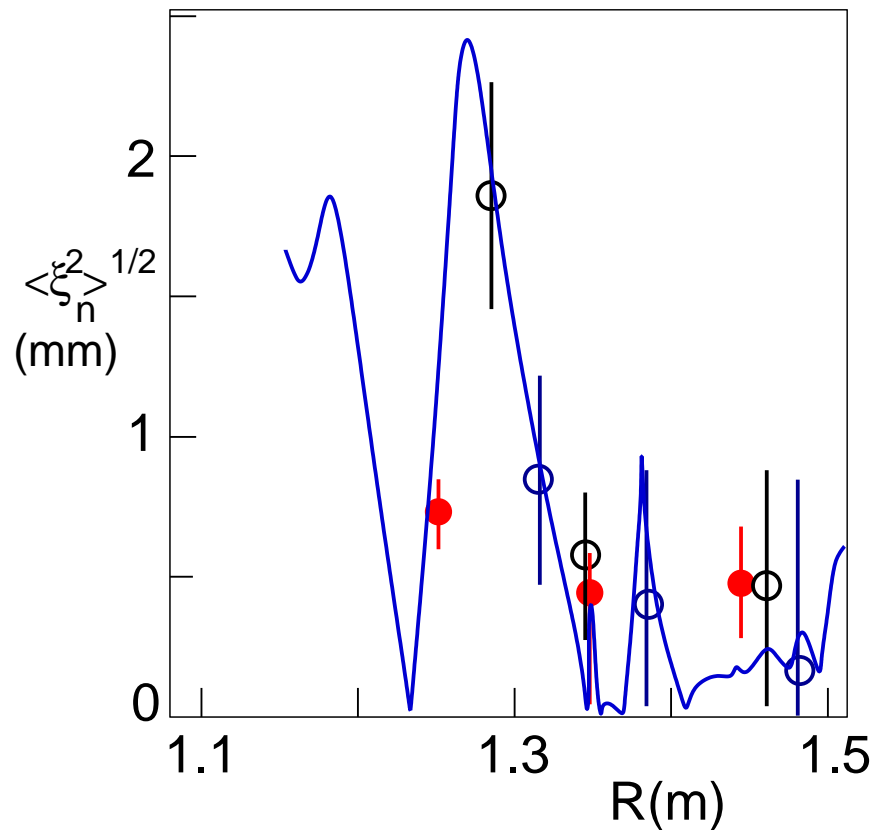


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## Reflectometer confirms localized BAAE structure

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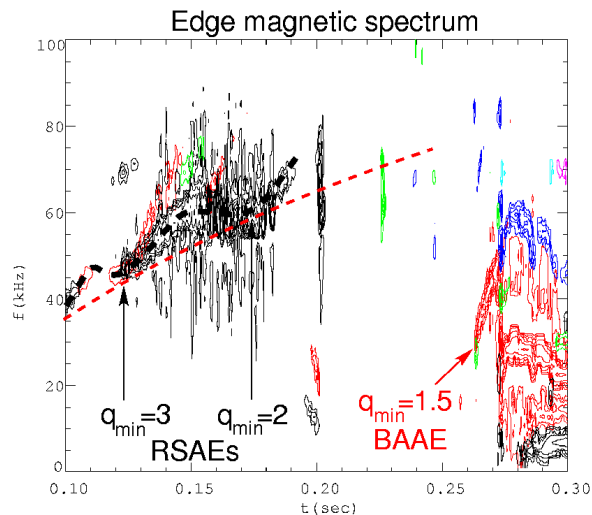


- Three plasmas with different densities  $n_e = 3 - 3.8 \times 10^{19} m^{-3}$ ,
  - each plasma has three point measurement
- Vertical axis refers to red solid circles #123816.
- Two other plasma measurements were normalized.
- Measurements are taken at signal maximum.
- Internal fluctuations level  $\delta n/n \sim 2 \times 10^{-3}$ .

## BAAEs complement RSAEs for MHD spectroscopy

JET

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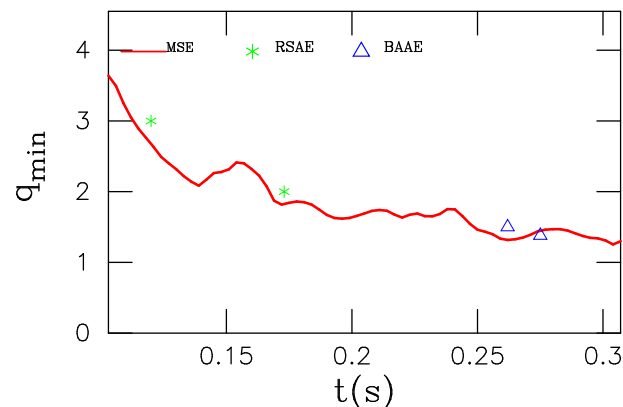
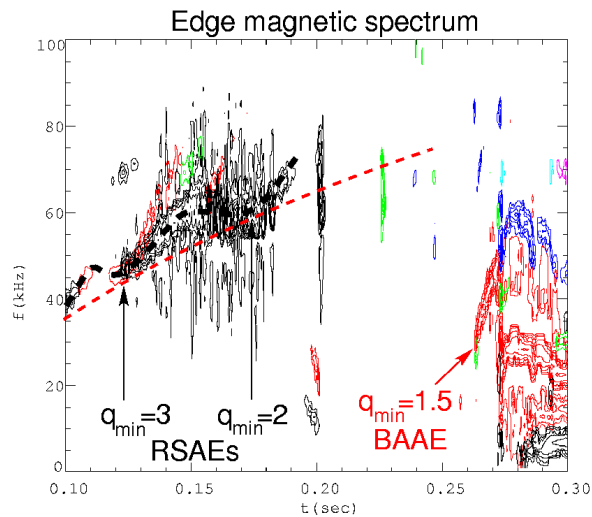


- In low- to medium- $\beta$  plasma BAAEs complement RSAE/TAEs.
- In high- $\beta$  plasma BAAEs may be the only tool for MHD spectroscopy.
- Zero BAAE frequency point (plasma frame) indicates rational  $q_{min}$ .
- BAAE activity is terminated by some IRE at  $t = 0.275s$ .  
Potential interplay of beam driven instabilities with internal  $m = 3/n = 2$  kink-like instability - similar to TAE/sawtooth nonlinear interplay (Bernabei'01, Sharapov'06).
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## *Discussion and Summary*

*JET*

*PPPL*

- Theory and numerical analysis show:
  - the existence of geodesic curvature induced gaps in the Alfvén/acoustic continuum below GAM frequency (van der Holst'00),
  - low-n global beta-induced Alfvén/acoustic eigenmodes - BAAE are found,
  - BAAEs exist in finite beta plasma within wider BAE gap.

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  - low-n global beta-induced Alfvén/acoustic eigenmodes - BAAE are found,
  - BAAEs exist in finite beta plasma within wider BAE gap.
- BAAEs are different from BAEs (Heidbrink-Turnbull-Chu-Huysmans) interpretation as BAAEs require compressibility effect, i.e. sound wave coupling:
  - frequency can sweep up from almost zero in reversed shear.
  - frequency is lower  $0 < \Omega < \sqrt{\gamma\beta/2}/q_{min}$  vs.  $\Omega = \sqrt{\gamma\beta(1 + 1/2q_{min}^2)}$  for BAE/GAM.
  - both low shear gap BAAEs can coexist (similar to RSAE/TAEs)

## Discussion and Summary

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- Theory and numerical analysis show:
  - the existence of geodesic curvature induced gaps in the Alfvén/acoustic continuum below GAM frequency (van der Holst'00),
  - low-n global beta-induced Alfvén/acoustic eigenmodes - BAAE are found,
  - BAAEs exist in finite beta plasma within wider BAE gap.
- BAAEs are different from BAEs (Heidbrink-Turnbull-Chu-Huysmans) interpretation as BAAEs require compressibility effect, i.e. sound wave coupling:
  - frequency can sweep up from almost zero in reversed shear.
  - frequency is lower  $0 < \Omega < \sqrt{\gamma\beta/2}/q_{min}$  vs.  $\Omega = \sqrt{\gamma\beta(1 + 1/2q_{min}^2)}$  for BAE/GAM.
  - both low shear gap BAAEs can coexist (similar to RSAE/TAEs)
- Kinetic modification of MHD theory is required for new global modes (Zonca'96, Mikhailovski'98):
  - damping is expected to be strong due to phase velocity of acoustic component close to thermal ion velocity.
  - dominant electron plasma is expected to be favorable for BAAE existence.

## Summary (continued)

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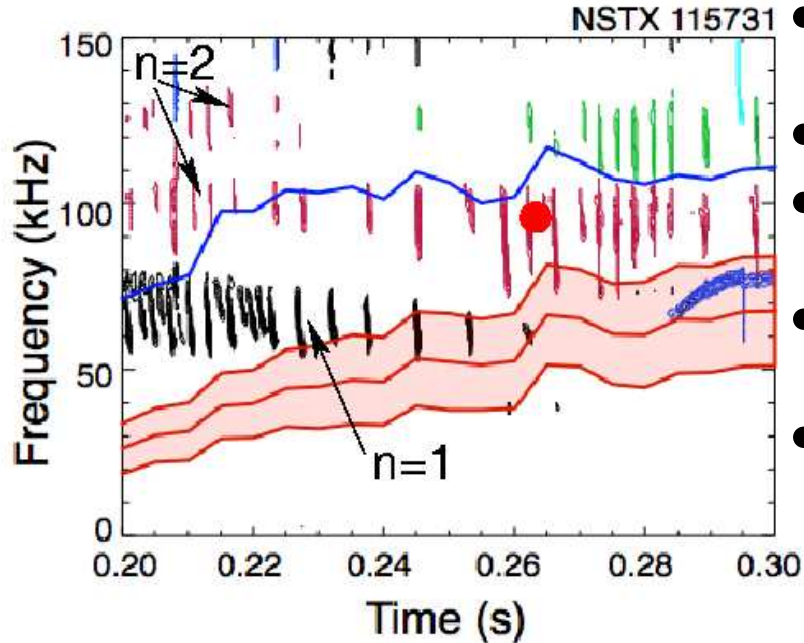
- NOVA analysis shows existence of global BAAEs in ICRH JET and NBI NSTX plasmas.
- Qualitatively NOVA predicts BAAE frequency evolution in agreement with observations on both tokamaks.
- In NSTX  $n = 2$  BAAE internal structure, frequency and their evolution are in agreement with NOVA.
  - MSE measurements on NSTX seem to validate theory and MDH ( $q_{min}$ ) spectroscopy via BAAEs.
  - Maybe useful for burning plasmas, ITER.
- For pure electron plasma (lowest  $f$ ) BAAE frequency is above the measured value in JET by factor  $\sim 1.77$ .
- Need to reconcile theory and experiment:
  - may imply local reversed shear with  $q_{min} = 1.5$  but strong indications exist for  $q_0 = 1$ ,
  - possible redistribution of the current drive due to:
    - \* MHD activity H-minority transport,
    - \* ICRH current drive,
    - \* runaway electrons in low density JET plasma.

Thank you !

## TAEs can transform to BAAE at strong drive

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- Pink shaded area is theoretical BAAE gap scaling with rotation and  $n = 2$ .
- Upper blue curve is core  $n=2$  TAE frequency.
- Red dot is NOVA BAAE frequency for  $n = 2$ : consistent with observations.
- At the start of the discharge TAE are unstable,  $f \simeq 100\text{kHz}$ .
- Later TAEs transform to EPM - BAAE:
  - Chen'94, Cheng '95, Gorelenkov '03.
  - Toroidal rotation is strongly sheared and may affect the mode localization.