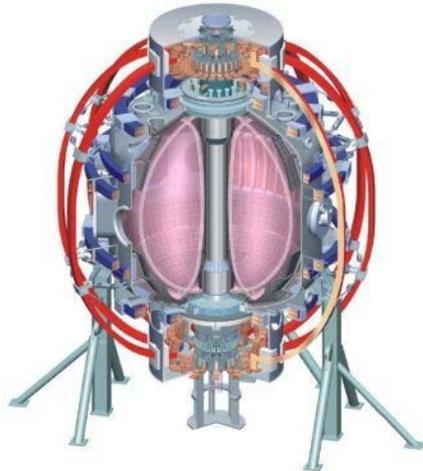


# Ideal Perturbed Equilibrium Code and Arbitrary Jump Conditions at Rational Surfaces

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**NSTX Monday Physics Meeting  
 B318, PPPL  
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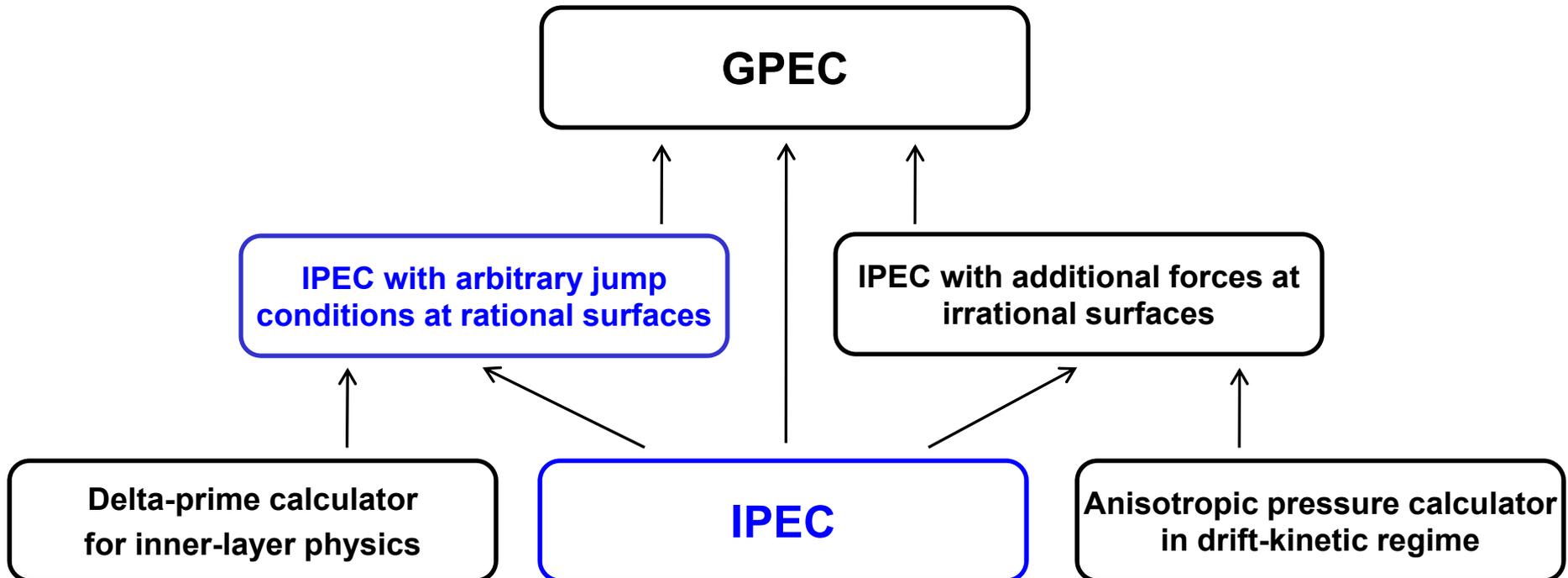


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# Motivation to perturbed equilibrium codes

- Perturbed equilibrium codes are efficient to study 3D field physics in tokamaks
  - IPEC solves ideal case
  - Future codes will include non-ideal effects



# Outline

- Features in Ideal Perturbed Equilibrium Code (IPEC)
  - How to use IPEC
  - Features in inputs and outputs
  - Important concepts
- Progress on IPEC with arbitrary jump conditions
  - Numerical implementations
  - Studies in simple cases

# How to run IPEC

- Copy
  - /p/nstxusr/nstx-users/jpark/ipec/ipec\_1.00/---, or
  - /p/nstxusr/nstx-users/jpark/ipec/ipec\_1.00/rundir/LinuxLahey64/---
- Run
  - /dcon, and then
  - /ipec
- Take a look
  - ipec\_singfld\_n#.out,
  - ipec\_xbnormal\_n#.out, etc
- IPEC is as easy as DCON to run

# How to make input

- Only two inputs : 2D equilibrium and 3D error-field
- Open and modify
  - **equil.in** for geqdsk

```
Eq_filename="./g132633.00608"
```

- **ipec.in** for error-field on the plasma boundary surface

```
harmonic_flag=t  
cosmn(2)=1.2e-4  
sinmn(3)=5.7e-5,...
```

```
jac_in="hamada"  
jsurf_in=0  
tmag_in=1
```

```
data_flag=t  
data_type="vac3d1" ("surfmn1")  
infile="./rwm_1kAt.dat"
```

```
jac_in="pest"  
jsurf_in=2  
tmag_in=0
```

- **ipec.in** for fixed boundary solution

```
fixed_boundary_flag=t
```

# Issues on input and output coordinates (I)

- IPEC supports many different coordinates
  - Magnetic coordinates are defined

$$\vec{B} \cdot \vec{\nabla} \theta \propto \frac{R^\alpha r^\beta}{B^\gamma B_p^\delta}$$

`power_rin, power_cin, power_bin, power_bpin,...`

- Surface weighted spectrum is supported

$$b_{mn} = \frac{\oint (\delta \vec{B} \cdot \hat{n})(\vartheta, \varphi) e^{-i(m\vartheta - n\varphi)} da}{\oint da}$$

`jsurf_in=1`

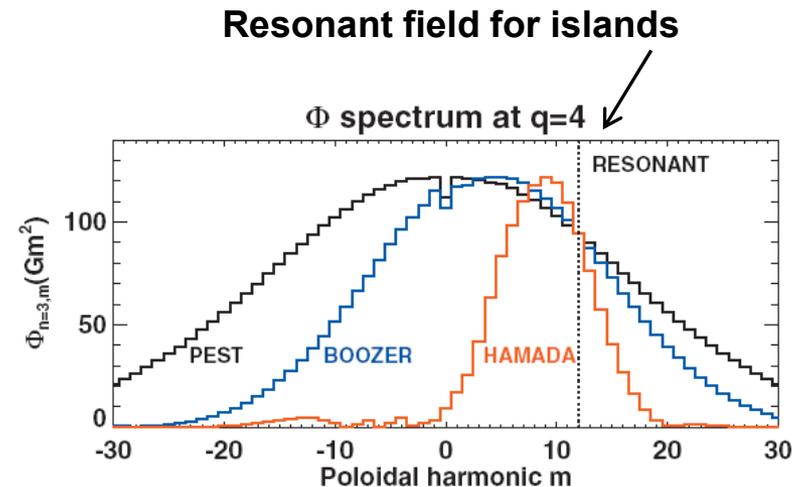
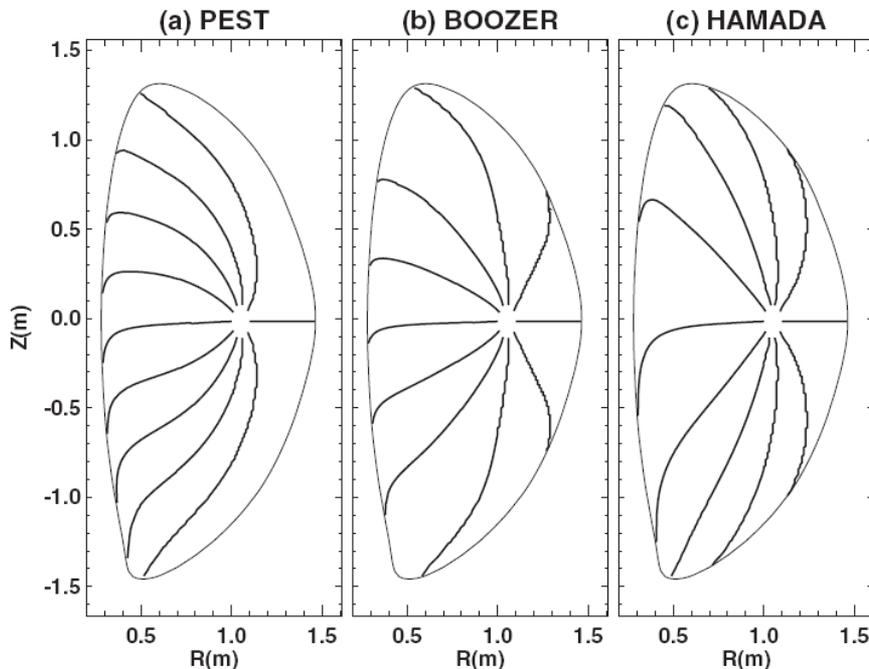
- Ordinary toroidal angle is supported

`tmag_in=0`

- Inputs and outputs can have different coordinates
- Working coordinates are different, and set by DCON

# Issues on input and output coordinates (II)

- Each coordinate has different strength
  - Hamada : best convergence, poor for inboard side
  - PEST and Equal arc : user-friendly, poor for outboard side
  - Boozer : mild convergence and accuracy everywhere
- Physics are independent on coordinates



# Output structures in IPEC (I)

- Output ascii files :

- **ipec\_response\_n#.out** : Plasma inductance, permeability, reluctance matrix, each eigenvalues and eigenvectors (**resp\_flag**)
- **ipec\_singcoup\_matrix\_n#.out** : Coupling matrix between external fields and total resonant fields, singular currents, islands (**singcoup\_flag**)
- **ipec\_singcoup\_svd\_n#.out** : The  $i^{\text{th}}$  important field for resonant fields, etc, and each singular values, coupling matrix (**singcoup\_flag**)

**When error-field is provided,**

- **ipec\_control\_n#.out** : Boundary response, amplification
- **ipec\_singfld\_n#.out** : Resonant fields, induced islands, Chirikovs (**singfld\_flag**)
- **ipec\_pmod\_n#.out** : Eulerian and Lagrangian  $|b|$  profiles (**pmodb\_flag**)
- **ipec\_xbnormal\_n#.out** : Normal  $\xi$  and  $b$  field profiles (**xbnormal\_flag**)
- **ipec\_brzphi\_n#.out** :  $(r,z,\phi)$  components of perturbed  $b$  fields in  $(r,z,\phi)$  coordinates (**brzphi\_flag**, **xrzphi\_flag**, **vbrzphi\_flag**, **vpbrzphi\_flag**, etc)

# Output structures in IPEC (II)

- Output binary files
  - `xbnormal.bin`, `ipec_pmodb.bin` : profiles for 1d xdraw (`bin_flag`)
  - `xbnormal_2d.bin`, etc : profiles for 2d xdraw (`bin_2d_flag`)
  - `pflux_re_2d.bin` : perturbed flux surfaces (`bin_2d_flag`)
  - `bnormal_spectrum.bin` : surfmn type plots for b normal fields
- Run
  - `/xdraw xbnormal`, etc
- Xdraw is a powerful tool to examine results promptly, by various interactive features, but
- Other tools (such as IDL) for ascii files are needed to make journal-quality figures

# Output : normal displacements and fields

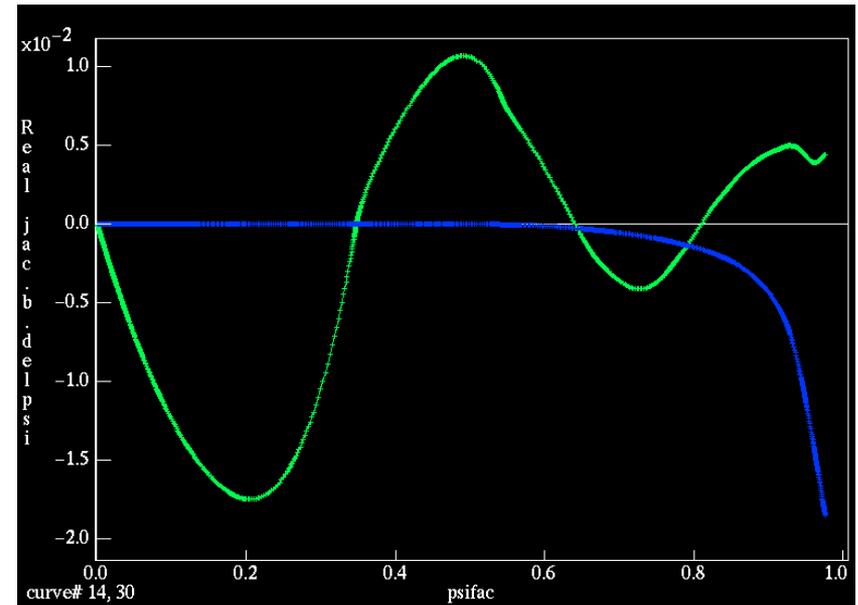
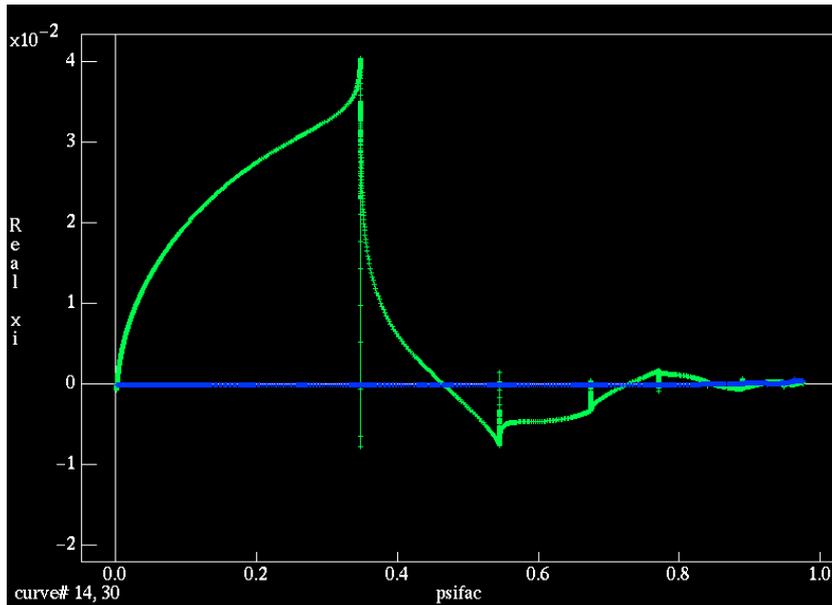
- See `ipec_xbnormal_n#.out` for Fourier harmonic profiles
- `/xdraw xbnormal`

g132633.00608 with n=1 EF currents

Normal displacement

$m=2, m=18$

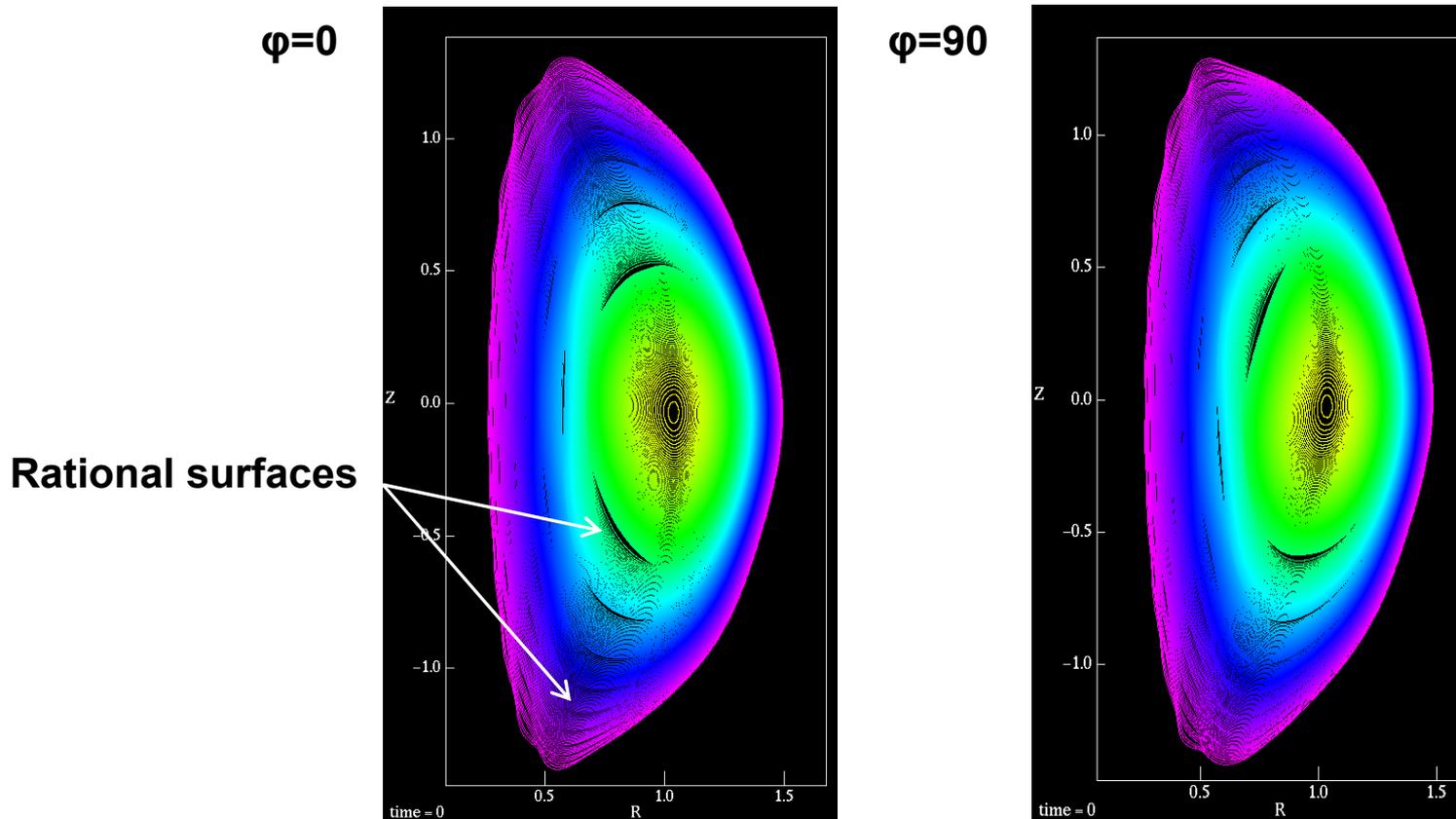
Normal b field



# Output : perturbed flux surfaces and fields

- See `ipec_xbnormal_n#.out` for fourier harmonic profiles
- `/xdraw xbnormal_2d`, or `/xdraw pflux_re_2d`

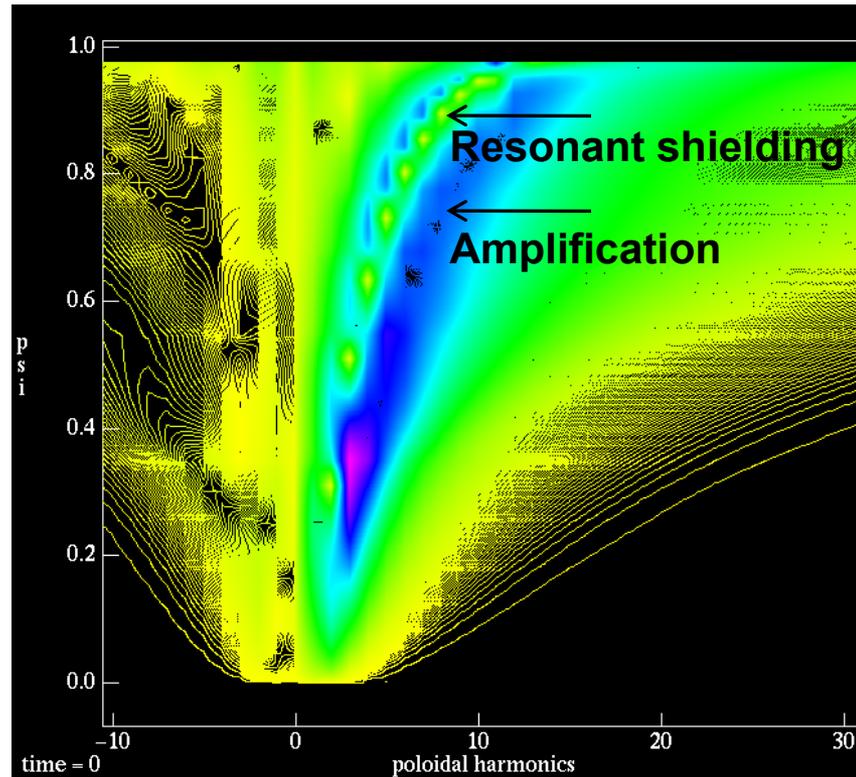
g132633.00608 with n=1 EF currents



# Output: surface weighted b field spectrum

- See `ipec_xbnormal_n#.out` for fourier harmonic profiles
- `/xdraw bnormal_spectrum`

g132633.00608 with n=1 EF currents

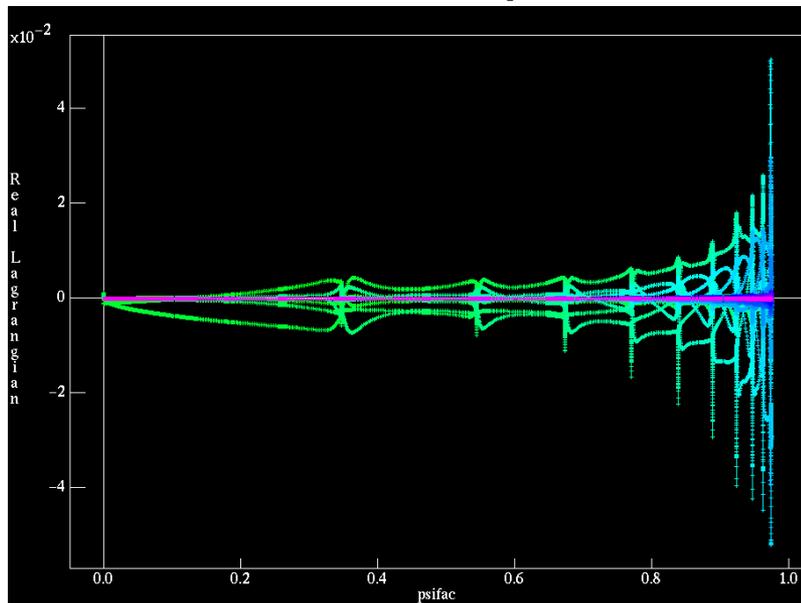


# Output: perturbed $|b|$ profiles

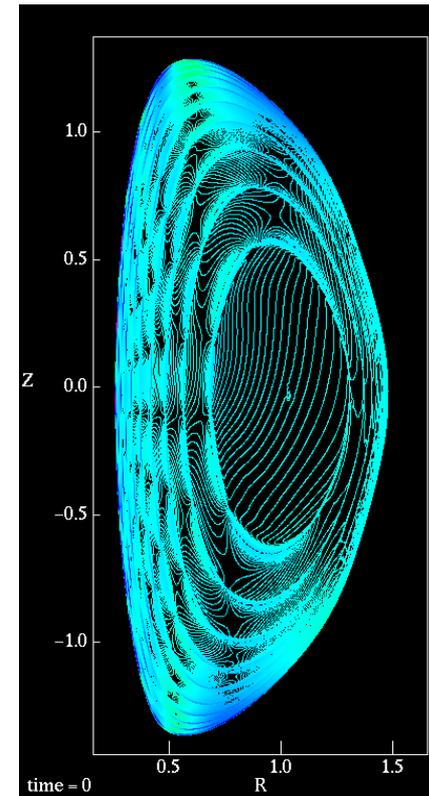
- See `ipec_pmodb_n#.out` for fourier harmonic profiles
- `/xdraw pmodb`, or `/xdraw pmodb_2d`

g132633.00608 with n=1 EF currents

Each harmonic profiles



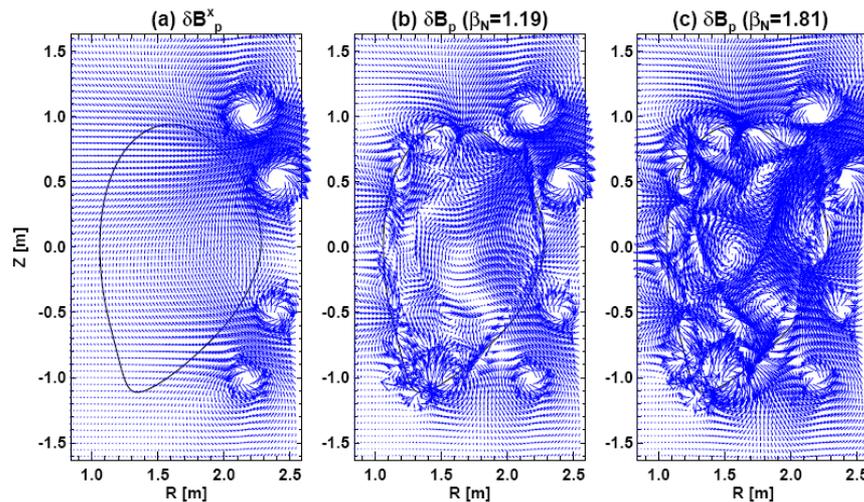
Contour



- NTV routines are separate and not yet implemented in IPEC

# Output : perturbed $\xi$ and $b$ vectors in (r,z,phi)

- See `ipec_*rzphi_n#.out` for data
  - Vacuum fields outside the plasma is not provided from IPEC
  - Use `ipec_pbrzphi_n#.out` and add your vacuum fields outside ( $I=0$ )
  - See `ipec_pbrzphi` (or `vpbrzphi`)\_n#.out to have plasma response
  - See `ipec_xrzphi_n#.out` to have plasma displacements



- Poincare tracing code (S. Hudson) is separate

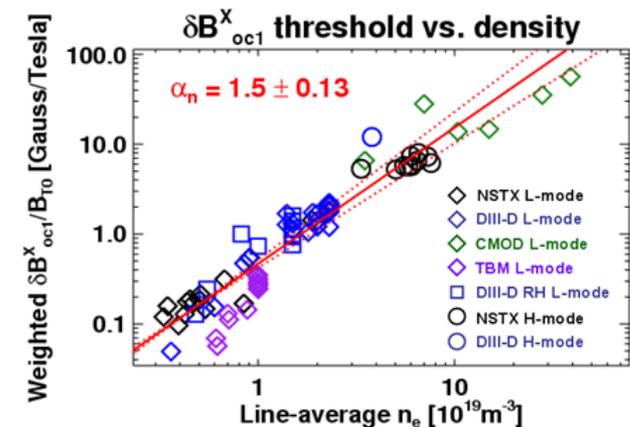
# Output : resonant fields, islands, Chirikovs

- See `ipec_singfld_n#.out` for resonant fields, islands, etc

q	psi	real(singflx)	imag(singflx)	real(singcur)	imag(singcur)	islandwidth	chirikov
2.000	3.46826953E-01	4.69650517E-04	-5.88121248E-04	1.17667469E+03	-1.47349436E+03	7.61589337E-02	7.72369004E-01
3.000	5.44035627E-01	1.82758806E-04	-2.42367945E-04	4.43245509E+02	-5.87815740E+02	3.85164222E-02	5.93357690E-01
4.000	6.73860934E-01	6.10155924E-05	-7.92386151E-05	1.41461753E+02	-1.83710966E+02	2.03607632E-02	4.21800633E-01
5.000	7.70403052E-01	4.94533020E-05	-6.13473827E-05	1.10785858E+02	-1.37431115E+02	1.59441126E-02	4.63529414E-01
6.000	8.39197447E-01	6.91010435E-05	-7.94839799E-05	1.49261365E+02	-1.71688974E+02	1.58277653E-02	6.43480303E-01
7.000	8.88391696E-01	6.79358284E-05	-7.34817892E-05	1.41171858E+02	-1.52696464E+02	1.33011570E-02	7.48965610E-01
8.000	9.23910435E-01	1.57830604E-04	-1.96006561E-04	3.14671694E+02	-3.90784265E+02	1.76860405E-02	1.52374141E+00
9.000	9.47124401E-01	2.88304755E-04	-3.85648019E-04	5.46143270E+02	-7.30543172E+02	1.97119845E-02	2.52142928E+00
10.000	9.62759965E-01	4.93641101E-04	-5.63773598E-04	8.84543509E+02	-1.01021223E+03	2.07356092E-02	3.66874048E+00
11.000	9.74063904E-01	2.23222979E-04	-1.45122837E-04	3.77801042E+02	-2.45617899E+02	1.06256207E-02	1.87998552E+00

- When `singcoup_flag=t`, overlap fields will also be provided
  - This is being used for locking thresholds

mode	real(ovf)	imag(ovf)	overlap(%)
1	-1.59566880E-04	2.09106744E-04	3.03126685E+01
2	-1.12919170E-04	1.25480940E-04	1.94538134E+01
3	2.22724911E-05	-4.13769485E-05	5.41529532E+00
4	2.69295167E-06	-2.98598016E-05	3.45507274E+00
5	7.35413851E-06	-5.78315646E-05	6.71830237E+00
6	-6.01999767E-05	7.27180964E-05	1.08792214E+01
7	5.92764183E-05	1.23037725E-06	6.83261137E+00
8	-5.60090623E-06	6.11843237E-05	7.08049261E+00
9	-3.00285513E-05	-3.81663356E-05	5.59652444E+00
10	2.09479454E-05	-1.07832051E-04	1.26591081E+01



# Output : response characteristics

- See **ipec\_response\_n#.out** for response matrices
  - **Surface and plasma inductance** : Virtual casing currents to fields
  - **Plasma permeability** : External fields to total fields
  - **Eigenvalues and eigenvectors for each matrix** show intrinsic properties of plasma response each matrix
- See **ipec\_singcoup\_matrix\_n#.out** for resonant coupling
  - **The  $i^{\text{th}}$  important field distributions** are complete set of distributions affecting resonant fields
  - **Singular values** show how important they are in error field correction problems
- Coupling for NTV is being planned

# Important notes for IPEC run

- IPEC results are robust enough for low- $\beta$  plasmas, but sensitivity exists in high- $\beta$ , H-mode plasmas, as DCON
- Boundary cutting issue :
  - If  $p' \neq 0$  in edge, results depend on the relative locations of rational surfaces
  - Put boundary at the first half in a q-interval (**sas\_flag=t, dmlim=0.1-0.5 in dcon.in**) to have robust answer
- Toroidicity issue:
  - NSTX is the most difficult,  $n=1$  is ok, but  $n \geq 2$  is often unreliable
  - Make sure plasma is far from stability limit for  $n \geq 2$
- Internal instability issue:
  - Refine equilibrium when DCON finds internal instability, or
  - Avoid  $q=1$  surface by setting **psi\_low > psi (q=1) in dcon.in**

# Summary and future work (I)

- IPEC is easy to run (/dcon, /ipec, /xdraw with equil.in, dcon.in, ipec.in)

I am still testing some features and preparing manuals, and these jobs will be done **by the end of June**

- Vacuum (Bio-Savart) part will be implemented **by the end of July**
- NTV and Poincare tracing routines are presently separate, but planned for implements **by this year**
- These new features will be version-controlled
- IPEC can be sensitive as DCON, but issues will be constantly studied

# Jump conditions are always necessary to solve force-balance equations

- DCON solves Euler-Lagrange equation for displacements

$$\vec{\Xi} = \left\{ \left( \vec{\xi} \cdot \vec{\nabla} \psi \right)_{mn} \mid m_{\min} \leq m \leq m_{\max}, m_{\max} - m_{\min} + 1 = M \right\}$$

$$\begin{bmatrix} \vec{\Xi} \\ \vec{F} \cdot \vec{\Xi}' + \vec{K} \cdot \vec{\Xi} \end{bmatrix}' = \begin{bmatrix} -\vec{F}^{-1} \cdot \vec{K} & \vec{F}^{-1} \\ \vec{G} - \vec{K}^+ \cdot \vec{F}^{-1} \cdot \vec{K} & \vec{K}^+ \cdot \vec{F}^{-1} \end{bmatrix} \begin{bmatrix} \vec{\Xi} \\ \vec{F} \cdot \vec{\Xi}' + \vec{K} \cdot \vec{\Xi} \end{bmatrix} \longrightarrow \frac{d\vec{u}}{d\psi} = \vec{L} \cdot \vec{u}$$

- 2M independent solutions exist depending on initial conditions
- F operator is singular at rational surfaces, and thus 2M independent solutions should be solved between rational surfaces

$\frac{d\vec{U}_1}{d\psi} = \vec{L} \cdot \vec{U}_1$ 
 $\frac{d\vec{U}_2}{d\psi} = \vec{L} \cdot \vec{U}_2$ 
 $\frac{d\vec{U}_3}{d\psi} = \vec{L} \cdot \vec{U}_3$

$\psi = 0$ 
 $\psi(q1)$ 
 $\psi(q2)$ 
 $\psi = 1$

# Ideal DCON and IPEC uses only M independent solutions to match boundary

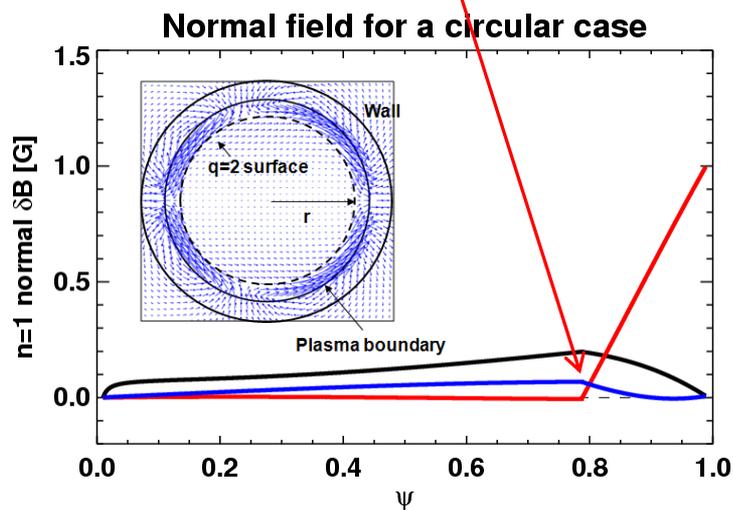
- DCON integrates equations from the magnetic axis – where only regular solutions are allowed  $\vec{\Xi} = \vec{0}$

- Only the half part starts from the axis  $\vec{U} = \begin{pmatrix} \vec{I} & 0 \\ 0 & \vec{I} \end{pmatrix}$

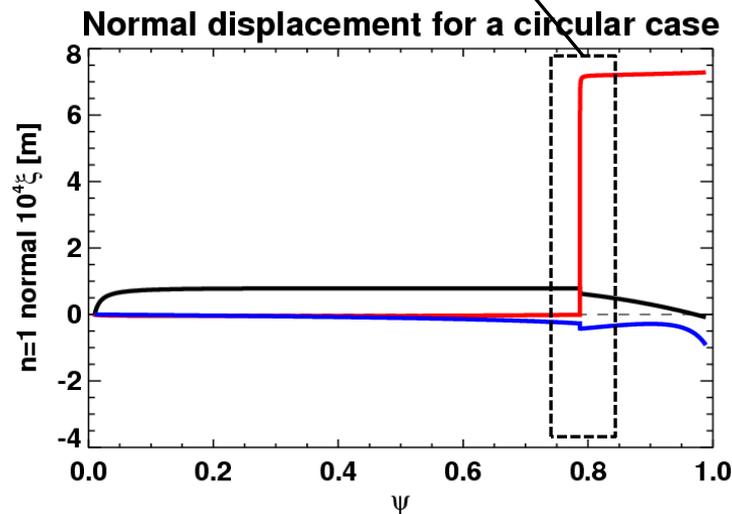
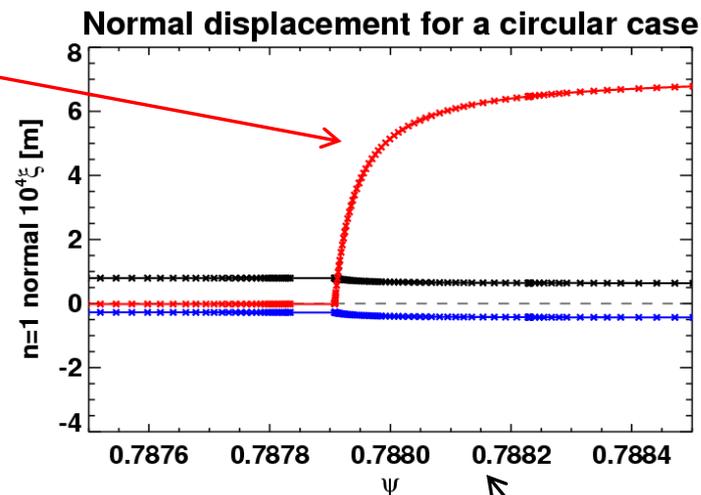
- At rational surfaces
  - **Non-resonant solutions continuously cross**
  - **Large resonant solutions are forced to be zero**
  - **Small resonant solutions are reinitialized**
  - **Gaussian eliminations (mixing initial conditions) are used to sort out resonant solutions**
- At the boundary, M independent solutions reach and can be determined by boundary conditions on the plasma boundary

# Ideal solutions have zero resonant fields by shielding currents suppressing islands

- Small resonant solution guarantees no resonant field at the rational surface



- Ideal DCON is very precise by preserving small solutions – which are easy to handle in the frame of displacements



# Outer-layer solutions can be determined differently with arbitrary jump conditions

- Outer-layer is more ideal
- Outer-layer equations are still E-L equations, and they can be differently determined if inner-layer gives different jump conditions
  - Full relaxation (smooth)
  - Magnetic island size (resonant field)
  - Delta-prime, etc
- Large solutions are now necessary, but their stronger singularity must be more carefully treated

# Stepped pressures may be necessary when large resonant solutions are retained

- Resonant solutions are

$$\begin{bmatrix} \xi \\ f\xi' \end{bmatrix} = c_s \begin{bmatrix} \text{Small} \\ z^{-1/2+\sqrt{-D_I}} \\ z^{1/2+\sqrt{-D_I}} (\sqrt{-D_I} - 1/2) f_0 \end{bmatrix} + c_l \begin{bmatrix} \text{Large} \\ z^{-1/2-\sqrt{-D_I}} \\ z^{1/2-\sqrt{-D_I}} (-\sqrt{-D_I} - 1/2) f_0 \end{bmatrix}$$

where

$$z = \psi - \psi_R$$

$$D_I = -\frac{1}{4} + \frac{2\pi f P'}{q' \chi'^3} \left\langle \frac{1}{|\nabla V|^2} \right\rangle \left[ 1 - \frac{2\pi f P'}{q' \chi'^3} \left\langle \frac{1}{|\nabla V|^2} \right\rangle \right] + \left\langle \frac{B^2}{|\nabla V|^2} \right\rangle \frac{P'}{q'^2 \chi'^4} \left\{ P' \left[ \left\langle \frac{1}{B^2} \right\rangle + \left( \frac{2\pi f}{\chi'} \right)^2 \left\langle \frac{1}{B^2 |\nabla V|^2} \right\rangle \right] + \frac{\chi''}{\chi'} \right\}$$

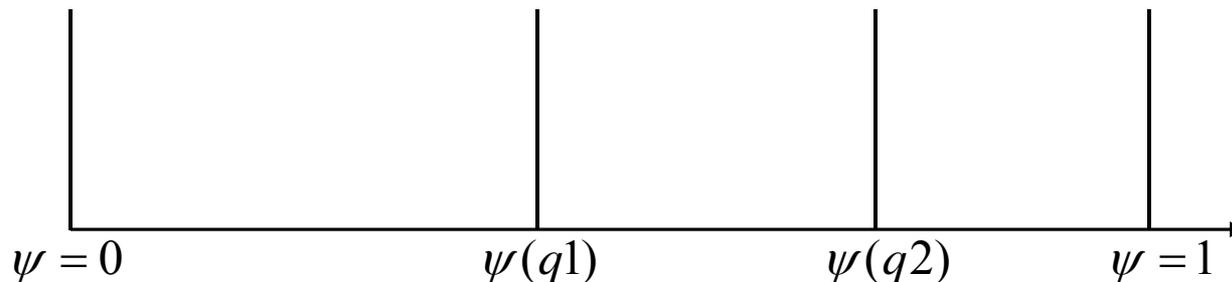
$$\delta B \propto z \xi$$

- Large solutions make resonant field singular unless  $D_I \geq -\frac{1}{4}$
- Stepped pressure  $P'=0$  ensures  $D_I = -\frac{1}{4}$  and resonant solutions can reach a constant

# Arbitrary jump conditions have been implemented in IPEC

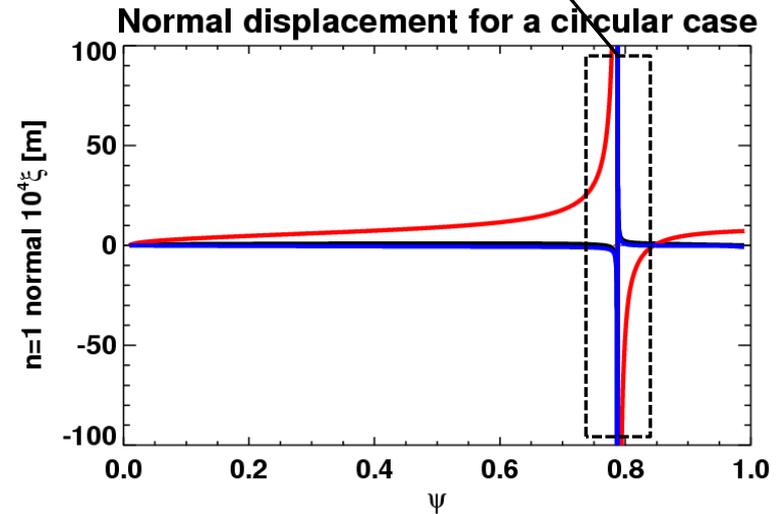
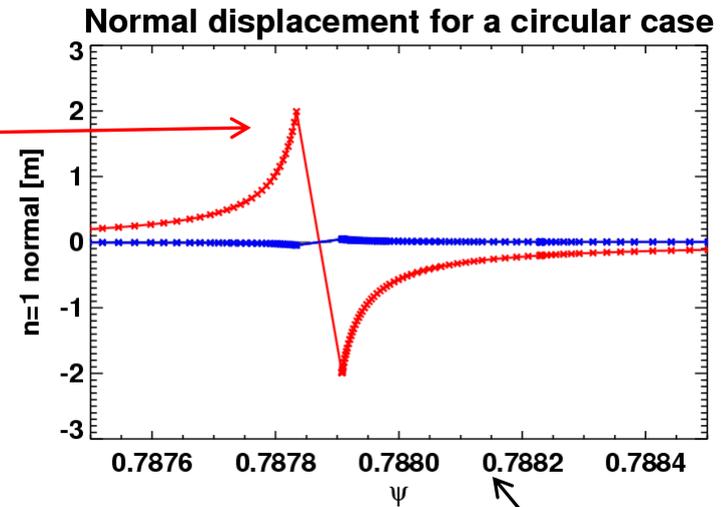
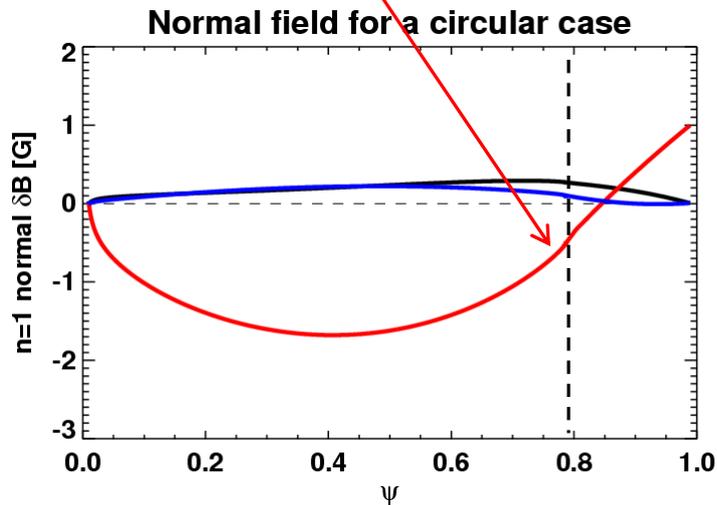
- Large solutions are also reinitialized and preserved as well as small solutions at each rational surface
  - Asymptotic power series are used to identify small and large solutions
- In each domain, one can determine  $M+1$  solutions using  $M$  edge boundary conditions and 1 jump conditions (Backward matching)
  - Displacement should change parity when crossing rational surfaces

$$\begin{bmatrix} u_{11}(\psi_R) & \cdots & u_{1M}(\psi_R) & l_{1M}(\psi_R) \\ \vdots & \ddots & \vdots & \vdots \\ u_{M1}(\psi_R) & \cdots & u_{MM}(\psi_R) & l_{MM}(\psi_R) \\ u_{L1}(\psi_L) & \cdots & u_{LM}(\psi_L) & l_{LM}(\psi_L) \end{bmatrix} \begin{bmatrix} c_1 \\ \vdots \\ c_M \\ c_l \end{bmatrix} = \begin{bmatrix} \xi_1(\psi_R) \\ \vdots \\ \xi_M(\psi_R) \\ \xi_l(\psi_L) \end{bmatrix}$$



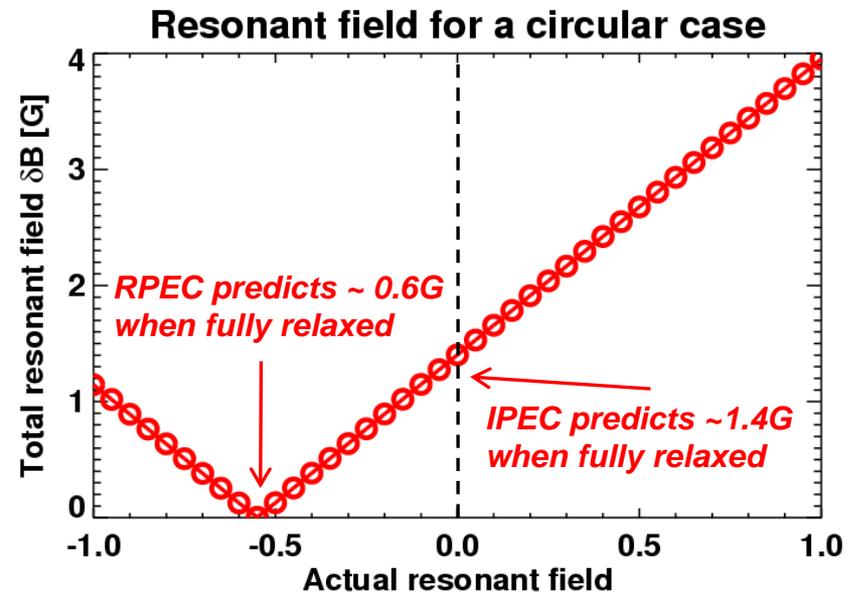
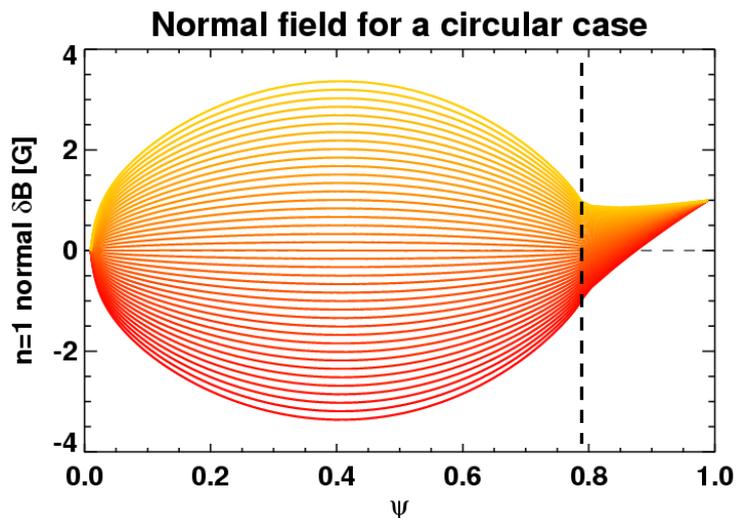
# Relaxed solutions have finite resonant fields with finite shielding currents

- Large solutions are preserved and cross the rational surface with jump conditions
- Singular currents can be removed and resonant fields are finite



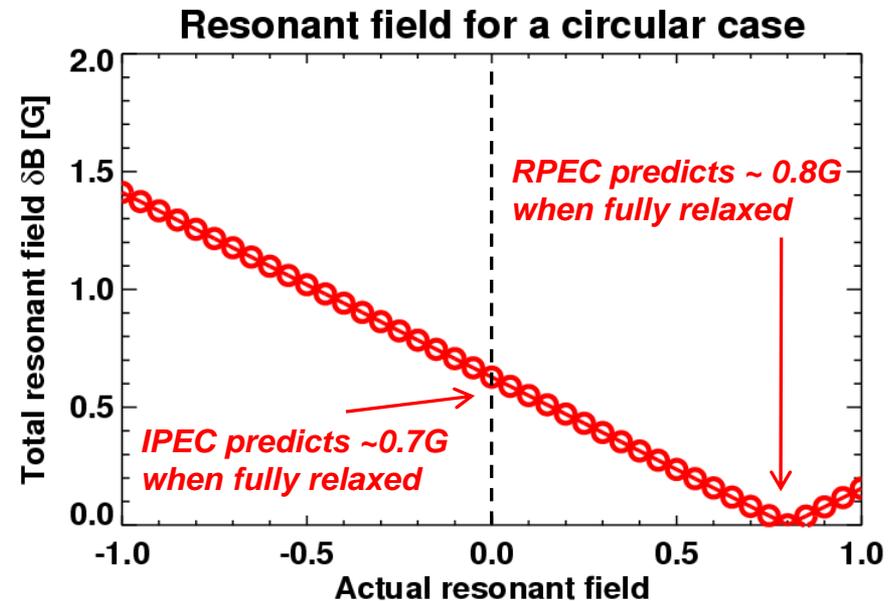
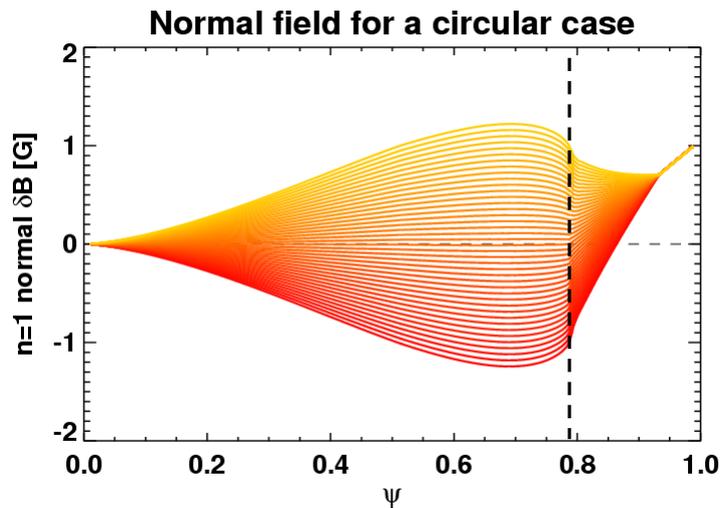
# Solutions with islands can be unstable

- Fixed boundary cylinder analysis with  $m=2, n=1$ 
  - Islands open with more singular currents unless crossing zero
  - This is unstable and  $dW < 0$  mode exists



# Solutions with islands can be stable

- Fixed boundary cylinder analysis with  $m=4$ ,  $n=2$ 
  - Islands open with less singular currents
  - This is stable and all modes  $dW > 0$



# Summary and future work (II)

- IPEC has been modified towards arbitrary jump conditions
- Initial test shows that islands can make plasmas unstable even in a very stable case like a cylinder
- Physics may be richer in strong toroidal and high pressure cases
- Thus, different options will be tested
  - **Full 2M matrix matching**
  - **E-L equations for normal b fields**
- Modified IPEC will be able to describe evolutions of outer-layer solutions and be coupled with inner-layer calculators