Overview of selected topics in SOL turbulence and blob-filament research

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Outline

- Introduction and blob theory 101
- Experimental comparisons
- Density shoulders
- Statistical properties
- Near SOL
- Other topics
- Ideas for future directions

Introduction & Motivation

- Coherent structures in the form of blob-filaments ("blobs") arise from edge/SOL turbulence: flux tubes filled with "extra" plasma.
- Filaments extended along B can deliver heat and particles to the divertor target.
- Blob-filaments propagate radially delivering heat and particles to the main chamber walls.
- Understanding the structure and characteristics of this turbulence is important. It may impact
 - recycling (divertor and main chamber)
 - deleterious PMI interactions: sputtering and erosion
 - spreading of the near SOL heat flux channel
 - impurity transport
 - divertor operation and detachment
 - RF interactions: scattering of ECH, LH waves; density at antenna, e.g. for ICRF coupling

Blob-filament theory

definition wall ulletbasic propagation mechanism **→** B • [Krasheninnikov, S. I. Phys. Lett. A 283, 36 (2001)] plasma circuit • F plasma blob η_{\parallel} Current tries to short out ⊥pol η_{sheath} polarization charges: affects blob velocity. plasma blob

Review: Krasheninnikov, D'Ippolito & Myra, J. Plasma Phys. 74, 679 (2008)

Blob velocity scaling

- Flux depends on blob velocity, size, density/pressure, frequency ullet
- Important quantities determining v_h include ۲
 - blob amplitude $\delta p/p$
 - blob scale size $\delta_{\rm b}$
 - parallel electrical connection \Leftrightarrow resistivity and magnetic geometry L_{\parallel}
- Simplest blob model: interchange driven, no parallel variation of blob ۲ pressure, sheath connected (negligible plasma resistivity and X-point effects)

 $v_b = 2c_s \frac{\rho_s^2}{\delta_b^2} \frac{L_{\parallel}}{R} \left(\frac{\delta p}{p}\right)$ balances curvature-drift current with sheath end-loss current

Disconnected inertial limit (strong parallel resistivity and/or X-point ۲ effects)

$$v_b \sim c_s \left(\underbrace{\frac{\delta_b}{R} \delta p}_{R} \right)^{1/2}$$

balances curvature-drift current with ion polarization current, i.e., inertia

... Blob velocity scalings

• Balancing curvature-drift, sheath end-loss & ion polarization current defines characteristic values:

$$v_* \sim c_s \left(\frac{\rho_s^2 L_{\parallel}}{R^3}\right)^{1/5} \qquad \qquad \delta_* \sim \rho_s \left(\frac{L_{\parallel}^2}{\rho_s R}\right)^{1/5}$$

• and dimensionless blob velocities and sizes

$$\hat{v} = v_b / v_* \qquad \qquad \hat{\delta} = \delta_b / \delta_*$$

• Dimensionless scaling

$$\hat{\mathbf{v}} = \begin{cases} \hat{\delta}^{1/2}, & \hat{\delta} << 1 & \text{inertial} \\ 1/\hat{\delta}^2, & \hat{\delta}^2 >> 1 & \text{sheath connected} \end{cases}$$

X-points and magnetic shear can disconnect filaments



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Experimental comparisons



- Scaling of v_b confirmed in basic laboratory experiments
- Validation in tokamaks more challenging

Not understood in database:

- V_{rad} distribution, amp. scaling. ...
- blob creation rate & size
- V_{pol} dependence amp.



Zweben PPCF 2016

... Experimental comparisons - NSTX

• Disconnection and smaller $V_{b,rad}$ predicted and observed near separatrix



... Experimental comparisons - TCV

C.K. Tsui, et al., PoP 2018





• comparisons with v scalings for $\hat{a} (= \delta)$ 1.6 (red) and 3.3 (blue)



For better validation and predictive capability:

- Simultaneous measurements of
 - 2D blob size and velocity (e.g. GPI)
 - internal blob plasma parameters $n_e, T_e, (T_i)$ (e.g. probes)
 - parallel structure and X-point or divertor connection
- Theoretical understanding (and numerical simulation) of
 - blob creation rate and location
 - background turbulence saturation level
 - blob size distribution
 - blob velocity distribution (blobs not isolated interact)

Many other interesting papers on blob dynamics: theoretical and experimental

- Thrysøe Blobs and **neutrals** PPCF 2016
- Hacker Blob generation rate PoP 2018
- Zhang & Krasheninnikov Blobs in framework of **nonlinear drift waves** PoP 2016
- Walkden Isolated filament motion for MAST NF 2015
- Stepanenkov Blobs in sheath-connected **arbitrary topology** PoP 2017
- Angus Blobs in **3D** PoP 2014
- Zweben NSTX 2D correlations PoP 2017
- Kube Blob amplitude and size scalings PoP 2016
- Allan Ion temperature of blobs in MAST PPCF 2016
- Birkenmeier Filament transport warm ions in AUG NF 2015
- Held **Finite Ti** blobs NF 2016
- Carralero Blobs and SOL heat transport NF 2018
- Avino X-pt effects on blobs PRL 2016
- Russell C-Mod blobs and flows PoP 2016
- Carralero High density transition of blob filaments in ASDEX NF 2014
- Fuchert ASDEX L- and H-mode GPI blobs PPCF 2014
- Kocan Intermittent SOL transport in ASDEX NF 2013
- Wiesenberger Unified **blob-hole transport scaling** laws PoP 2017
- Simon Edge turbulent transport scaling PPCF 2014

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Density shoulders in the SOL

- Density shoulder: increase of λ_n in the far SOL at high n_e
 - observed in many tokamaks; consistent with convective (not diffusive) transport
- Impacts plasma (blob-filament) interaction with main chamber walls



... Density shoulders in the SOL

- Alternative shoulder formation mechanisms:
 - ionization in main SOL [Lipschultz PPCF 2005 C-Mod]
 - divertor neutral processes, e.g. recycling [Wynn NF 2018]
- Broaden $\lambda_n \sim L_{\parallel} \Gamma_{\perp} / \Gamma_{\parallel} \Rightarrow$ increase Γ_{\perp} or reduce Γ_{\parallel}
 - CX friction may clog the flow [Militello PPCF 2016]
- Recent JET [Wynn NF 2018], TCV [Vianello NF 2017] and ASDEX-U [Carralero JNM 2017] experiments suggest that changes in Λ_{div} are not sufficient to give a shoulder. Also
 - H- and L- mode shoulders may be different
 - Divertor target configuration, fueling/seeding important
- Whether or not larger Λ_{div} causes shoulders, it is still expected to affect blob-filament propagation and main chamber PMI interaction.

Positive feedback loop: upstream-to-divertor power flow

"Thermal transport catastrophe and the tokamak edge density limit," D. A. D'Ippolito and J. R. Myra, Phys. Plasmas **13**, 062503 (2006).

- Blob velocity and transport rate increases with divertor collisionality
- 2 self-sustaining states:
 - connected, slow moving blobs with warm X-pt
 - disconnected, fast moving blobs with cold X-pt
- May be related to ASDEX-U observation of SOL width broadening at detachment [Sun PPCF 2017]

- C-Mod N₂ seeding [LaBombard PSI 2018]: no relationship between divertor conditions and upstream profiles
 - But all blob-filaments were electrically disconnected in these datasets by X-point shear



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Blob statistical properties

- Surface and atomic physics interactions are nonlinear should not use mean profiles
- Use statistical description for main chamber PMI and divertor studies
- Garcia PRL 108 265001 (2012); PoP 23 052308 (2016) and others
- Single point probe measurements of blob arrivals in far SOL described by
 - Poisson process (exponential distribution of waiting times between uncorrelated pulses)
 - prescribed "pulse" shape (the blob)
 - exponential distribution of amplitudes
- PDFs characterized by intermittency parameter

$$\gamma = \frac{\tau_d}{\tau_w} = \frac{\text{duration time} = \delta_b / v_b}{\text{avg. waiting time}}$$



- "Pulse" shape, mean amplitude, $\gamma \Rightarrow rms/mean$, normalized time correlation function, fluctuation spectrum S(ω), skewness, kurtosis ...
- Quite successful in comparisons with C-Mod and TCV data [Garcia et al.]

Application: Ensemble-averaged profiles

[D'Ippolito PoP 2002, Militello PoP 2008]

- If blobs:
 - non-interacting (e.g. $\gamma \ll 1$, far SOL)
 - velocity distribution known
 - blob shape and size distribution known
 - amplitude distribution and decay rate known (e.g. parallel losses)
 - waiting time (blob frequency) known

then \Rightarrow radial profile of n_e , Γ , ... in SOL





$$n(x) = \frac{2\pi n_{b0}}{\tau_{w}} \frac{2R}{L_{||}} \frac{\int dr_{b} r_{b}^{4} f(r_{b}) e^{-\mu x r_{b}^{2}}}{\int dr_{b} r_{b} f(r_{b})}$$

e.g. for $u_x \sim 1/r_b^2$, Gaussian shape, constant amp, decay rate ~ μ , size dist. $f(r_b)$

D'Ippolito Myra, and Krasheninnikov 2002

Application: Blob-hole dynamics



• However, GPI did not show inward "hole" motion in time-lag correlations

Motion of holes

- Inward motion expected theoretically
 - inward impurity transport
 - turbulence spreading into pedestal
- Probes [e.g. Boedo PoP 2014]
 - Skewness S < 0 (~ inside separatrix) for holes
 - S > 0 (~ outside) for blobs
 - Inward/outward V_{rad} for holes/blobs
- Recent GPI analysis (Zweben)
 - Negative S not seen
 - Inward motion of minima *on average*
 - Some *individual* minima/maxima have positive/negative V_{rad}



Average Vrad per discharge



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Near SOL, SOL width, turbulence

• Midplane SOL width λ_q from

$$\frac{\mathbf{q}_{\perp}}{\lambda_{\mathbf{q}}} = \frac{\mathbf{q}_{\parallel}}{\mathbf{L}_{\parallel}}$$

 L_{\parallel} = parallel scale length, or if SL entire field line

• Parallel flux is "classical" (e.g. collisional, sheath ...)

$$q_{\parallel} \equiv g p c_s$$

- Focus on turbulent heat (or particle, $p \rightarrow n$) flux at the separatrix
 - QL theory

$$q_{\perp} = \left\langle \widetilde{p} \, \widetilde{v}_{x} \right\rangle$$

• If pressure advection dominates

$$\widetilde{\mathbf{p}} = \frac{\mathrm{i}\widetilde{\mathbf{v}}_{\mathrm{X}}}{\omega}\nabla\mathbf{p}$$

• and e.g. for interchange-type modes $\omega = i\gamma_{mhd}$ giving a (near SOL) diffusive flux $q_{\perp} = -\frac{1}{\gamma_{mhd}} \langle \tilde{v}_x^2 \rangle \nabla p \implies \text{scaling for } \lambda_q$

Drift-interchange turbulence on NSTX

• Instead of assuming interchange turbulence, *define* v_{turb}

$$\widetilde{p} = \frac{i\widetilde{v}_x}{v_{turb}} \nabla p \qquad \qquad \frac{\widetilde{p}}{p} \to \frac{\widetilde{I}_{GPI}}{I_{GPI}} = \frac{V_{rad}}{v_{turb}\lambda_p}$$

and test relationship to drift and interchange frequencies



Myra, Russell, and Zweben, PoP 2016

- Assuming drift-interchange + additional caveats \Rightarrow scaling for $\lambda_{q, \text{ turbulence}}$
 - small but not negligible for NSTX
 - could dominate Eich and Goldston HD at large I_p and large R

Relation of QL turbulent flux to blob flux

- Turbulence Gaussian towards core \rightarrow intermittent in the far SOL
 - Skewness S increases into the SOL (many experiments)
 - Intermittency parameter $\gamma_{int} = \tau_d / \tau_w$ decreases into the SOL
 - $S = 2/\gamma_{int}^{1/2}$ for Poisson statistics, exponential pulse amplitudes and exponential pulse shape [Garcia, 2016]
- Consistent with notion that
 - Large amplitude fluctuations form outward propagating blobs
 - Blobs accelerate and hence separate into the SOL



• Number of blobs/frame decreases into the SOL; in the same region the velocity increases

Status

- Need to understand physics and scaling of
 - p_b the peak pressure in the blob, radial birth location
 - f_p the packing fraction, or equivalently (assuming δ_b/v_b is known) the blob frequency f_b or waiting time
 - blob size distribution
 - know the most stable blob size for SOL propagation
 - blob generation rate:
- Work on blob generation rate by [Fuchert PPCF 2016]
- Drift wave dynamics and blob formation [Krasheninnikov PLA 2016]; [Zhang PoP 2016]
 - relates blob amplitudes, sizes, separations and poloidal speeds

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Neoclassical drift effects and turbulence

XGC1 simulation analysis (DIII-D) [Keramidas Charidakos PoP 2018]

- Ions cross the separatrix due to banana motion; and can be subsequently lost to the divertor
- Electrons cross the separatrix due to turbulence; maintain ambipolarity
 - controls density SOL width
- Ion orbit losses set up electrostatic potential $\Phi(\theta)$ near separatrix
 - complex interaction between $\Phi(\theta)$, E_r shearing and turbulence



... XGC1 analysis: filamentary structures in SOL

Myra, Russell ... TTF 2018

 5 flux surfaces starting near the separatrix and moving out: Contours → {0.005, 0.01, 0.022, 0.05, 0.10} to show blobs (not holes)



- High field side (HFS) is populated only near the separatrix
 - Blob-filaments propagate out only on low field side (LFS)
 - Consistent with turbulent flux at separatrix
- Turbulent near separatrix; fades in SOL
- "Neoclassical" drifts dominate ion motion
 - Local quasi-neutrality \Rightarrow ion entrainment in electron filaments
 - Not axisymmetric

Divertor localized modes



Scotti NF 2018

- Observed experimentally on NSTX-U, MAST, C-Mod
- Understanding of anomalous divertor transport needed: are these modes important?
- Modeling in progress



Ideas for future work (experiment and theory)

- Basic blob physics
 - blob creation rate and location; packing fraction
 - turbulence saturation level; blob amplitude
 - blob size and velocity distributions
- Hole dynamics and impurity transport
- Drift orbits, turbulence and SOL width:
 - Coexistence and transitions
 - Scaling of near and far SOL turbulent/blob flux with I_p, R, ...
- Density shoulder formation and relation to disconnection, divertor regime
- Divertor transport: effect of fluctuations, divertor modes ...
- Neutral and atomic physics
 - divertor; instabilities and interaction with filaments
 - fueling and edge profiles, sources for simulation

Extras

Regimes for Zweben PPCF 2016 blob database



- L-mode SC near sep and disconnected in far SOL
- H-mode SC near sep and marginally disconnected in far SOL

Zweben PPCF 2016: theoretical challenges



• Velocity distribution (with negative Vrad)



Vrad and Vpol scaling with amplitude and distance from separatrix

... QL turbulent flux to blob flux

• For isolated blobs

 $q_{\perp} = p_b v_b f_p$



with peak pressure p_b , velocity v_b , packing fraction $f_p \sim \gamma$

- QL flux is $q_{\perp} = \langle \widetilde{p} \, \widetilde{v}_x \rangle$
- For blobs identify $\tilde{p} \sim p_b$, $\tilde{v}_x \sim v_b$ and at the birth region $f_p \sim 1$
- Packing fraction (or intermittency parameter) is key

 $\gamma = \frac{\tau_{d}}{\tau_{w}} = \frac{\text{duration time} = \delta_{b} / v_{b}}{\text{avg. waiting time}}$

proportional to blob generation frequency $1/\tau_{\rm w}$