Understanding and controlling the plasma kink mode response on DIII-D

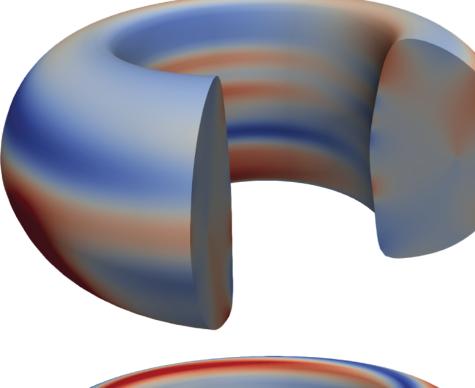
by Jeremy M. Hanson¹, A.F. Battey¹, T.C. Luce^{2,3}, G.A. Navratil¹, E.J. Strait², and F. Turco¹

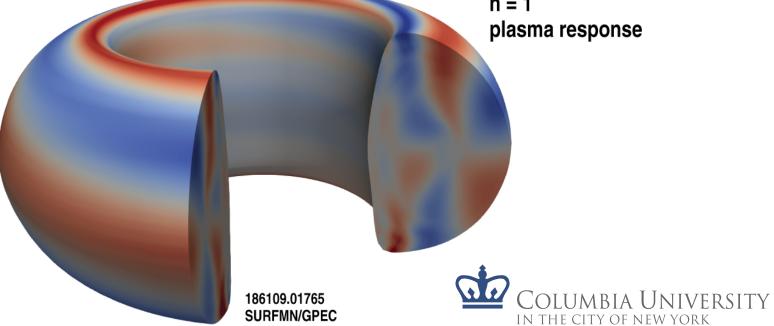
1. Columbia University 2. General Atomics 3. ITER Organization

Presented at **Princeton Plasma Physics** Laboratory

February 10, 2025









Intrinsic error field

n = 1

Stable tokamak kink response is useful for model validation and plasma control

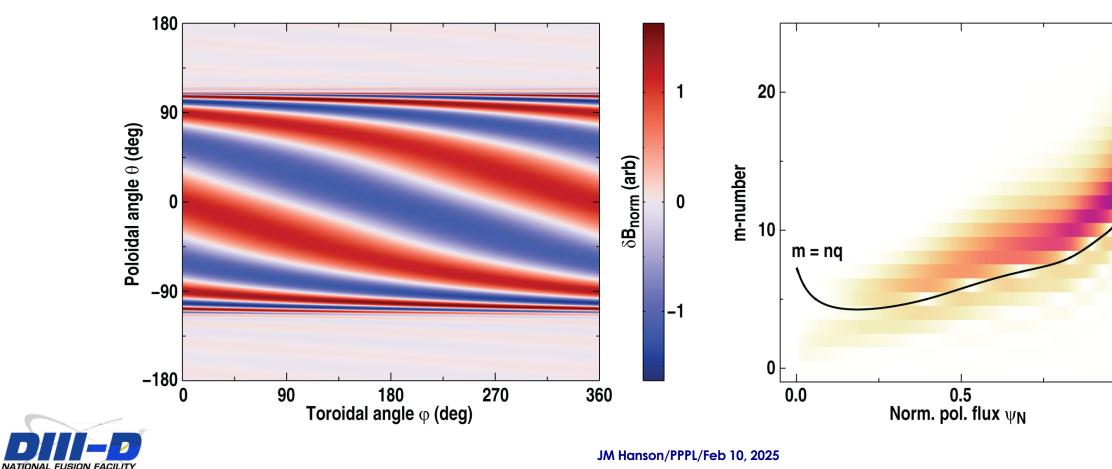
The kink mode is a long-wavelength, helical distortion with m > nq

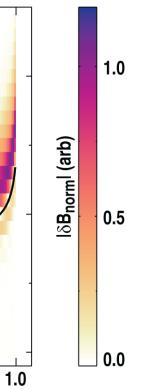
- Often stable but weakly damped \rightarrow can be driven by applied 3D fields

Ideal MHD least-stable n = 1 mode DCON, 176078, t = 2.45 s

Stability boundary sets a key pressure limit

2





Stable tokamak kink response is useful for model validation and plasma control

The kink mode is a long-wavelength, helical distortion with m > nq

- Often stable but weakly damped \rightarrow can be driven by applied 3D fields
- Stability boundary sets a key pressure limit

Stable response is an observable for simulation comparisons

- Measure using applied perturbations
- Can be predicted by MHD codes

And a lever for plasma control

- Error field correction, ELM suppression





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- Identify stabilizable regimes

– Inform control models

Outline

1. Background

- DIII-D, stable kink excitation, measurements, feedback

2. Resistive contribution to plasma response at low torque

– In the ITER baseline regime

3. Simulating feedback-controlled error field correction

- Integration of error field and MHD response codes



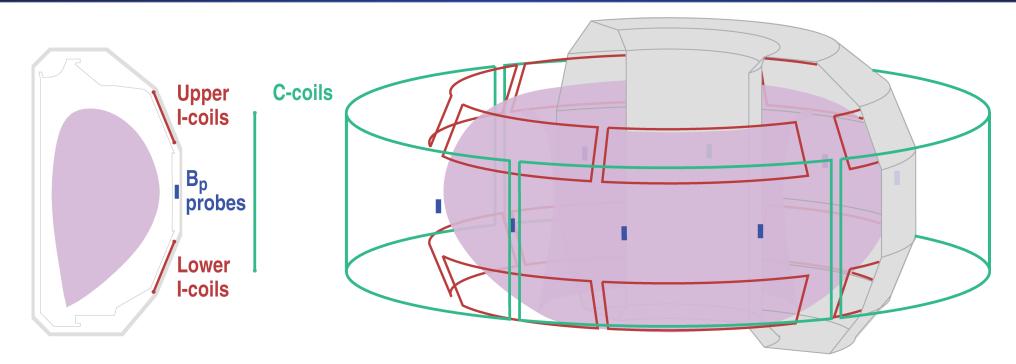
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DIII-D is well equipped for 3D response studies



- 12 internal coils (I-coils) with adjustable lower/upper phasing
- 6 external coils (C-coils)
- >100 pair-difference magnetic sensors
- Plasma rotation can be controlled with balanced NBI
- High disruption tolerance

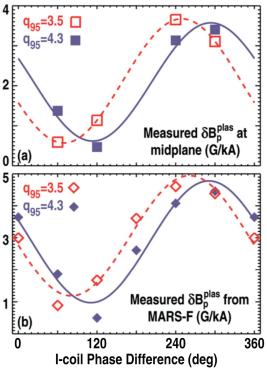


Applied 3D fields can be used to excite and control the plasma kink mode

•	Driven stable plasma response consistent with kink mode excitation predictions ¹	;
	 Varied poloidal structure of applied field 	
	 Compared two different q₉₅ values 	



MJ Lanctot, et al., Phys. Plasmas 17 (2010) 030701. 1. 2. EJ Strait, et al., Phys. Plasmas 11 (2004) 2505.



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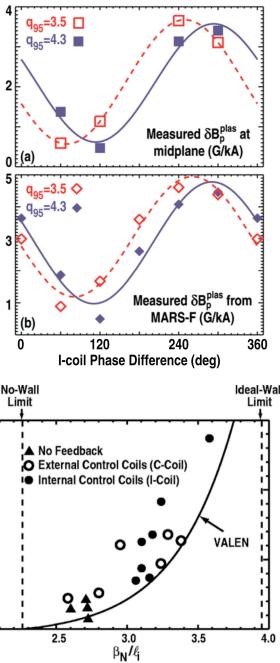
- Driven stable plasma response consistent with kink mode excitation predictions¹
 - Varied poloidal structure of applied field
 - Compared two different q_{95} values

- Unstable resistive wall mode (RWM) growth rates consistent with simulated dispersion relation²
 - Kink mode interaction with wall eddy currents = RWM
 - Feedback control enables approach to ideal wall β -limit



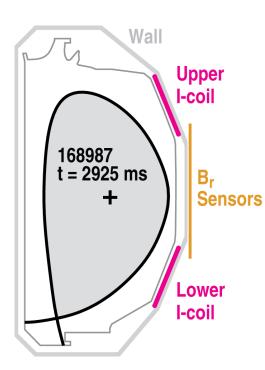
 $\gamma \tau_w$

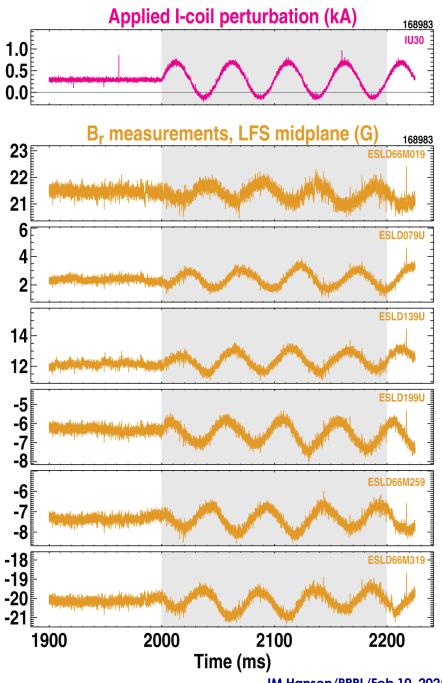
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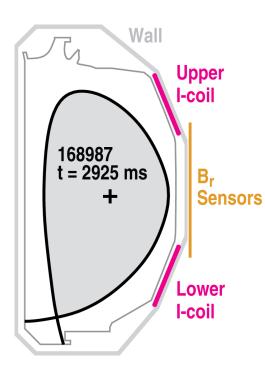
Synchronous analysis yields plasma response to rotating 3D perturbation

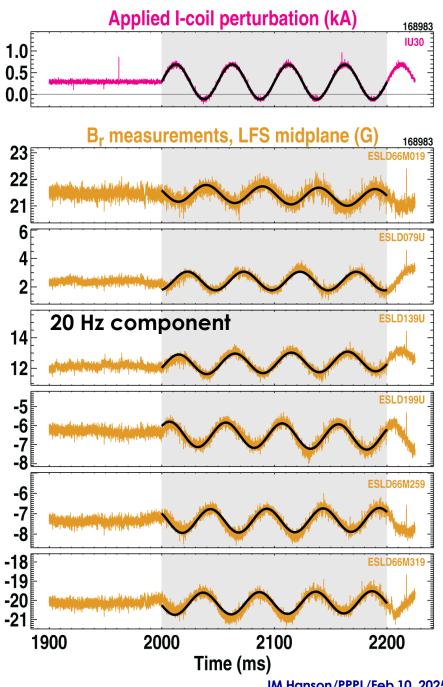






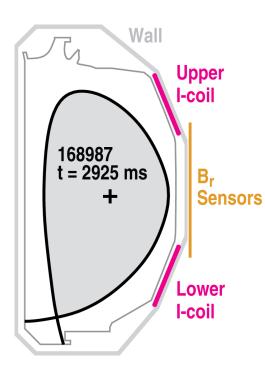
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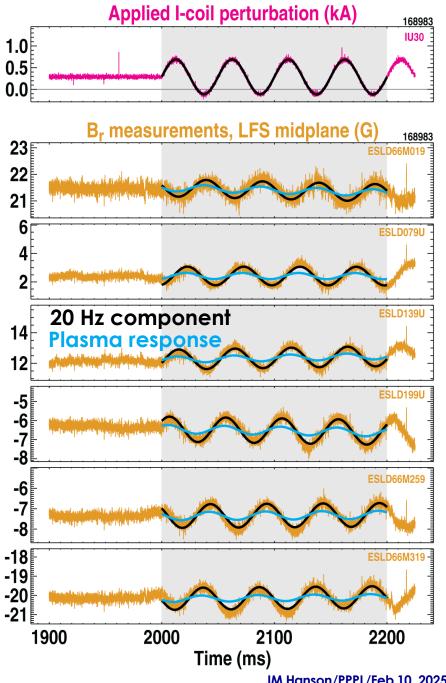






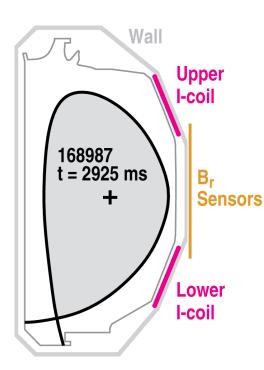
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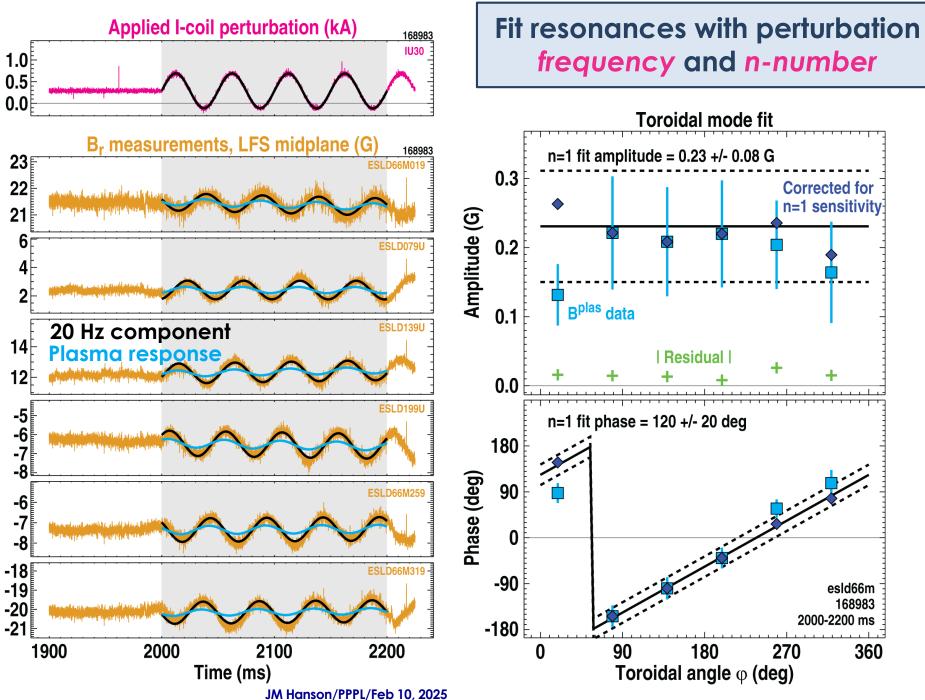




Synchronous analysis yields plasma response to rotating **3D** perturbation







Single-mode model describes tokamak n=1 plasma response and kink stability

Simple response model has a single mode^{1,2}

$$\tau_{\rm w} \frac{{\rm d}B_s}{{\rm d}t} = \gamma \tau_{\rm w} B_s + M_{sc}^* I_c$$



JM Hanson/PPPL/Feb 10, 2025

1. H. Reimerdes, et al., Phys. Rev. Lett. 93 (2004) 135002. 2. R. Fitzpatrick, Phys. Plasmas 21 (2014) 092513. 3. M. Okabayshi, et al., Phys. Plasmas 8 (2001) 2071. 4. C. Paz-Soldan, et al., Nucl. Fusion 54 (2014) 073013.



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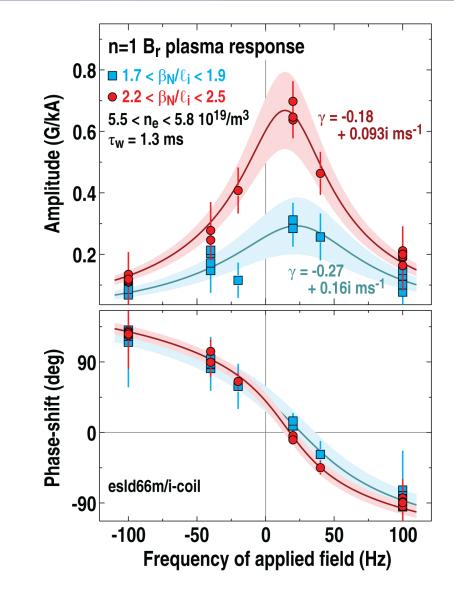
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In the Fourier domain

$$\frac{B_{s,\text{plas}}}{I_c} = \frac{B_s - B_{s,\text{vac}}}{I_c} \approx M_{sc}^* \frac{1 + \gamma \tau_w}{(i\omega\tau_w - \gamma\tau_w)(i\omega\tau_w + 1)}$$

- Model consistent with measured n=1 response
 - Peak at **resonance** with mode frequency



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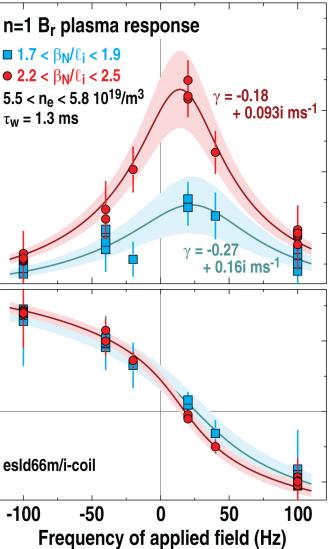
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- Model consistent with measured n=1 response
 - Peak at **resonance** with mode frequency
- Also compatible with RWM feedback and error field control experiments^{3,4}



- **0.8** | **1.7** < β_N/ℓ_i < 1.9 Amplitude (G/kA) 0.6 τ_w = 1.3 ms 0.4 Phase-shift (deg) 90 esld66m/i-coil -90 -100

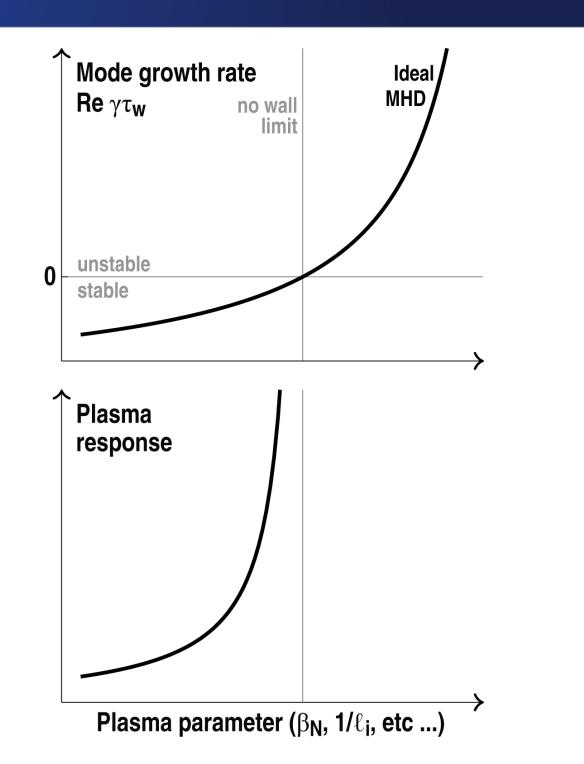


1. H. Reimerdes, et al., Phys. Rev. Lett. 93 (2004) 135002. 2. R. Fitzpatrick, Phys. Plasmas 21 (2014) 092513. 3. M. Okabayshi, et al., Phys. Plasmas 8 (2001) 2071. 4. C. Paz-Soldan, et al., Nucl. Fusion 54 (2014) 073013.

Plasma response depends on mode stability



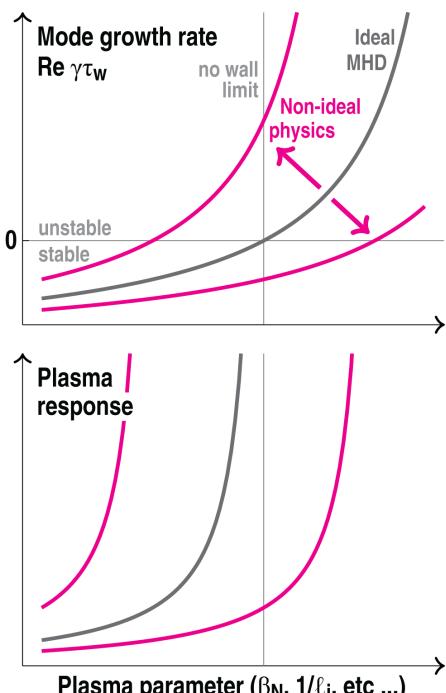
- In ideal MHD this is the no-wall limit



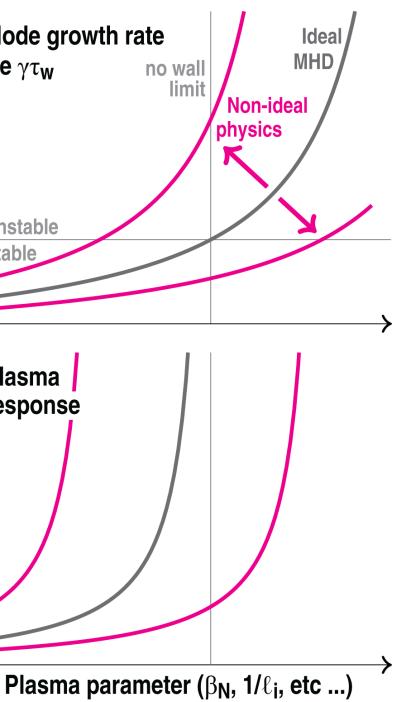


Plasma response depends on mode stability

- Increasing response indicates lacksquareapproaching stability limit
 - In ideal MHD this is the no-wall limit
- Non-ideal physics modifies limits and response
- Response is a useful tool for stability control and model validation

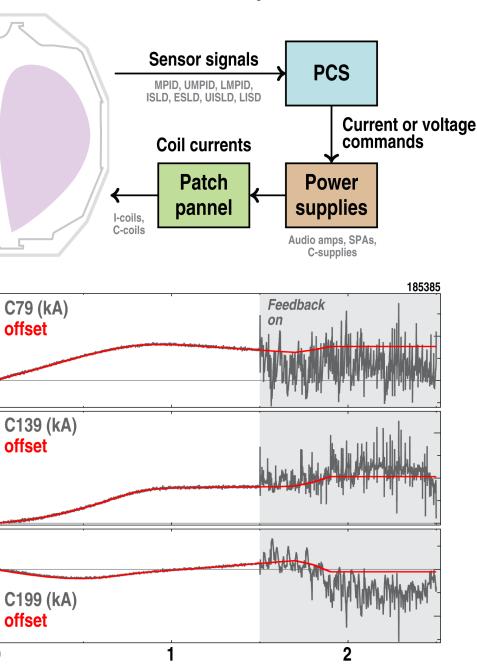


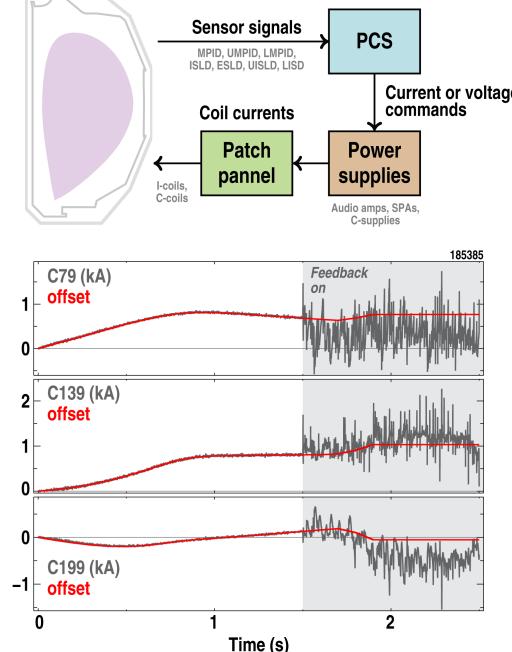




Magnetic feedback enables error field optimization

- Measure perturbed field with magnetic sensors, feed back with 3D coils
- Feedback optimizes error field correction
 - Sum with feedforward correction derived from Ohmic COMPASS scans
- Important for sensitive plasmas
 - High $\beta_N \rightarrow$ strong plasma response, possible unstable RWMs
 - Low input torque → locking likely



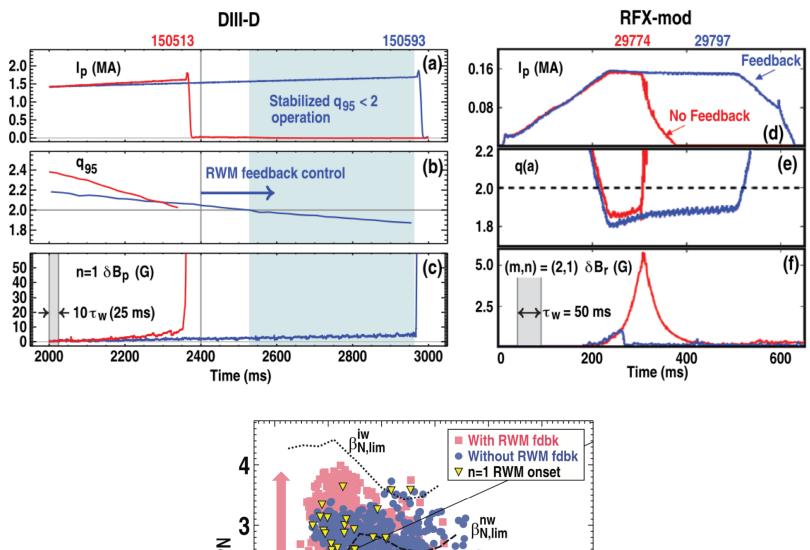




3D control system

Magnetic feedback helps expand the tokamak operating space

• To below q_{edge} = 2 - In DIII-D and RFX-MOD¹

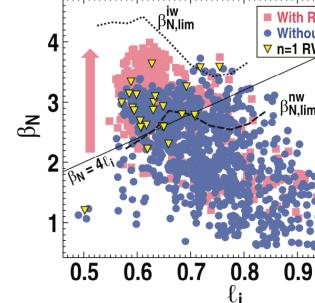


• To higher β_N

– In DIII-D high-q_{min} scenario²

1. JM Hanson, et al., Phys. Plasmas 21 (2014) 072107. 2. JM Hanson, et al., Nucl. Fusion 57 (2017) 056009.





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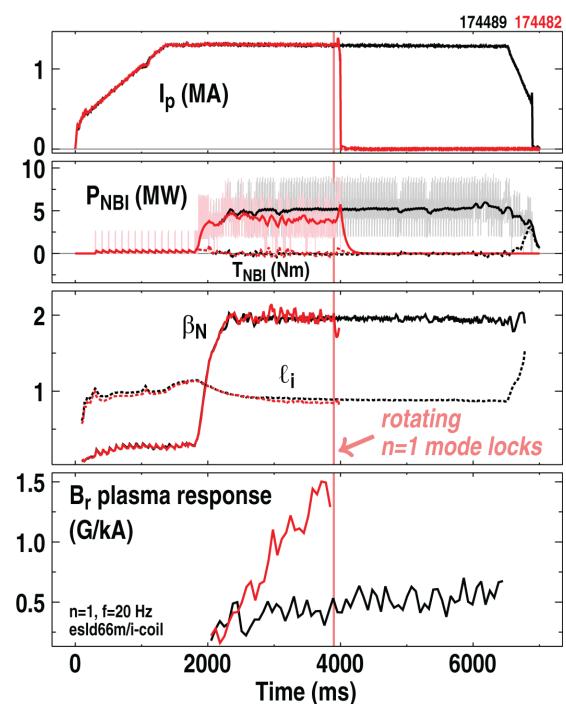
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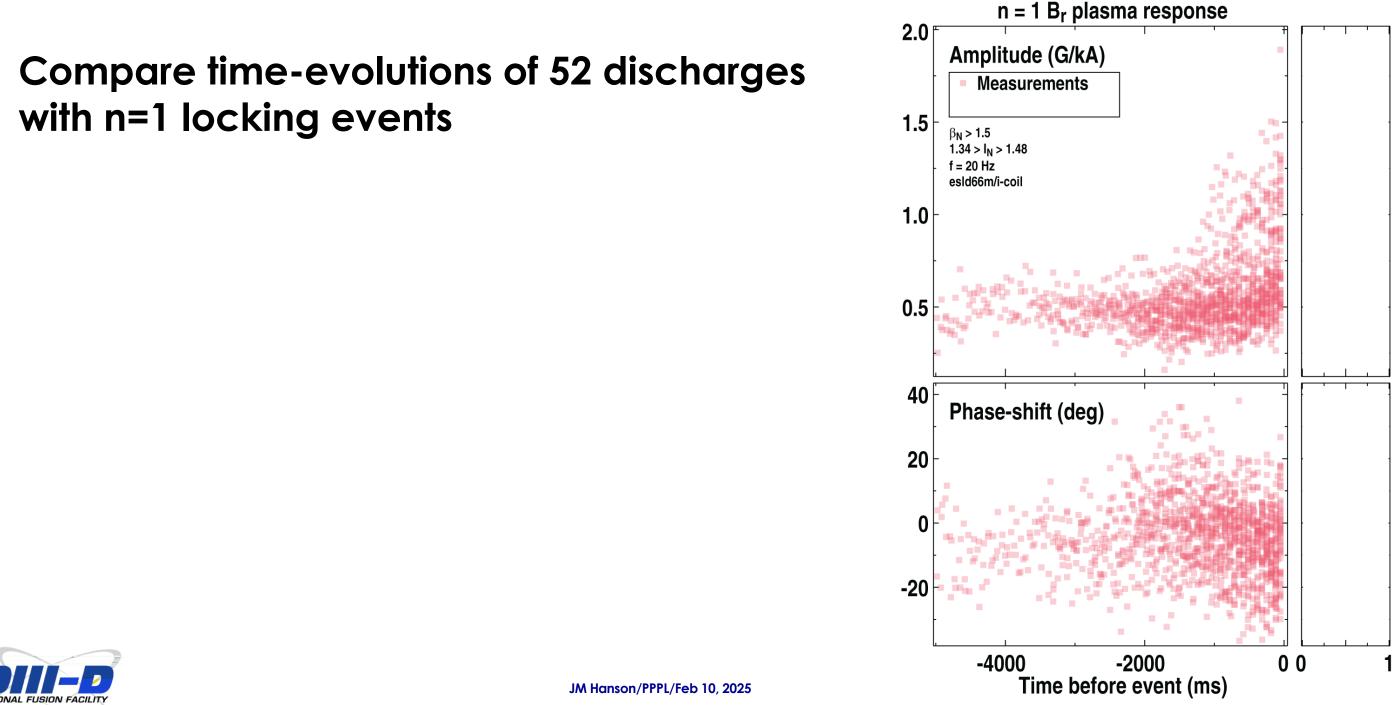
DIII-D ITER baseline demonstration discharges present stability challenge

- Combination of low torque, low q₉₅≈3 and H-mode edge leads to stability challenges¹
 - Below ideal MHD limits, but ...
 - − ... current profile shape near q=2 and 3 → tearing
- Low frequency plasma response increase prior to tearing mode locking





Many shots exhibit response increase before locking



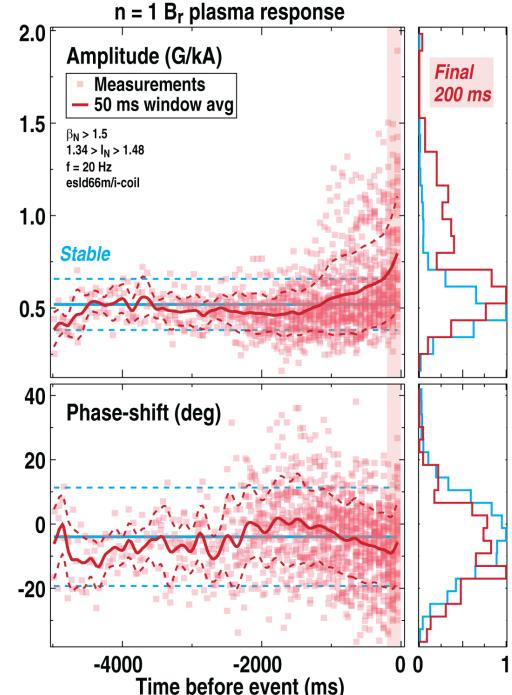




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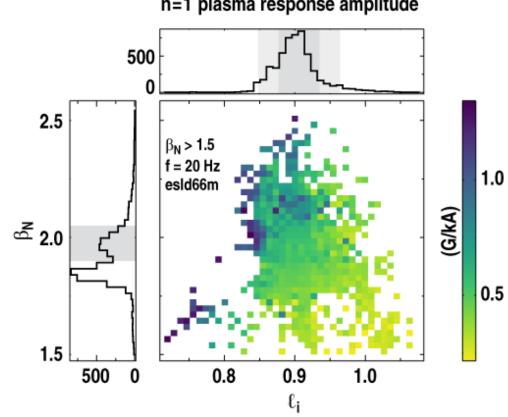
- Compare time-evolutions of 52 discharges with n=1 locking events
- Average evolution shots increasing response amplitude
 - Exceeds baseline of 60 stable shot
- Not all cases show increased response
 - $-\beta_{N}$ influences response but not TM stability in this regime¹
 - Modes sometimes lock more quickly than 100 ms response timescale







Plasma response is sensitive to MHD equilibrium



Examine large ITER demo shot dataset

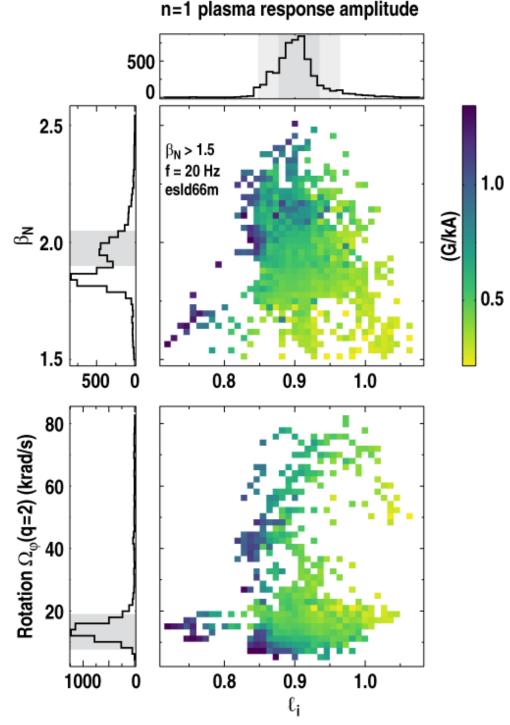
- 148 shots with $\beta_N > 1.5$, $I_N \approx 1.4$
- 5525 100 ms time-intervals
- Sensitivity to β_N and ℓ_i suggests ideal MHD link





n=1 plasma response amplitude

Plasma response is sensitive to MHD equilibrium



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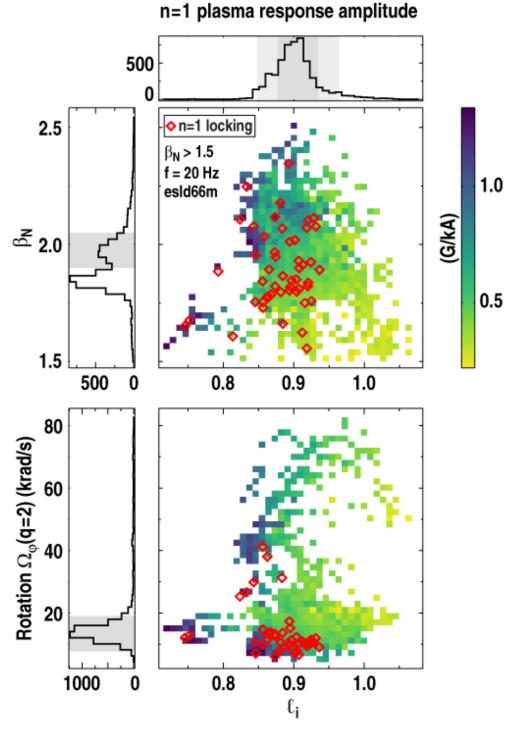
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- Strongest response at low ℓ_i and low rotation





n=1 plasma response amplitude

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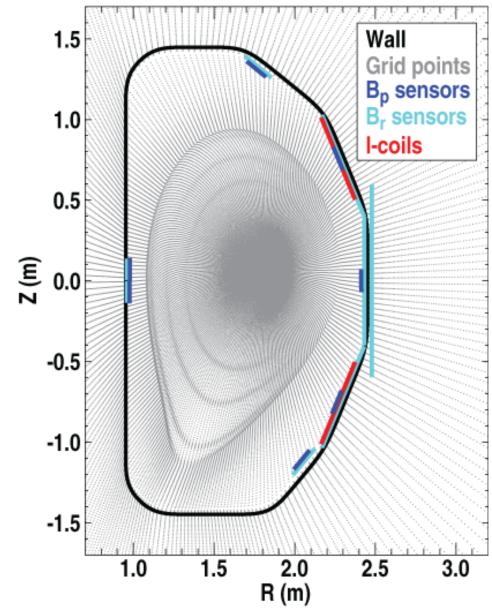
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- 5525 100 ms time-intervals
- Sensitivity to β_N and ℓ_i suggests ideal MHD link
- Strongest response at low ℓ_i and low rotation
- Higher incidence of locking in this regime
 - -52 n = 1 mode locking events
 - All followed by disruption, sometimes at reduced I_{p}



MARS-F code simulates toroidal mode plasma response

- MARS-F solves linearized, perturbed MHD model with resistive wall and coils¹
 - Solves for perturbed field over a large domain
 - Can apply rotating perturbations with coils
 - Compare with magnetic sensor measurements by averaging predictions over sensor locations
- A variety of plasma physics contributions can be included
 - Plasma rotation
 - Single-fluid resistivity
 - Kinetic contributions





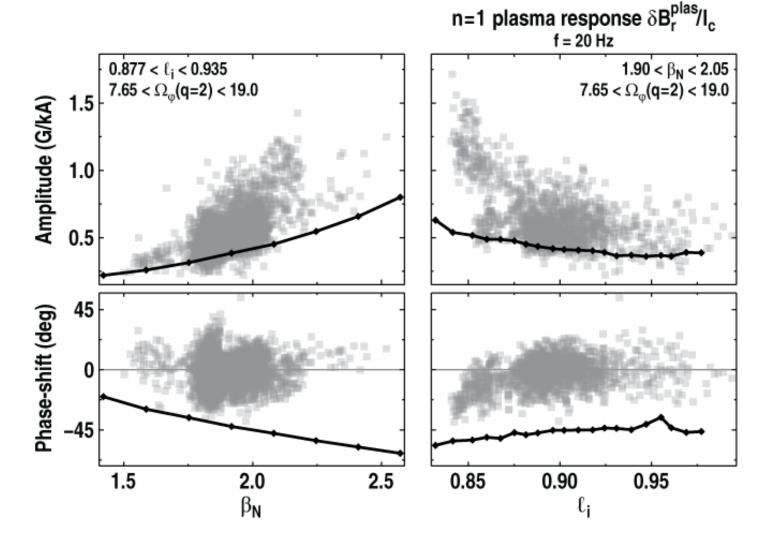


MARS DIII-D model

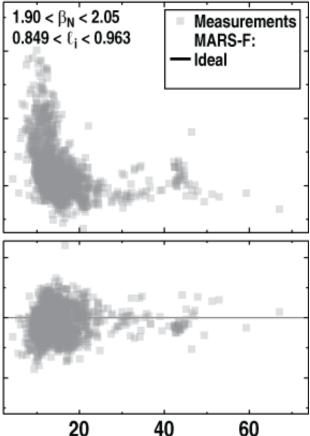
1. Y. Liu, et al., Phys. Rev. Lett. 84 (2000) 907.

Resistive response simulations compatible with experimental dependencies

 Ideal MHD qualitatively consistent with β_N and ℓ_i dependencies



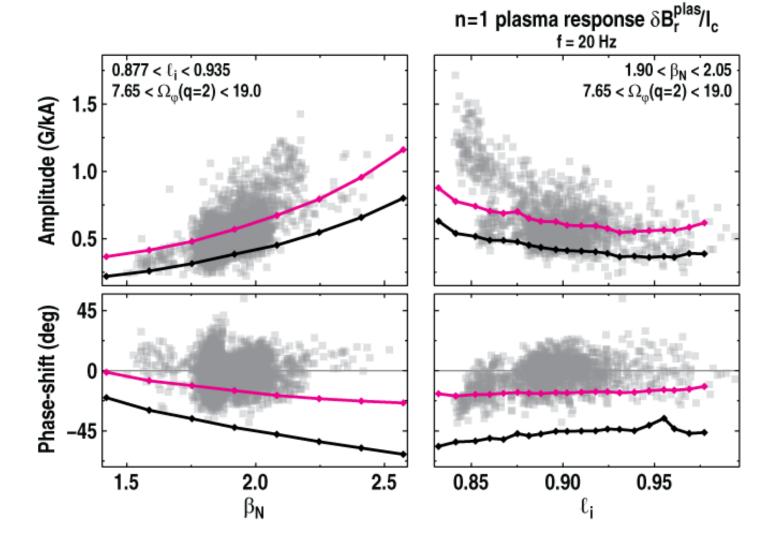




 $\Omega_{\omega}(q=2)$ (krad/s)

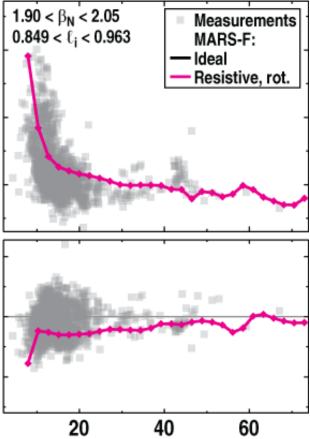
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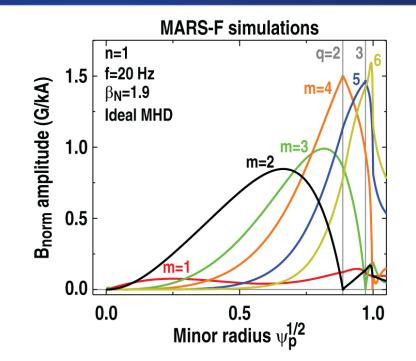
- Including resistivity and rotation improves agreement
 - Modeling consistent with large change at low rotation





 $\Omega_{\omega}(q=2)$ (krad/s)

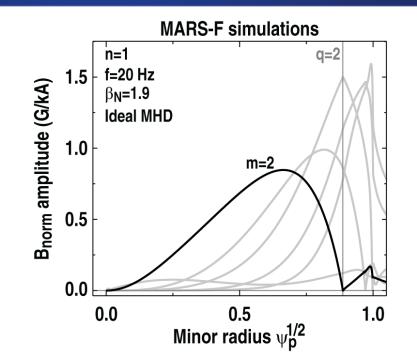
 Ideal MHD: pitch-resonant fields screened at rational surfaces







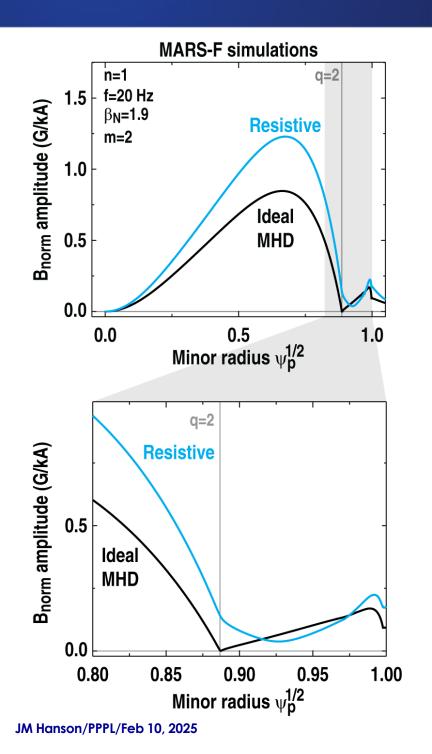
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- Ideal MHD: pitch-resonant fields screened at rational surfaces
 - Most visible at q=2
- Screening currents decay in the presence of resistivity
 - Permitting non-zero resonant components

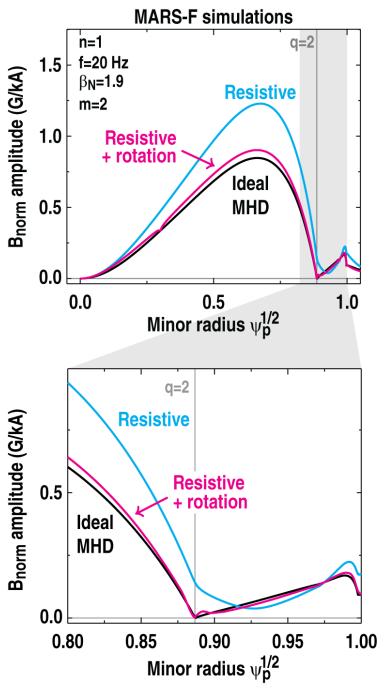






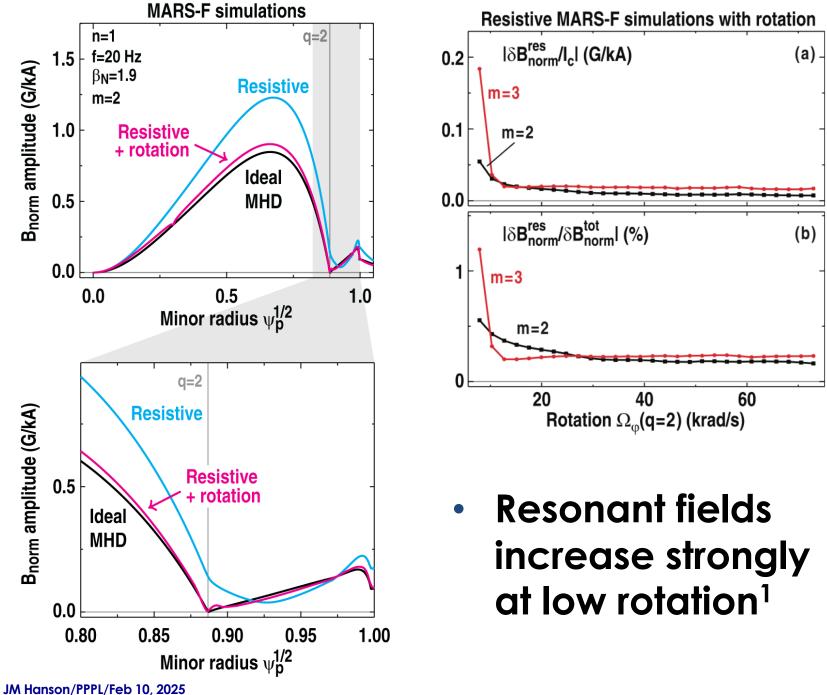
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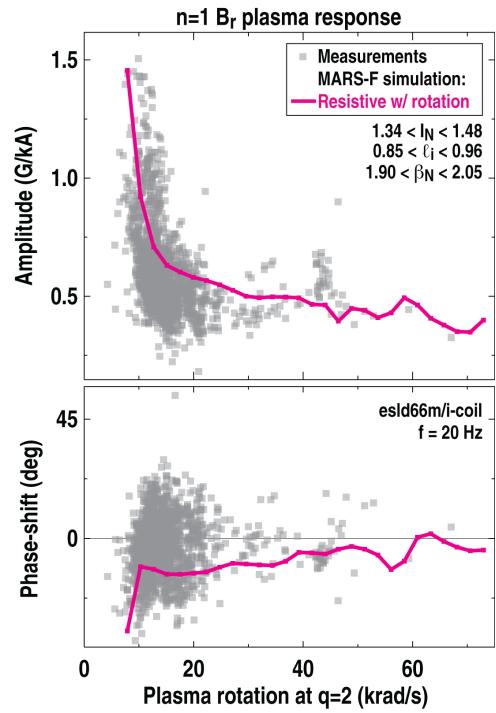




1. JM Hanson, et al., Phys. Plasmas 28 (2021) 042502.

What have we learned about the transition to locking?

- Is the applied perturbation driving tearing? lacksquare
 - Simulations show resonant field increase ...
 - \dots but q = 2 surface rotates faster than perturbation: ≈1 kHz vs 20 Hz
 - This regime has (2,1) stability and locking issues even without perturbations¹





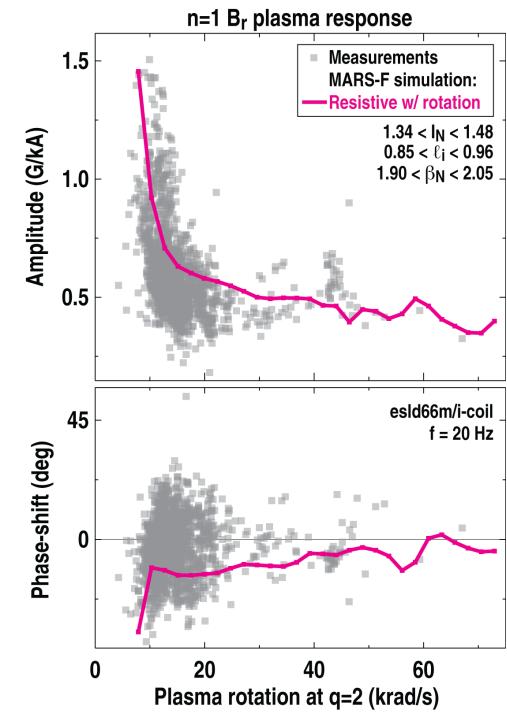


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Kink stability weakens at low rotation

- Below no-wall limit, so don't expect instability
- But increasing error field amplification likely ____ contributes to locking and disruption



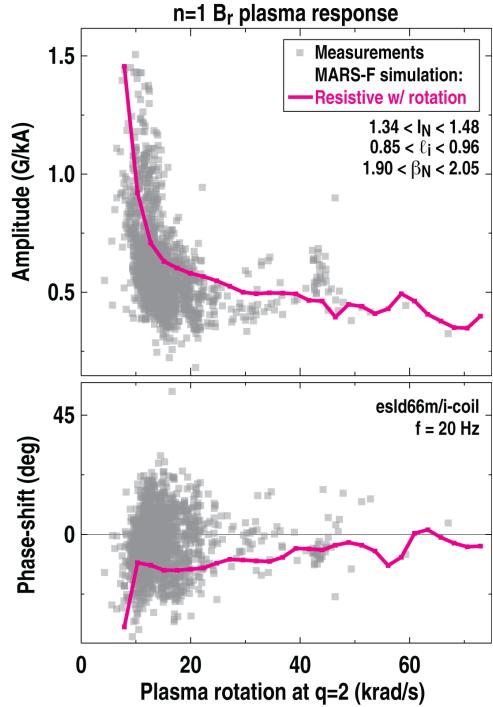




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 - Below no-wall limit, so don't expect instability
 - But increasing error field amplification likely contributes to locking and disruption
- **Ultimate solution:** current profile optimization for (2,1) tearing stability²







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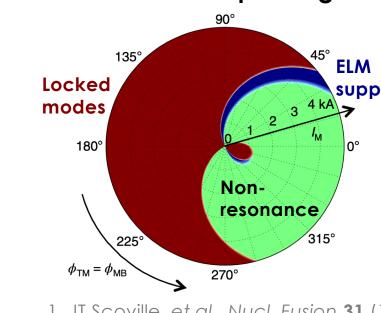


Correcting plasma error field response improves tokamak performance

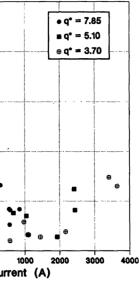
- Small equilibrium coil deviations from axisymmetry lead to field errors
 - Implicated in MHD instability onsets and locking
 - Observed in many tokamaks
 - Critical threshold $\delta B/B_0 \sim 10^{-4}$
- Low density locked mode threshold improved with applied n=1 field¹
 - Attributed to error field correction
- Now understand that managing 3D field can have many benefits
 - Rotation optimization², ELM suppression³ ...



DIII-D Low density locked mode threshold¹ • q* = 7.85 ∎ q* = 5.10 ⊕ q* = 3.70 (10¹³ cm⁻³) 2.5 Density -3000 -2000 3000 n=1 Coil Current (A)





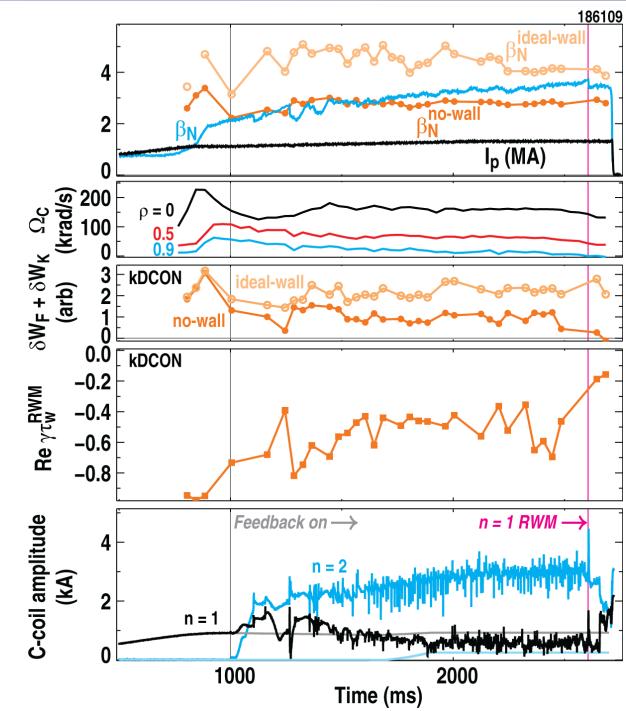


KSTAR simulated 3D operating window³

1. JT Scoville, et al., Nucl. Fusion 31 (1991) 875. 2. AM Garofalo, et al., Nucl. Fusion 42 (2002)1335. 3. J-K Park, et al., Nature Phys. 14 (2018) 1223.

Advanced tokamak discharge reaches stability limit with n=1 + n=2 feedback control

- β_N ramped to 3.8, above ideal n=1 no-wall limit¹
- Born-locked n=1 mode before disruption
 - Growth time: 1 ms $\approx \tau_w$
 - Coincides with approach to kinetic MHD marginal point





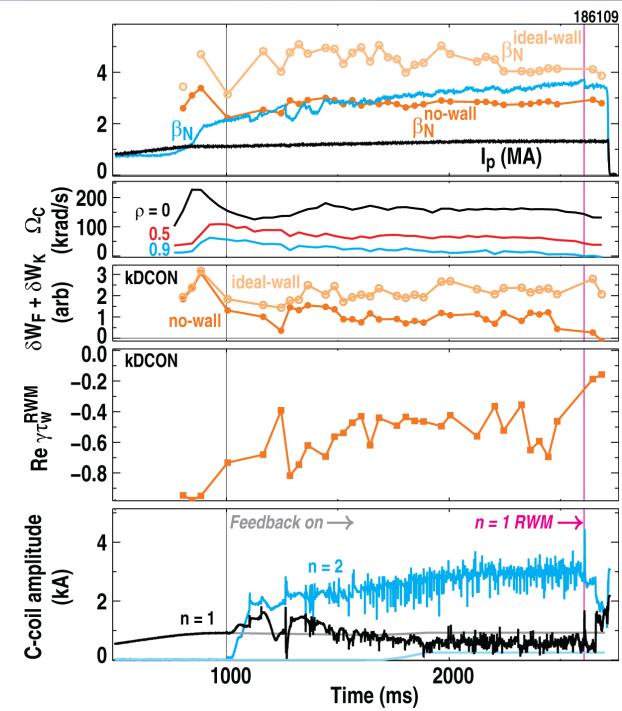


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n=1 + n=2 C-coil feedback

- Fast RWM control and EF correction
- Baseline evolution indicates EF response





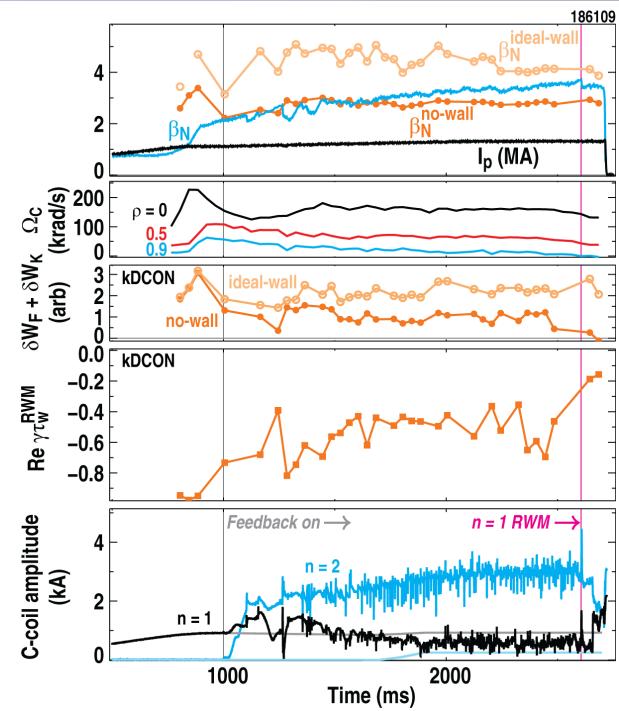


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Can we understand what the feedback did?



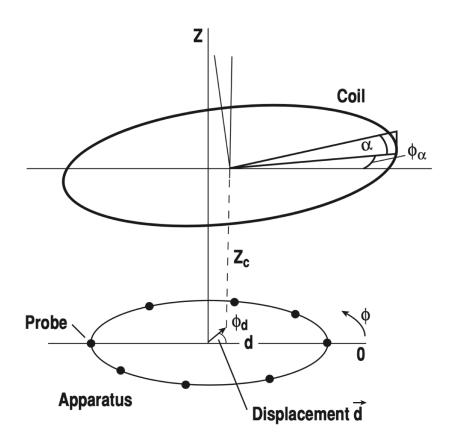




Recent developments enable error field correction simulations

Perturbed coil and measurement apparatus¹

- DIII-D n=1 and n=2 EF source model is well established¹
 - Based on in-vessel coil asymmetry measurements
- Can simulate plasma response with DCON²/GPEC³
 - Critical to include: *plasma selects EF harmonics*
 - Ideal MHD + kinetic model allows stable response prediction above no-wall limit



- Idea: apply codes to simulate DIII-D high β feedback
 - Multi-harmonic plasma response "collapsed" into single-mode model



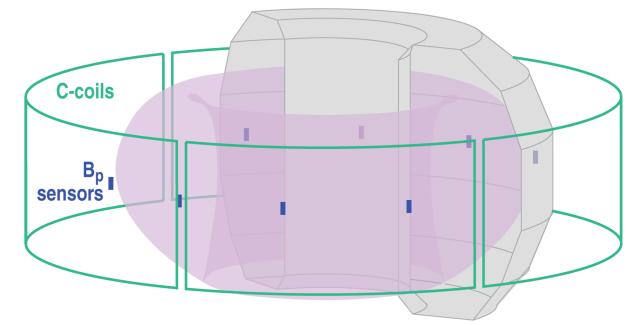
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- JL Luxon, et al., Nucl. Fusion 43 (2003) 1813.
- 2. AH Glasser, Phys. Plasmas 23 (2016) 072505.

3. J-K Park and NC Logan, Phys. Plasmas 24 (2017) 032505.

Simplifying assumptions lead to tractable model for feedback error field correction

- 1. Single toroidal arrays of coils and sensors, single-n
 - Represent toroidal modes with complex scalars ${\rm B}_{\rm s}$ or ${\rm I}_{\rm c}$
 - Consistent with toroidal mode fitting in feedback algorithm



2. Fields are small enough that plasma response is linear

- Perturbed fields proportional to applied coil currents $B_s = A_{sc}I_c$
- Not always true: EF can change plasma rotation, changing A_{sc}

3. Plasma response is fast relative to equilibrium time scale

- − Response time $\tau_w \approx 2.5$ ms << $\tau_E \approx 100$ ms
- Leads to time-independent model



Model equations are straightforward¹

 Sensor mode has vacuum and plasma response contributions, from intrinsic EF and coils

$$B_{s} = B_{s}^{\text{EF,v}} + B_{s}^{\text{EF,p}} + (A_{sc}^{v} + A_{sc}^{p})I_{c}$$
SURFMN GPEC

Proportional gain feedback law as in real-time algorithm

$$I_c = G_{cs}B_s$$

- Feedforward commands, sensor baselining, and sensor vacuum compensation also included
 - But omitted here for simplicity



1. H. Reimerdes, et al., Fusion Sci. Tech. 59 (2011) 572.

Solving yields predicted coil currents and sensor fields

Closed loop solution

$$I_{c} = \frac{G_{cs}}{1 - A_{sc}G_{cs}}B_{s}^{\text{EF}}, \text{ with } B_{s}^{\text{EF}} = B_{s}^{\text{EF},v} + B_{s}^{\text{EF},p} \text{ and } A_{sc} = A_{s}^{v}$$
$$B_{s} = \frac{1}{1 - A_{sc}G_{cs}}B_{s}^{\text{EF}}$$

Infinite gain limit yields "optimal" current for nulling EF sensor field

$$I_c^{\text{opt}} = -\frac{1}{A_{sc}} B_s^{\text{EF}}$$



$\frac{v}{sc} + A_{sc}^{p}$

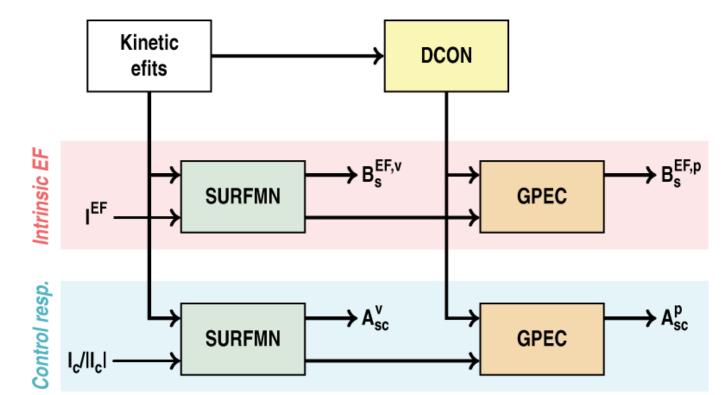
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- 2.

Modular simulation codes enable flexible workflow

- **SURFMN:** 3D vacuum fields, given coil currents¹
- **DCON:** MHD mode spectrum and stability²
- **GPEC:** MHD spectrum response to 3D fields³
- **Exploit model linearity**
 - Fields from different code runs can be summed
 - Extract amplification terms A_{sc} with unit amplitude coil currents
 - Sensor fields from integrals over sensor areas



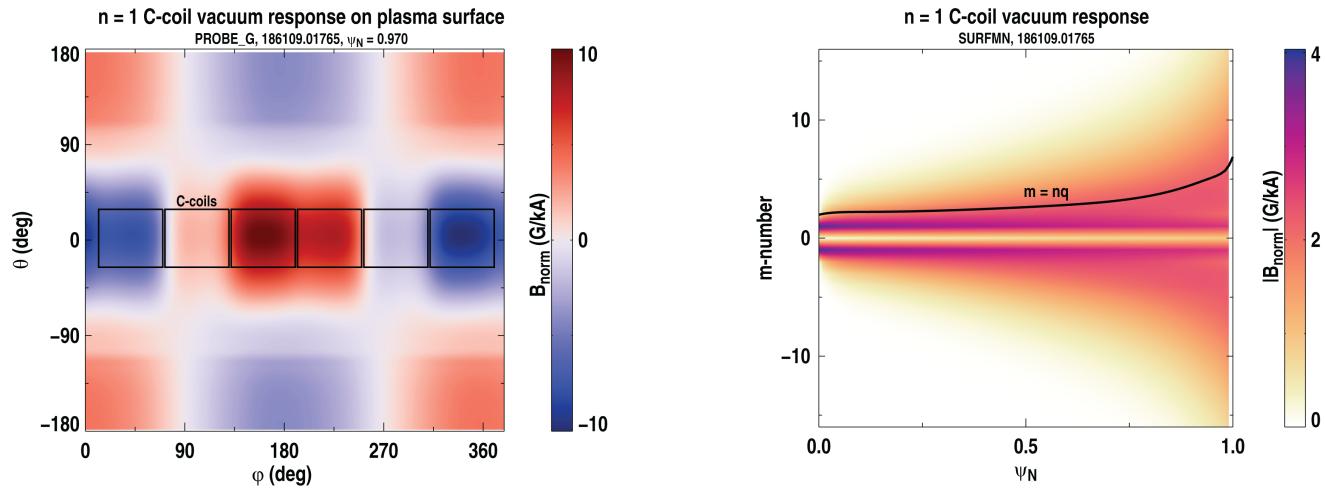




Simulation workflow for feedback error field correction terms

1. MJ Schaffer, et al., Nucl. Fusion 48 (2008) 024004. AH Glasser, Phys. Plasmas 23 (2016) 072505. 3. J-K Park and NC Logan, Phys. Plasmas 24 (2017) 032505.

C-coil n=1 vacuum field has a broad m-spectrum

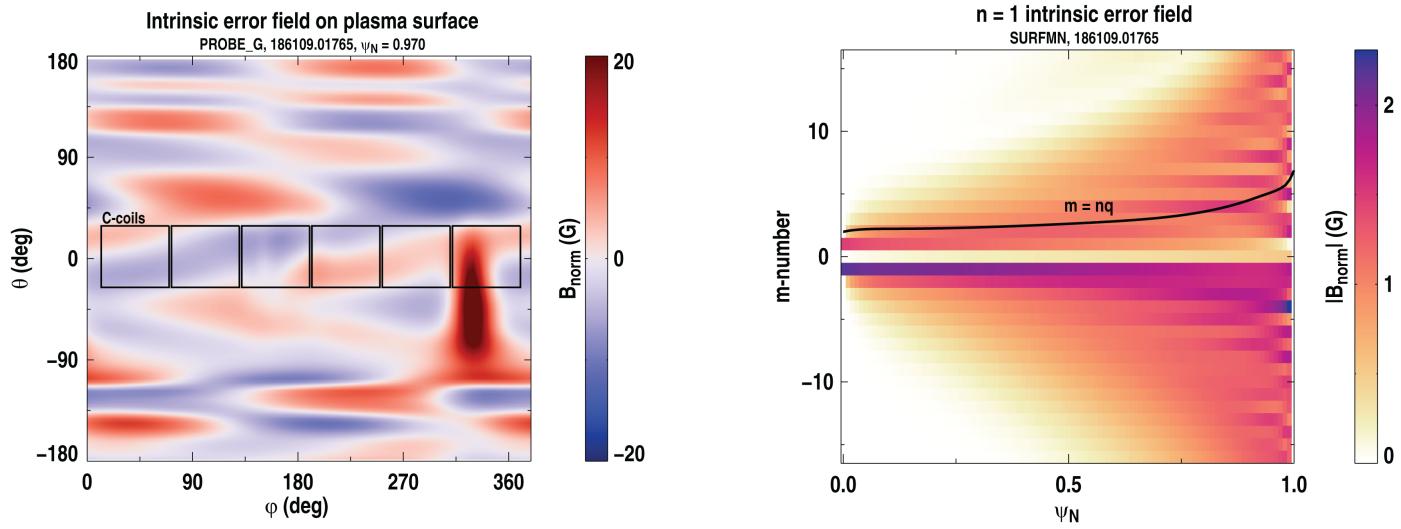


SURFMN predicts n=1 C-coil perturbation can couple to pitch-resonant and nonresonant harmonics





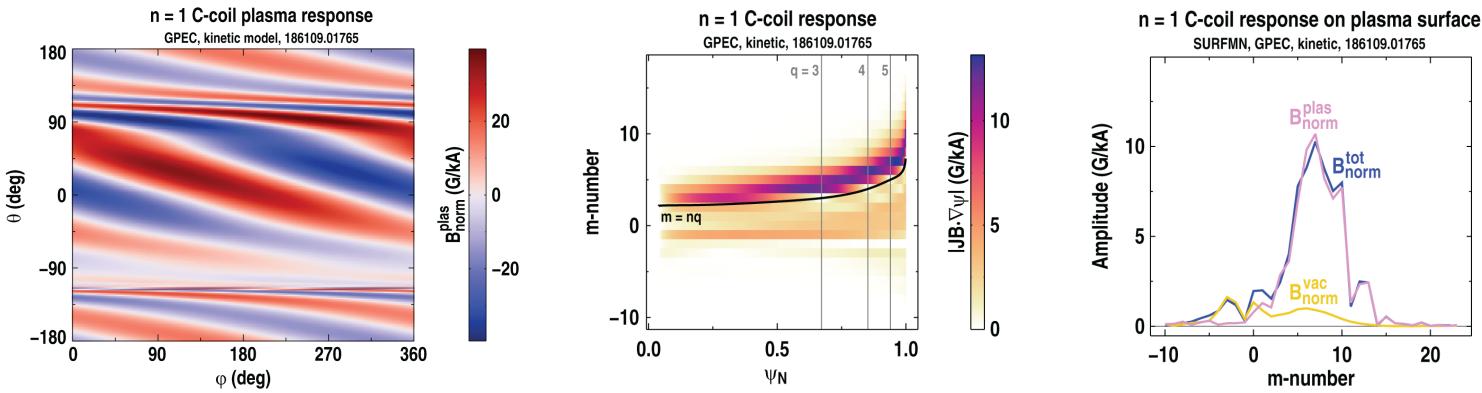
Intrinsic error field spectrum is also broad



- SURFMN predicts broad intrinsic error field spectrum
 - Anti-resonant peak in right-handed plasma



Kink-like plasma response to C-coil field predicted

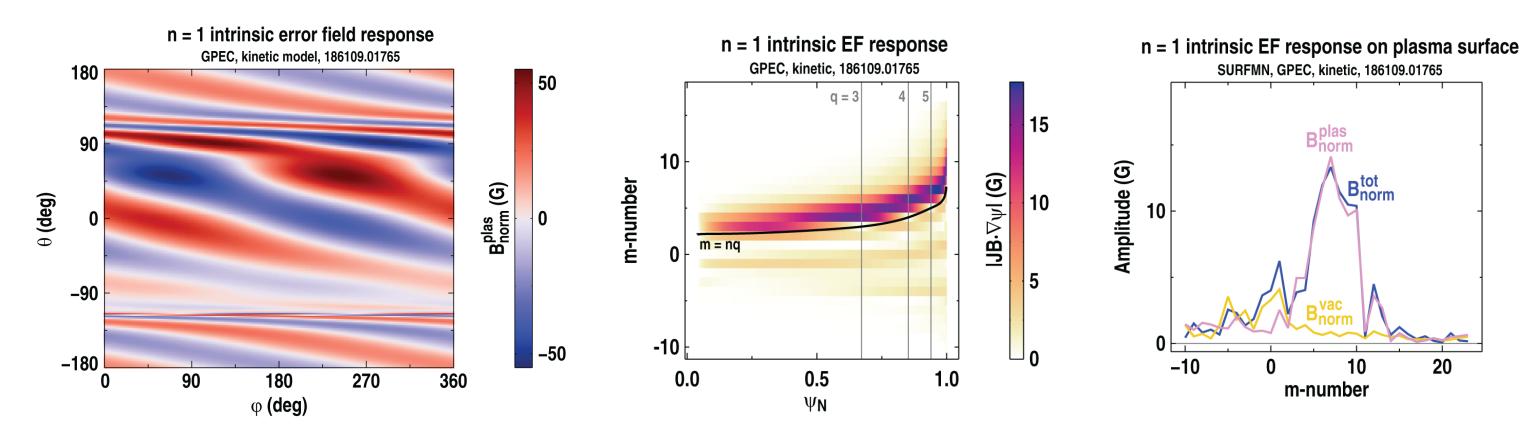


Plasma response dominates over vacuum field ullet





Intrinsic error field response prediction is also kink-like

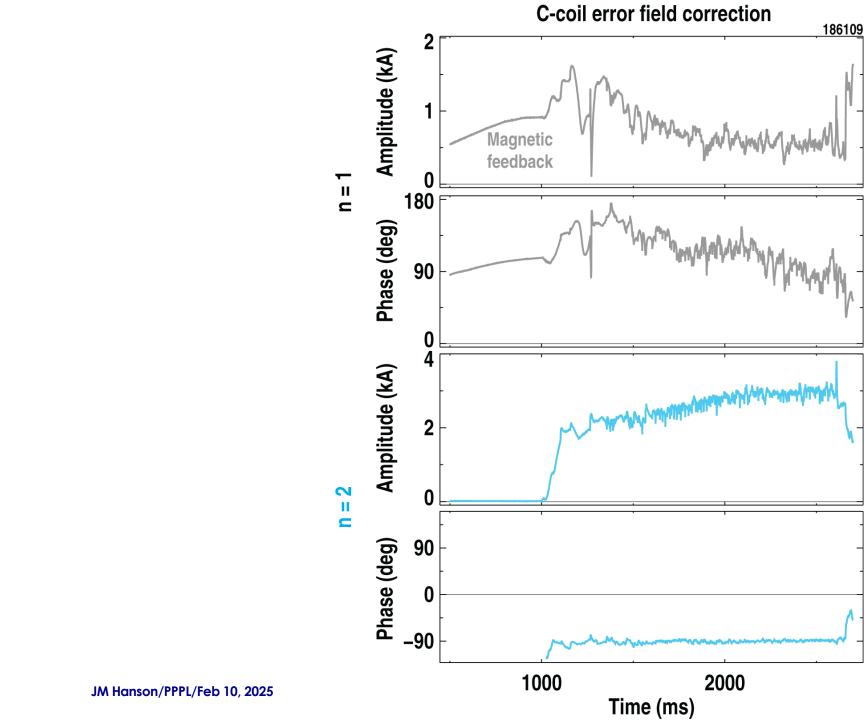


 Coils that couple to the plasma kink mode should be excellent for error field control





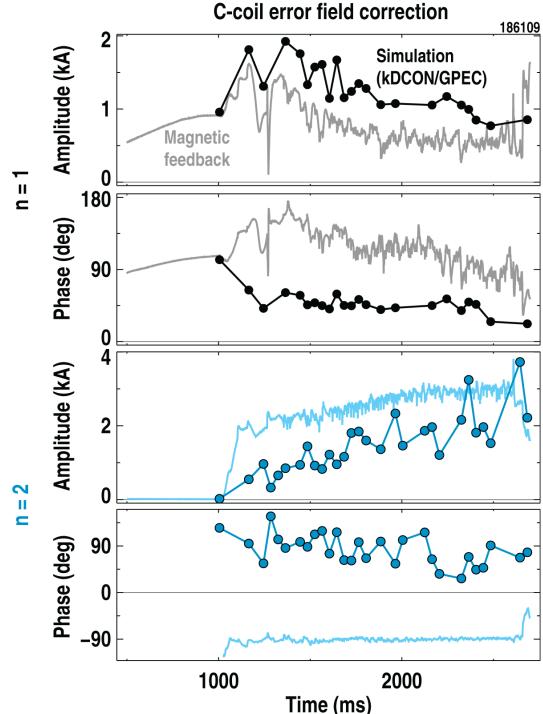
Can simulations predict feedback baseline evolution?





Can simulations predict feedback baseline evolution?

- Predict n=1 amplitude to within 50%, phase to within 90°
- n=2 predictions are more scattered
 - Large phase disagreement

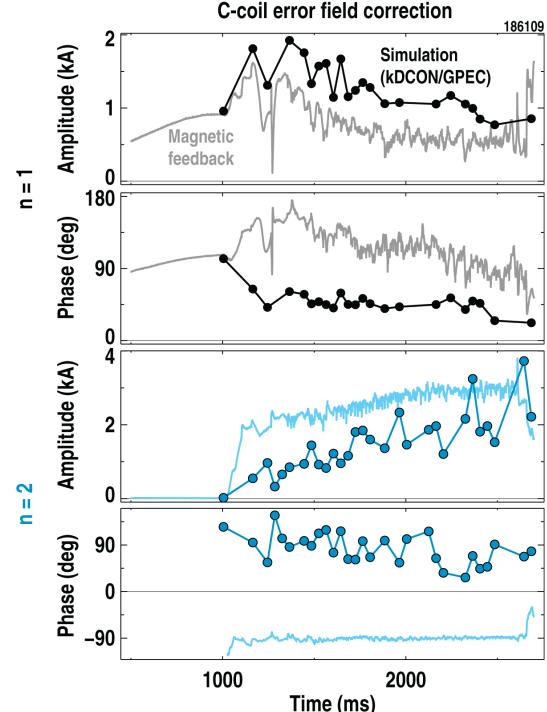




Can simulations predict feedback baseline evolution?

- Predict n=1 amplitude to within 50%, phase to within 90°
- n=2 predictions are more scattered
 - Large phase disagreement
- Several possible explanations for discrepancies
 - Sub-optimal feedback? shot did exhibit rotation braking and instability
 - EF source model inaccuracies?
 - Plasma response calculation? 180 phase shift would be surprising





Conclusions

- Kink mode response is a key observable for validating simulations and lever for plasma control
 - Easily driven with applied 3D fields
- **Resistive response linked to mode locking in low-torque regime**
 - Measurements consistent with resistive MHD simulations
 - Simulations show weakening shielding at rational surfaces as rotation slows
- New simulations facilitate error field correction predictions
 - Link error field source model with plasma response simulation
 - Validation effort ongoing





Extra slides

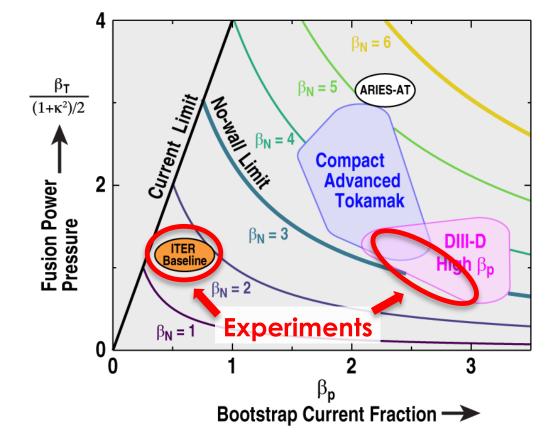


DIII-D can explore different operating regimes of interest for fusion

- High current scenarios have good confinement, but high current drive need
 - Example: ITER baseline scenario (IBS)¹
 - Peaked current profile → high no-wall limit
- Increasing β_p is associated with higher bootstrap and non-inductive fractions
 - Higher degree of profile self-organization²
 - Broad current profile \rightarrow high with-wall limit³
 - Advanced tokamak candidate for compact fusion pilot⁴
- Stability challenges differ, will show results from different regimes



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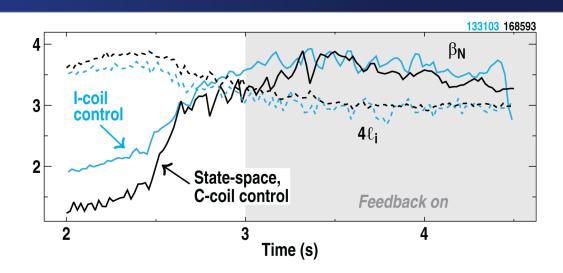
1. EJ Doyle, et al., Nucl. Fusion **50** (2010) 075005. 2. PA Politzer, et al., Nucl. Fusion 48 (2008) 075001. 3. JM Hanson, et al., Nucl. Fusion 57 (2017) 056009. 4. RJ Buttery, et al. Nucl. Fusion 61 (2021) 046028.

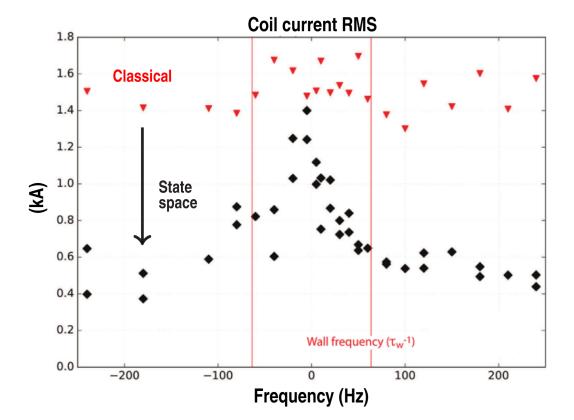
Simulations inform control strategies

- Evaluated state-space control approach¹
 - Incorporating reduced-order VALEN model
- Accessed $\beta > \beta^{no-wall}$ using external coils
- Led to reduced power requirement
 - Compared with proportional gain



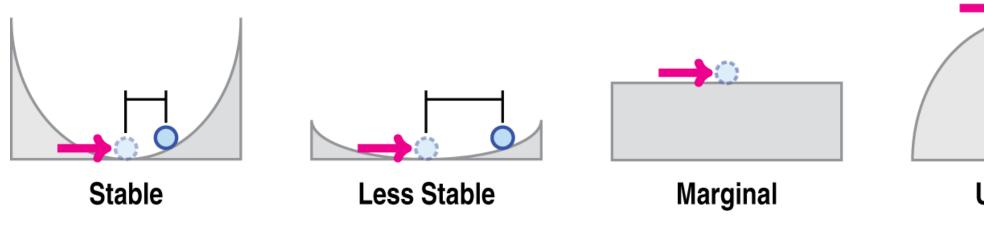






Perturbative experiments can help assess stability

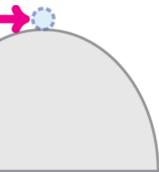
Consider a family of mechanical systems



- Apply a small perturbation
- If the system is stable, we can **measure a finite response**
- **Response contains information** about the proximity to marginal

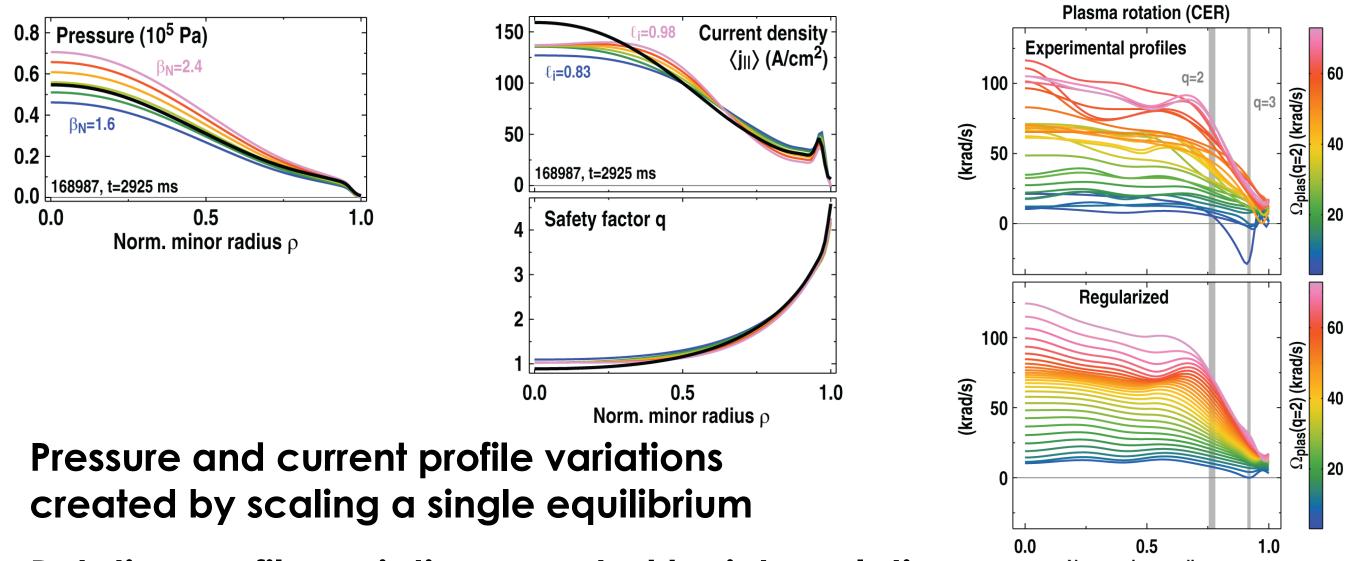






Unstable

Profile variations allow investigations of β_N , ℓ_i , and rotation dependencies



- Rotation profile variations created by interpolating experimental profiles



Norm. minor radius p