

Electron Transport in NSTX and ETG

J.-H. Kim, W. Horton

Institute for Fusion Studies, University of Texas at Austin

S.M. Kaye, R. Bell, L.P. LeBlanc

Princeton Plasma Physics Laboratory

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Overview

- Heat fluxes and diffusivities of the Electron Temperature Gradient(ETG) mode are compared with NSTX TRANSP data.
- GS2 Linear Analysis for the wide range of the poloidal wavenumber k_y .
- Linear Analysis based on 3-field fluid model
- Test simulations of the fluid model and Heat flux estimation
- Conclusion and future direction

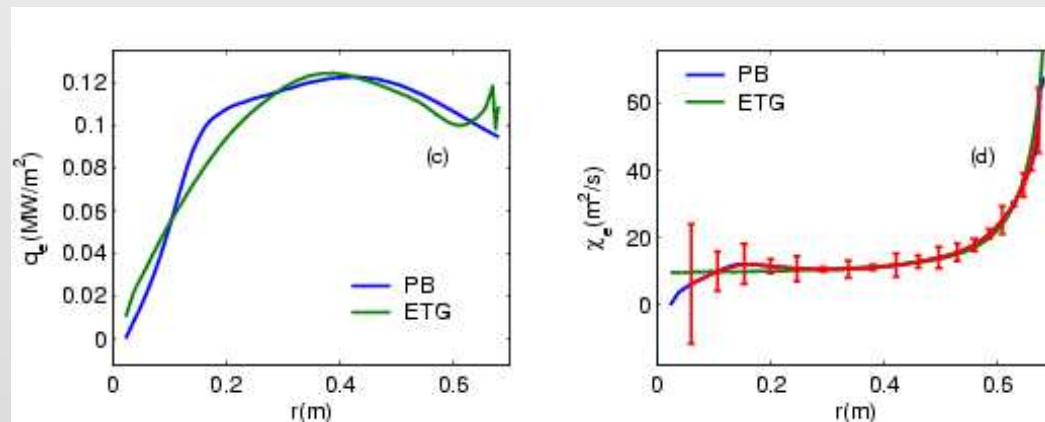
Electron Temperature Gradient(ETG)

For ETG mode,

$$q_e = C_e^{em} n_e T_e q \frac{c^2 v_e}{\omega_{pe}^2 R} \left(\frac{R}{L_{Te}} - \frac{R}{L_c} \right) \quad \beta_e > \beta_{e,cr} \quad (1)$$

$$q_e = C_e^{es} n_e T_e q^2 \left(\frac{\rho_e^2 v_e}{L_{Te}^2} \right) \left(\frac{R}{L_{Te}} - \frac{R}{L_c} \right) \quad \beta_e < \beta_{e,cr} \quad (2)$$

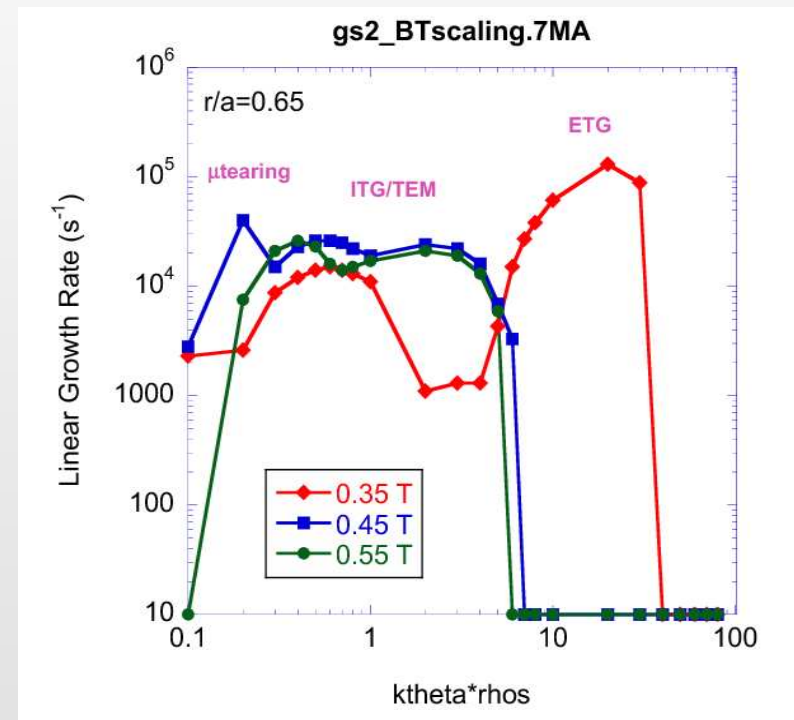
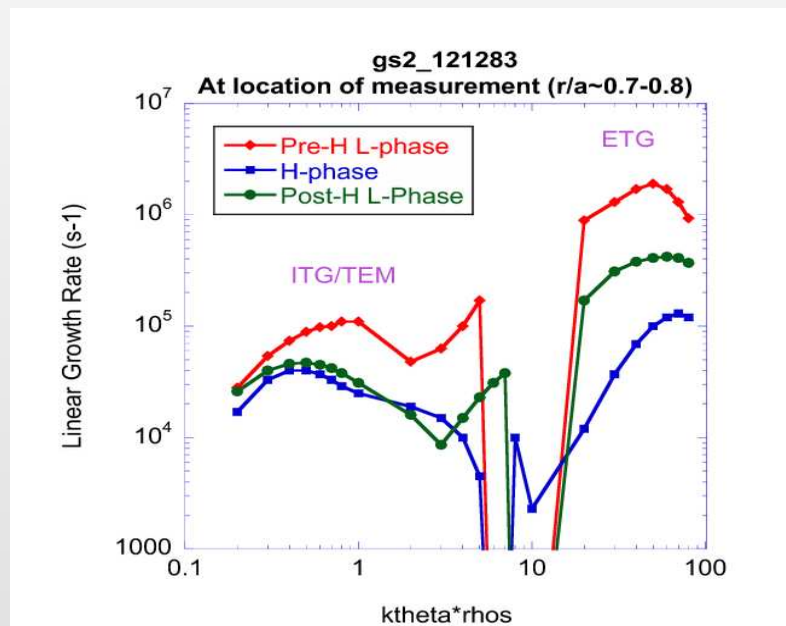
where $\beta_{e,cr} = (L_{Te}/qR)^2$ NSTX 106194 with 2.4 MW HHFW power



From Horton et. al. Nuclear Fusion 2005

GS2 Linear Stability

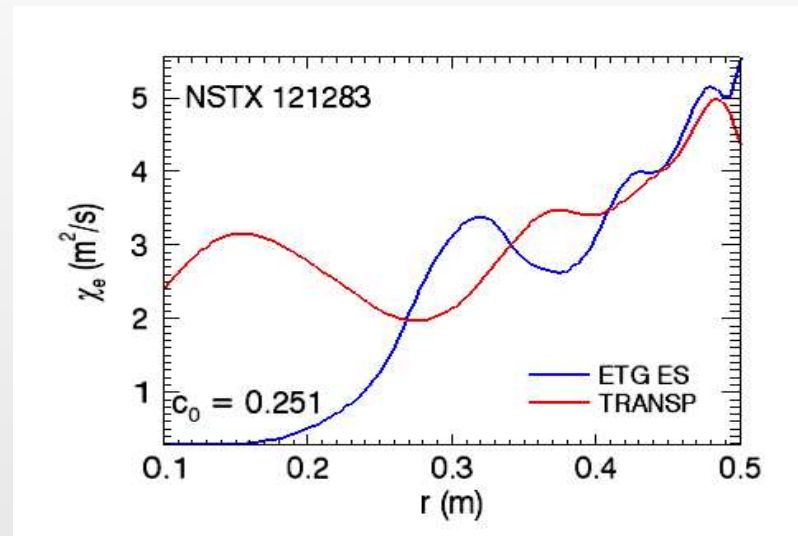
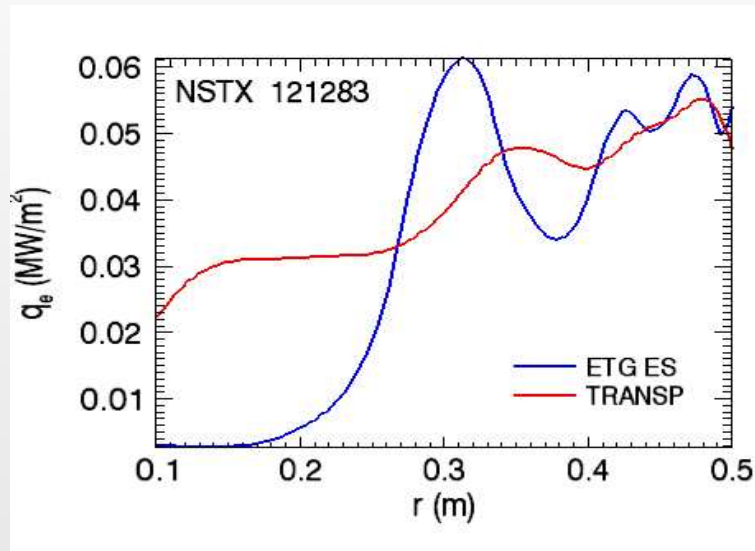
At $r/a = 0.65$, the wavenumber $k_y \rho_e \sim 1$ corresponding to ETG mode are the most unstable.



On the contrary, at the core, the modes $k_y \rho_e \sim 1$ are stable or less unstable compared to the mode, $k_y \rho_s \sim 1$ (microtearing, TEM). At high B_T , the ETG modes are suppressed.

NSTX 121283

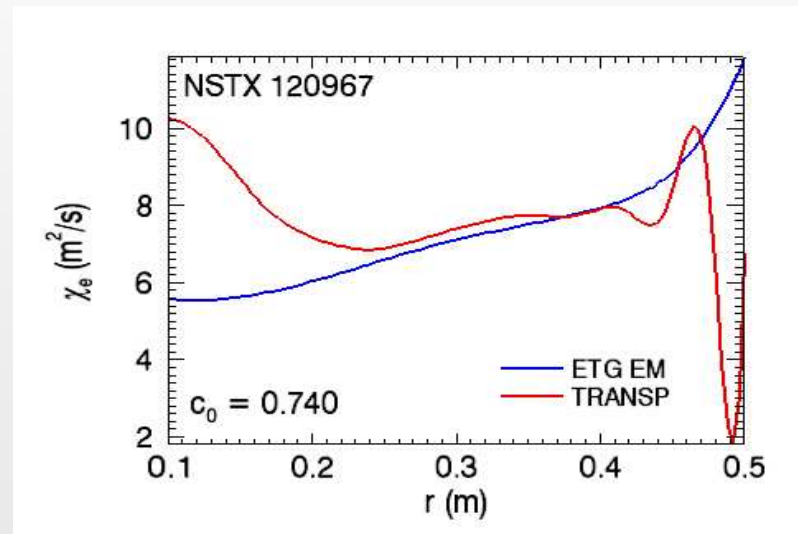
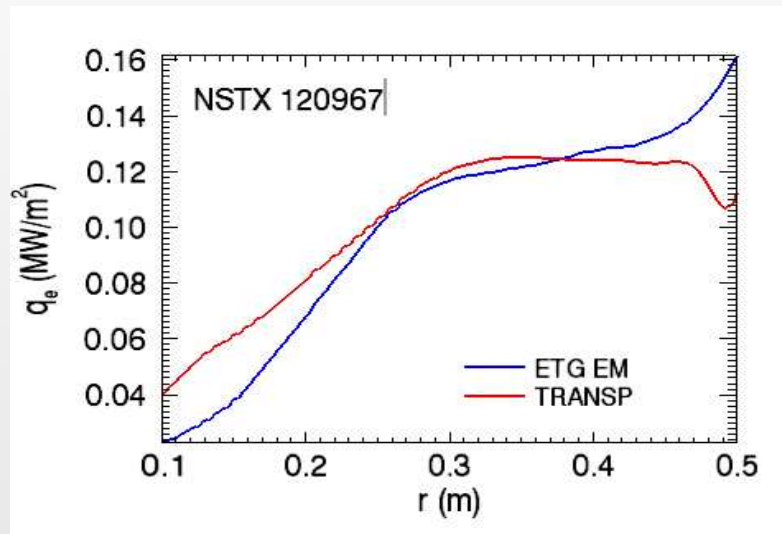
- $\beta_e < 0.1$ at the outer region



- The **electrostatic** scaling does have good agreement with the TRANSP q_e and χ_e in the pre-H mode and H-mode.

NSTX 120967

- $\beta_e \sim 0.1$ at the outer region

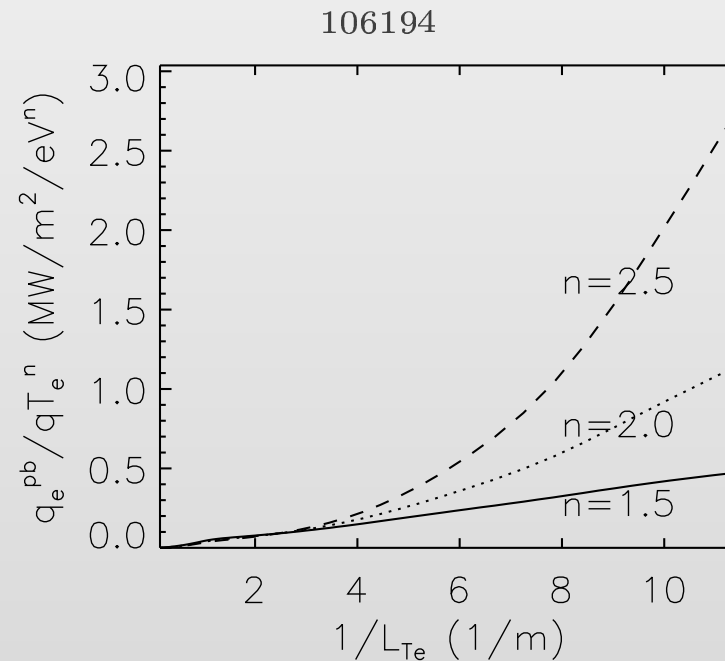


The [electromagnetic](#) scaling does have good agreement with the TRANSP ones.

Temperature Dependence

The scaling law shows that heat flux is proportional to

- ETG EM $\sim T_e^{3/2}$
- Bohm $\sim T_e^2$
- gyro-Bohm $\sim T_e^{5/2}$
- ETG ES $\sim T_e^{5/2}$



ETG Model

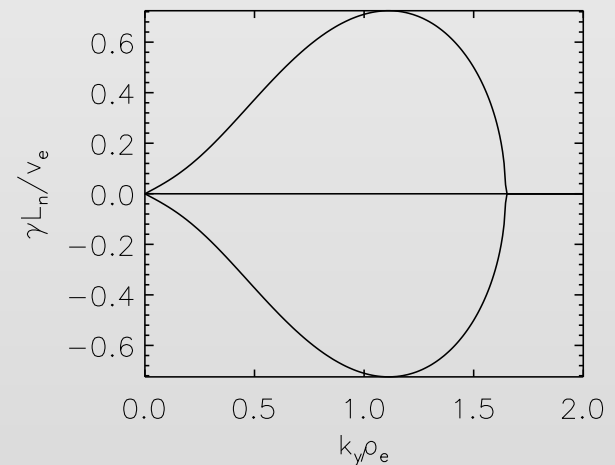
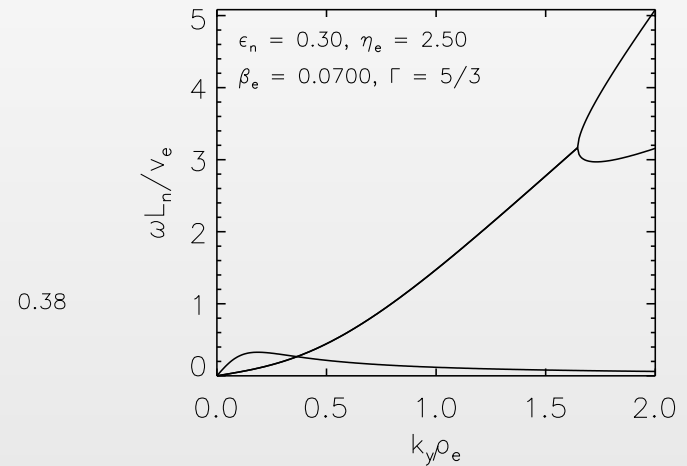
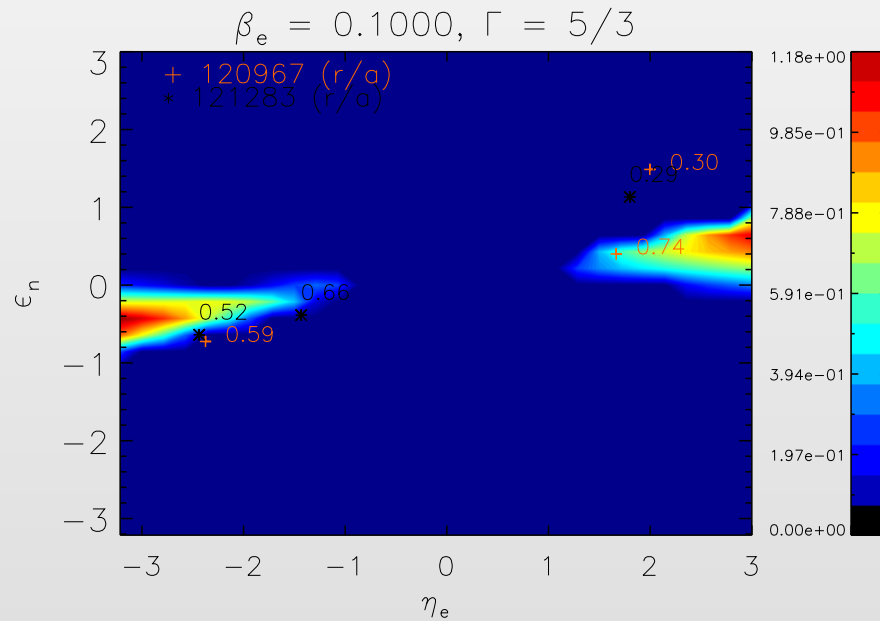
The non-dimensional ETG three-field model is derived by Horton (NF, 2005) in a locally toroidal geometry,

$$\begin{aligned}
 (1 - \nabla_{\perp}^2) \frac{\partial \hat{\phi}}{\partial t} &= [1 - 2\epsilon_n + (1 + \eta_e)\nabla_{\perp}^2] \frac{\partial \hat{\phi}}{\partial y} + 2\epsilon_n \frac{\partial \delta \hat{T}_e}{\partial y} \\
 &\quad + [\hat{\phi}, \nabla^2 \hat{\phi}] + \partial_{\parallel}^{nl} \nabla_{\perp}^2 \hat{A} - \mu \nabla^4 \hat{\phi} \\
 \left(\nabla^2 - \frac{\beta}{2} \right) \frac{\partial \hat{A}}{\partial t} &= \frac{\beta}{2} (1 + \eta_e) \frac{\partial \hat{A}}{\partial y} + 2\partial_{\parallel}^{nl} \hat{\phi} - \partial_{\parallel}^{nl} \delta \hat{T}_e - [\hat{\phi}, \nabla^2 \hat{A}] - \frac{\eta}{\mu_0} \nabla^2 \hat{A} \\
 \frac{\partial \delta \hat{T}_e}{\partial t} &= S_{\text{rf}}(r) - s_{\text{edge}}(r) - [\eta_e - 4\epsilon_n(\Gamma - 1)] \frac{\partial \hat{\phi}}{\partial y} - 2\epsilon_n(2\Gamma - 1) \frac{\partial \delta \hat{T}_e}{\partial y} \\
 &\quad - (\Gamma - 1) \partial_{\parallel}^{nl} \nabla^2 \hat{A} - [\hat{\phi}, \delta \hat{T}_e] + \chi_{\perp} \nabla_{\perp}^2 \delta \hat{T}_e + \chi_{\parallel} \left(\partial_{\parallel}^{nl} \right)^2 \delta \hat{T}_e
 \end{aligned}$$

where $T_e/T_i = 1.0$ is assumed and $(x, y)/\rho_e \rightarrow (x, y)$,
 $tv_e/L_n \rightarrow t$.

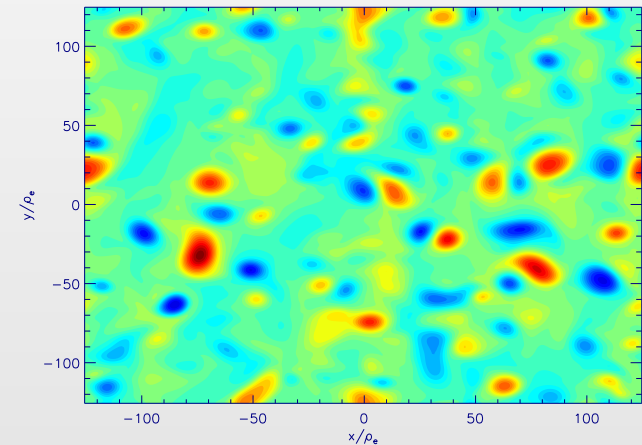
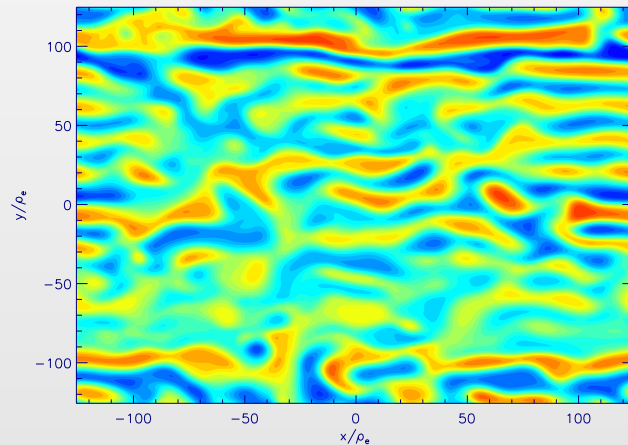
Linear Analysis

$0.3 < k_y \rho_e < 1.0$ modes have the maximum growth rates.
The most unstable modes are electrostatic in nature.



Nonlinear Simulation(TEST)

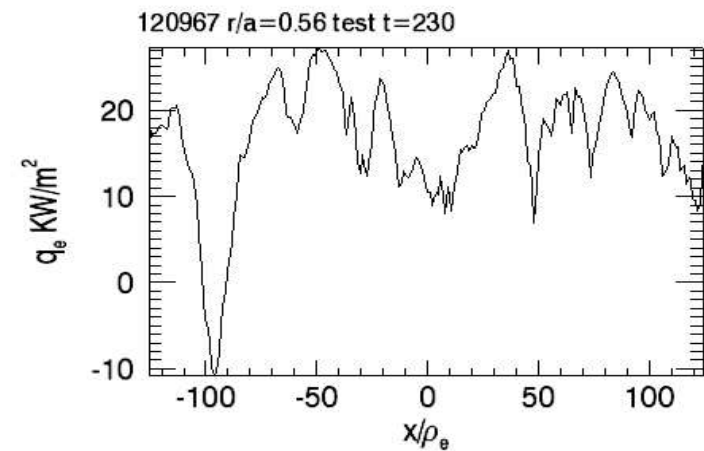
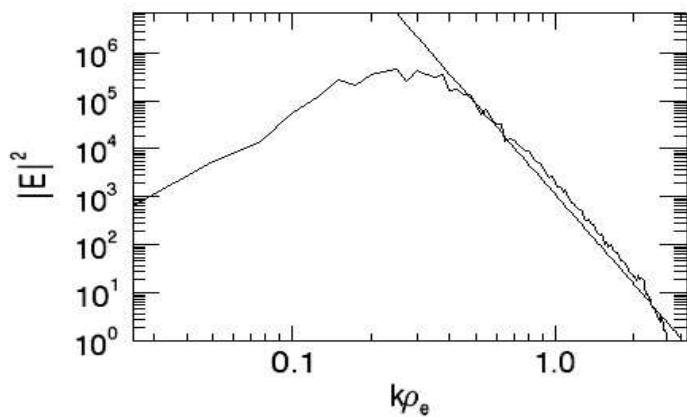
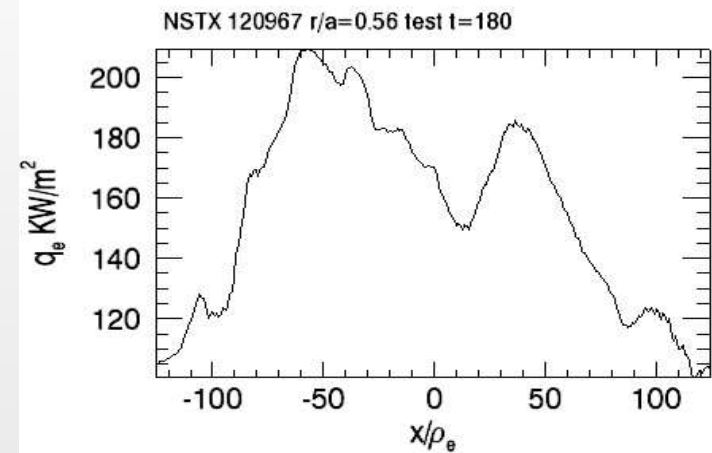
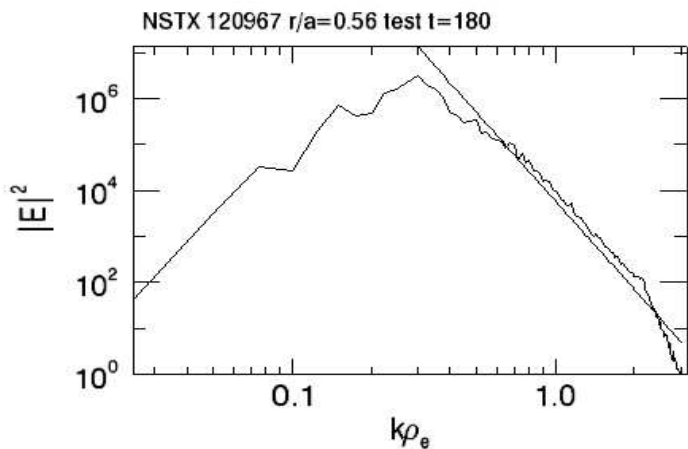
- Using the pseudospectral method, the set of the electrostatic equations are solved with $\eta_e = -2.7$, $\epsilon_n = -0.8$ and $\beta_e = 0.19$ from $120967 r/a = 0.56$. The box size is $[240\rho_e, 240\rho_e]$.



- The left figure is taken when the linear growth is fully suppressed. The streamer-like structure formation is almost finished.
- The right figure is taken when the structure is decay by setting $\eta_e = 0$.

Nonlinear Simulation II(TEST)

For each case, ($\% q_e = 120 \text{ KW/m}^2$ from TRANSP)



Conclusion and Future Direction

- The linear analysis indicates that the ETG modes $k_y \rho_e \sim 1$ are unstable at the outer region.
- The electron heat transport in NSTX 106194 and 120967, or 121283 can be explained by the electromagnetic or electrostatic ETG modes.
- We need to refine the β_e criterion for the transition between EM and ES.
- We are going to investigate the effect of the inverted profile $\nabla n_e \cdot \nabla T_e < 0$ on the transport.