

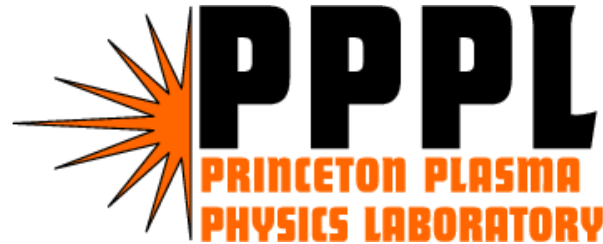
TORIC Modeling of HHFW in NSTX

P.T. Bonoli, J.C. Wright,
and J. Lee



C.K. Phillips, T. Brecht,
E. Valeo, R. Bell, J. Hosea,
B. LeBlanc, G. Taylor,
J. R. Wilson, and the NSTX
HHFW Team

D. L. Green
and P. Ryan



R.W. Harvey **COMPX**

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Outline

- TORIC code development:
 - Elimination of numerical instability in code at high mode resolution.
 - “Radial” parallelization of block tri-diagonal matrix solver.
- Sensitivity analysis of ion/electron absorption partition in HHFW.
- Progress in implementing realistic SOL models in full-wave solvers (AORSA)
- Future work

Numerical instability at high poloidal mode number eliminated

$$\frac{4\pi i}{\omega} \vec{J}_e^{(2)} = -\frac{c^2}{\omega^2} \nabla_{\perp} \times [\lambda_0 (\nabla_{\perp} \times \vec{E}_{\perp})] + i \frac{c^2}{\omega^2} \{ \nabla_{\perp} \times [\xi_0 (\vec{b} \cdot \nabla) E_{\parallel} \vec{b} + \vec{b} (\vec{b} \cdot \nabla) [\xi_0 \vec{b} \cdot (\nabla_{\perp} \times \vec{E}_{\perp})]] \},$$

$$\lambda_0 = \frac{1}{2} \frac{\omega_{pe}^2}{\Omega_{ce}^2} \frac{v_{te}^2}{c^2} [-x_e Z(x_e)],$$

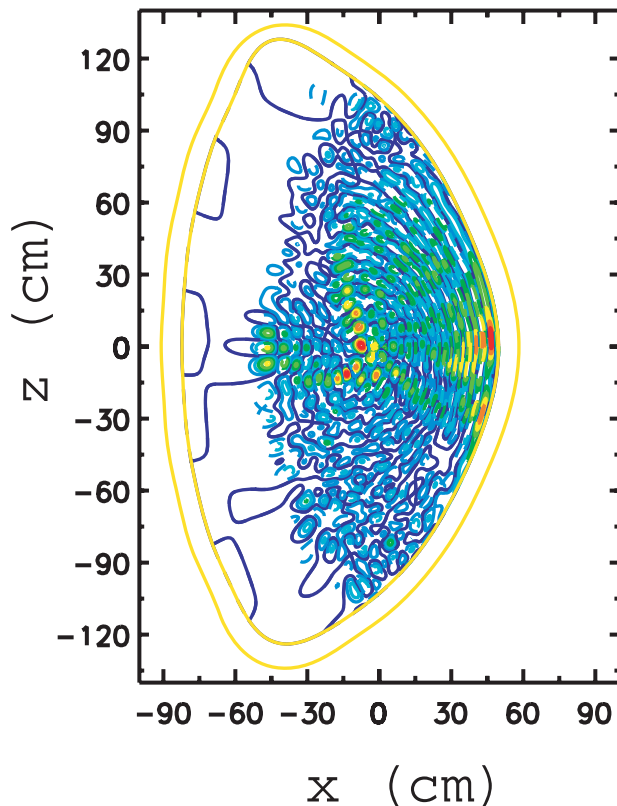
$$\xi_0 = \frac{1}{2} \frac{\omega_{pe}^2}{\omega \Omega_{ce}} \frac{v_{te}^2}{c^2} [x_e^2 Z'(x_e)].$$

- **Second order electron FLR current:**
 - **First term (above) corresponds to TTMP**
 - **Second term is cross-term between TTMP and ELD**
- **Found instability was caused by error in the cross-term.**
- **Also resolves instability seen in the LHRF TORIC solver.**

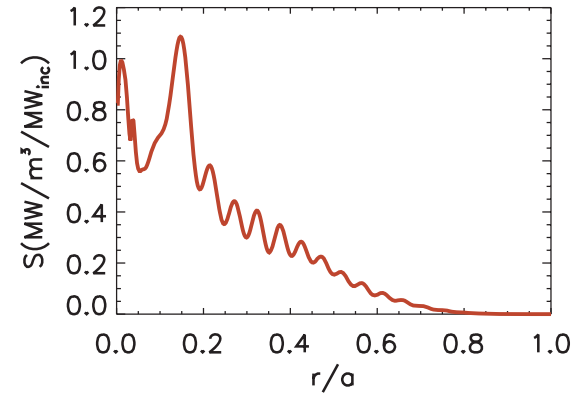
High resolution simulations of HHFW heating in NSTX can now be carried out using TORIC

[$N_r=480$, $N_m=511$]

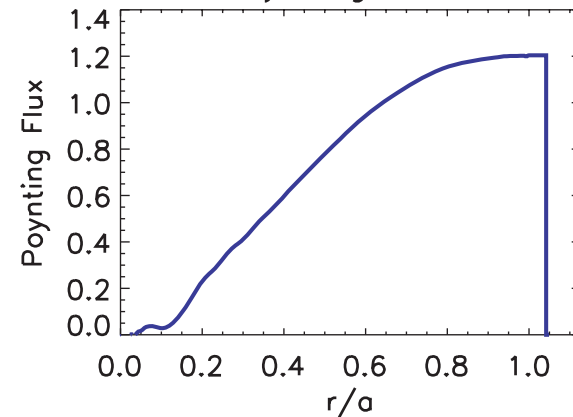
$\text{Re} \{ E_+ \}$



Electron Absorption



Poynting Flux



- Prior to correction, Poynting flux and power absorption behaved poorly as N_m was increased beyond 63-127.

New TORIC parallel solver is 2-4 times faster

Uses a 3-D parallelization of block-tridiagonal matrix system

Current Solver

$$\underline{L}_i \cdot \vec{x}_{i-1} + \underline{D}_i \cdot \vec{x}_i + \underline{R}_i \cdot \vec{x}_{i+1} = \vec{y}_i$$

$$\underline{D}_{i+1} = \underline{D}_{i+1} - \underline{L}_{i+1} \times \underline{D}_i^{-1} \times \underline{R}_i$$

$(6N_m) \times (6N_m)$

- Serial (Radial direction [$i=1.. N_\psi$]: Thomas Algorithm)

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2-D Parallel (Poloidal m modes :
Scalapack matrix calculation $\rightarrow 6N_m \times 6N_m$)

New Solver

$$Ax \equiv \begin{bmatrix} a_1 & b_1 & & & & \\ c_2 & a_2 & b_2 & & & \\ & c_3 & a_3 & b_3 & & \\ & & \ddots & \ddots & \ddots & \\ & & & & b_{n-1} & \\ & & & & c_n & a_n \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \\ x_3 \\ \vdots \\ x_n \end{bmatrix} = \begin{bmatrix} r_1 \\ r_2 \\ r_3 \\ \vdots \\ r_n \end{bmatrix} \equiv r.$$

- 1-D Parallel (Radial direction: combination of Divide-and-Conquer and Odd-even cyclic reduction Algorithms) \rightarrow # P1 groups

P.Garaud, Mon.Not.R.Astron.Soc,391(2008)1239-1258

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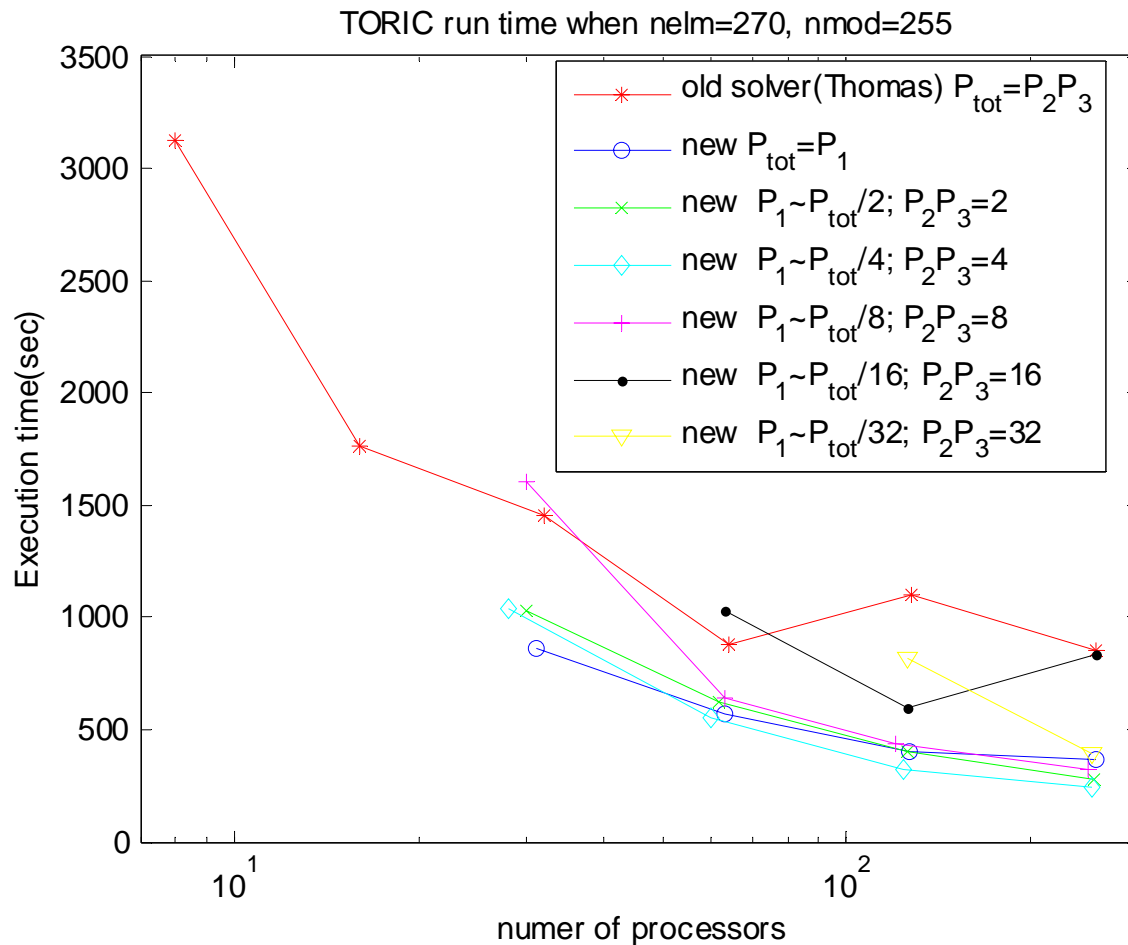
2-D Parallel (Poloidal m modes :
Scalapack matrix calculation $6N_m \times 6N_m$)

\rightarrow #P2*P3 processors

= 3-D grid (P_tot=P1*P2*P3)

New solver is crucial for efficient parallel performance of TORIC in TRANSP

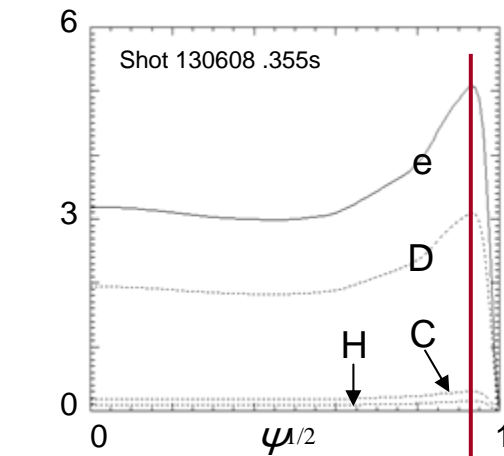
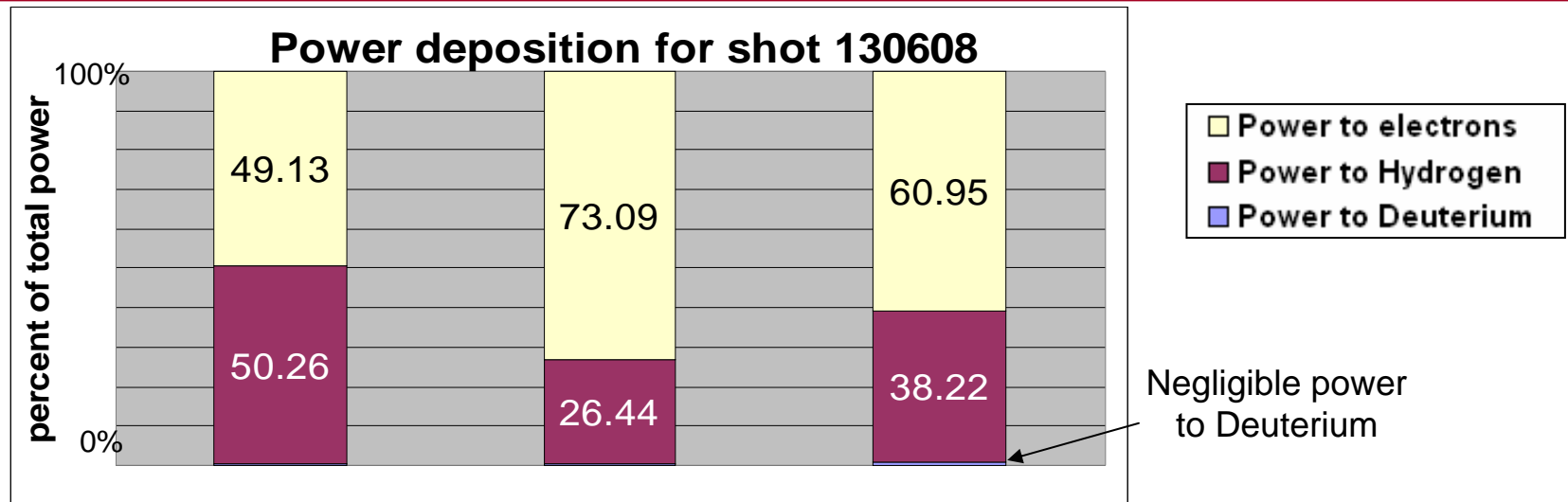
- Optimized grid $\rightarrow P_1 = P_{\text{tot}}/4$ and $P_2 P_3 = 4$



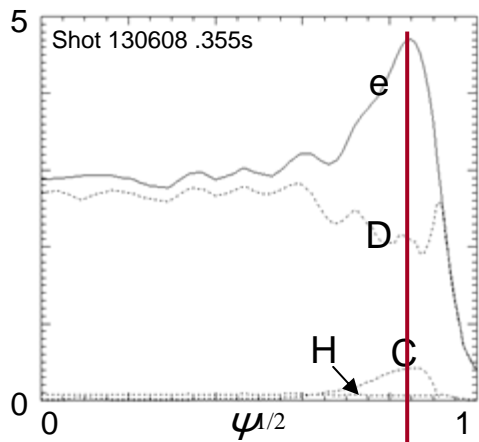
J. Lee, MIT

Will provide sufficient mode resolution in parallel implementation in TRANSP

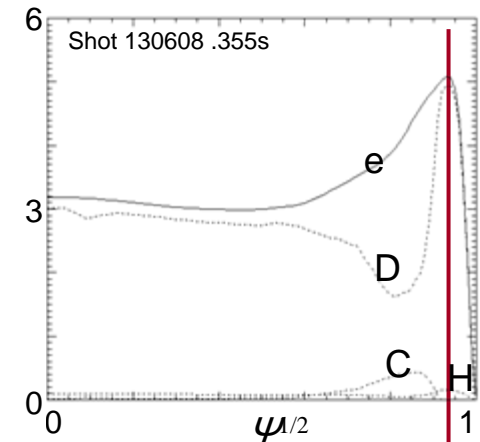
Profile shape, especially near the edge, affects power split between ions and electrons



Menard smoothing
 $n_i \propto n_e$ and n_c from \bar{z}_{eff}



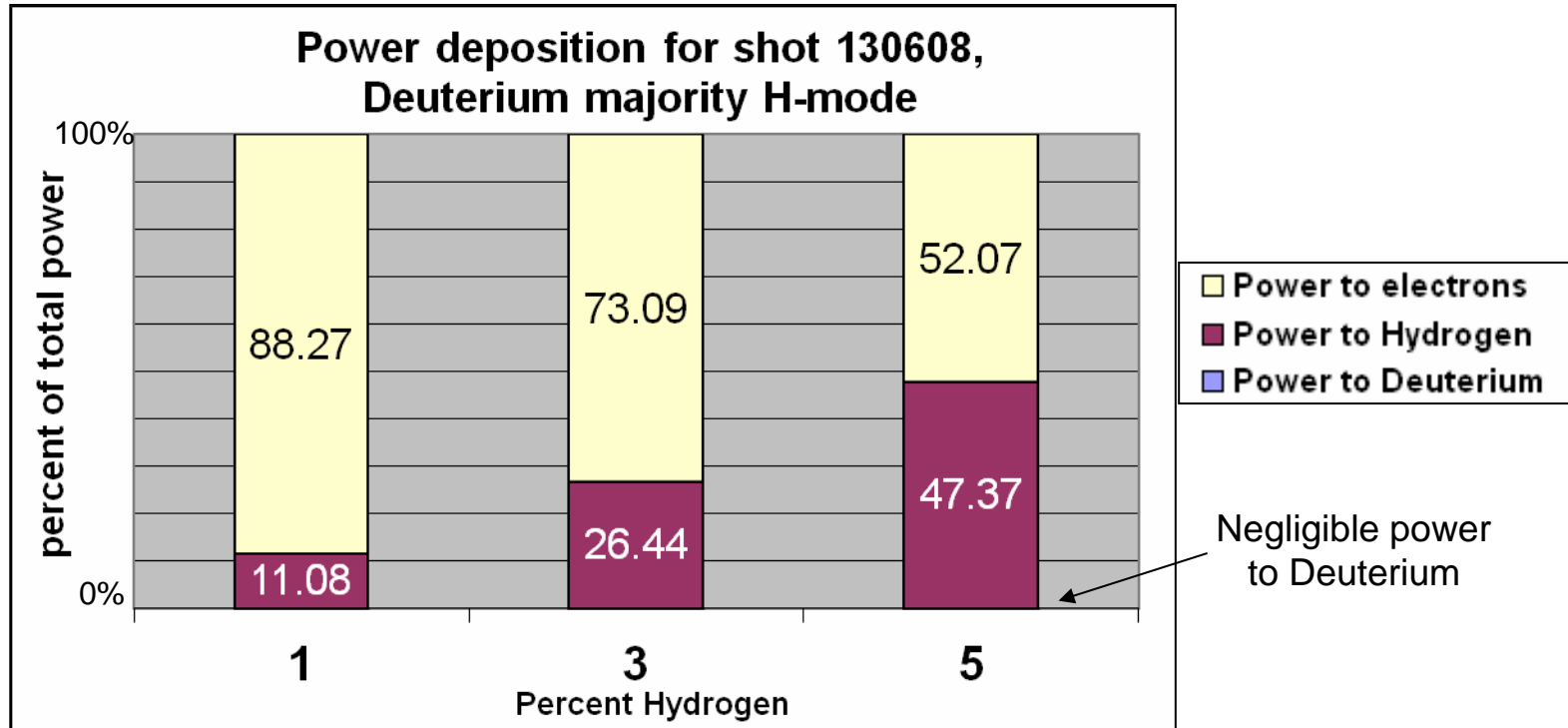
R. Bell spline fit



R. Bell smoothing

n_C from CHERS
 n_H/n_D from $H\alpha/D\alpha$
 n_D from \bar{z}_{eff} & quasineutrality

Hydrogen absorption increases as H concentration increases in H-mode plasmas

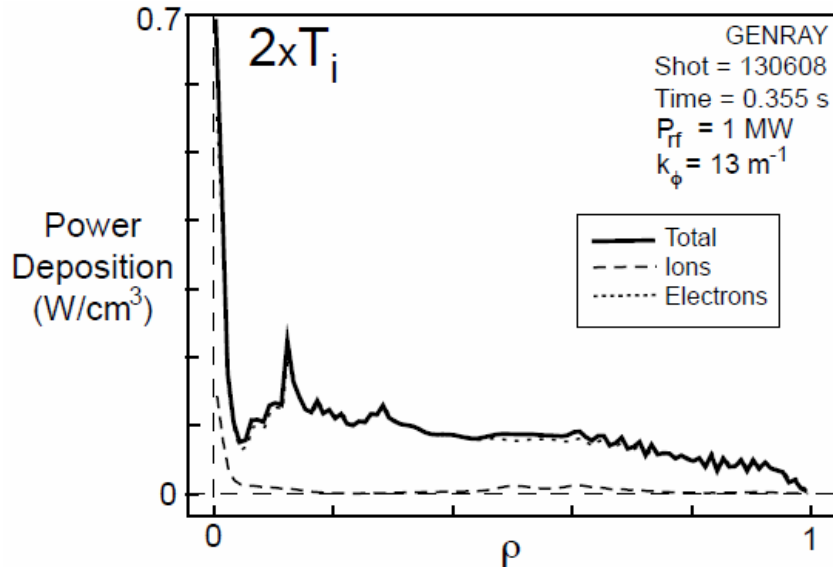


This level of sensitivity necessitates careful measurement and control of H concentration in experimental plasmas

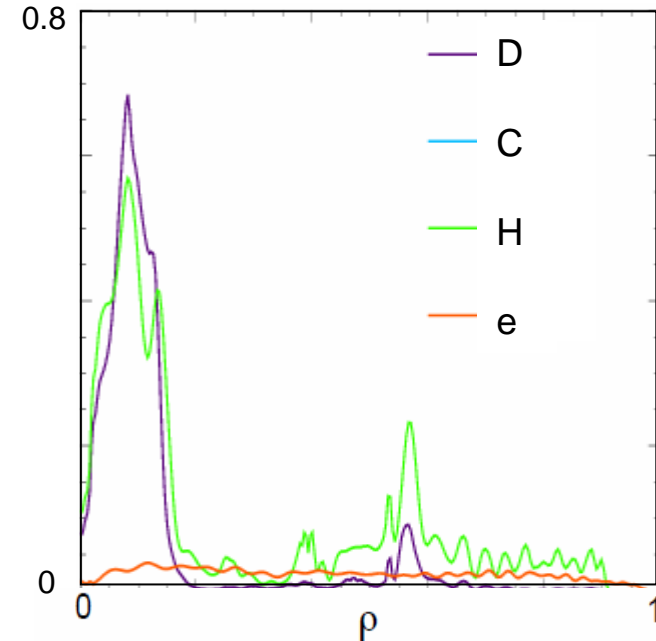
Comparable trend found in L-mode simulations

GENRAY and TORIC predict electron absorption decreases as T_i increases

GENRAY



TORIC



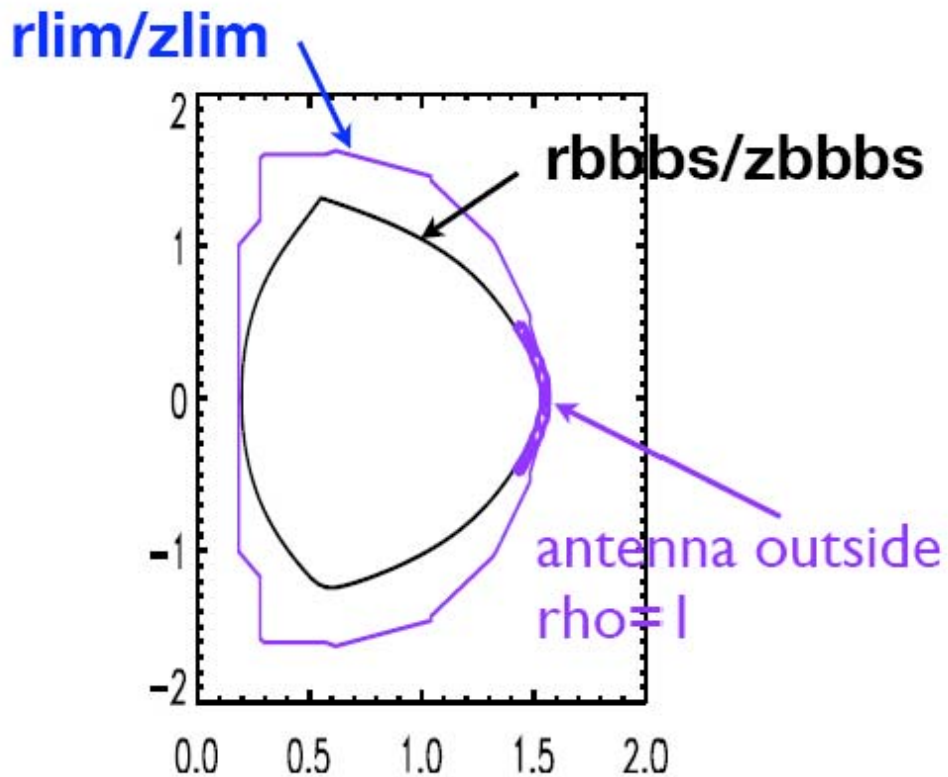
Code:	Power to:	CHERS	CHERSx2	CHERSx4
TORIC	e	71.77	39.4	7.91
	H	26.85	27.36	4.43
	D	1.38	33.23	87.66
GENRAY	e	99.3	84.7	45.3
	H	0.7	9.1	20.2
	D	0	6.1	31.4

TORIC predicts stronger dependence on T_i than GENRAY

>> recent TORIC code fix may resolve the discrepancy

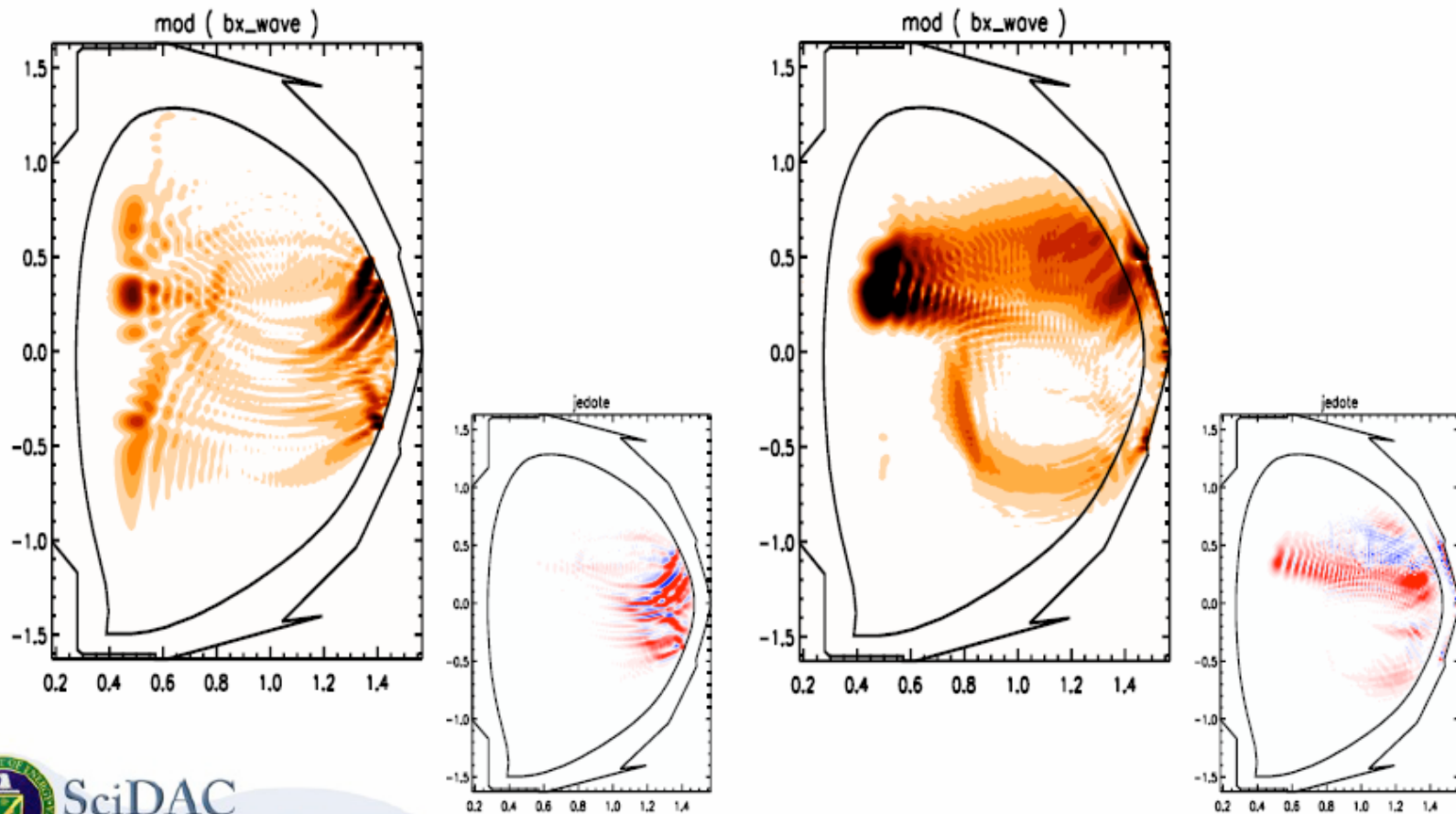
AORSA boundary has been extended to the limiter

- combining Thompson scattering, charge exchange and reflectometer data with **SOLPS** results gives an estimate of the 2D profiles
- edge gradients are large



Wave fields depends on location of the launcher and SOL profiles

- NSTX HHFW, shot 128797@400ms, D ~1keV, maxwellian, 1MW@30MHz, nPhi = -18, 128x256 using smooth profiles



Conclusions

- RF power deposition on ions and electrons is sensitive to:

- Edge density profiles

High density nearest the edge of the plasma causes more ion absorption than when density rises away from the edge.

- Hydrogen concentration

Hydrogen absorbs more RF power as its concentration increases in both H- and L- mode plasmas.

- Ion temperature

Ion absorption increases significantly with rising ion temperatures. (even modest temperatures of 3 keV)

- Future work:

- Revisit dependence of power partitioning on T_i

- Explore effect of limiter boundary on wave propagation

- Carry out 3D field reconstructions for wave propagation in SOL

- Improve algorithms in TRANSP for more accurate, faster simulations