

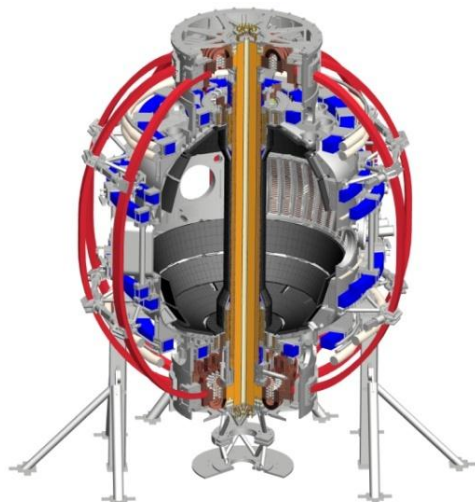
NSTX Upgrade Cryo-pumping Design Progress, Particle Control Plans

John Canik

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Motivation

- Particle control is needed to meet NSTX-U programmatic goals
 - Avoid density limit, radiative collapse during long-pulse (5s) discharge
 - Reduce collisionality to access new core physics
 - Control n/n_G for non-inductive scenarios
- Several PAC recommendations concern particle control
 - Perform cryo-pump design study as complement to Li efforts
 - PAC 29-4
 - PAC 29-5b
 - PAC 29-10
 - Consider alternatives to Li ELM-free scenario: Type-I or small-ELM
 - PAC 29-40
 - PAC 29-42
- Milestone R(12-2): Project deuterium pumping capabilities for NSTX-U using lithium coatings and cryo-pumping
 - Use existing discharges to assess persistence of pumping by Li coatings, project to NSTX-U pulse lengths
 - Develop cryo-pump design, analyze which scenarios and densities can be pumped with stationary deuterium inventory

Outline

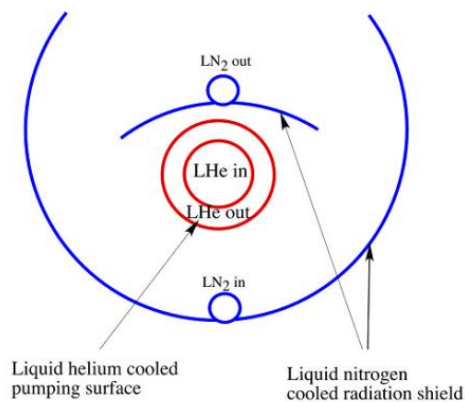
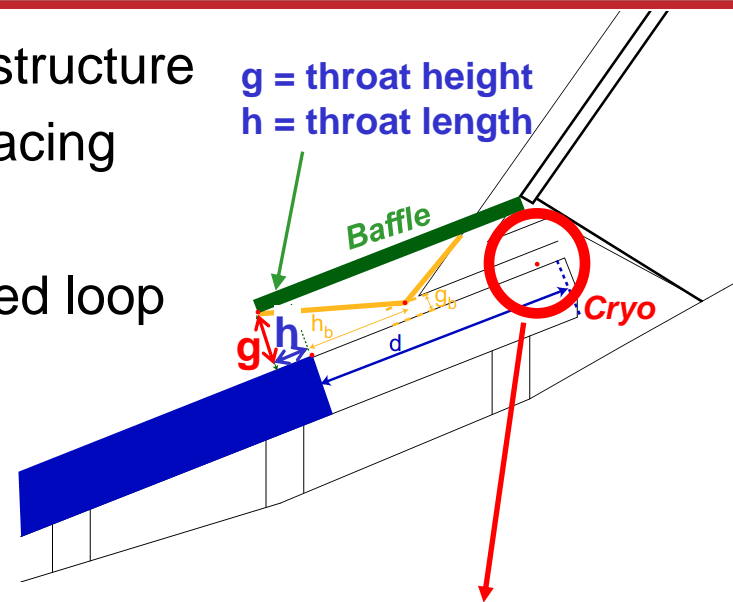
- Progress in cryo-pump design
 - Pumping model developed for use in plenum geometry design
 - Performance and flexibility of optimized system
- Analysis of use of lithium coatings for long-pulse
 - Time-dependent recycling characteristics in Li ELM-free plasmas
 - Long-pulse, ELMy plasmas with partially passivated lithium
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 - Near term analysis
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General design similar to DIII-D lower outer cryo-pump system is taken as starting point for design analysis

- Plenum location studied: under new baffling structure near secondary passive plates, possibly replacing some outer divertor plates and tiles
- Pumping capacity of a toroidal liquid He cooled loop (Menon, NSTX Ideas Forum 2002)
 - $S=24,000$ l/s @ $R=1.2$ m
 - Need plenum pressure of 0.83 mtorr to pump beam input (10MW~20 torr-l/s)
- Pumping rate:

$$I_{pump} = P_{pl} S = \frac{I_0}{S + C} S$$

- P_{pl} = plenum pressure
- I_0 = neutral flux into plenum
- C = throat conductance
- To optimize, need $C(g,h)$, $I_0(g,h)$



Cross-section of the pump (10 cm outer dia.)

Projected divertor parameters combined with semi-analytic pumping model are used to calculate pumping rates

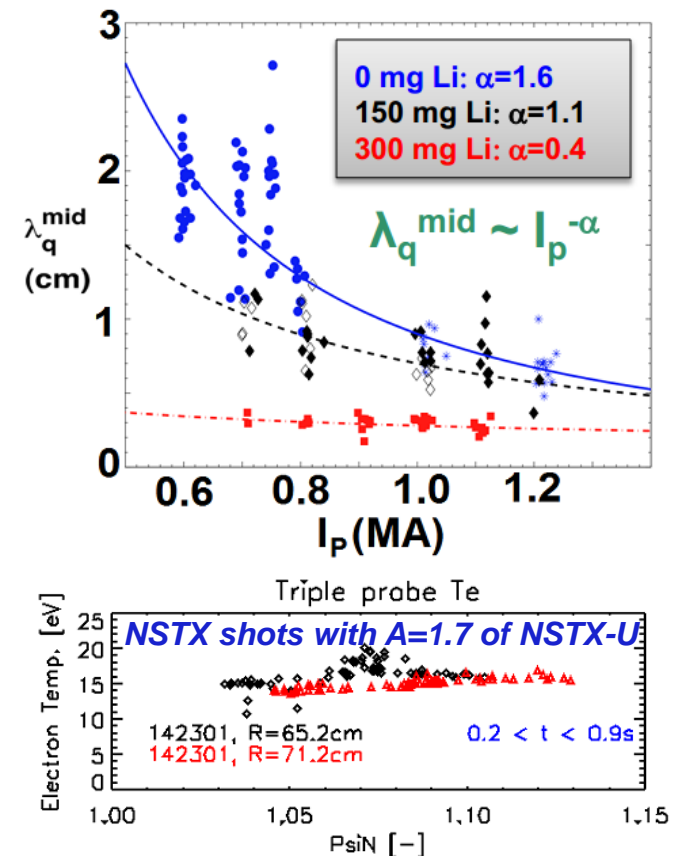
- Analytic model developed for DIII-D pumping studies (Maingi, NF '99)
 - Predicts plenum pressure, validated with DIII-D data
 - Requires divertor n , T , Γ profiles
 - Uses first-flight neutral model (insufficient for detached divertor)

- Heat flux, angle of B wrt PFC surface (α), and plasma temperature are sufficient to calculate n , Γ :

$$\Gamma_{\perp} = q_{\perp} / 7T$$

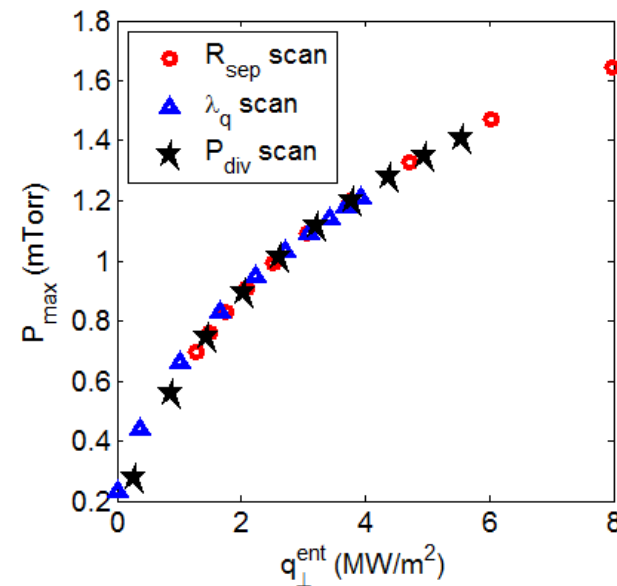
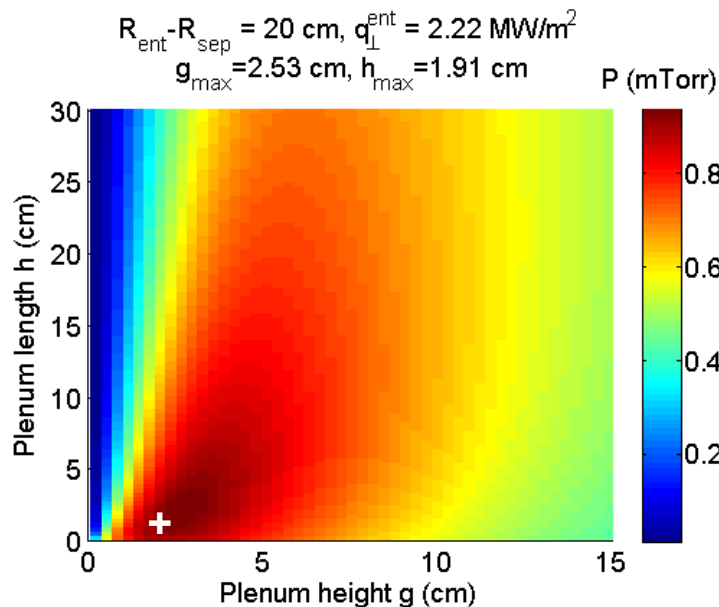
$$n = \Gamma_{\perp} / (\sin \alpha \sqrt{2T/m})$$

- Recent experiments yield scaling of SOL heat flux width
 - No-lithium scaling used here, but all trend towards $\lambda_q \sim 3\text{mm}$ at $I_p = 2\text{MA}$
 - $P_{\text{div}} = 5\text{ MW}$ assumed (1/2 of 10 MW input)
- Langmuir probes show $T_e \sim 15\text{-}20\text{ eV}$ in far SOL
 - $T_e \sim$ independent of I_p
 - $T_e \sim 15\text{ eV}$ assumed (NSTX-U-like discharges)



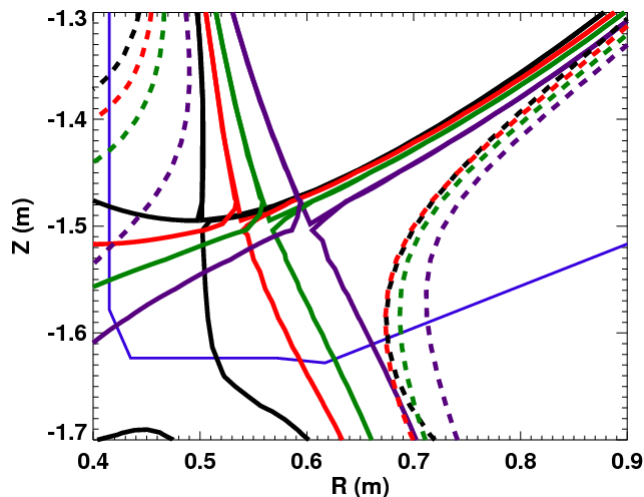
Pressure projections are used to optimize plenum geometry parameters

- Exponentially decaying heat flux footprint imposed, with $T_e=15$ eV
- Plenum entrance height, length are varied to maximize pressure
- Pressure in optimized plenum depends primarily on heat flux at pump entrance
 - Varied through R_{OSP} , flux expansion or $P_{tot} \Rightarrow$ profile effects not important
 - Reaching $P \sim 0.8$ mTorr (to pump 10 MW NBI) requires $q_{\perp}^{ent} \sim 2$ MW/m²
- Optimal plenum entrance for $P=0.8$ mTorr: height $g \sim 2.5$ cm, length $h \sim 2$ cm

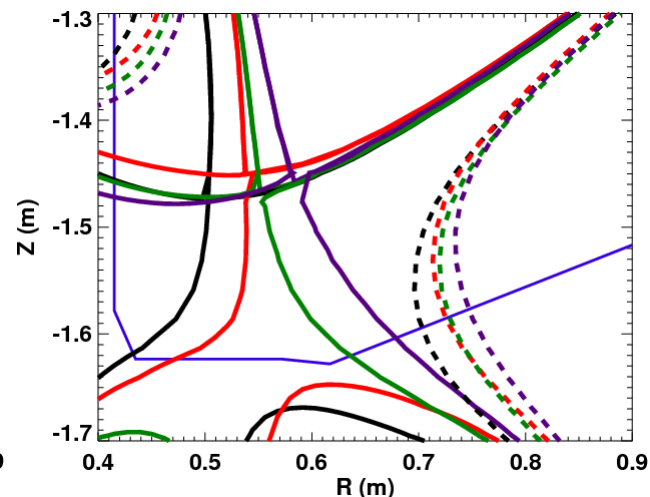


Equilibria with variety of R_{OSP} , flux expansion are used to map heat flux profiles, assess candidate pump entrance locations

- Standard and snowflake divertors considered
 - Four R_{OSP} each
 - $\psi_N=1.0, 1.03$ shown
 - Movement of $\psi_N=1.03$ strike line is much less than that of R_{OSP}

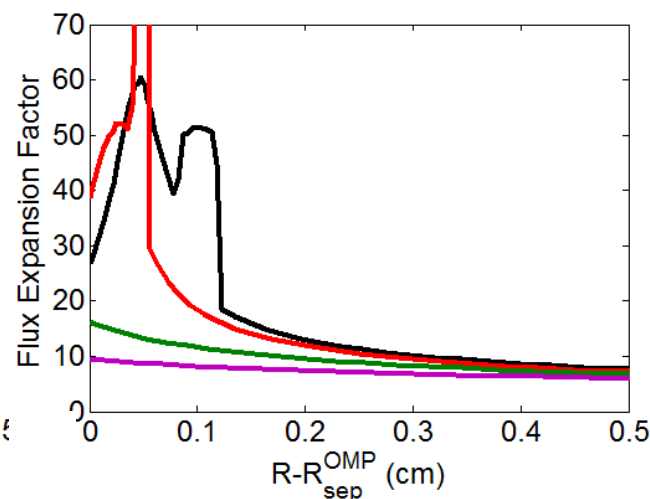
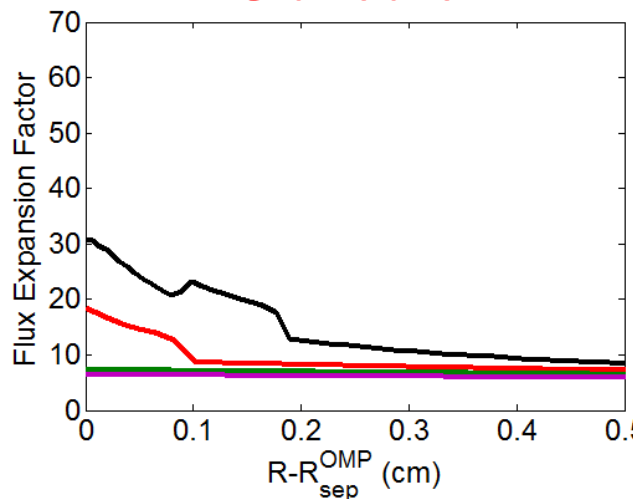


Standard



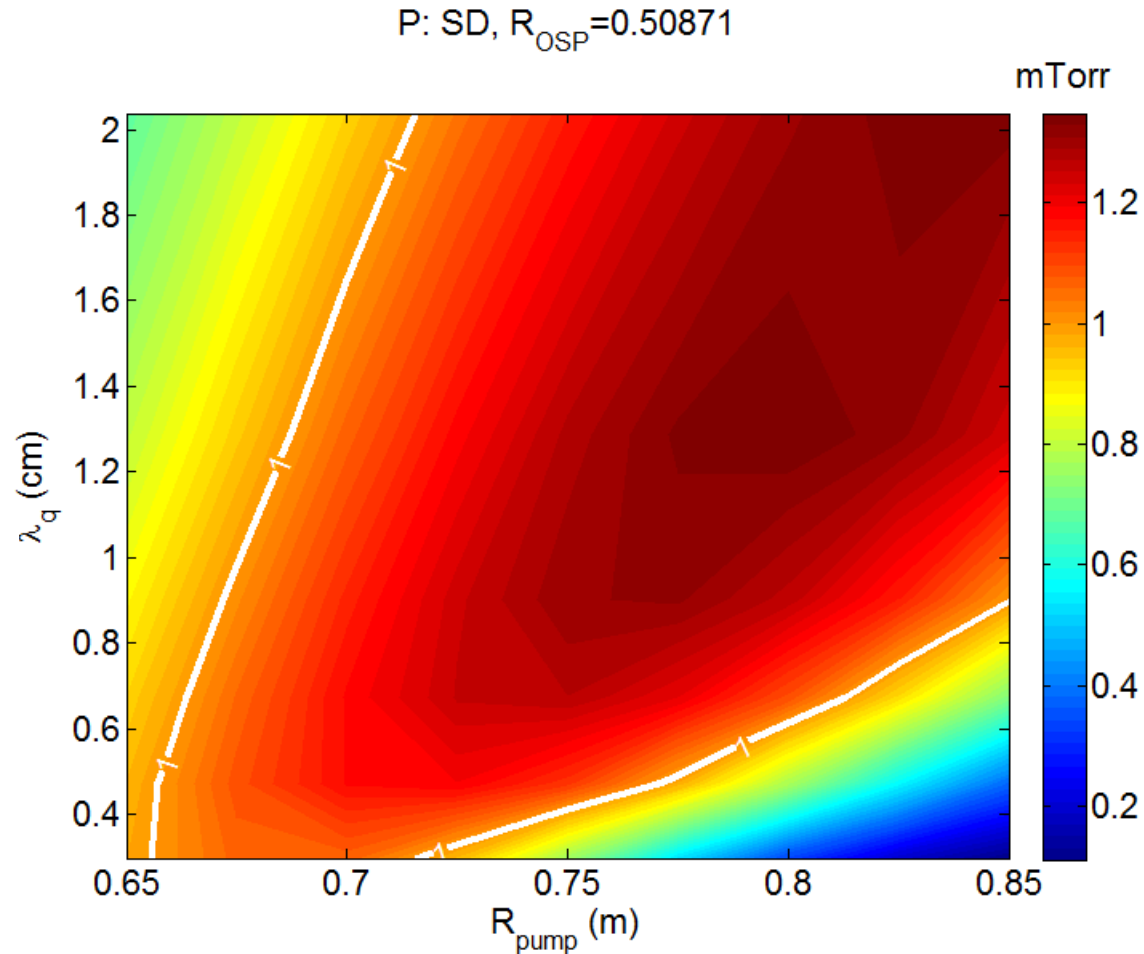
Snowflake

- Flux expansion, flux surface geometry used to convert midplane heat flux profile (from scaling) to divertor heat flux
 - As R_{OSP} is increased, flux expansion is decreased



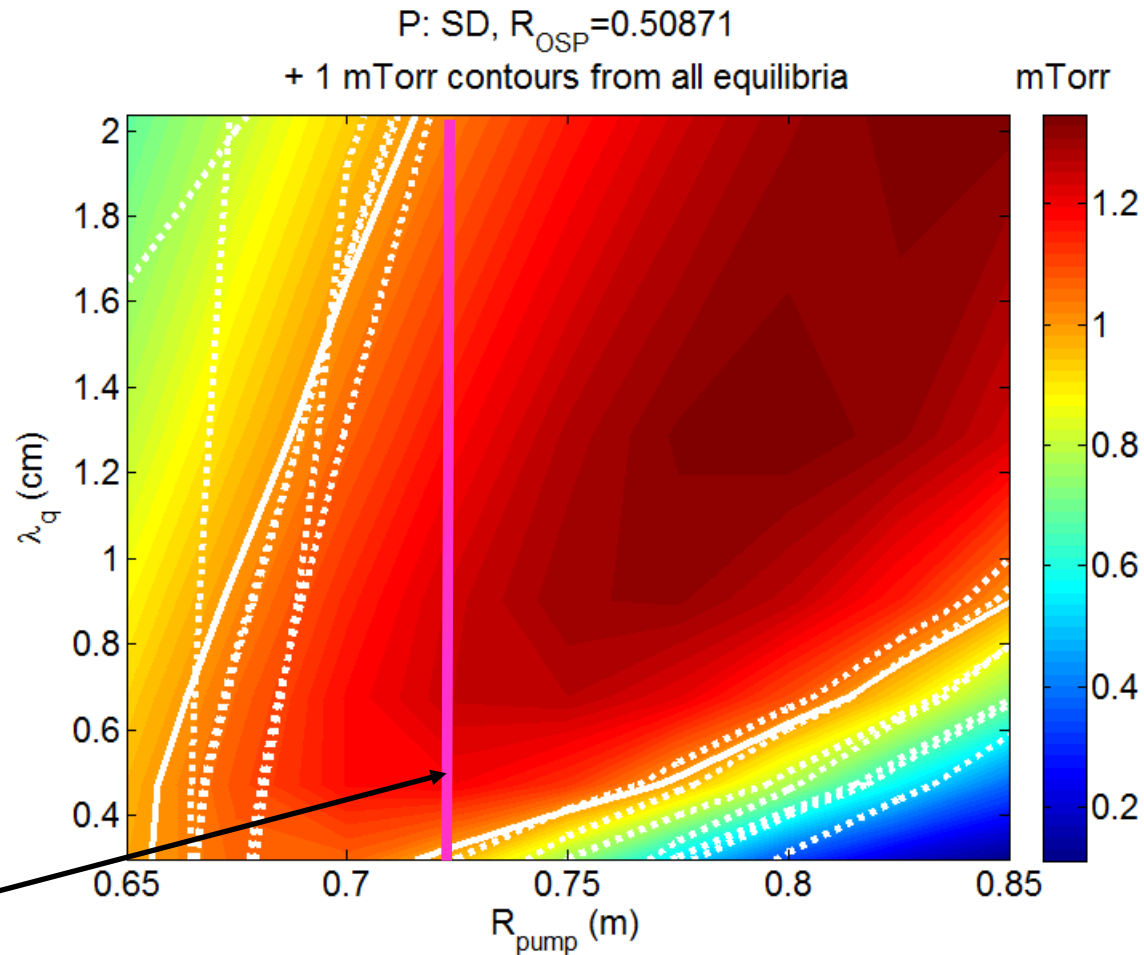
Realistic equilibria, heat flux scaling, and empirical T_e^{SOL} are used to project plenum pressure for candidate R_{pump}

- Analytic model for plenum pressure with optimized entrance parameters
- Pressure is non-monotonic with R_{pump} due to field geometry
 - At low R_{pump} , α is lower, so n/Γ_{\perp} is increased \Rightarrow more neutrals ionized before reaching pump



Realistic equilibria, heat flux scaling, and empirical T_e^{SOL} are used to project plenum pressure for candidate R_{pump}

- Analytic model for plenum pressure with optimized entrance parameters
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 - At low R_{pump} , α is lower, so n/Γ_{\perp} is increased \Rightarrow more neutrals ionized before reaching pump
- Optimizing position for narrowest SOL gives $R_{\text{pump}} \sim 0.7$
 - Narrow SOL gives least flexibility in moving R_{OSP} to improve pumping
 - $R_{\text{pump}} = 0.72$ gives high P for wide range of SOL width



$R_{\text{pump}}=0.72$ gives n_e control for range of I_p , equilibria

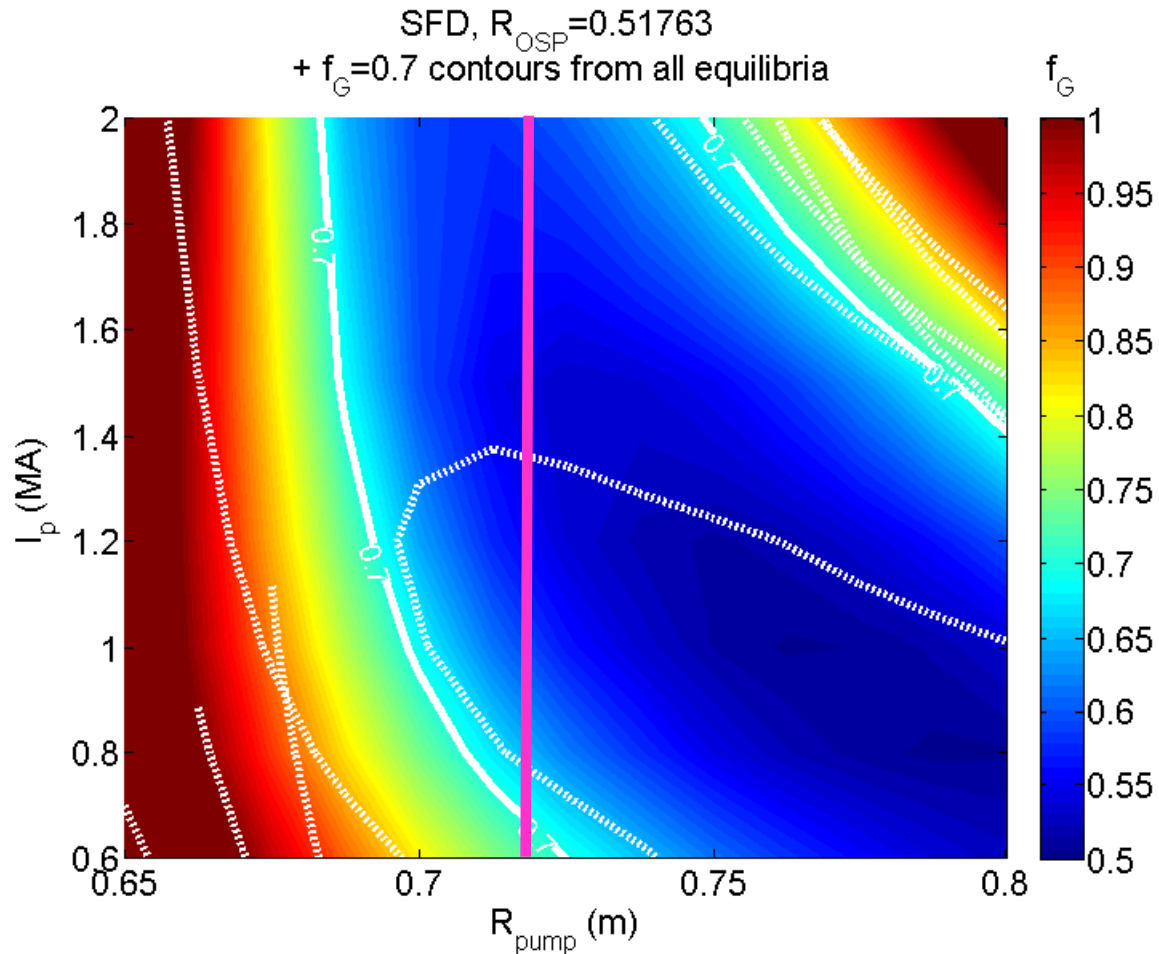
- Modified 2-pt model used to estimate n_e^{sep}

$$T_{OMP} = \left(T_{DIV}^{7/2} + \frac{7}{4\kappa_{0e}} q_{\parallel}^{\text{sep}} L \right)^{2/7}$$

$$n_{OMP} = f_{cal} \frac{2n_{DIV} T_{DIV}}{T_{OMP}} \frac{B_{OMP}}{B_{DIV}}$$

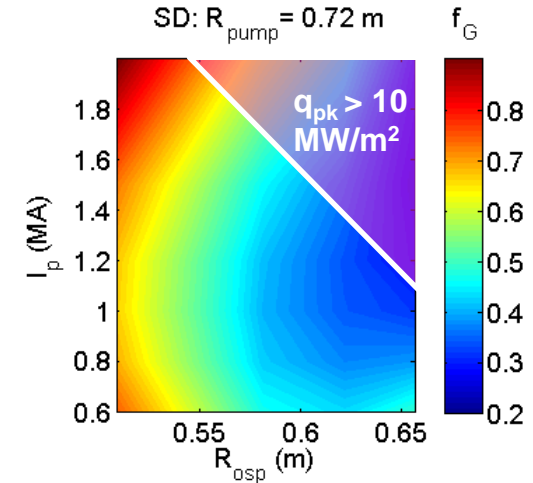
- $q_{\parallel}^{\text{sep}}$ from I_p scaling, T_e^{div} varied
- Final n_e^{sep} : pumping=NBI input

- $\bar{n}_e/n_e^{\text{sep}} \sim 3$ used to estimate $f_G = n/n_G$
 - Consistent with NSTX data

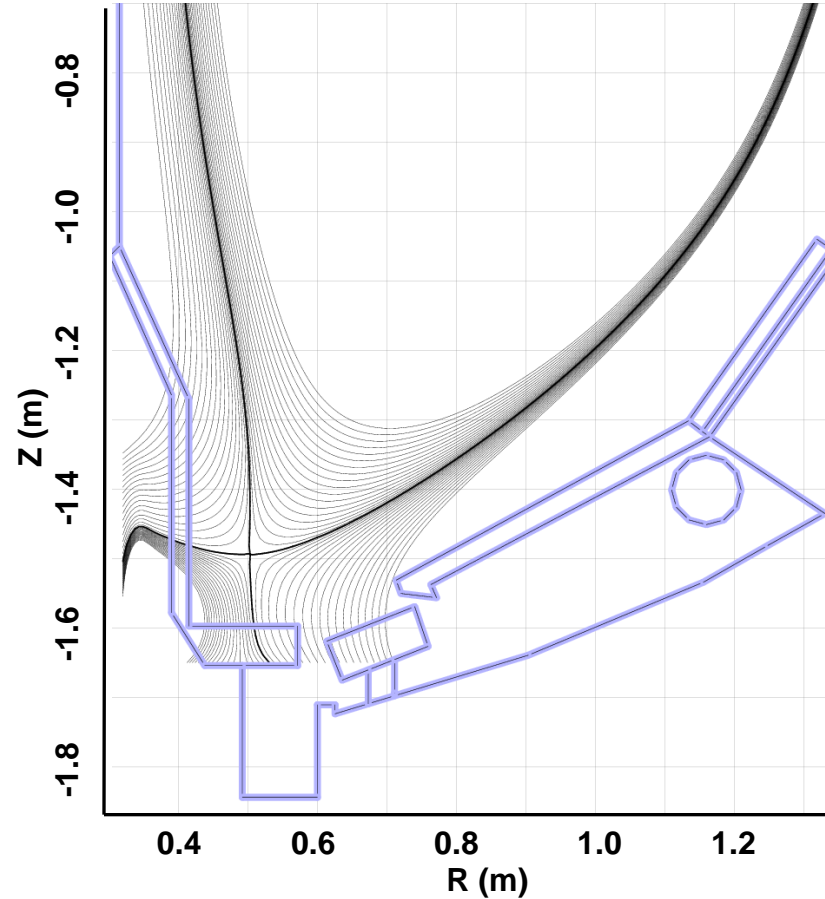
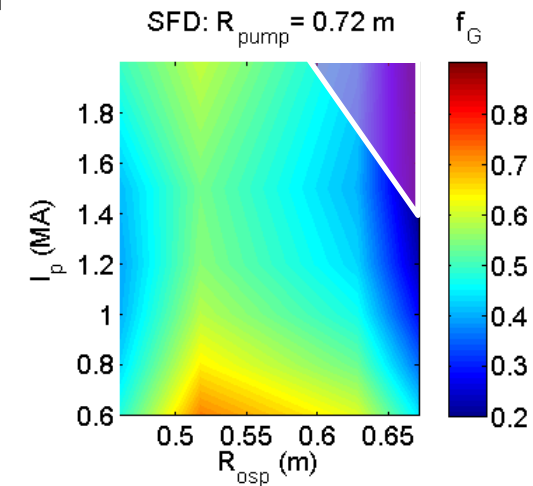


Optimized plenum geometry capable of pumping to low density for a range of R_{OSP} , I_p

- Equilibrium f_G down to < 0.5
 - Moving R_{OSP} closer to pump allows lower n_e , but limited by power handling



- High flux expansion in SFD gives *better* pumping with SOL-side configuration
 - More plasma in far SOL near pump
 - More room to increase R_{OSP} at high I_p

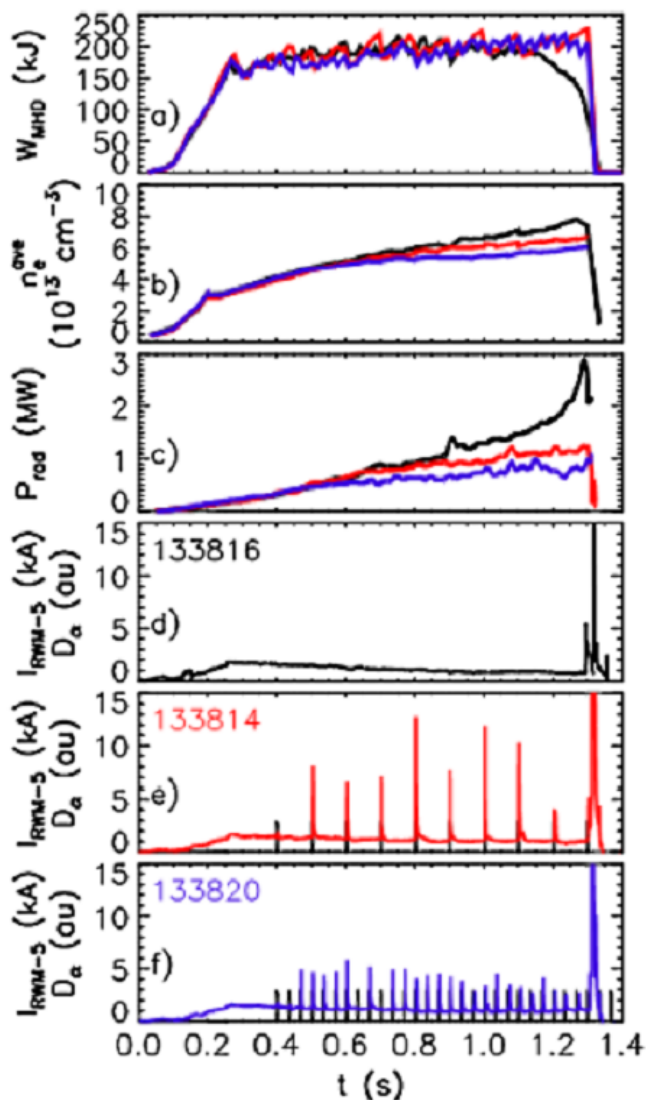


SOLPS geometry to be used in future calculations

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Scenarios with Li coatings and ELMs trend towards stationary D and C inventory—but how do they extrapolate?

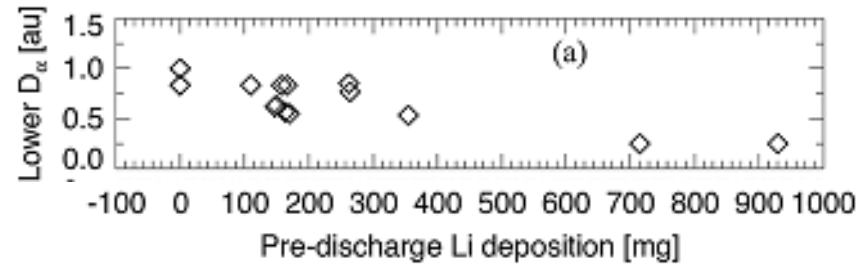


J. Canik - PRL 104, 045001 (2010)

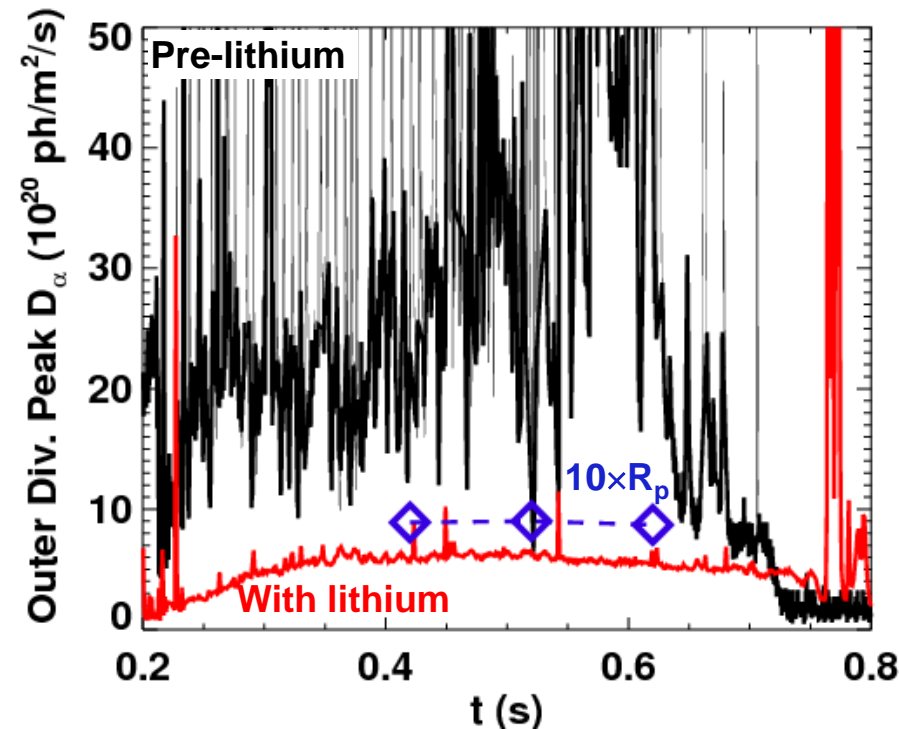
- Li coatings + triggered ELMs come closest to achieving externally controllable stationary D inventory, Z_{eff}
 - Use of triggered ELMs for particle control will be assessed in upgrade
- How do these parameters project to NSTX-U parameters?
 - Up to 5x longer pulse
 - Up to 2x higher NBI fueling
- How persistent is D pumping by Li?
 - Part of Milestone R(12-2)
 - Project deuterium pumping capabilities for NSTX-U using lithium coatings and cryo-pumping

Low-recycling conditions with lithium coatings last throughout NSTX discharges

- Heavily lithium coated, low-recycling, ELM-free discharges studied
 - Most thoroughly analyzed 2008 pre- to post-lithium discharges
- Peak D_α emission at outer divertor does not increase toward the end of the discharge
 - And in fact often decreases
 - Without lithium, recycling increases throughout shot
 - Inferred PFC particle recycling coefficient $R_p \sim$ constant



R Maingi, PRL 107, 145004 (2011)



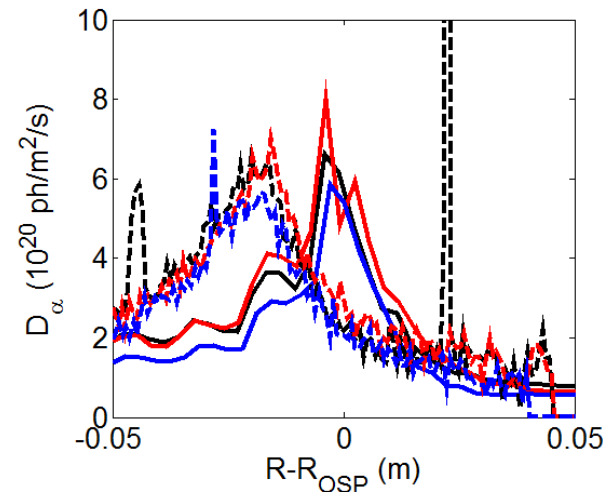
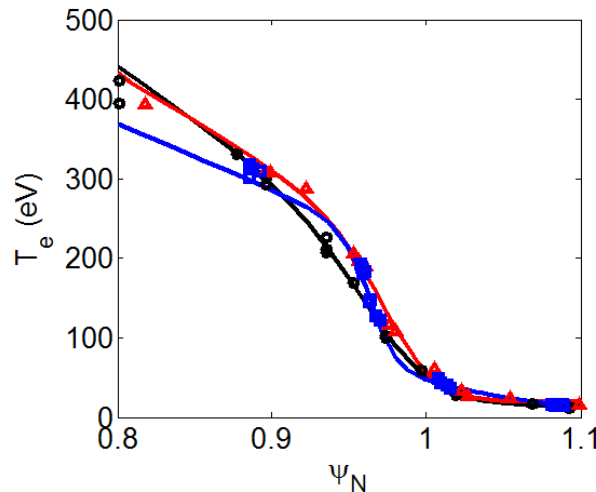
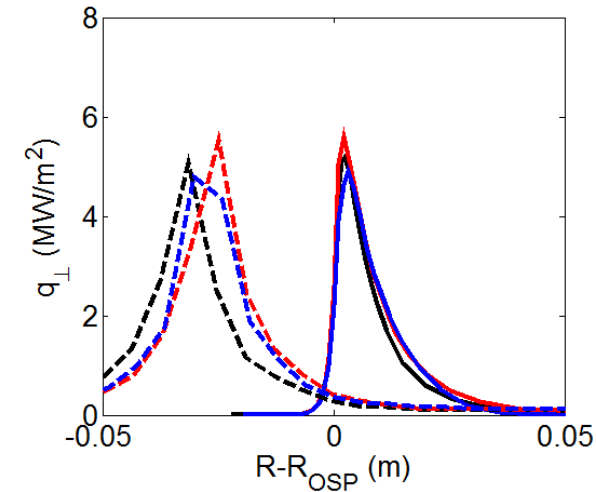
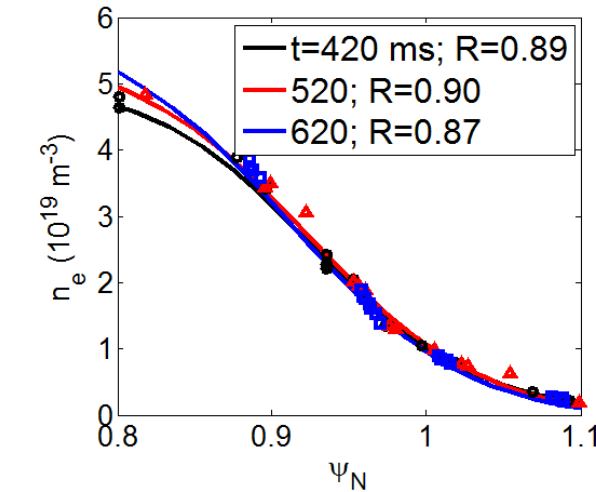
SOLPS modeling indicates recycling coefficient remains low throughout low- δ discharge

- Measurements show little change during shot
 - Points/dashed lines are measurements
 - SOL n_e , T_e , Peak heat flux, D_α all \sim constant

- Constraints in modeling*:
 - Fitted n , T profiles
 - Peak q_{div} (T_e^{sep})
 - Peak D_α (R_p)

- Inferred R_p remains low
 - 0.89, 0.90, 0.87
 - Without Li: $R_p=0.98$

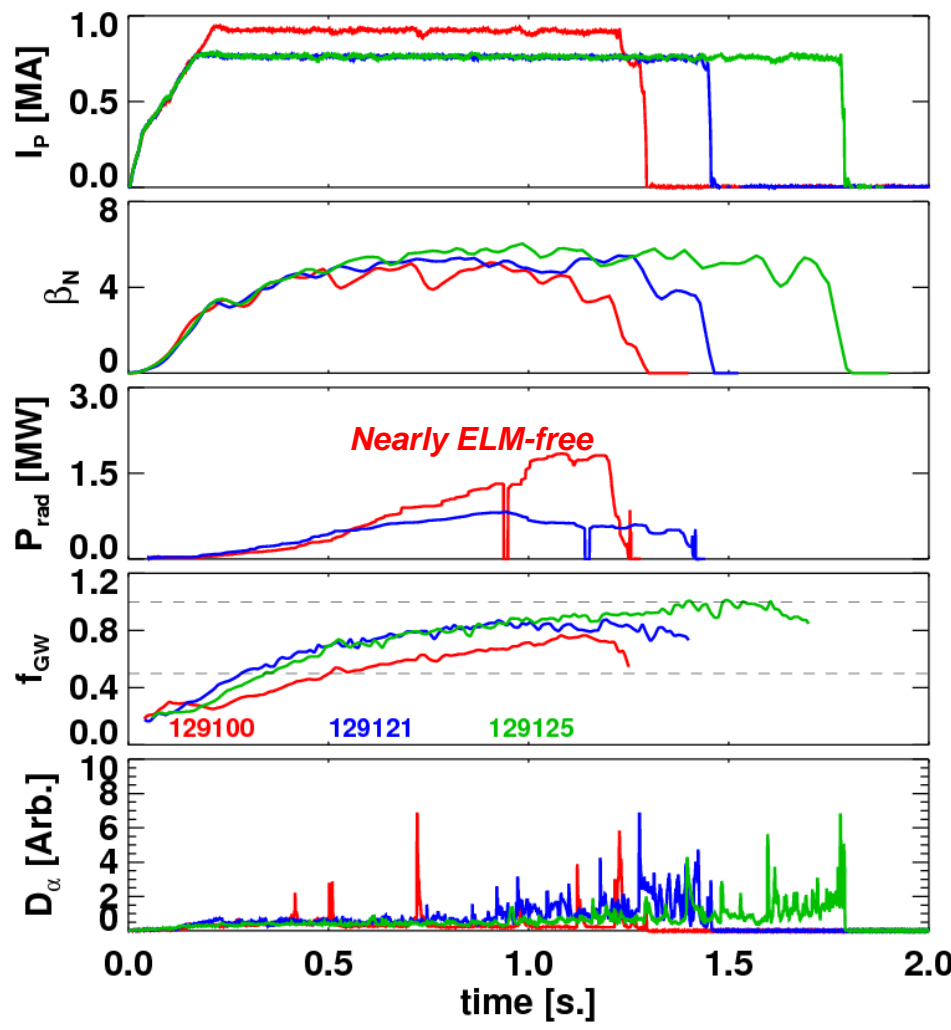
\Rightarrow Li pumping appears to persist over these pulse lengths (~ 1 s)



*J Canik, JNM 415, S409 (2011)

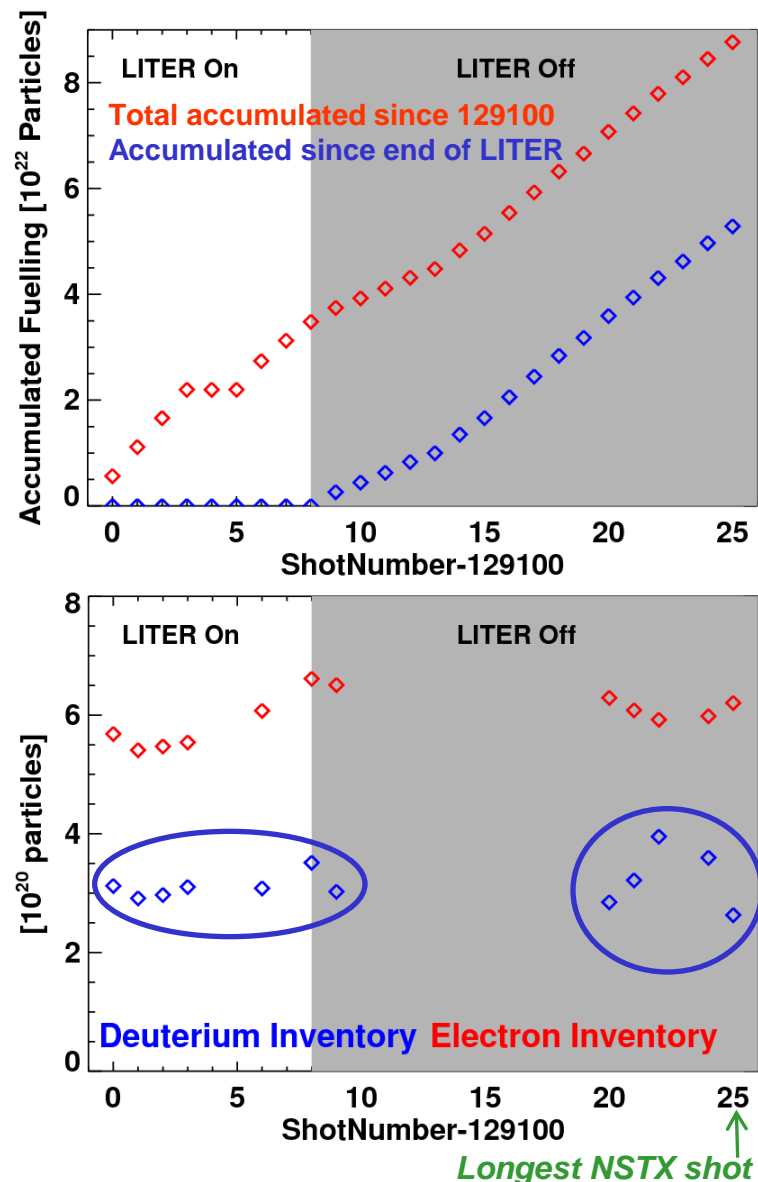
Long-pulse discharges following the shut-off of LITER maintained high performance while approaching stationary n_e

- **129100**: 900 kA shot just before LITER ran out.
- **129121**, **129125**: long pulse optimization sequence
 - **129125**: longest NSTX shot
- Density time trace (at high f_G) ~stationary after 1 s
 - No large $n=1$ MHD modes
 - Equivalent to beginning of an NSTX-U shot
- Without LITER, ELMs returned
 - Mostly small
 - Radiated power reduced



Experiments following LITER shut-off show D inventory control for many shots

- LITER operated for ~90 discharges prior to lithium running out
 - ~20 shots taken without LITER
 - Integrated discharge time ~25 s
 - Accumulated fuelling $\sim 5 \times 10^{22}$ particles
 - ~5 NSTX-U discharges (10 MW, 5s)
 - Plasma not held constant
 - He GDC performed between shots
 - Fairly constant D inventory maintained throughout sequence
- ⇒ May be possible to tailor lithium deposition to provide long-pulse pumping while maintaining ELMs for impurity control



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Next steps for particle control analysis during Upgrade outage

- Cryo-pumping design
 - Confirm plenum optimization using SOLPS (B2-EIRENE)
 - More comprehensive treatment of neutral transport (beyond first-flight)
 - Can treat radiative/detached divertor
 - Investigate design details of chosen plenum geometry
 - Is clearing area currently occupied by divertor tiles feasible?
 - Prepare for engineering design
- Lithium persistence for long-pulse (with ELMs)
 - Further modeling with 2D fluid codes (UEDGE/SOLPS/OEDGE)
 - Recycling analysis for high- δ , longer pulse ELM-free discharges
 - Analysis of long, ELMy discharge
 - Extrapolation to NSTX-U
 - Longer pulse, higher NBI particle input
- Begin studying compatibility/interaction of cryo and lithium pumping
 - Could lithium coat the pumping surface?
 - What plenum pressure can be achieved with SOL modified by lithium coatings (e.g., λ_q , P_{div}/P_{tot} , n_e)

Plans for years 1 and 2 of NSTX-U operation

- Validate physics design of cryo-pump
 - Measure plasma parameters at likely pump entrance location
 - Document Γ , T_e as I_p , P , flux expansion, etc are varied
 - Perform engineering design (begin during outage with incremental funding)
- Particle control with lithium coatings
 - Develop ELMy scenarios with lithium coatings
 - Assess ELM triggering with thick lithium coatings
 - Perform experiments with controlled scans lithium deposition amounts (including none), document recycling and ELM characteristics
 - Test passivation of lithium with D_2 glow for control of pumping properties
 - Optimize lithium deposition (ELMs vs. pumping), combine with impurity control techniques (snowflake, gas puff, etc.) towards long-pulse
 - Test persistence of lithium coatings
 - Measure recycling characteristics as power, ion flux, pulse length are varied
 - Use rapid SGI gas pulses to measure SOL pump-out vs time within shot
 - Later stages: measure impurity behavior with Li on Mo tiles

Long term plans (NSTX-U years 3-5)

- Install cryo-pump as part of long-pulse divertor
 - Present plan: cryo in upper divertor, liquid lithium on lower divertor
 - Contingent on available resources
- Explore performance of pumping system
 - Document pumping rates as P , I_p , R_{OSP} are varied
 - Test pumping of high flux expansion divertor
 - Assess n/n_G achievable with pumping in various conditions, and develop low-density, high-performance scenario
 - Develop long-pulse, density controlled plasmas for range of n/n_G
 - Compare to lithium-based pumping

Summary

- A promising cryo-pump design point has been identified that is compatible with standard and snowflake divertors
 - Based on semi-analytic pumping model
 - Divertor profile projections based on NSTX data
- D inventory control using lithium coatings appears to extrapolate to 5s NSTX-U discharges
 - Estimated recycling coefficient remains low at $R \sim 0.9$ for 1 s discharges
 - D inventory in ELMy plasmas controlled ~ 25 s following LITER shutoff
- Studying compatibility and interplay between cryo-pumping and lithium is an important research goal for NSTX-U

BACKUP

Analytic pumping model* used to optimize pumping chamber

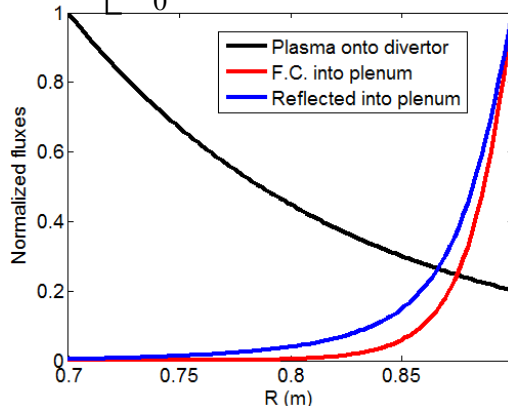
- Uses first-flight model for neutral flux into pump plenum
- Requires knowledge of divertor plasma profiles
- Validated against DIII-D experiments

$$P_{pl} = \frac{I_0}{S + C} \Rightarrow P_{pl} = \frac{I_0}{S + C} \frac{C}{C_{eff}} \quad \leftarrow \text{Plenum pressure corrected for penetration of neutrals through long duct (verified using EIRENE)}$$

$$I_0 = \int_{R_{min}}^{R_{max}} \Gamma_0(r) F(r) T(r) 2\pi R_m dr \quad \leftarrow \text{Neutral current into plenum}$$

$$F(r) = \frac{1 - \cos \Theta_{max}(r)}{2}; \Theta_{max}(r) = \tan^{-1} \left(\frac{g}{R - R_{ent}} \right) \quad \leftarrow \text{Solid angle of plenum entrance}$$

$$T(r) = \exp \left[\frac{-1}{v_0} \int_R^{R_{max}} n_e(r) \langle \sigma v \rangle_{EH}(r) dr \right] \quad \leftarrow \text{Transmission of neutrals through plasma}$$



**Origin of neutrals making it into plenum tends to be localized to near-entrance region
Dominantly due to solid angle factor**

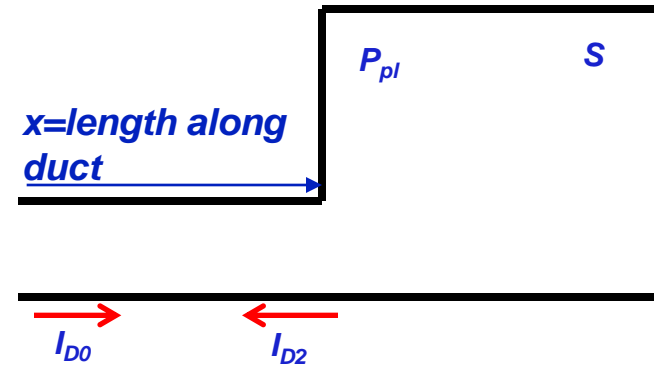
*R. Maingi, Nucl. Fus. 39 (1999) 1187

Penetration of neutrals through a long throat is accounted for to correct the conductance

- $I_{D0} = I_{D0}(x)$ = current of “fast” atomic deuterium entering from plasma

If fast atoms are turned into thermal molecules on collision with the wall, then:

$I_{D0}(x) = I_{D0}(0) \cdot F(x)/F(0)$, where F is the solid angle factor evaluated along x



- I_{D2} = current of thermal molecules leaving
- I_{D2} = volume integral of sources (I_{D0}), sinks ($P_{pl}S$)

$$\Rightarrow I_{D2}(x) = I_{D0}(x) - P_{pl}S$$

- Pressure is $\Delta P = \int_0^h I(x)\sigma(x)dx$, $\sigma = \frac{3}{4\bar{v}} \frac{H}{A^2}$, $\frac{1}{C} = \int_0^h \sigma(x)dx$
- So plenum pressure is

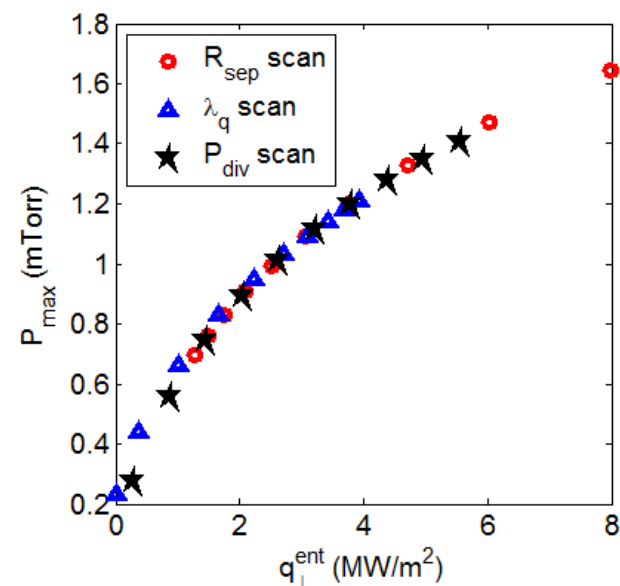
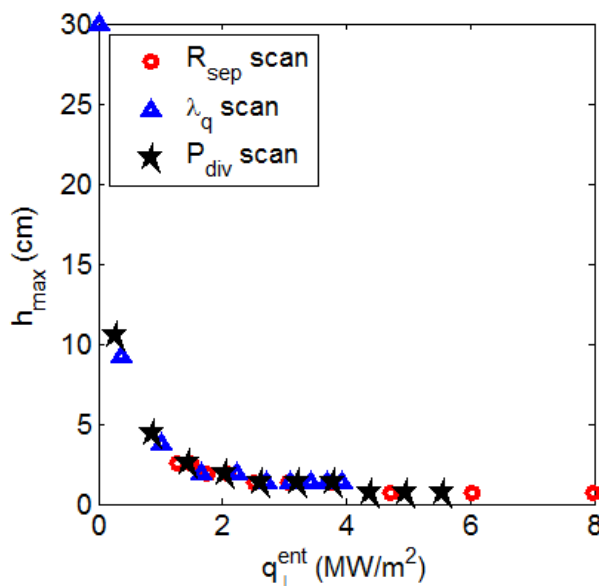
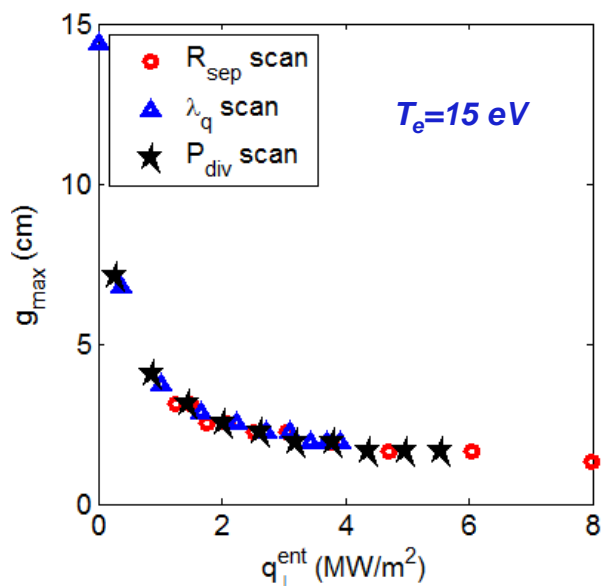
$$P_{pl} = \int_0^h I_{D2}(x)\sigma(x)dx = \int_0^h I_{D0}(x)\sigma(x)dx - \int_0^h P_{pl}S\sigma(x)dx$$

$$= I_{D0}(0) \int_0^h \frac{F(x)}{F(0)} \sigma(x)dx - \frac{P_{pl}S}{C} = \frac{I_{D0}(0)}{C_{eff}} - \frac{P_{pl}S}{C} = \frac{I_{D0}(0)}{S+C} \frac{C}{C_{eff}}$$

$$C_{eff} = \int_0^h \frac{F(x)}{F(0)} \sigma(x)dx$$

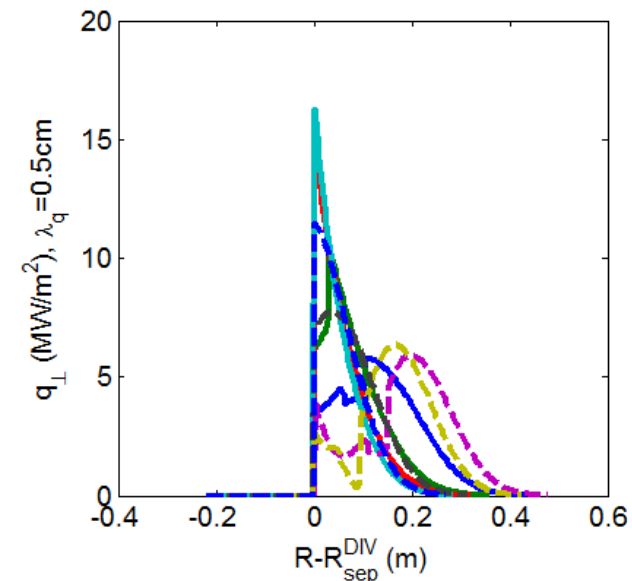
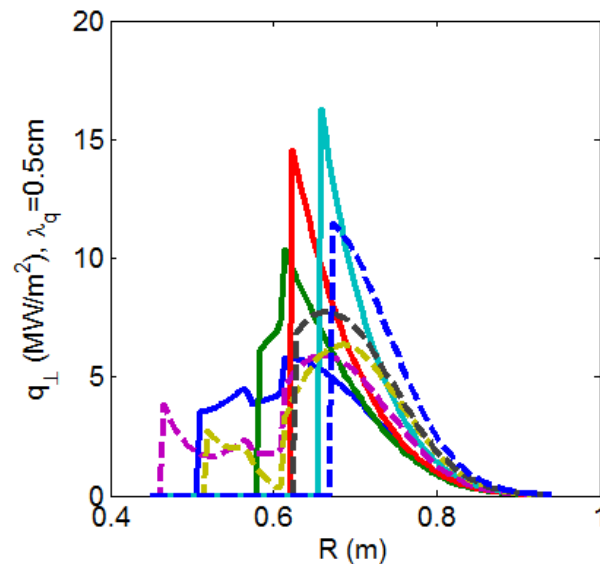
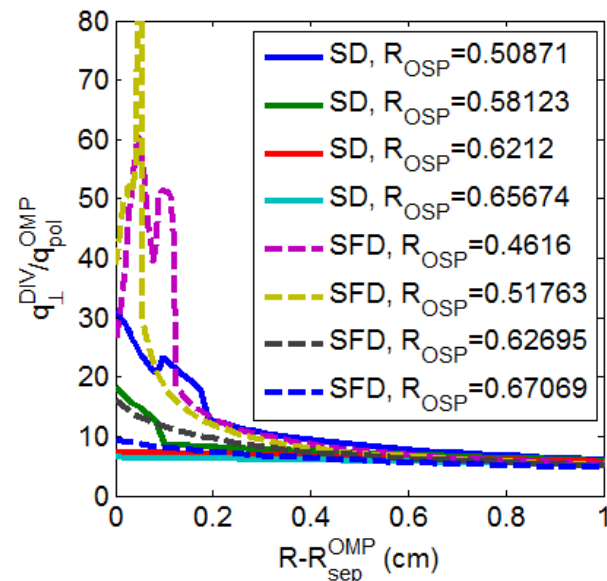
For given pump entrance position, heat flux at pump entrance orders the “optimal” geometry parameters

- Optimal throat height/length depend mainly on heat flux near entrance
 - Doesn't matter if it's varied by moving the OSP, changing flux expansion, or changing total power
 - T_e affects maximum pressure achievable, but only weakly affects g/h
- Optimizing for $P=0.8\text{mTorr}$ at $T_e=15.0\text{ eV}$ gives $g\sim 2.5\text{ cm}$, $h\sim 2\text{ cm}$ at $q\sim 2\text{MW/m}^2$



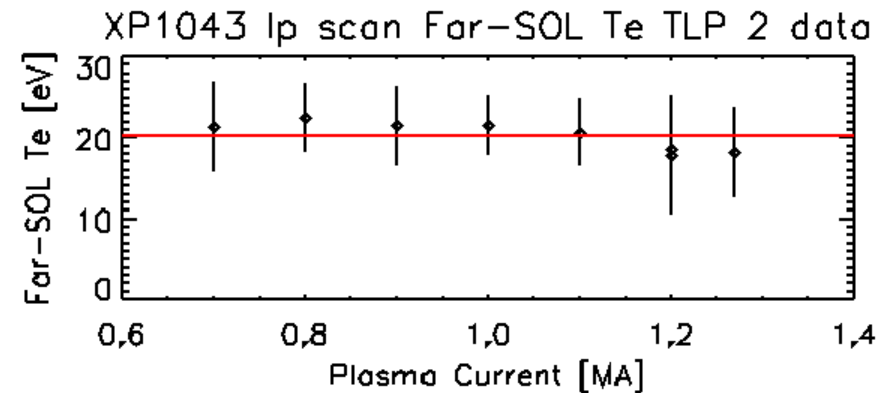
Projecting heat flux profiles

- Exponential poloidal heat flux profile imposed at midplane
 - $P=5$ MW (e.g., 1/2 of 10 MW goes to outer divertor)
 - $\lambda_q^{\text{OMP}} \sim 0.3\text{-}2.0$ cm
- Mapped along field lines to divertor
 - Total geometric heat flux reduction factor shown on left
 - Example heat flux profiles showing for $\lambda_q^{\text{OMP}}=5\text{mm}$
 - Heat flux high at $R=0.7$, significantly lower at 0.8

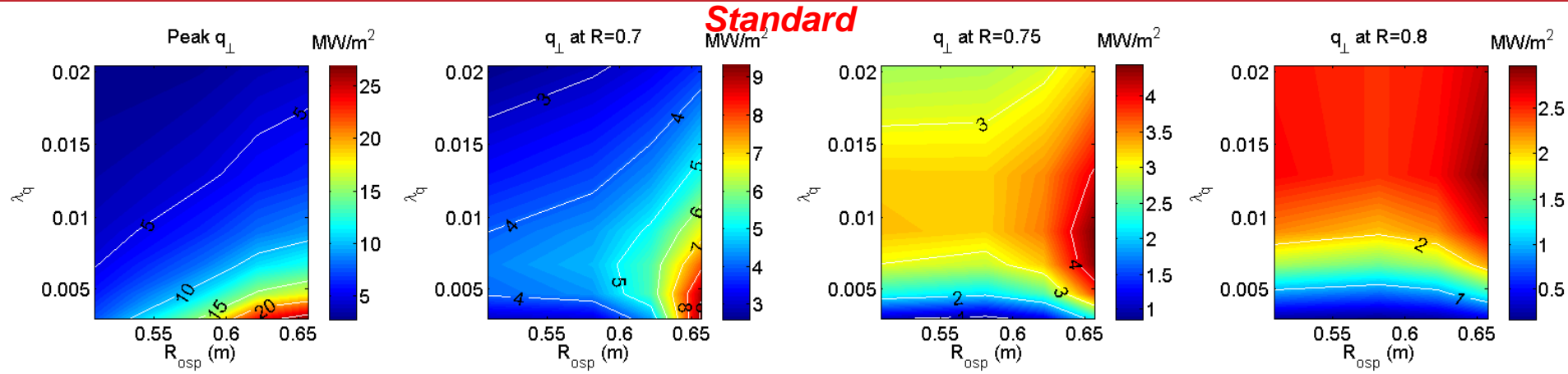


Far-SOL T_e does not vary significantly with I_p

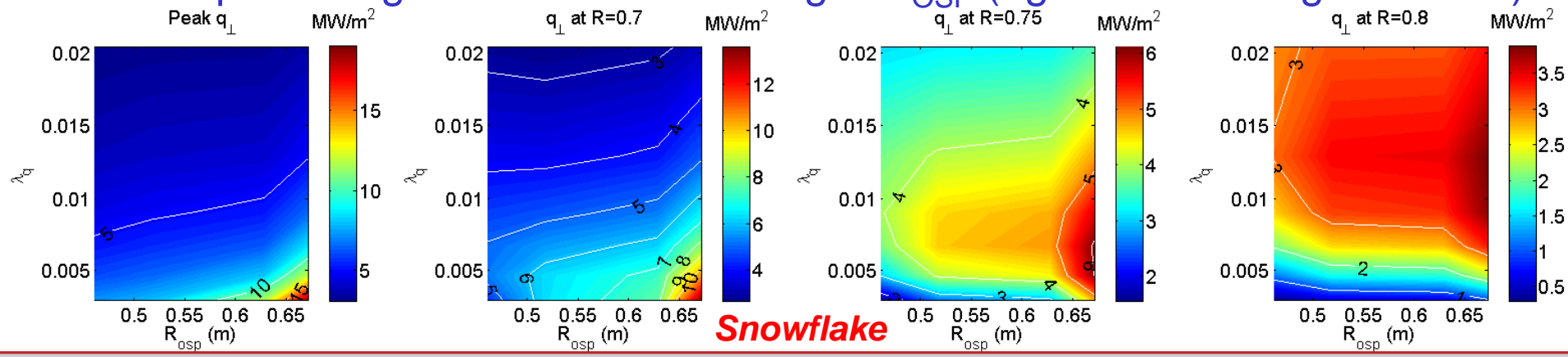
- Triple probe used to determine T_e (avoids turbulence issue)
- “Far-SOL” defined as beyond second sep.
- Variation in T_e not statistically significant



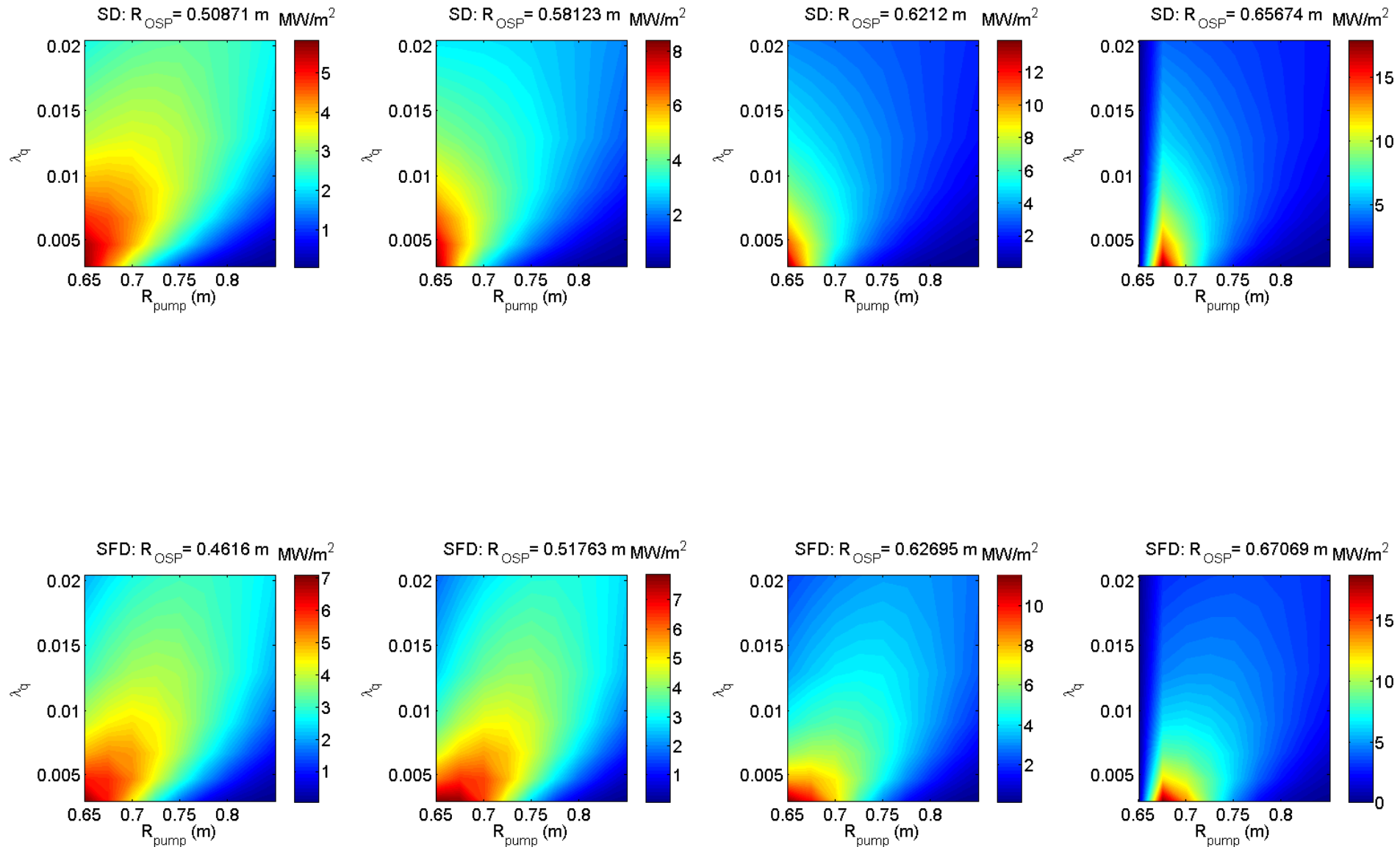
Heat flux projections show plenum entrance at $R \sim 0.7-0.75$ m likely to provide sufficient pumping



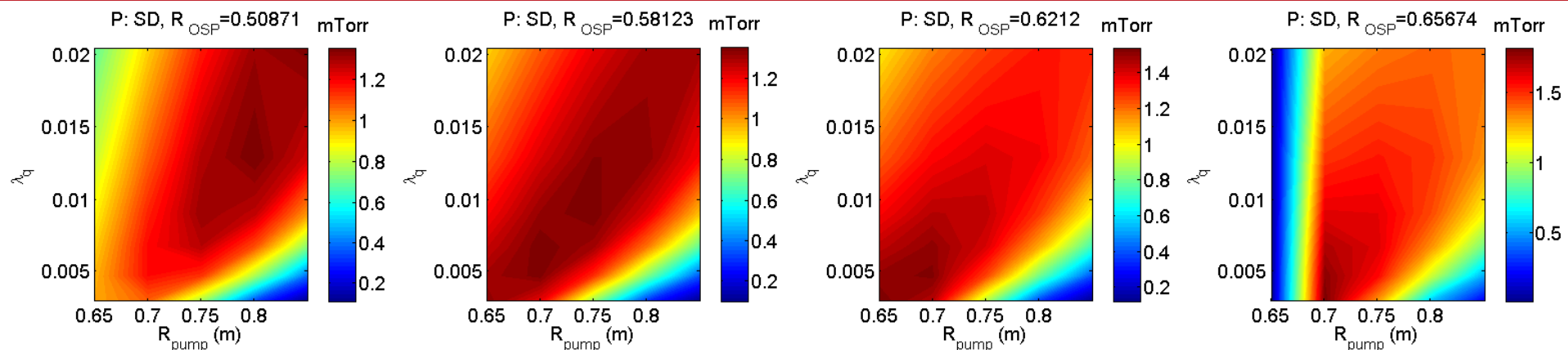
- Power handling: peak heat flux < 10 MW/m²
 - Restricts R_{OSP} for narrow SOL (wider range for SFD)
- Pumping: q_{\perp} entrance $> \sim 2$ MW/m²
 - Requires larger SOL widths for larger R_{OSP} (again wider range for SFD)



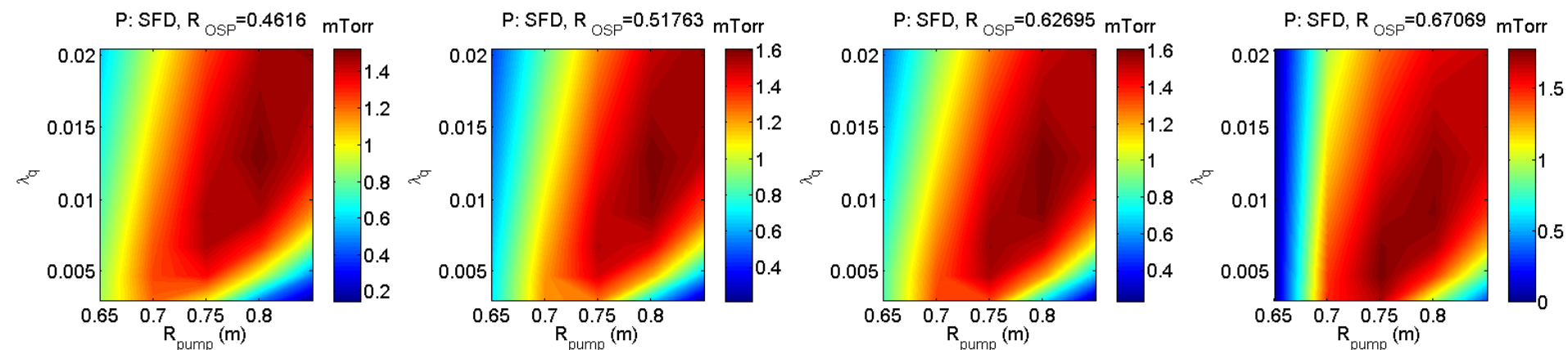
Heat flux at potential plenum entrances for 8 equilibria



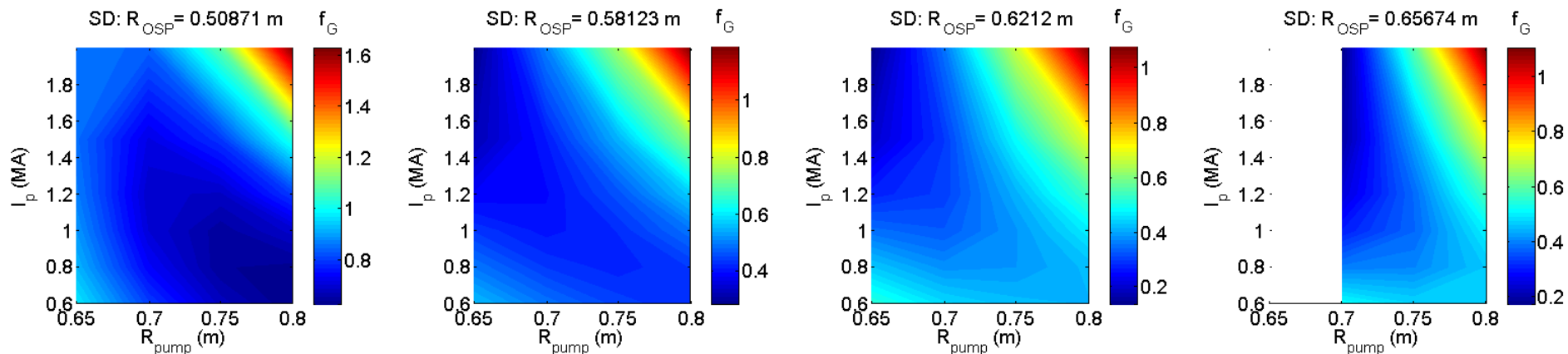
Projections show plenum entrance at $R=0.72$ can give >1 mTorr for wide range of SOL width, equilibria



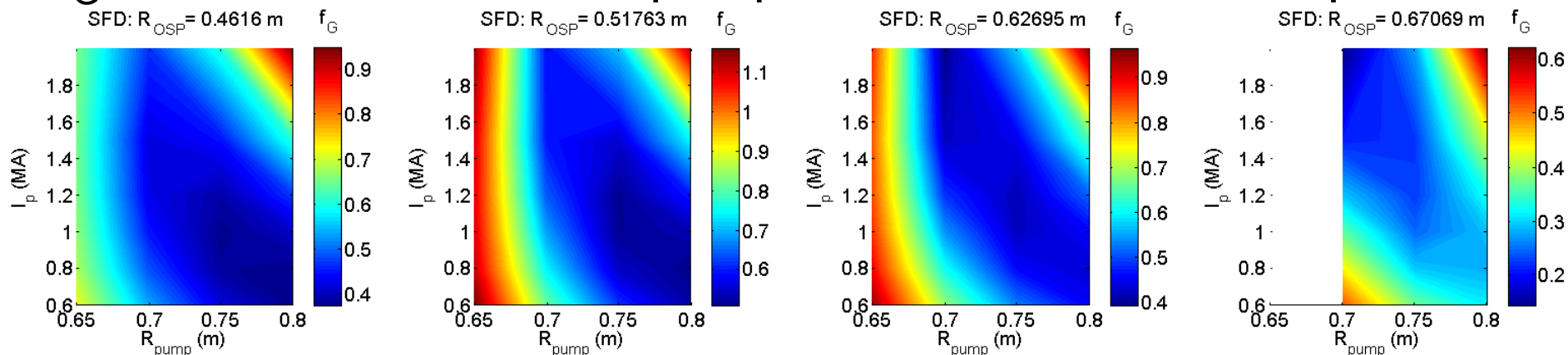
- Heat flux profiles, T_e^{div} , and optimized entrance parameters used in analytic model for plenum pressure
- Optimizing position for narrowest SOL gives $R_{pump} \sim 0.72$
 - Narrow SOL gives least flexibility in moving R_{OSP} to improve pumping



$R_{\text{pump}}=0.72$ supports low Greenwald fraction for range of I_p , equilibria

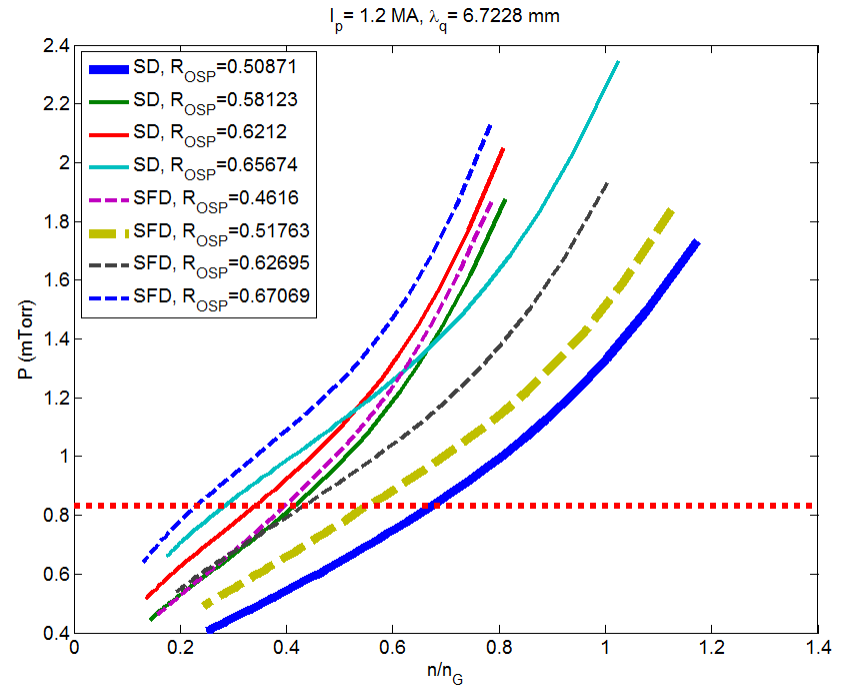
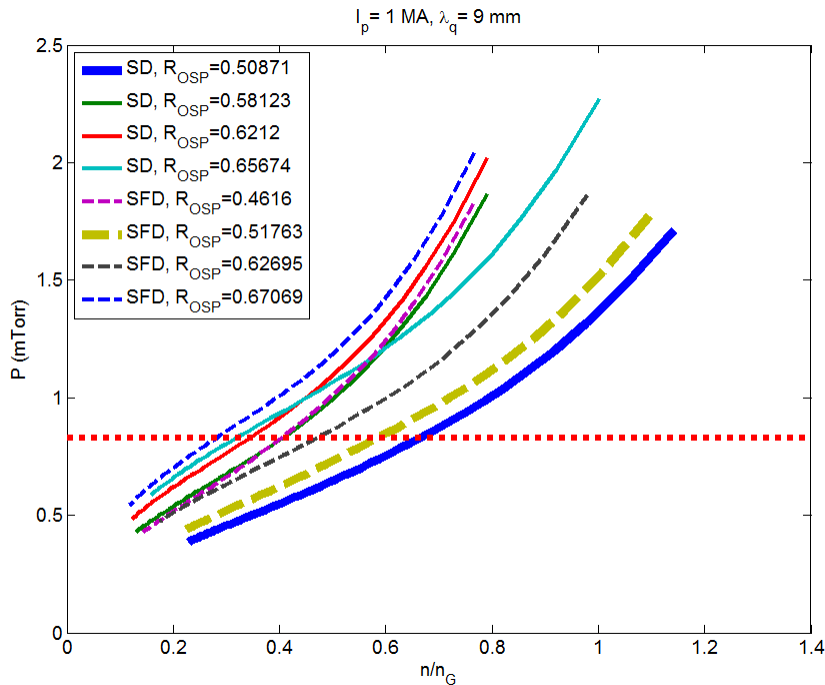


- $q_{\parallel}^{\text{sep}}, T_e^{\text{div}}$ used in modified 2-pt model used to estimate n_e^{sep}
 - $q_{\parallel}^{\text{sep}}$ from I_p scaling, T_e^{div} varied
- $n_e/n_e^{\text{sep}} \sim 3$ assumed to estimate f_G
- f_G shown is that at which pumped flux balances NBI input



Estimating achievable n/n_G

- n/n_G varied by scanning T_e^{div}
- To pump beams, need $P \sim 0.8$ mTorr
- f_G shown is where the pumping balances beam input
 - Minimum achievable n_e -> could puff to increase



The Basic Two-Point Model

- Begins from the fluid equations and simplifies...
- Provides simple relations for upstream and target (PFC) plasma parameters
- Varying levels of complexity can be implemented
 - Fluid reconstruction via generalized 2-point (e.g. OSM/OEDGE code)
 - Coupling with Monte Carlo neutrals and impurities (e.g. DEGAS 2/EIRENE/DIVIMP)
- **Start with the basics**

Assume:

$$T_e = T_i \quad \& \quad p = p_e + p_i$$

$$\frac{d}{dx} \left[\left(\frac{1}{2} m_i v^2 + 5kT \right) n v - \kappa_{0e} T_e^{5/2} \frac{dT_e}{dx} \right] = Q_R + Q_E$$

Assume:

**Conduction Dominates
Neglect Sources**

$$\frac{d}{dx} \left[-\kappa_{0e} T_e^{5/2} \frac{dT_e}{dx} \right] = 0$$



$$T_u = \left[T_t^{7/2} + \frac{7}{4\kappa_{0e}} q_t L \right]^{2/7}$$

Simple Extensions Attempt to Capture More Physics

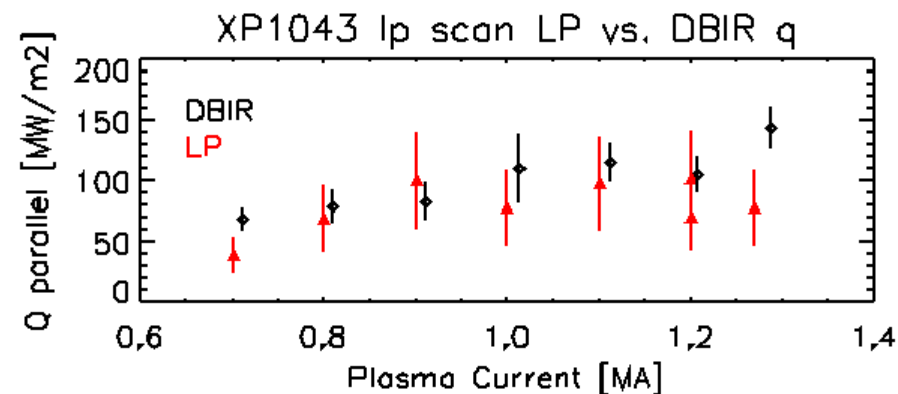
- Volumetric loss terms can be included via f_{power} term
- Term can be estimated with interpretative modeling in lieu of better div. Bolom. Coverage
- Comparison of nominal LP and DBIR results are encouraging
- Two values of f_{power} used following: 0 and 0.5

Radiation and charge-exchange

$$q_{rad} + q_{cx} = f_{power} q_0$$

$$(1 - f_{power}) q_0 = q_t = \gamma n_t c_{st} k T_t$$

$$\frac{T_t}{T_u} \propto (1 - f_{power})^2$$



Updated Upstream Density

- Force balance in the ST requires modification to 2-PM
 - Typical formulation assumes “straight” flux tubes
 - 1.5m OMP vs. 0.5m target results in significant variation
- Flux-tube definition allows conversion of magnetic field to area
- Not yet consistently applied everywhere in calculations

$$F_u = F_t$$

$$P_u A_u = P_t A_t$$

$$N_u T_u = N_t T_t (1 + M^2) \frac{A_t}{A_u}$$

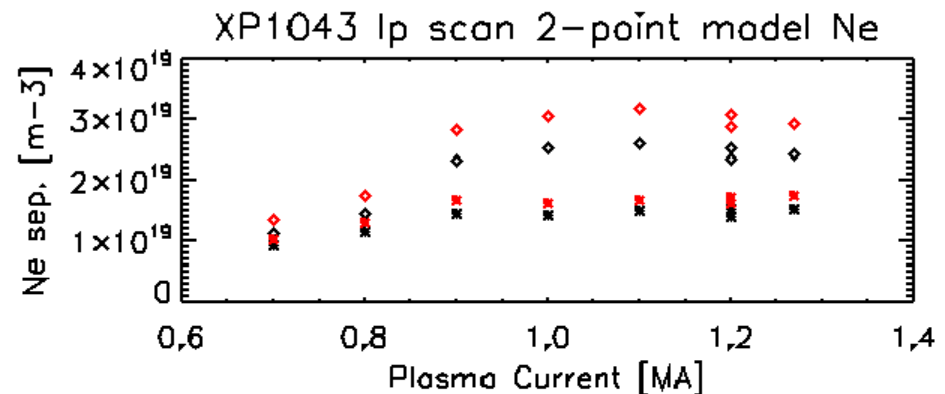
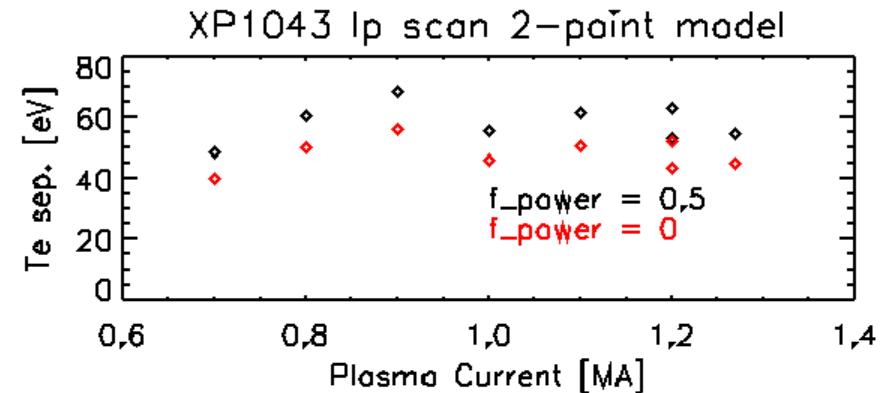
$$M = v/c_s \geq 1 \quad \text{Mach No. at sheath}$$

$$B A = \Psi_0 = \text{const.} \rightarrow \frac{A_t}{A_u} = \frac{B_u}{B_t}$$

$$N_u = \frac{N_t T_t (1 + M^2)}{T_u} \frac{B_u}{B_t}$$

Upstream Quantities Determined via 2-Point Model

- Parallel connection length calculated from EFIT02
 - q_{peak} used to locate nominal Ψ_N value for integration
 - Solution not sensitive to variance in length (robust model from target)
- Interpolated MPTS density at the upstream temperature shown for comparison
- Uncertainty not yet propagated in calculations to determine significance in discrepancy



Wall fuel uptake is ~zero during long-pulse H-modes in pumped machines

- FY09 Joule Milestone Report
- Close balance is observed between particle input (beams+puff) and pumping
 - True for both DIII-D and C-Mod
 - Motivates pumping the beam input as figure of merit

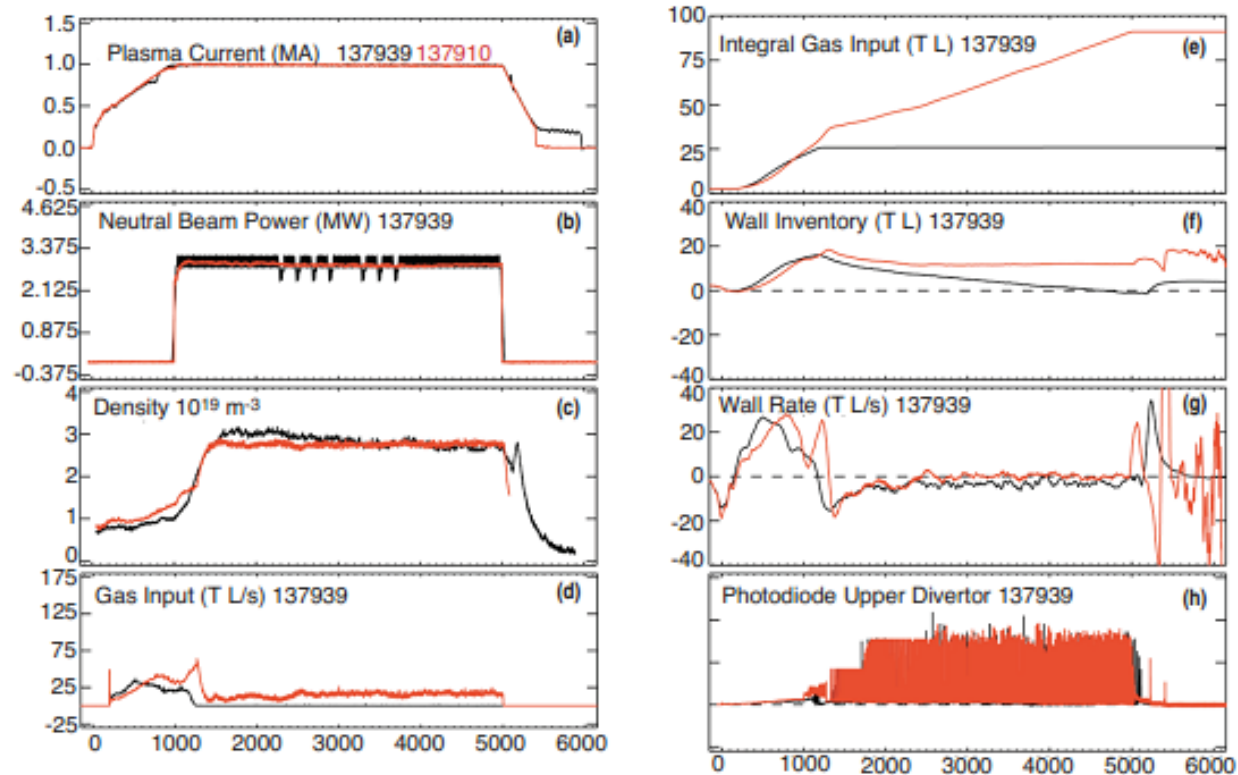


Fig. 2. ECH heated (red) H-mode plasma is compared with a NBI (black) discharge. The plasma current, neutral beam or ECH power, electron density, gas input, integral of the gas input; wall inventory and wall rate from the dynamic particle balance; along with the photodiode signal are compared with a neutral beam heated H-mode DIII-D shot. Note that in both cases, the wall flux is quite large in the L-mode period, but during the ELMing H-mode, the wall flux is very close to zero.