

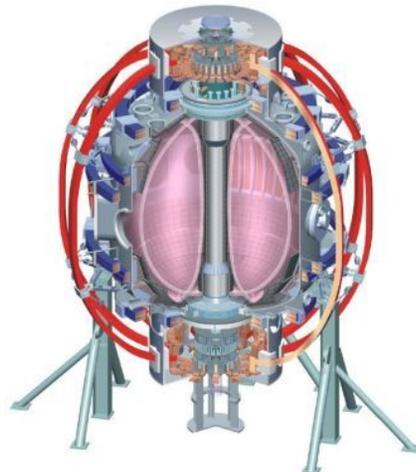
# Recent progress in transport and turbulence research at NSTX

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# NSTX accesses a broad range of parameter space to address many turbulence and transport issues

- Low aspect ratio, high  $\beta$  (high  $\nabla P$ ), and strong  $E \times B$  flow shear in NSTX stabilize “traditional” electrostatic low-k turbulence, such as Ion Temperature Gradient (ITG) and Trapped Electron Mode (TEM)
- Ion thermal transport is close to neoclassical (collisional) in NSTX H-modes  
→ electron and ion thermal transport is largely decoupled
- With relatively small magnetic field Electron Temperature Gradient (ETG) turbulence can cause significant transport at  $\rho_e$ -scales (high-k)  
→ desire for high-k turbulence measurements to correlate with electron transport and validate with nonlinear ETG simulations
- Achievable range of  $\beta_T \leq 40\%$  can also lead to significant EM contribution  
→ micro-tearing turbulence can cause electron transport through magnetic flutter
- Must still consider low-k turbulence, for example in L-mode (ITG/TEM) and in H-mode pedestal region (ITG/TEM/KBM)  
→ desire for large scale (low-k) turbulence measurements such as BES

# Overview

- First nonlinear gyrokinetic simulations of micro-tearing turbulence for “high beta” NSTX H-mode plasmas
- Parametric dependence of high-k turbulence measured by a microwave scattering diagnostic in “low beta” plasmas
- First low-k turbulence measurements from a newly implemented BES diagnostic
- Summary

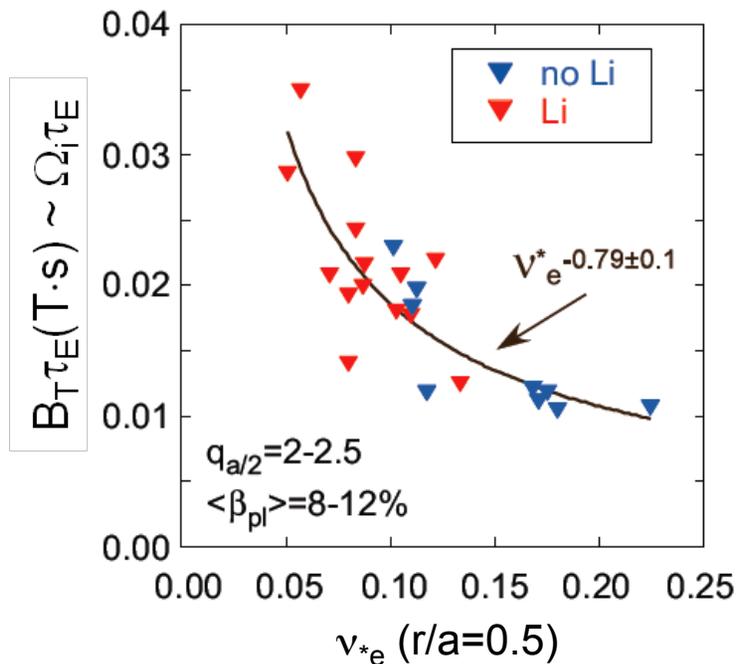
See website for recent APS presentations:

[http://nstx.pppl.gov/DragNDrop/Scientific\\_Conferences/APS/APS-DPP\\_11/](http://nstx.pppl.gov/DragNDrop/Scientific_Conferences/APS/APS-DPP_11/)

# NSTX H-mode thermal confinement scaling shows strong dependence on collisionality

- Dimensionless collisionality ( $\nu_*$ ) scan holding  $q$ ,  $\beta$ ,  $\rho_* = \rho_i/a$  constant **without** or **with** Lithium wall conditioning

$$\nu_{*e} = \frac{\nu_{ei} / \epsilon}{\epsilon^{1/2} \nu_{Te} / qR}$$



← NSTX (Kaye et al., Nucl. Fusion 2007; APS 2011)

$$\Omega \tau_E^{th} \sim \nu_{*e}^{-0.8}$$

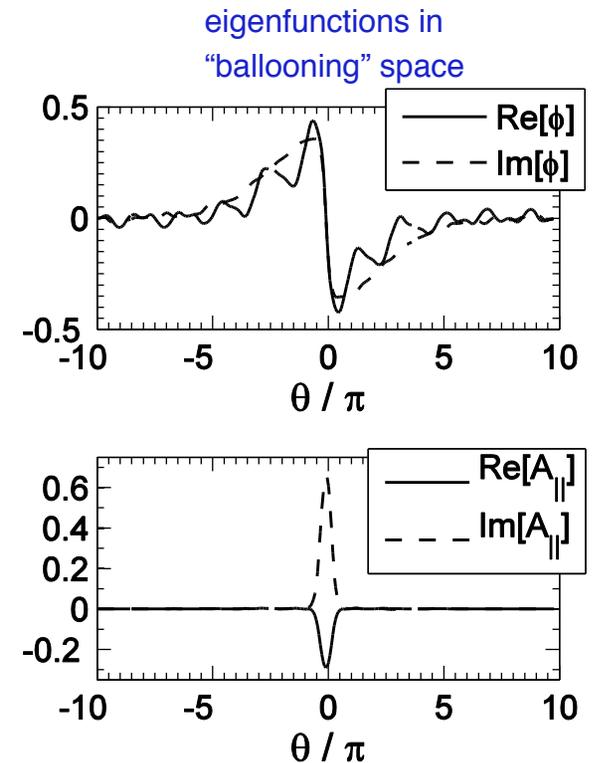
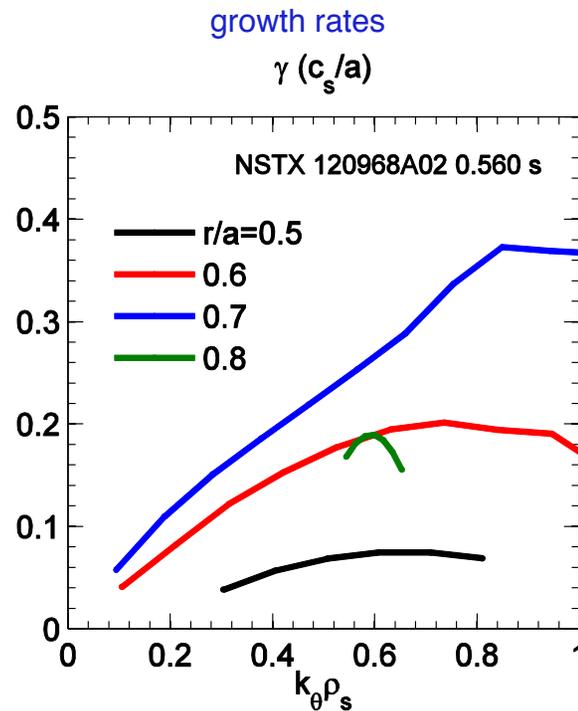
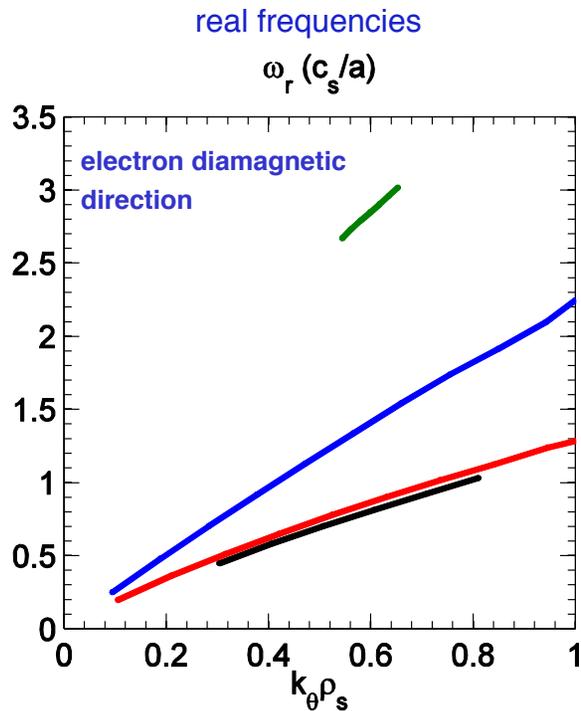
ITER (PIPB, Doyle et al., Nucl. Fusion 2007)

$$\Omega \tau_E^{th,04(2)} \sim \nu_{*e}^{-0.2}$$

- Ion transport is neoclassical, consistent with strong toroidal flow and flow shear
- What is the cause of anomalous electron thermal transport?
- Will favorable  $\tau_E$  scaling hold at lower  $\nu_*$  envisioned for next generation ST (high heat flux, CTF, ...)?

# Microtearing modes found to be unstable in many high $v_*$ discharges

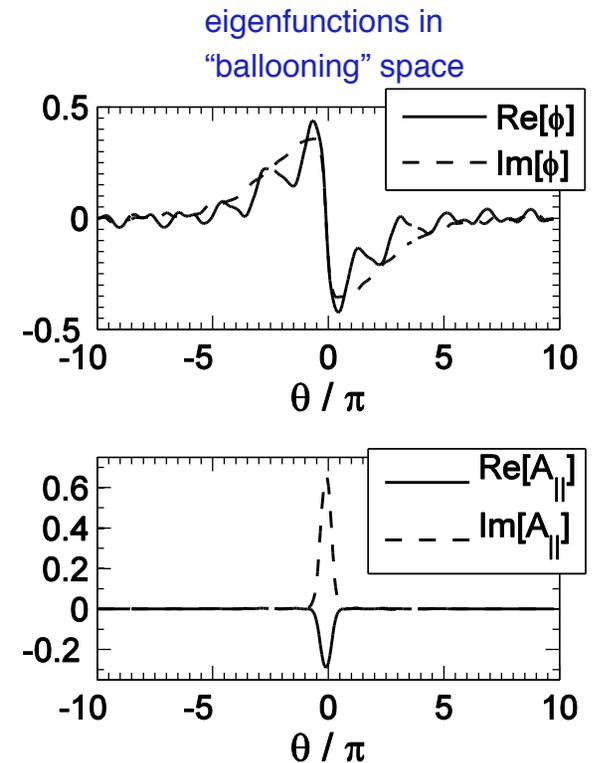
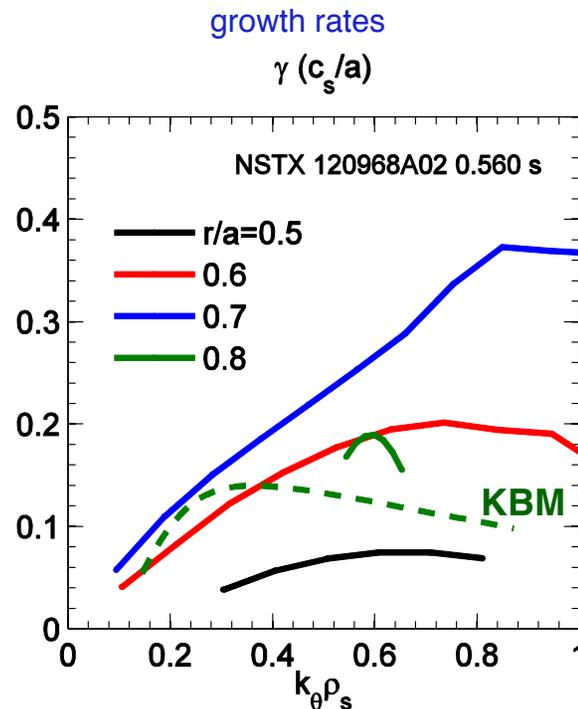
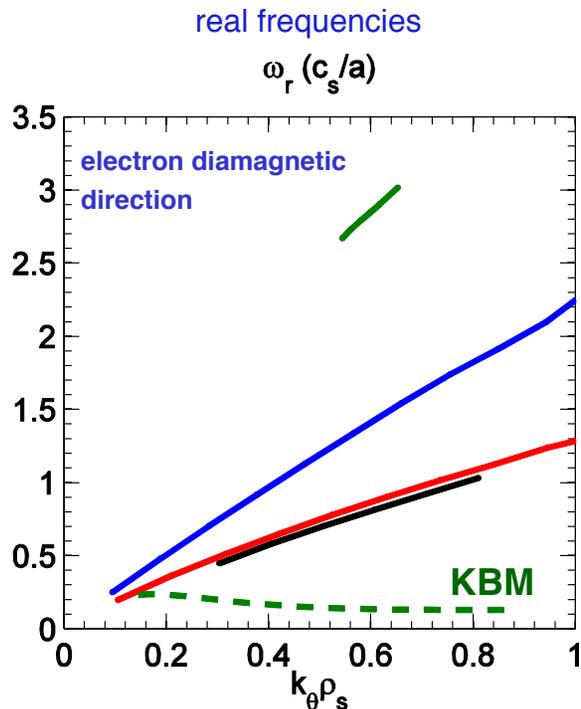
- Microtearing dominates over  $r/a=0.5-0.8$ ,  $k_\theta \rho_s < 1$  ( $n \approx 5-70$ )
- Real frequencies in electron diamagnetic direction,  $\omega \approx \omega_{*e} = (k_\theta \rho_s) \cdot (a/L_n + a/L_{Te}) \cdot (c_s/a)$
- ETG mostly stable due to larger  $Z_{\text{eff}} \approx 3$ ,  $(R/L_{Te})_{\text{crit,ETG}} \sim (1 + Z_{\text{eff}} T_e/T_i)$



Linear GYRO simulations [Candy & Waltz, Phys. Rev. Lett. (2003); <https://fusion.gat.com/theory/Gyro>]  
with kinetic ions and electrons, fully electromagnetic, collisions, local general equilibrium

# Microtearing modes found to be unstable in many high $v_*$ discharges

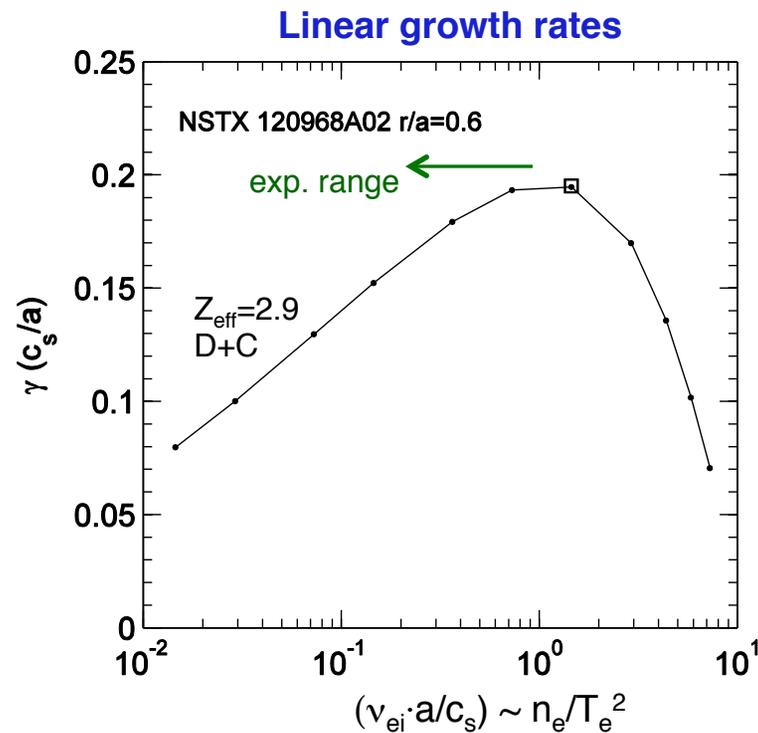
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- KBM competes farther out ( $r/a \geq 0.8$ ) where  $\alpha_{\text{MHD}} = -q^2 R \beta'$  much larger (larger  $q$ ,  $a/L_n$ )



Linear GYRO simulations [Candy & Waltz, Phys. Rev. Lett. (2003); <https://fusion.gat.com/theory/Gyro>]  
with kinetic ions and electrons, fully electromagnetic, collisions, local general equilibrium

# A distinguishing feature of the microtearing mode is the non-monotonic dependence on $v^{e/i}/\omega$

- Peak  $\gamma$  occurs for  $v^{e/i}/\omega = Z_{\text{eff}} \cdot v_{ei}/\omega \sim 1-6$ , similar to slab calculations (Gladd et al., 1980)
- $\gamma$  decreases with  $v_e$  in experimental range  $\rightarrow$  **qualitatively consistent with observed confinement scaling**



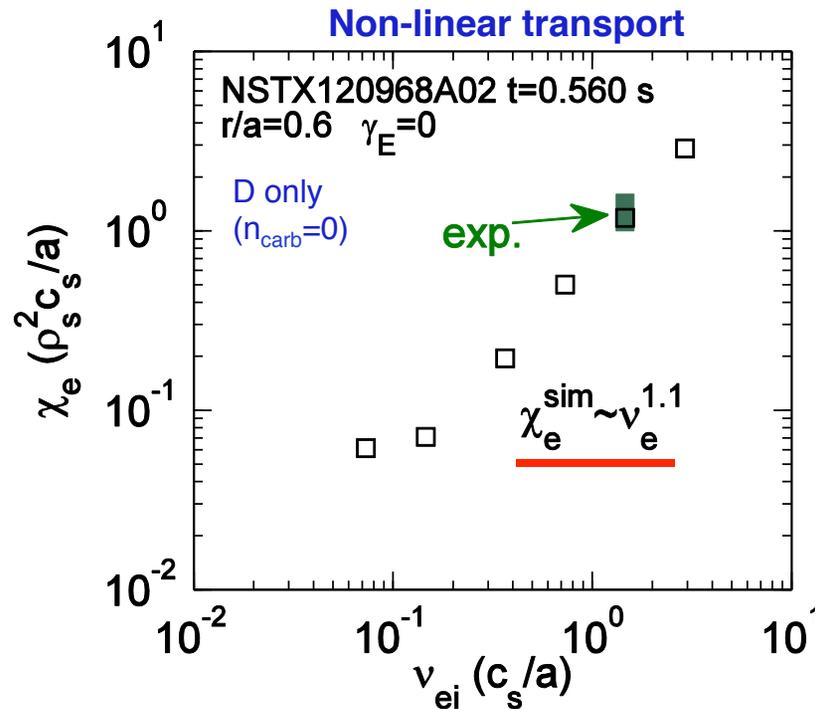
$$v^{e/i} = Z_{\text{eff}} v_{ei} \propto Z_{\text{eff}} \frac{n_e}{T_e^{3/2}}$$

\* Guttenfelder et al., *Scaling of linear microtearing stability for a high collisionality NSTX discharge*, submitted to Phys. Plasmas (Oct, 2011)

# Predicted nonlinear transport comparable to experiment, scales with $\nu_e$ similar to experimental confinement scaling\*

Nonlinear GYRO simulations with: local general equilibrium, kinetic ions and electrons, collisions, electromagnetic

Very fine radial resolution,  $\Delta x = 0.2\rho_s$  required to resolve rational surfaces



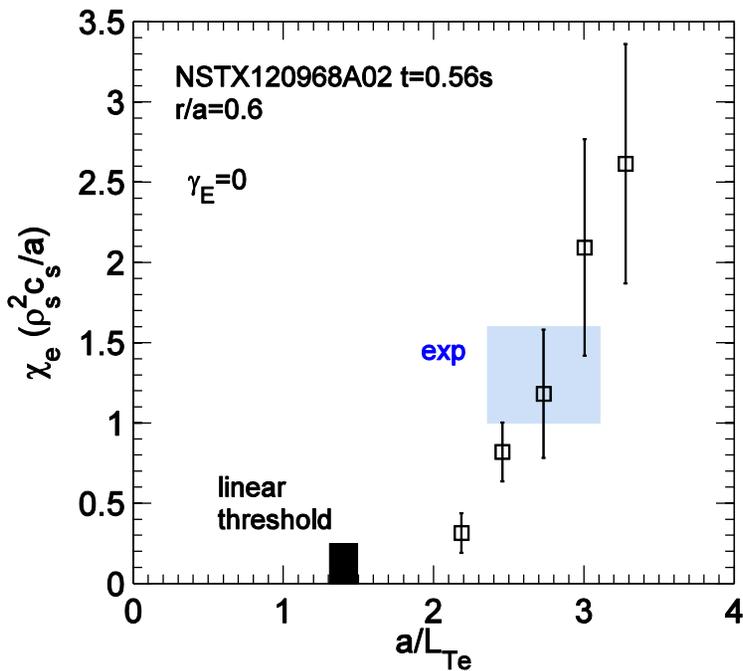
NSTX experimental scaling  
 $\Omega_i \tau_E \sim \nu_*^{-0.95}$

- As transport drops,  $a/L_{Te}$  will increase (for fixed heat flux), at some point ETG (TEM?) should become important
- This transition likely to determine limit of “favorable”  $\nu_*$  scaling

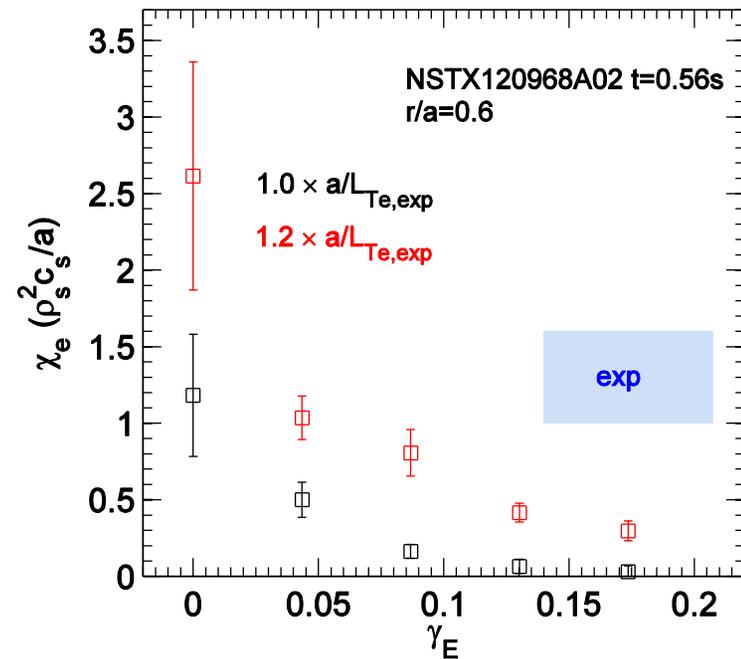
\* Guttenfelder et al., *APS-DPP invited talk T12.06, Salt Lake City (2011)*

# Predicted transport “stiff” with $\nabla T_e$ , susceptible to suppression via $E \times B$ shear

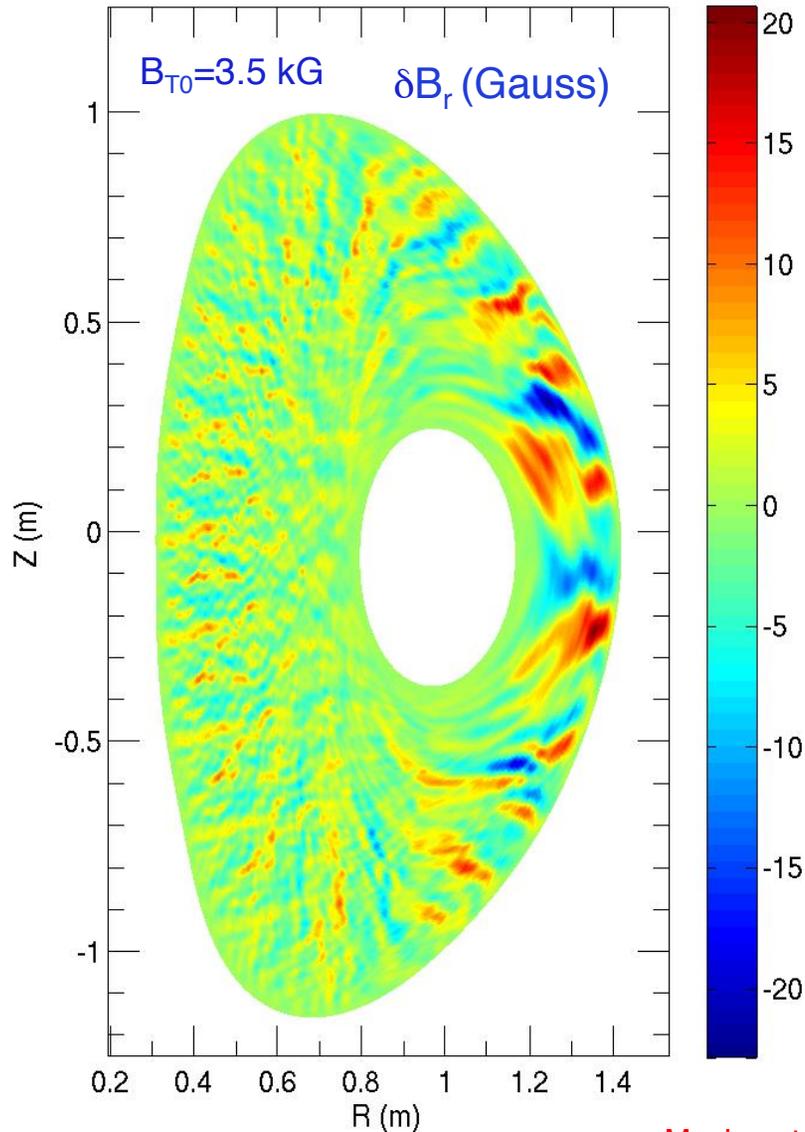
- Complicates simple interpretation from  $\chi_{e,\text{sim}} \sim v_e^{1.1}$  scaling
- Useful to characterize scaling of threshold gradient
- Transport reduced when increasing  $\gamma_E$  to local experimental value, **partially recovered with increase in  $\nabla T_e$**



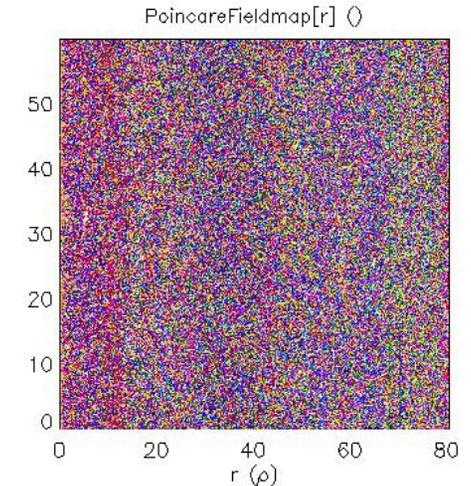
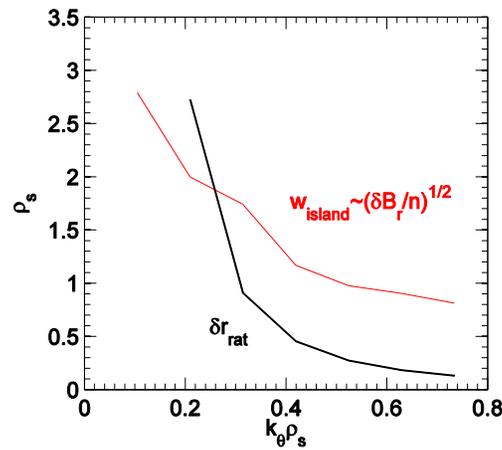
Error bars represent rms variation in transport, not standard error



# ~98% of transport due to magnetic “flutter” contribution



- Large magnetic perturbations,  $\delta B_r/B \approx 0.15\%$
- $w_{\text{island}}(n) > \delta r_{\text{rat}}(n)$ , island overlap  $\rightarrow$  **perturbed field line trajectories are stochastic\***
- $\chi_{e,EM} = 1.25 \rho_s c_s^2/a$  close to *collisionless* Rechester-Rosenbluth\* ( $\lambda_{\text{mfp}} = 12 \text{ m}$ ,  $L_c \approx 2.5 \text{ m}$ )



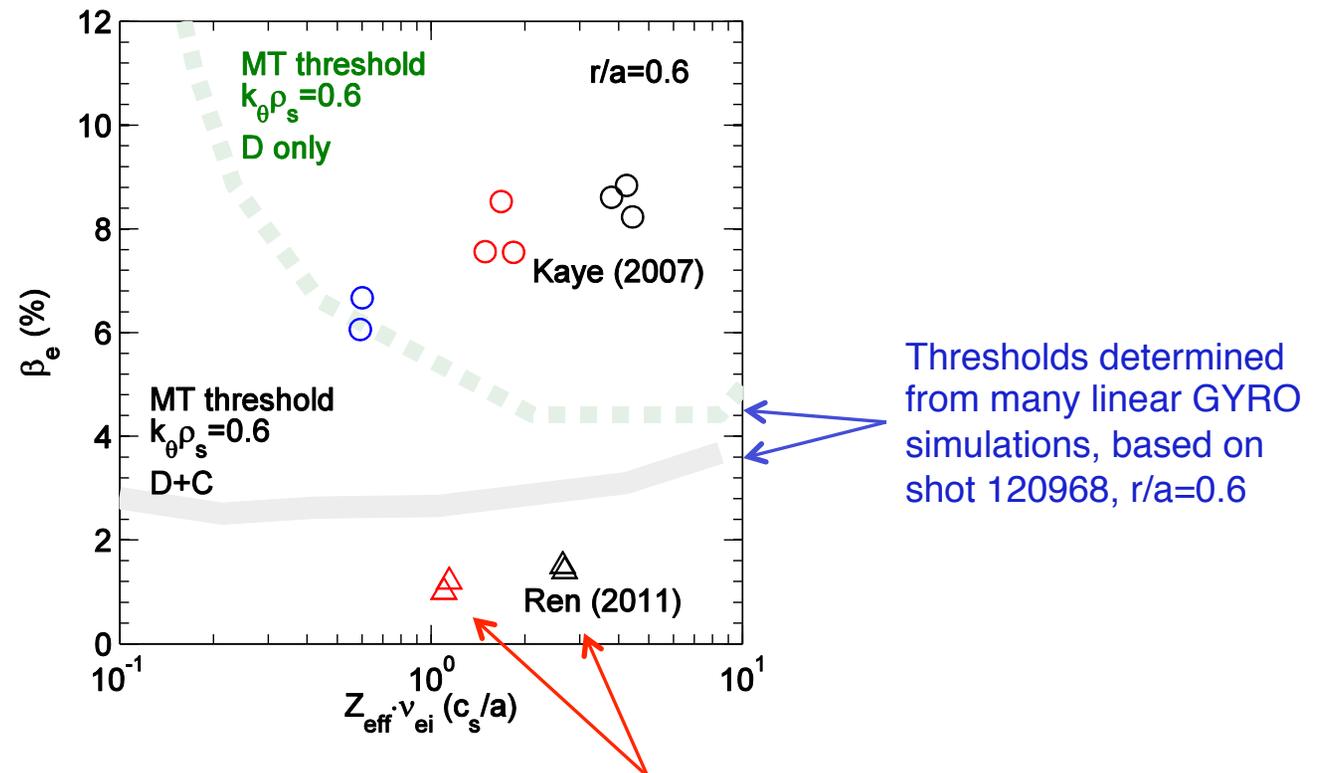
- New polarimetry system proposed for NSTX-Upgrade (UCLA, Jie Zhang et al. RSI, 2010) may be able to detect  $(\delta B/B)_{\text{internal}} \leq 0.1\%$  ( $1-2^\circ$  or  $\sim 0.3^\circ$  rms)

E. Wang et al., PoP (2011).

A.B. Rechester & M.N. Rosenbluth, PRL (1978)

# NSTX has studied electron transport for a range of beta and collisionality

- Microtearing exhibits a threshold in  $\beta_e$  (or  $a/L_{Te}$ ) that depends on  $\nu_{ei}$ ,  $Z_{eff}$ , etc...
- Distinguishes earlier scaling studies (Kaye, 2007) at higher beta compared to more recent studies (Ren, 2011)

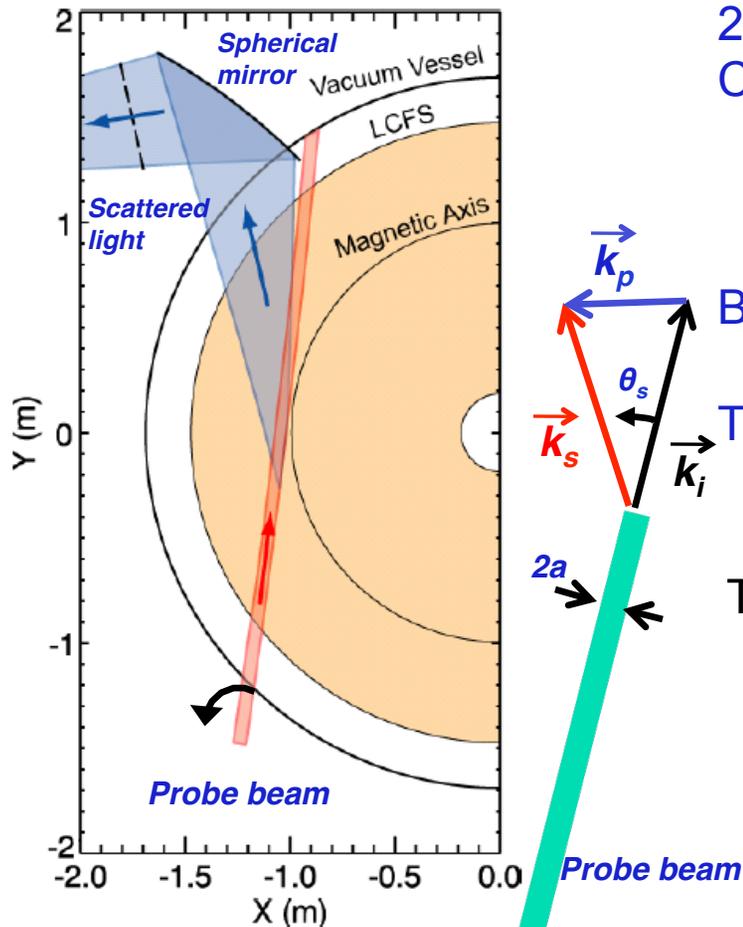


- Following studies investigate ETG turbulence at “low  $\beta_e$ ” where microtearing is predicted to be stable

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  - Predicts experimental level of electron thermal transport
  - Scaling of transport with collisionality ( $\chi_{e,sim} \sim \nu_e$ ) consistent with confinement ( $\Omega_i \tau_E \sim \nu_*^{-0.8}$ )
- Parametric dependence of high-k turbulence measured by a microwave scattering diagnostic in “low beta” plasmas
  - **Collisionality dependence of high-k turbulence in H-mode**
  - Density gradient stabilization of ETG turbulence in H-mode, partially validated with non-linear simulations
  - Suppression of ETG turbulence in reverse shear L-mode plasmas with e-ITB, partially validated with non-linear simulations
- First low-k turbulence measurements from a newly implemented BES diagnostic
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# High-k microwave scattering system capable of measuring electron-scale turbulence



D.R. Smith, PhD thesis, 2009

280 GHz microwave is launched as the probe beam. Coherent scattering by plasma density fluctuations occurs when the three-wave coupling condition is satisfied:

$$\vec{k}_s = \vec{k}_p + \vec{k}_i$$

Bragg condition determines  $k_p$ :

$$k_p = 2k_i \sin(\theta_s/2)$$

The scattered light has a frequency of:

$$\omega_s = \omega_p + \omega_i$$

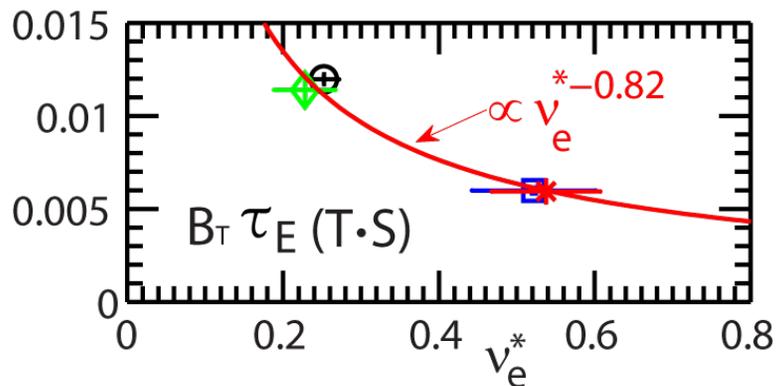
with  $\omega_s$  and  $\omega_i \gg \omega_p$

The scattering system characteristics are:

- Frequency bandwidth: 5 MHz
- Heterodyne receiver: Wave propagation direction resolved
- Measurement:  $k_r$  spectrum
- Wavenumber resolution:  $0.7 \text{ cm}^{-1}$  ( $2/a$  with  $a \approx 3 \text{ cm}$ )
- **Wavenumber range ( $k_r$ ):  $5\text{-}30 \text{ cm}^{-1}$  ( $\sim 5\text{-}30 \rho_s^{-1}$ )**
- Radial resolution:  $\pm 2 \text{ cm}$
- Tangential resolution: 5-15 cm
- Radial range: R=106 – 144 cm
- Minimal detectable density fluctuation:  $|\delta n_e(k)/n_e|^2 \approx 2 \times 10^{-11}$

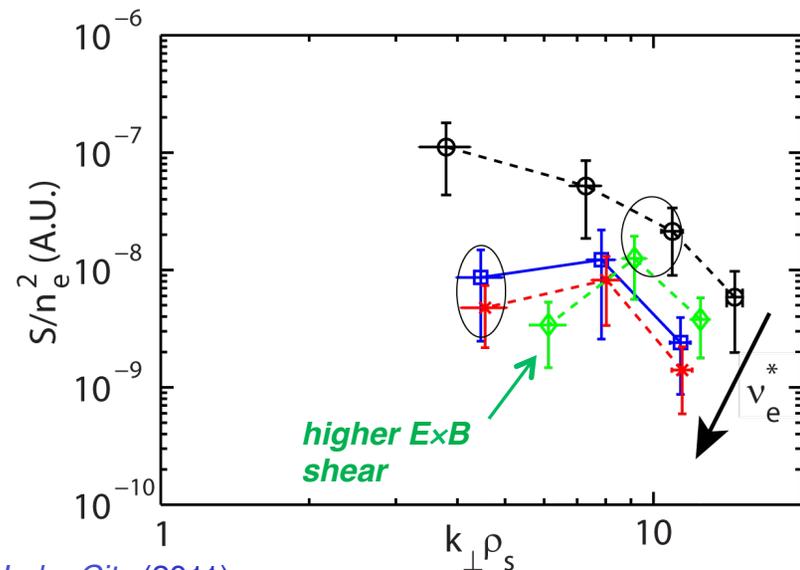
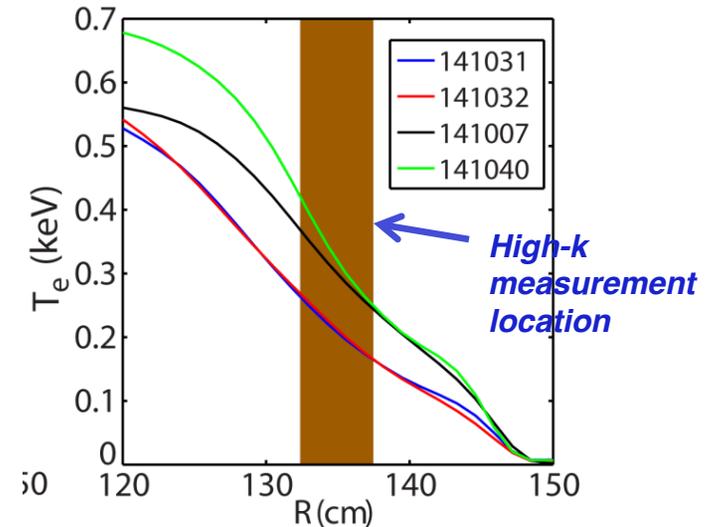
# More recent experiment to study $v_{*e}$ scaling of electron scale turbulence using high-k scattering

- Factor  $\sim 2.5$  change in  $v_{*e}$ ; local  $\rho_e$ ,  $\beta_e$ ,  $q_{95}$  are well matched around high-k measurement region ( $R=130-140$  cm)
- Confinement scaling  $\Omega_i \tau_E \sim v_{*e}^{-0.82}$ , similar to previous scaling  $\Omega_i \tau_E \sim v_{*e}^{-0.95}$



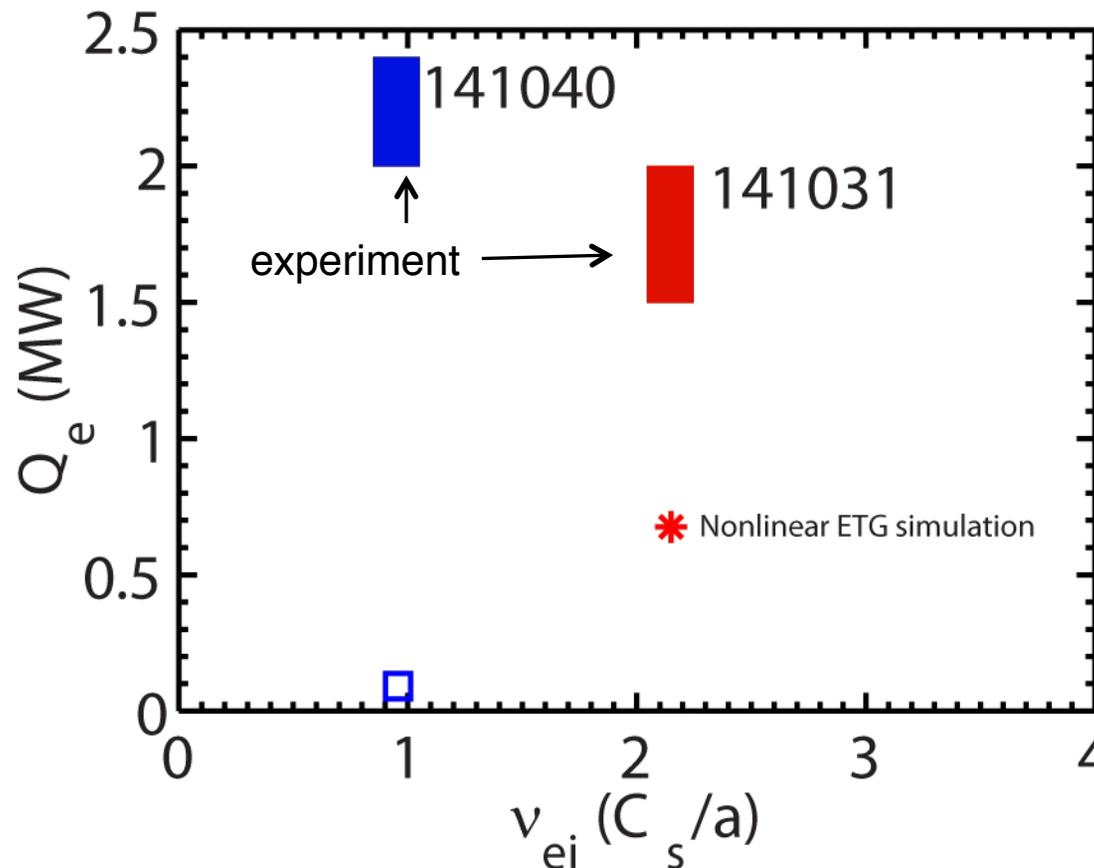
- High-k turbulence intensity decreases with increasing  $v_{*e}$
- Trend holds at all  $k_{\perp} \rho_s$  except for one case where local  $E \times B$  shear is  $\sim 2\times$  larger

\* Ren et al., APS-DPP invited talk T12.02, Salt Lake City (2011)



## Local nonlinear ETG simulations show large deviation from experimental transport

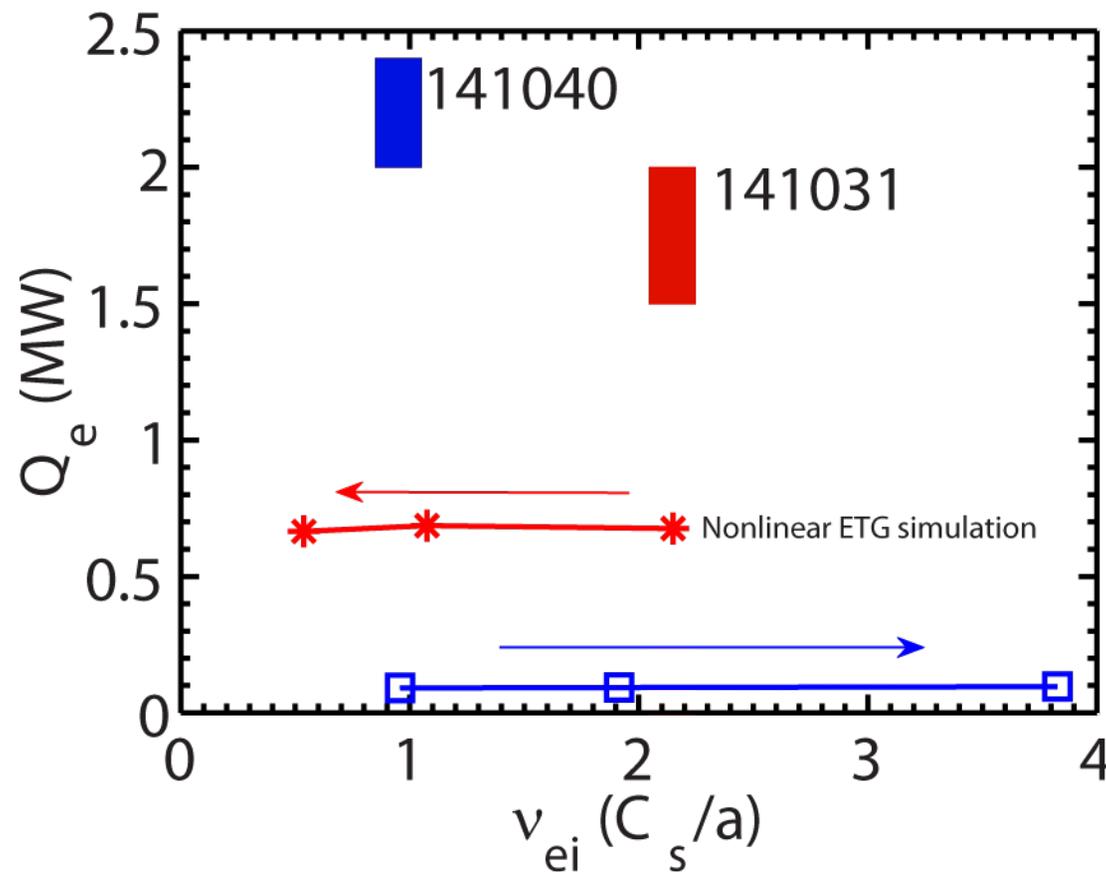
- In these lower beta discharges ETG is locally unstable (no microtearing)
- Predicted heat flux much smaller than experiment
- Can not be accounted for by sensitivity in  $a/L_{Te}$



Nonlinear GYRO ETG simulations with: local general equilibrium, kinetic ions and electrons, collisions, electromagnetic, flow and flow shear

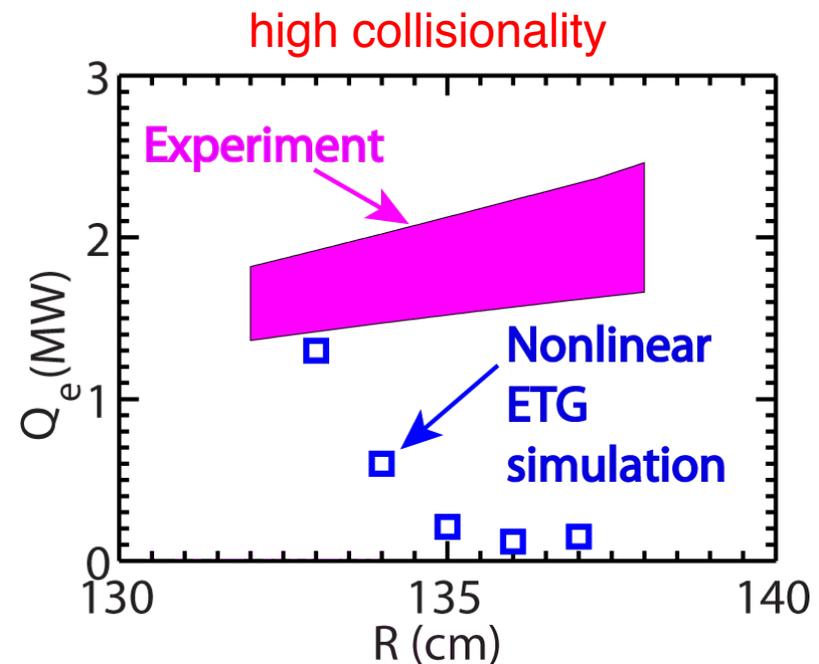
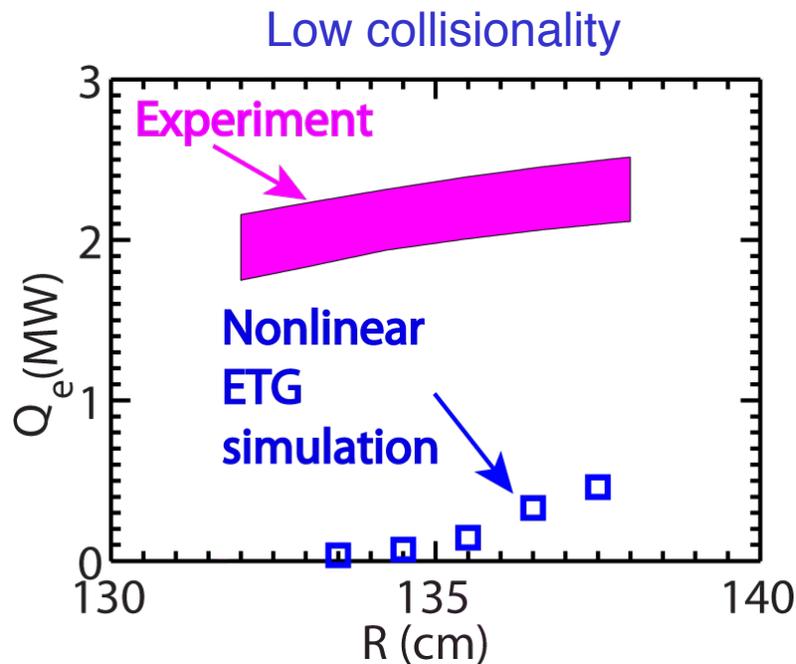
# Local nonlinear ETG simulations show weak collisionality dependence

- No dependence of predicted ETG transport on  $\nu_e$
- Can not explain confinement scaling through ETG dependence on  $\nu_e$  alone



# Profile variations lead to significant variation in local transport predictions

- Radial variation in other parameters (notably  $a/L_n$ ,  $q$  &  $s$ ) cause dramatic change in predicted transport over high-k measurement region ( $\Delta R \sim 4$  cm)
- Large discrepancy remains for the low collisionality shot
- Match with experimental  $Q_e$  found at inner radii for the high collisionality shot

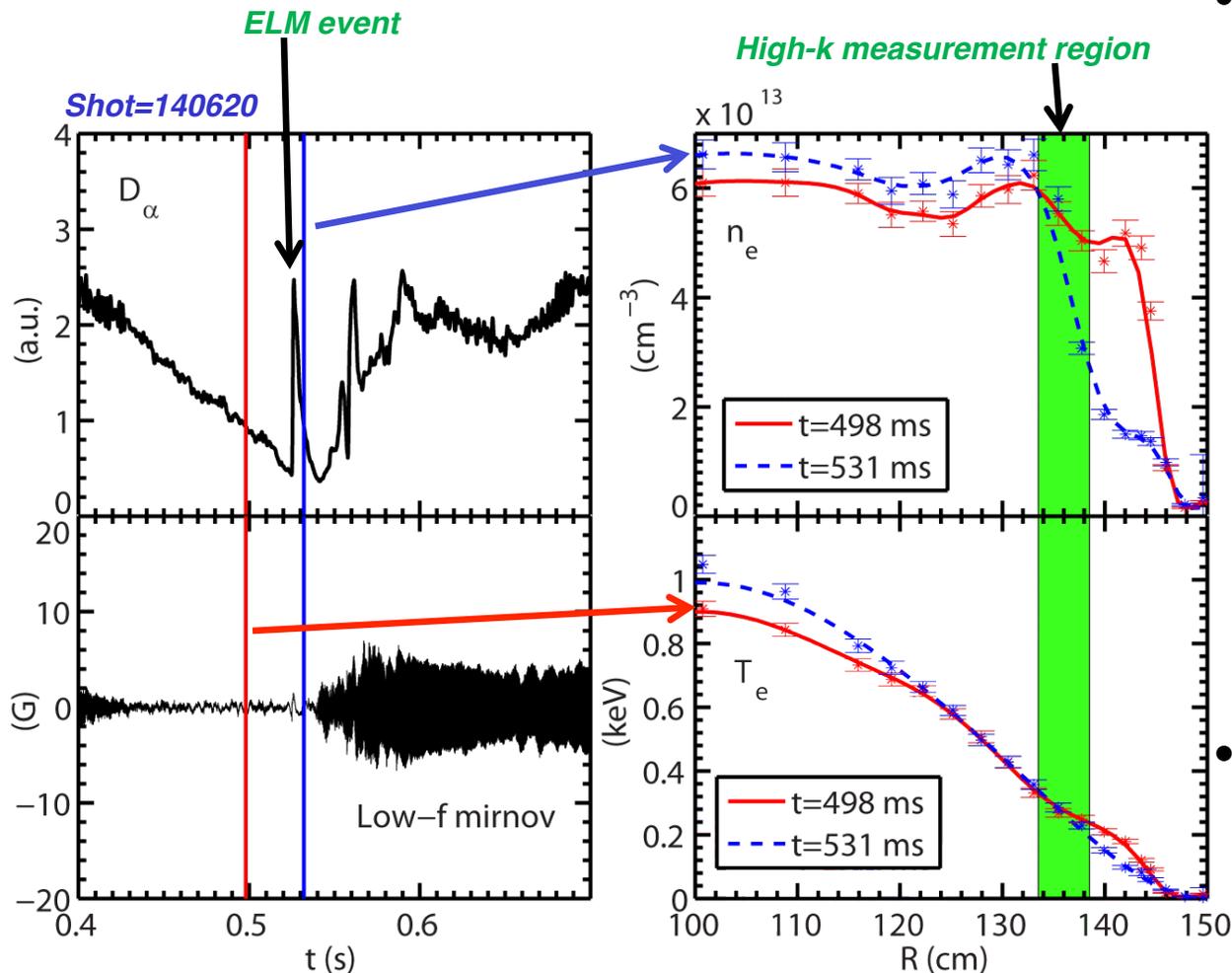


- Pursuing “global” ETG simulations with profile variations
- ITG also found to be unstable at different radii with  $\gamma_{lin} > \gamma_E \rightarrow$  ion scale ( $\sim 7 \rho_i$ ) turbulence spreading may also be important

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# Using increased density gradient induced by a large ELM as a tool for local turbulence studies

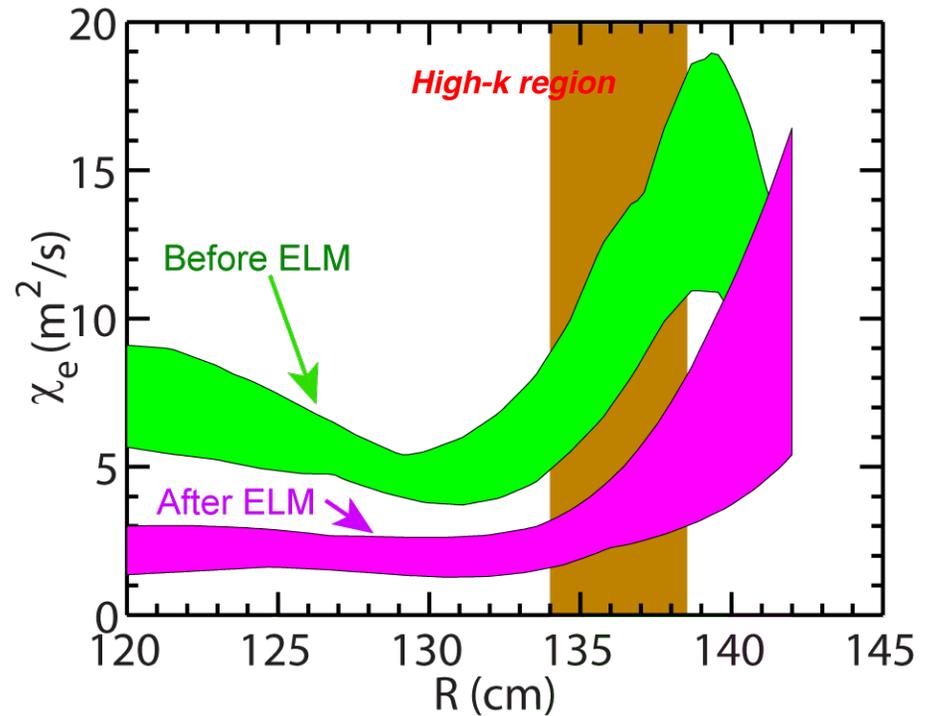
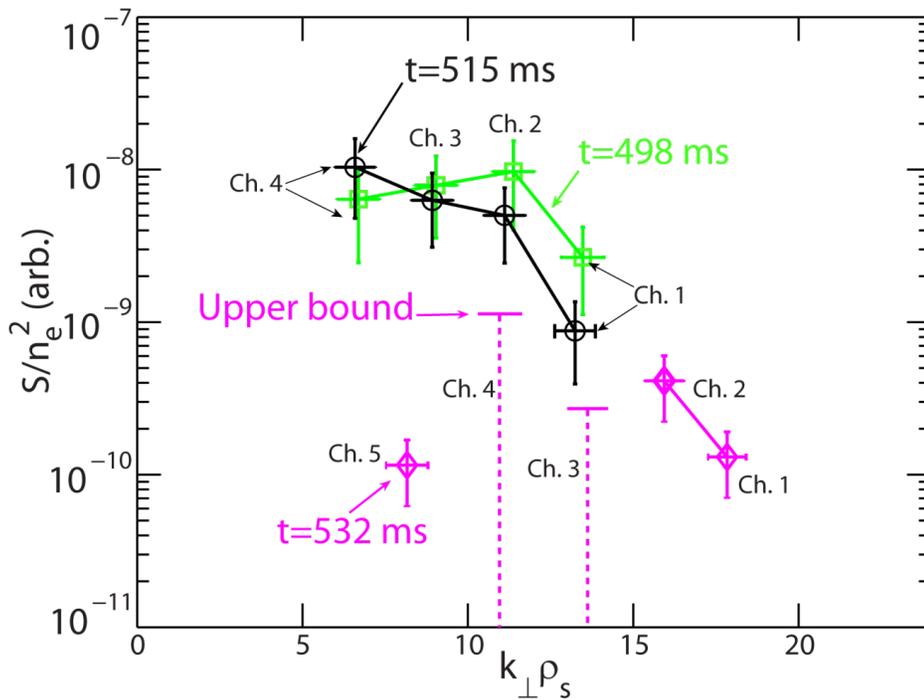


- After the ELM event:
  - Large density gradient developed in the high-k measurement region.
  - Electron temperature gradient also increases
  - Electron density has only a moderate decrease
  - Electron temperature remains essentially constant
- No large MHD mode appears before and right after the ELM event

# Correlation between reduced measured turbulence intensity and improved plasma thermal confinement\*

- Significant decrease in spectral power observed for  $k_{\perp}\rho_s \lesssim 10$

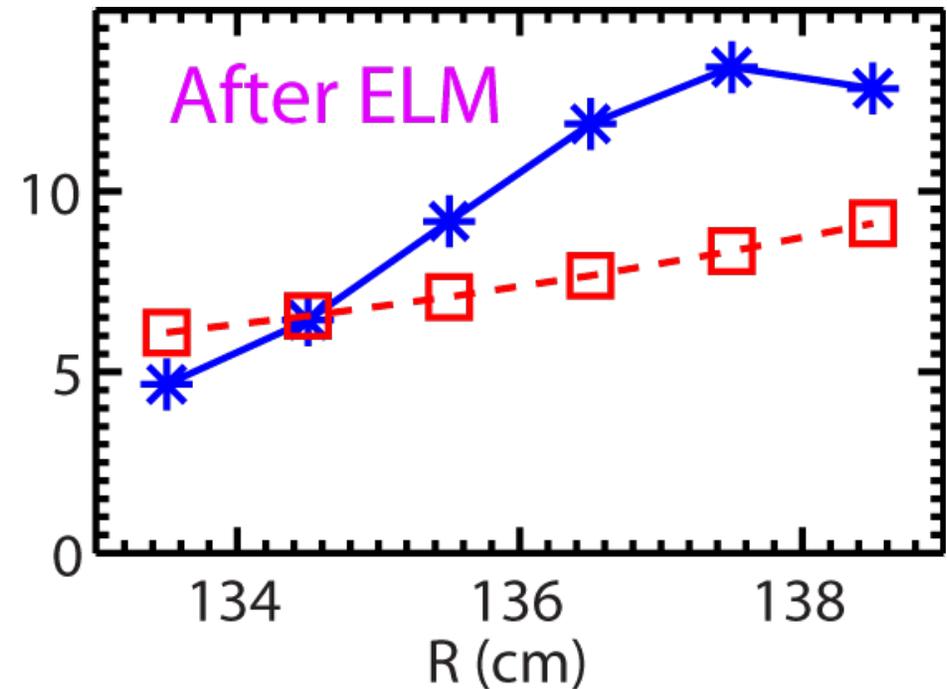
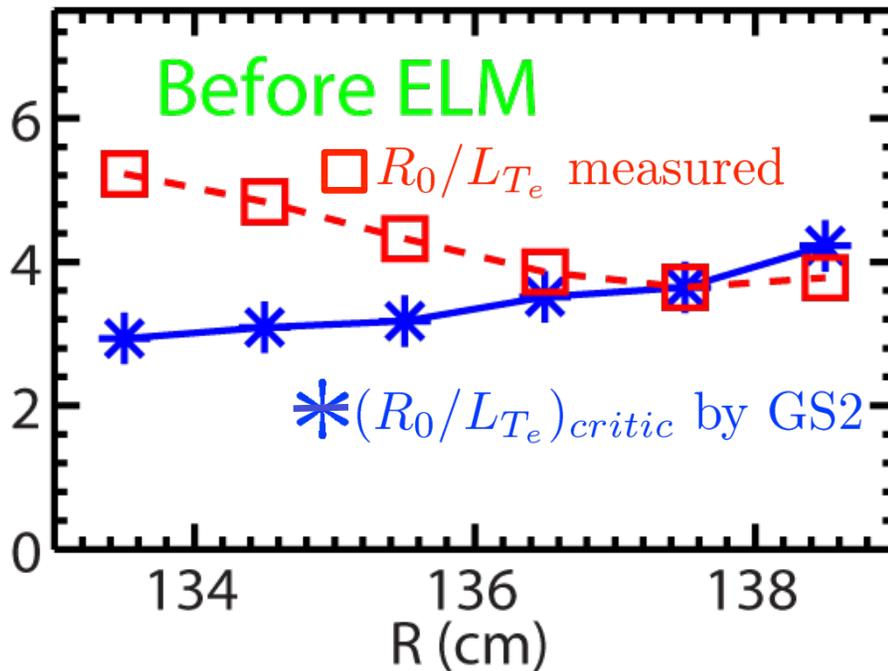
- Electron thermal diffusivity is decreased by a factor of  $\sim 2$  after the ELM event



\*Y. Ren et al., PRL 106, 165005 (2011)

# Threshold gradients for ETG modes are much higher after the ELM

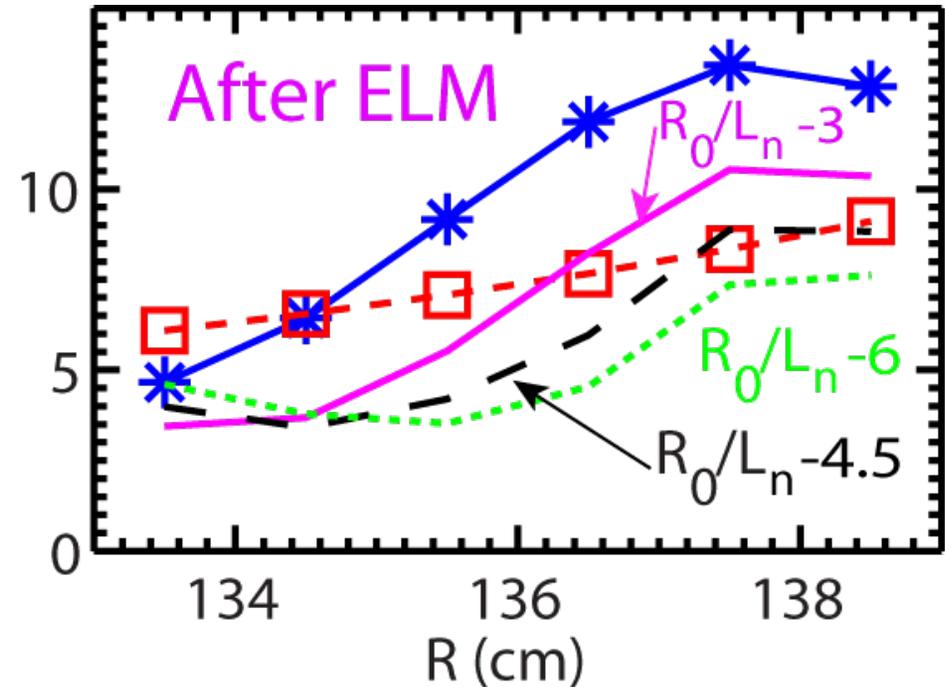
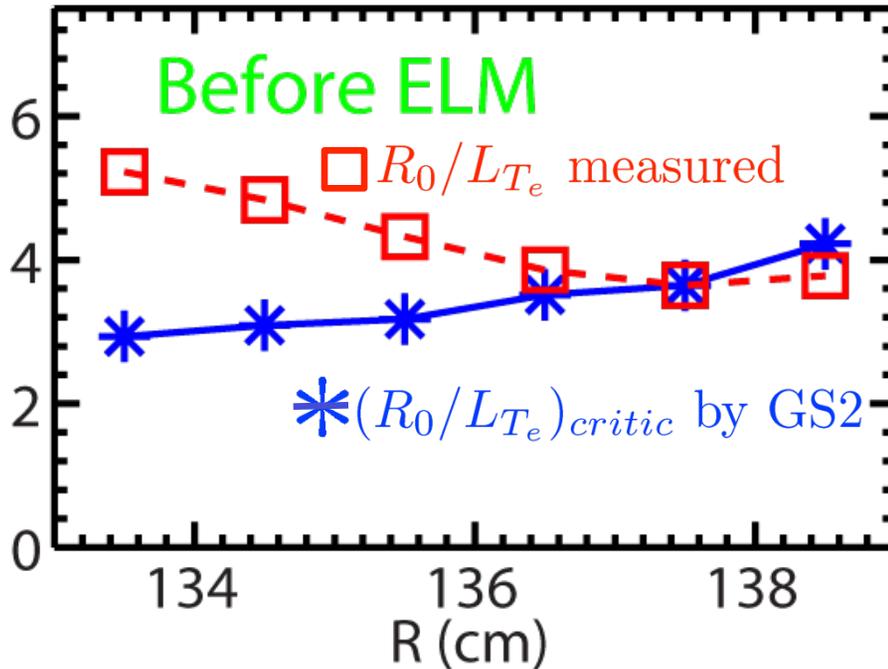
- Before ELM, ETG is largely unstable
- After ELM, ETG is largely stable



- Stability analysis performed with GS2 code (Kotschenreuther et al., 1995)

# Increase in ETG threshold gradient is due to large density gradient

- Before ELM, ETG is largely unstable
- After ELM, ETG is largely stable



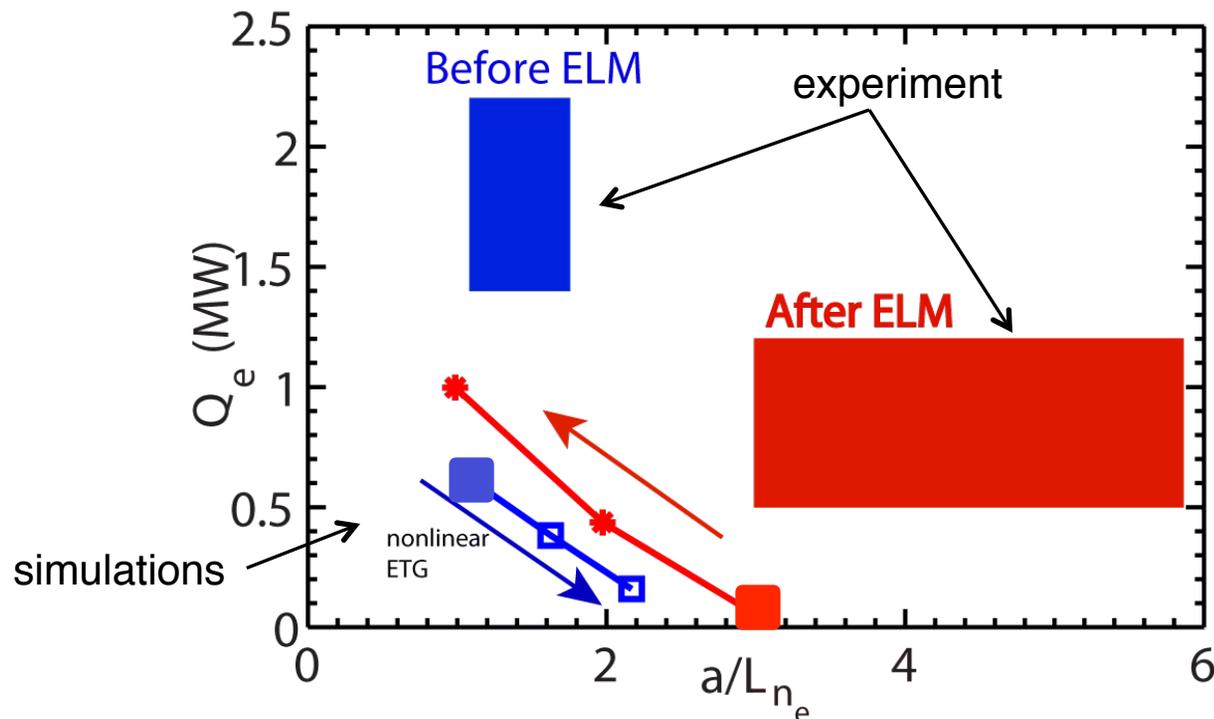
- Manually decreasing  $R/L_{ne}$  brings down critical gradient as expected from linear theory (e.g. Jenko et al, 2001)

$$(R_0/L_{T_e})_{crit} = \max\left\{ \left(1 + Z_{eff} \frac{T_e}{T_i}\right) (1.33 + 1.99\hat{s}/q) f(\epsilon, \kappa, \delta, \dots), \underline{0.8R_0/L_{ne}} \right\}$$

- Stability analysis performed with GS2 code (Kotschenreuther et al., 1995)

# Nonlinear ETG simulations reproduce observed dependence of electron transport on density gradient

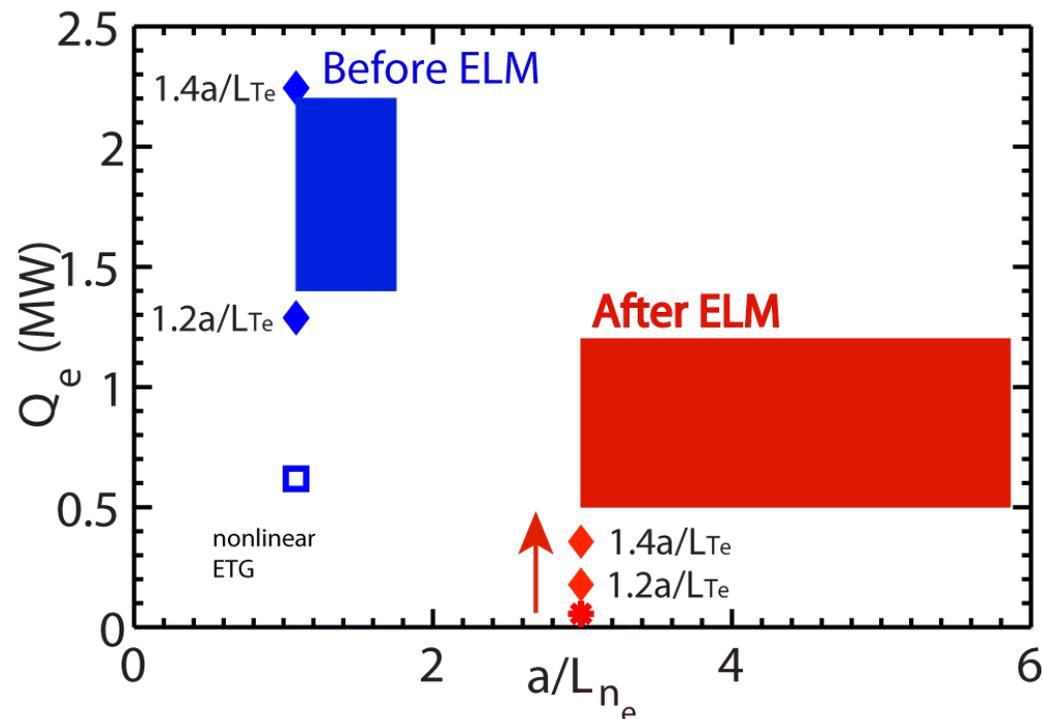
- Experimental  $Q_e$  is found to decrease after the ELM event with large density gradient
- The same trend is found from nonlinear ETG simulations, but does not agree quantitatively



Nonlinear GYRO ETG simulations with: local general equilibrium, kinetic ions and electrons, collisions, electromagnetic, flow and flow shear

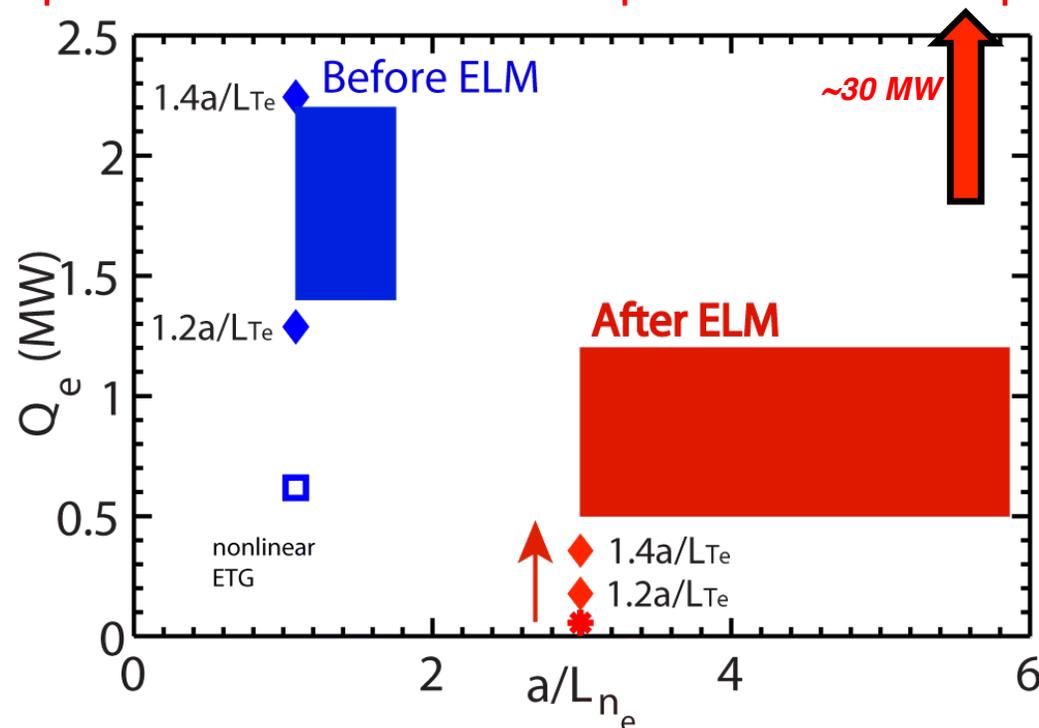
## Predicted $Q_e$ sensitive to temperature gradient

- Before ELM, a 20-30% increase in  $a/L_{Te}$  is able to match the experimental  $Q_e$
- After ELM, increasing  $a/L_{Te}$  by 40% after still cannot match experimental  $Q_e$



## Trapped Electron Mode (TEM) destabilized by large density gradient may contribute to transport

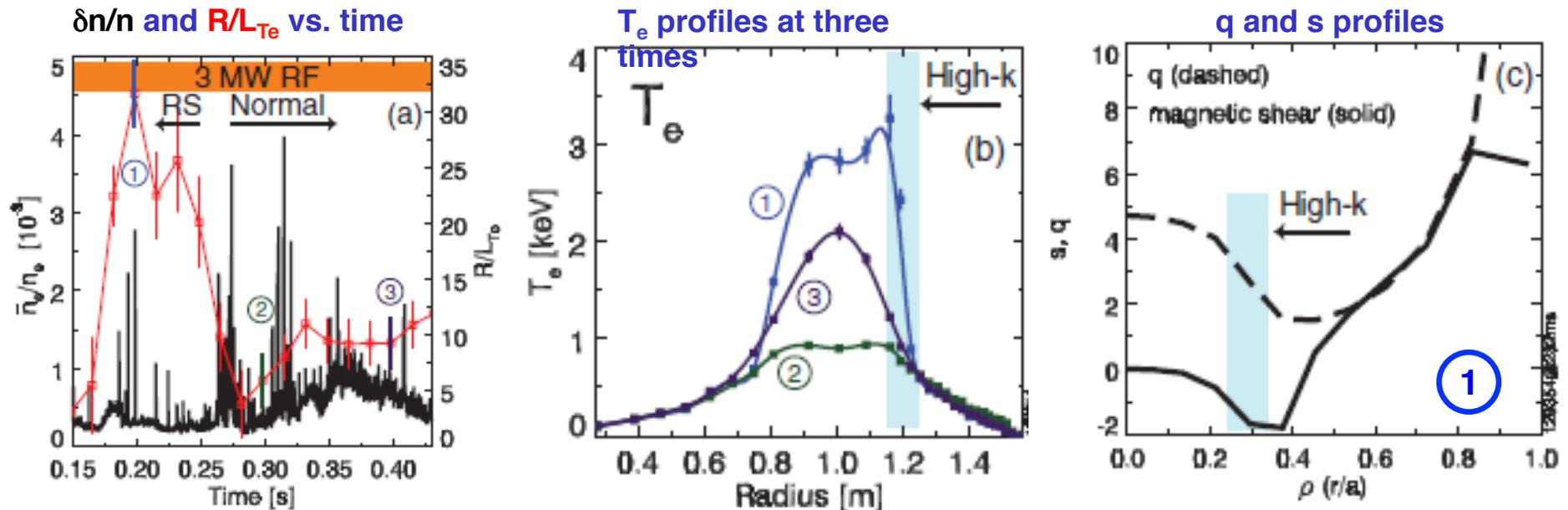
- Before ELM, a 20-30% increase in  $a/L_{Te}$  is able to match the experimental  $Q_e$
- After ELM, increasing  $a/L_{Te}$  by 40% after still cannot match experimental  $Q_e$
- Large TEM-induced transport ( $\sim 30$  MW) is predicted after ELM without  $E \times B$  shear stabilization
- Using experimental  $E \times B$  shear almost completely suppresses transport  
→ does not require much residual transport to match experimental  $Q_e$



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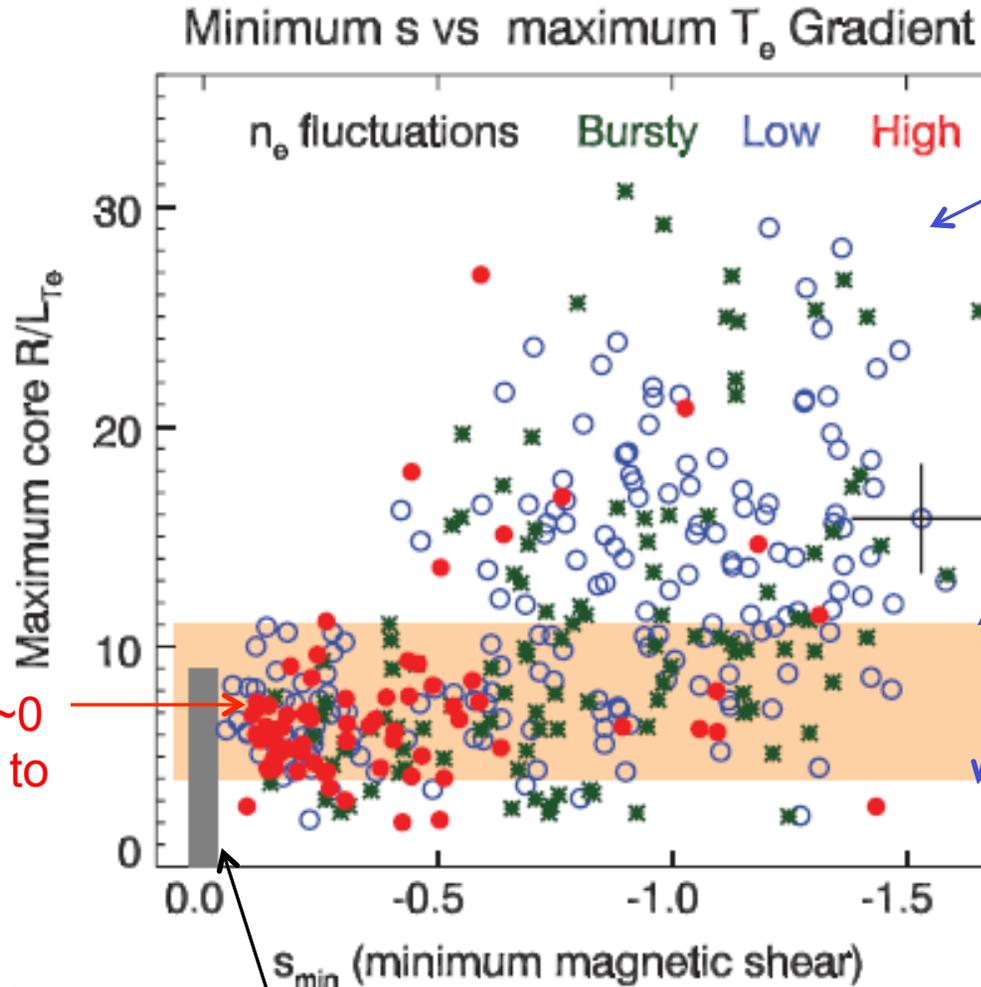
# Electron internal transport barrier (e-ITB) occurs in L-mode with large reverse magnetic shear, $s \ll 0$



- ① e-ITB occurs only during reversed shear portion of discharge, even in the absence of  $E \times B$  shear  
Very low, or bursty, high-k fluctuations in e-ITB
- ② Current is suddenly redistributed by MHD leading to monotonic  $q$  profile
- ③  $\rightarrow$  near zero or positive  $s$ , larger high-k fluctuations, smaller maximum gradient

# Largest gradients and weak high-k turbulence correlated with largest negative shear

Many discharges  
 2 MW NBI  
 0-3 MW RF  
 $1-4 \times 10^{19} \text{ m}^{-3}$   
 $Z_{\text{eff}}=1.1-4 \text{ (D, He)}$



Highest  $R/L_{Te}$  occur for  $s \ll 0$

Low or bursty high-k fluctuation intensity

Well above linear thresholds (supercritical)

For weak shear,  $s \sim 0$   $R/L_{Te}$  limited close to linear threshold

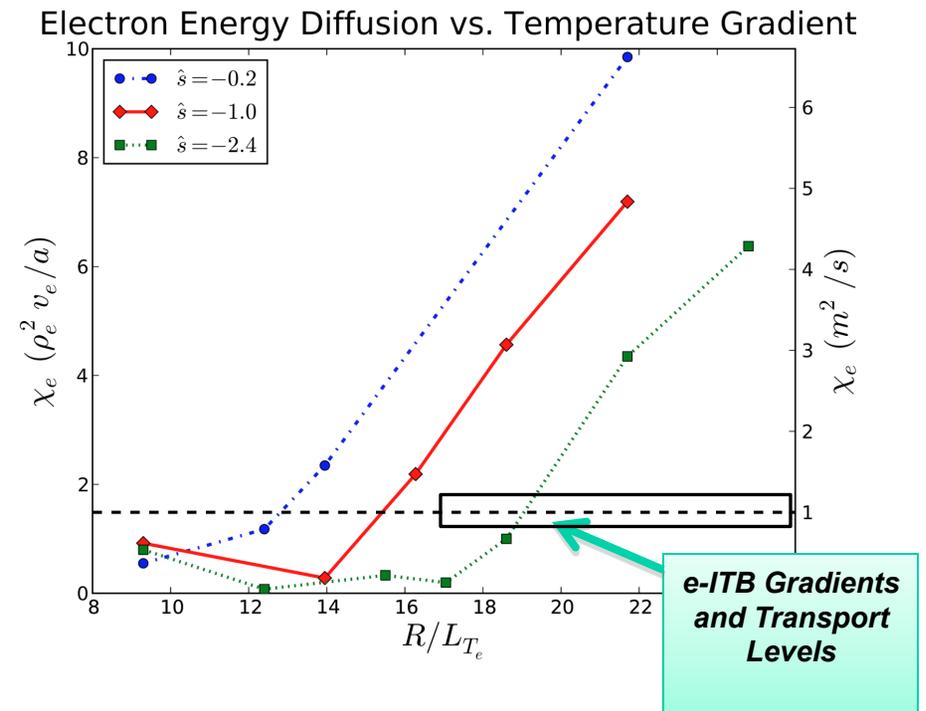
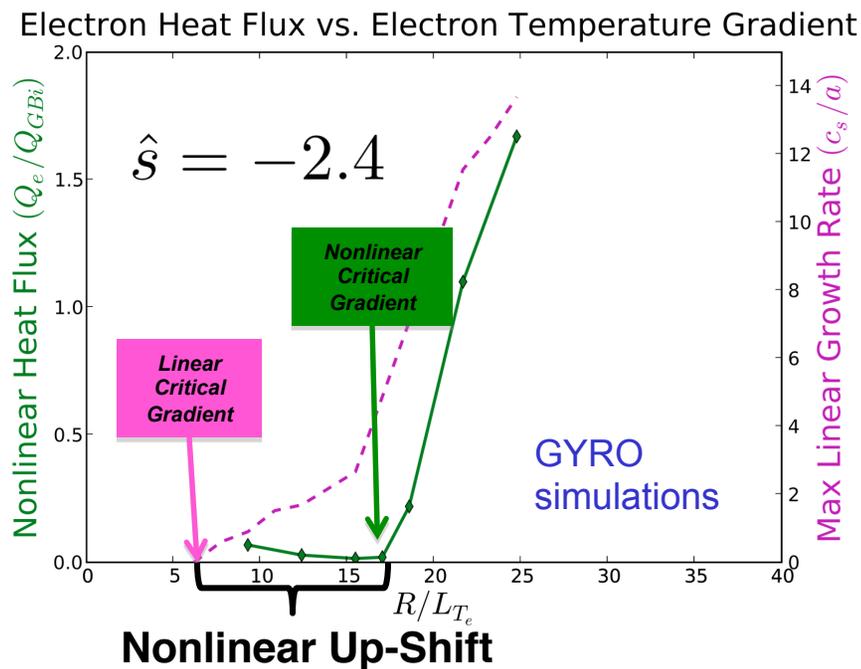
Large high-k fluctuation intensity

Characteristic H-mode values

H. Yuh, Phys. Rev. Lett. 106, 055003 (2011)

# Large negative magnetic shear causes strong upshift in effective non-linear ETG threshold

- **Nonlinear threshold** much larger than **linear threshold** for large negative shear in agreement with supercritical ETG gradients observed in experiments
- Nonlinear threshold increases with reverse shear magnitude,  $s < 0$
- **Magnitude and scaling of predicted transport consistent with experiment**



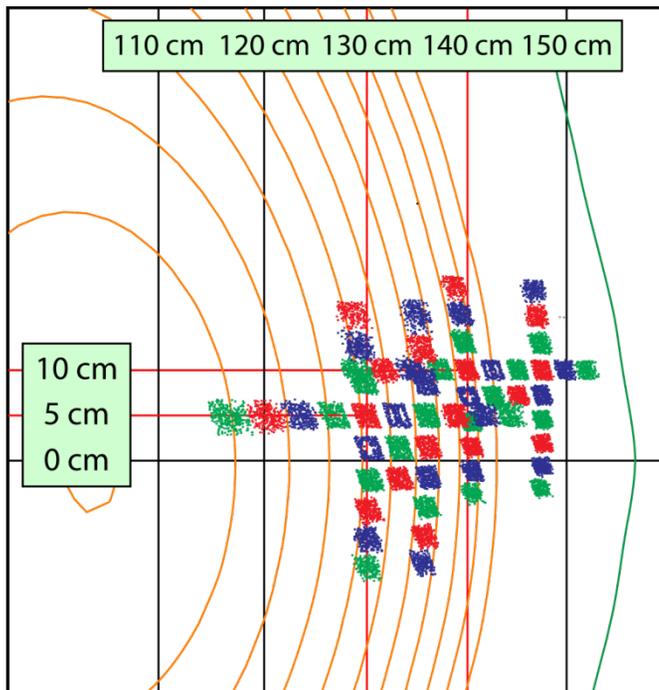
J.L. Peterson, Ph.D. Thesis, PPPL (2011); APS invited talk T12.03, Salt Lake City (2011)

# Overview

- Motivation
  - Improve understanding of core anomalous electron thermal transport in NSTX
  - Attempt to validate with turbulence simulations
- First nonlinear gyrokinetic simulations of micro-tearing turbulence for “high beta” H-mode plasmas
  - Predicts experimental level of electron thermal transport
  - Scaling of transport with collisionality ( $\chi_{e,sim} \sim \nu_e$ ) consistent with confinement ( $\Omega_i \tau_E \sim \nu_*^{-0.8}$ )
- Parametric dependence of high-k turbulence measured by a microwave scattering diagnostic in “low beta” plasmas
  - Collisionality dependence of high-k turbulence in H-mode
  - Density gradient stabilization of ETG turbulence in H-mode, partially validated with non-linear simulations
  - Suppression of ETG turbulence in reverse shear L-mode plasmas with e-ITB, partially validated with non-linear simulations
- First low-k turbulence measurements from a newly implemented BES diagnostic
  - Decrease in low-k turbulence from L-H mode transition, from edge to core
  - Poloidal correlation lengths in the edge pedestal region correlated with  $q$ ,  $s^{-1}$ ,  $\nabla n$
- Summary

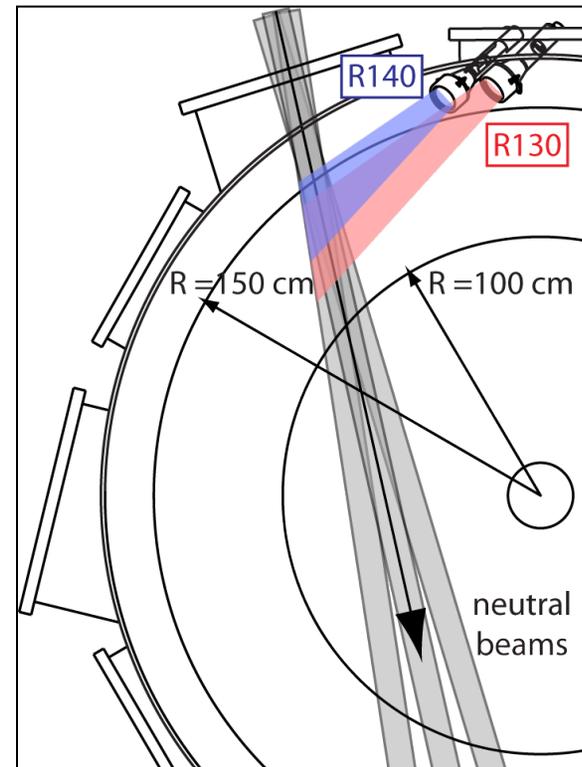
# Beam Emission Spectroscopy (BES) diagnostic recently commissioned, obtained data routinely during FY10-11

- Presently 32 detection channels
- 56 sightlines in radial and poloidal arrays spanning core to SOL
- 2 MHz sampling
- $k_{\perp} \rho_i \leq 1.5$  & 2-3 cm spot size
- Field-aligned optics with high throughput (etendue = 2.3 mm<sup>2</sup>-ster)



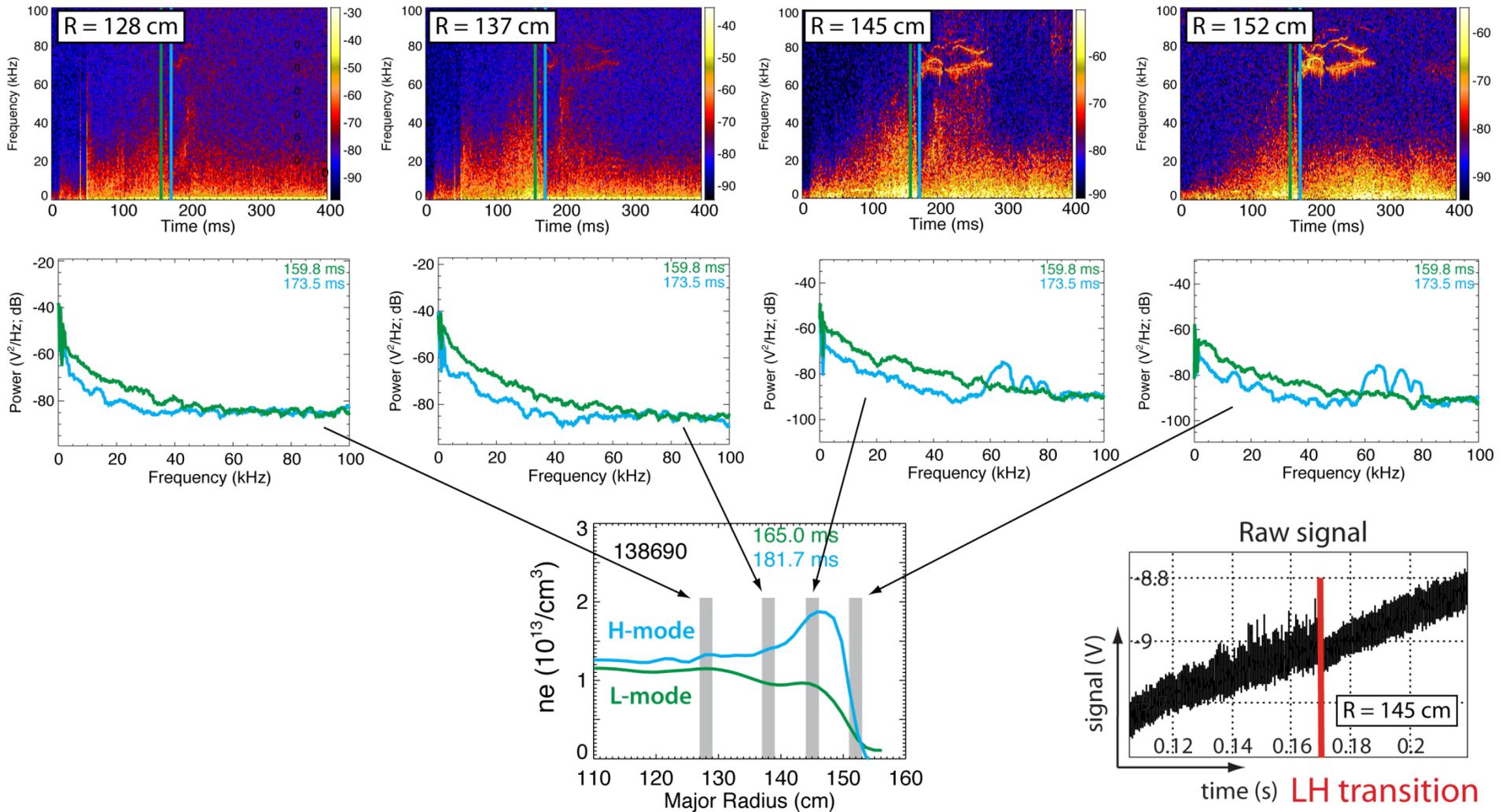
$$\frac{\delta I_{D\alpha}}{I_{D\alpha}} = \frac{\delta n}{n} \times C(E_{NB}, n, T_e, Z_{eff})$$

$\delta I_{D\alpha}$  ← neutral beam  $D_{\alpha}$  emission  
 $\frac{\delta n}{n}$  ← density fluctuation  
 $C \approx 1/2$



\*D.R. Smith et al., Rev. Sci. Instrum (2010)

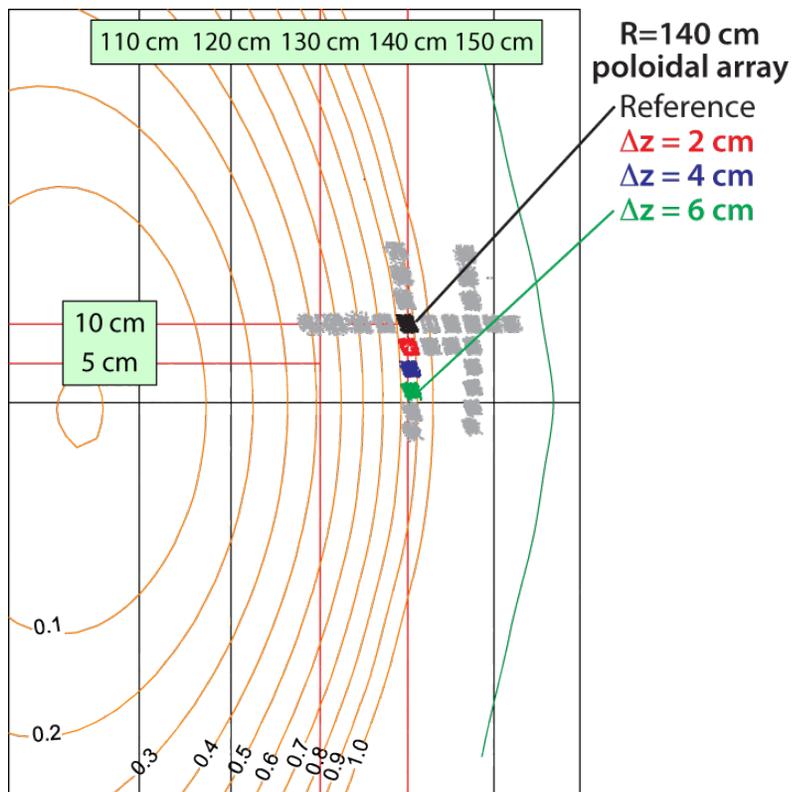
# Decrease in low-k turbulence observed at L-H transition from edge to core



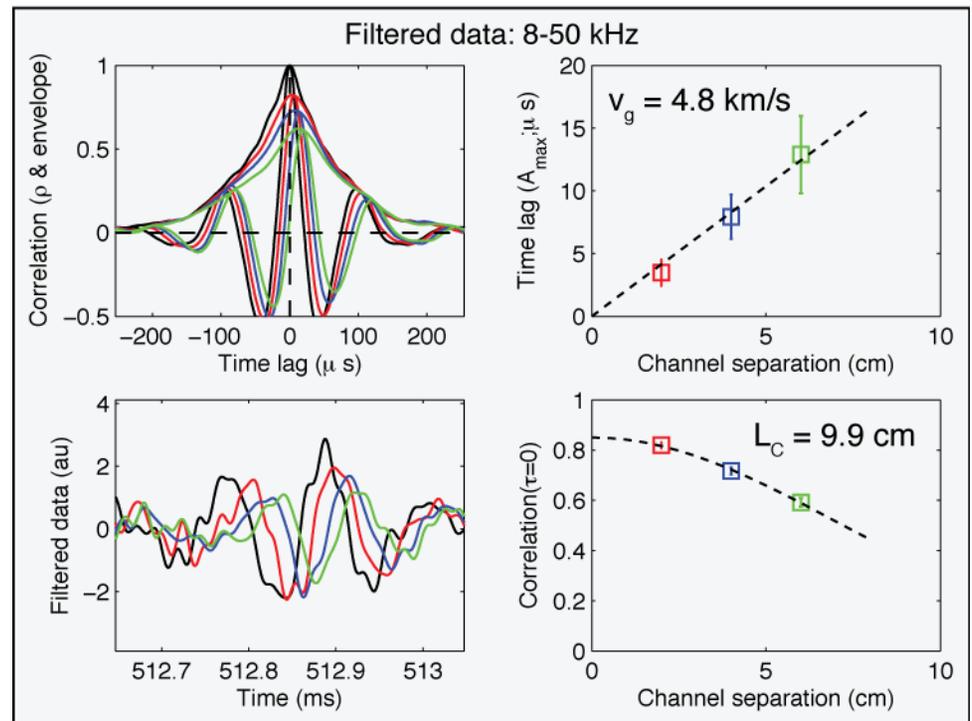
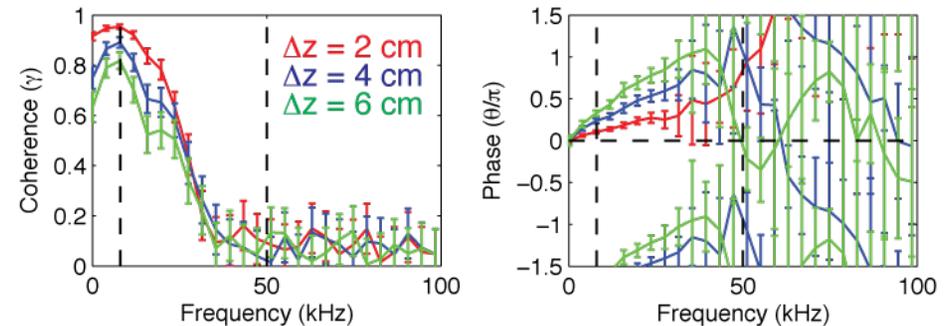
R. Fonck, G. McKee, D. Smith, and I. Uzun-Kaymak (UW-Madison) and B. Stratton (PPPL)

# Poloidal correlation length ( $L_C$ ) obtained from BES poloidal array at $R=140$ cm ( $r/a \approx 0.8-0.95$ )

- Correlation functions and coherence spectra indicate:  
 $k_\theta \approx 0.2-0.4 \text{ cm}^{-1}$  and  $k_\theta \rho_i \approx 0.1-0.25$
- Cross-phase and time-lags show eddy advection in ion diamagnetic direction



140991 @ 500-525 ms



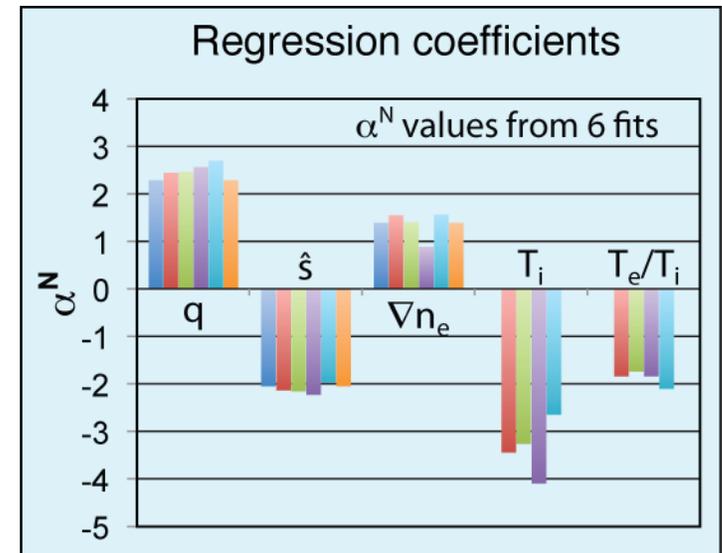
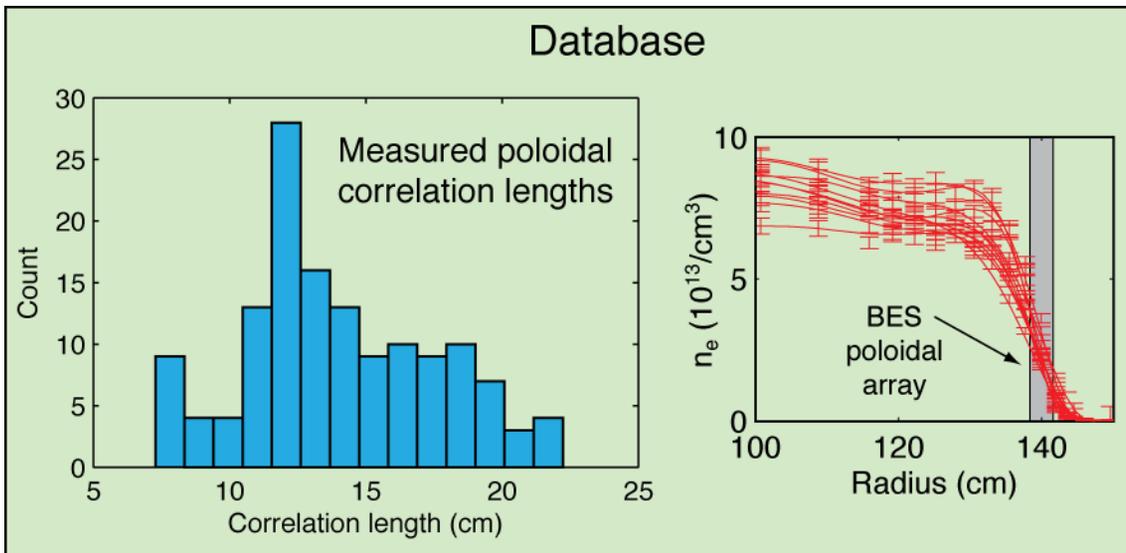
## Poloidal $L_c$ in pedestal region are 7-22 cm, appear to be correlated with $q$ , $1/s$ , $\nabla n_e$

- Poloidal  $L_c$  database for ELM-free, MHD-quiescent H-mode contains 130 entries from 29 shots (fixed  $B_{T0} = 4.4$  kG), in the pedestal region
- Regression analysis attempts to fit scaling of  $L_c$  to different parameters, e.g.  $n_e$ ,  $T_e$ ,  $T_i$ ,  $\nabla(n_e, T_e, T_i)$ ,  $a/L_{(n_e, T_e, T_i)}$ ,  $q$ ,  $\hat{s}$ ,  $\nu$ ,  $\beta$ , etc...

$$L_c = \bar{L}_c + \sum \alpha_i^N \frac{x_i - \bar{x}_i}{\sigma_{x_i}}$$

scaling coefficient
plasma parameters

- Best fits find poloidal  $L_c$  increases with  $q$ ,  $\nabla n_e$ , decreases with  $s$ ,  $T_i$ ,  $T_e/T_i$



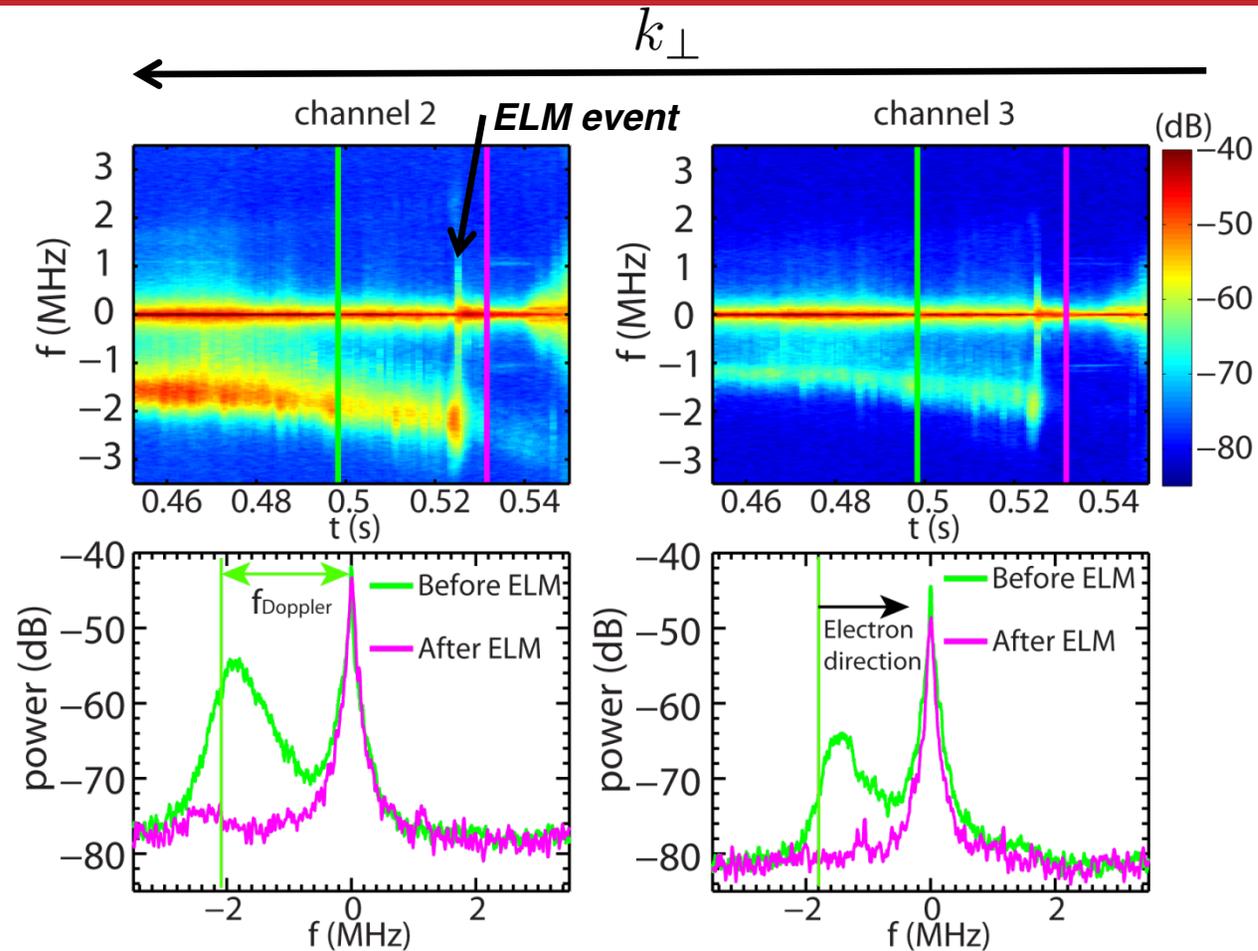
D.R. Smith et al., APS-DPP, Salt Lake City (2011)

## Summary: NSTX is making progress towards understanding anomalous electron thermal transport

- First nonlinear gyrokinetic simulation of micro-tearing turbulence in NSTX H-mode “high beta” plasmas reproduce magnitude and  $v_*$  scaling of electron transport
- In “lower beta” plasmas, where microtearing is stable:
  - High-k scattering intensity increases with decreasing  $v_*$  → ETG appears to be relevant at limited radial locations, doesn’t directly account for  $v_*$  scaling, requires more work to include profile effects, ITG may also be important
  - High-k turbulence stabilized by large density gradient → scaling reproduced by non-linear ETG simulations, TEM may become important at large  $a/L_{ne}$
  - e-ITB occurs at large negative magnetic shear in L-mode plasmas → large non-linear threshold reproduced by ETG simulations
- The newly implemented BES diagnostic provides first low-k turbulence measurements in NSTX
  - Decrease in low-k turbulence during L-H transition, from core to edge
  - Poloidal correlation lengths increase with  $q$ ,  $a/L_n$ , decrease with  $s$ ,  $T_i$ ,  $T_e/T_i$
- Please note, I have not covered many transport and turbulence topics: (1) core impurity and momentum transport, (2) transport driven by energetic particle modes, GAEs, CAEs, (3) many other edge/pedestal/SOL measurements → see NSTX website!



# Significant decrease in scattering signal power observed after the ELM



- All five channels saw decreased scattering power after the ELM event
- Interpretation has to take into account the change of wavenumber measured by each channel due to the increase density gradient & refraction after the ELM event

# Linear microtearing instability

- High- $m$  tearing mode around a rational  $q(r_0)=m/n$  surface ( $k_{\parallel}(r_0)=0$ )  
(Classical tearing mode stable for large  $m$ ,  $\Delta' \approx -2m/r < 0$ )
- Driven by  $\nabla T_e$  with time-dependent parallel thermal force\*  $\Rightarrow$  requires e-i collisions

## Conceptual linear picture

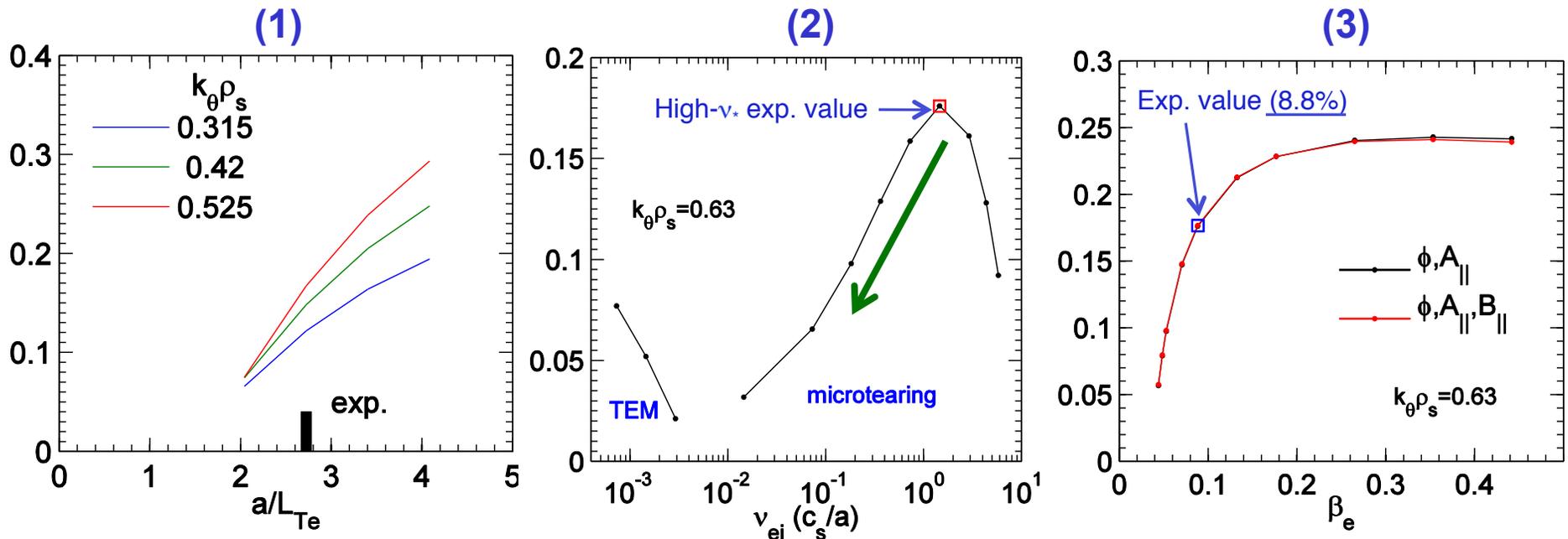
- Imagine helically resonant ( $q=m/n$ )  $\delta B_r$  perturbation  $\delta B_r \sim \cos(m\theta - n\varphi)$
- $\delta B_r$  leads to radially perturbed field line, finite island width  $w = 4 \left( \frac{\delta B_r}{B} \frac{rR}{n\hat{s}} \right)^{1/2}$
- $\nabla T_e$  projected onto field line gives parallel gradient  $\nabla_{\parallel} T_{e0} = \frac{\vec{B} \cdot \nabla T_{e0}}{B} = \frac{\delta B_r}{B} \nabla T_{e0}$
- Parallel thermal force ( $R_{T\parallel} \sim -\alpha(\omega) n_e \nabla_{\parallel} T_e$ ) drives parallel electron current that reinforces  $\delta B_r$  via Amperes's law  $k_{\perp}^2 \rho_s^2 \hat{A}_{\parallel} = \frac{\beta_e}{2} \hat{j}_{\parallel}$ ,  $B_r = ik_{\theta} A_{\parallel}$
- **Instability requires sufficient  $\nabla T_e$ ,  $\beta_e$ ,  $v_e$ , and time dependence ( $\omega$ ) important**

\*e.g. Hazeltine et al., Phys. Fluids 18, 1778 (1975); Gladd et al., Phys. Fluids 23, 1182 (1980);  
D'Ippolito et al., Phys. Fluids 23, 771 (1980); M. Rosenberg et al., Phys. Fluids 23, 2022 (1980).

# Microtearing Modes Found to be Unstable in High $v_*$ Discharge

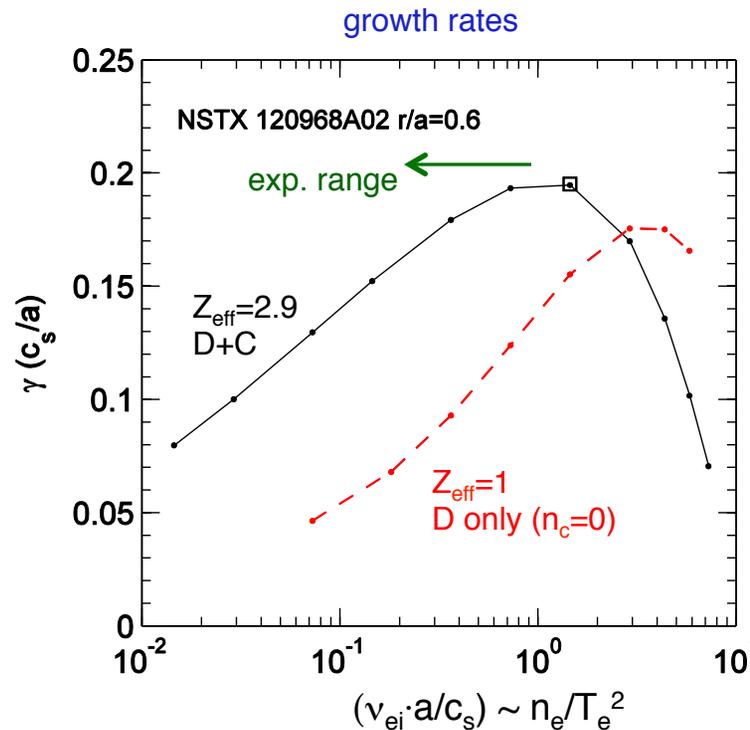
- Focusing on high- $v_*$  NSTX discharge, part of dimensionless scaling experiments where favorable scaling found  $\Omega\tau_E \sim v_*^{-0.95}$  [S.M. Kaye et al., Nucl. Fusion 47, 499 (2007)]
- Microtearing dominates  $k_\theta \rho_s < 1$  ( $n \approx 5-70$ ) in outer half-radius ( $r/a=0.5-0.8$ )
- ETG stable due to higher  $Z_{\text{eff}}=2.5-3.0$   $(R/L_{Te})_{\text{crit,ETG}} \sim (1+Z_{\text{eff}}T_e/T_i)$
- Microtearing exhibits threshold in  $\nabla T_e$ ,  $v_e$ ,  $\beta_e$
- Growth rates decrease with  $v_e < v_{e,\text{exp}}$  (consistent with experimental  $v_*$  scan)

Linear growth rates ( $\gamma \cdot a/c_s$ ) for NSTX 120968  $t=0.56$  s  $r/a=0.6$  with  $B_T=0.35$ T,  $I_p=0.7$  MA,  $P_{\text{NBI}}=4$  MW,  $n_e \approx 6 \times 10^{19}$  m $^{-3}$ ,  $T_e(0) \sim 1$  keV



# A distinguishing feature of the microtearing mode is the non-monotonic dependence on $v^{e/i}/\omega$

- Peak  $\gamma$  occurs for  $v^{e/i}/\omega = Z_{\text{eff}} \cdot v_{ei}/\omega \sim 1-6$ , similar to slab calculations (Gladd et al., 1980)
- $\gamma$  decreases with  $v_e$  in experimental range, qualitatively consistent with confinement scaling
- In addition to shifting peak in  $v^{e/i}/\omega$ ,  $Z_{\text{eff}}$  can enhance instability through shielding potential (from adiabatic ion response,  $\delta n_i \sim -Z_{\text{eff}} \delta \varphi / T_i$ )



$$v^{e/i} = Z_{\text{eff}} v_{ei} \propto Z_{\text{eff}} \frac{n_e}{T_e^{3/2}}$$

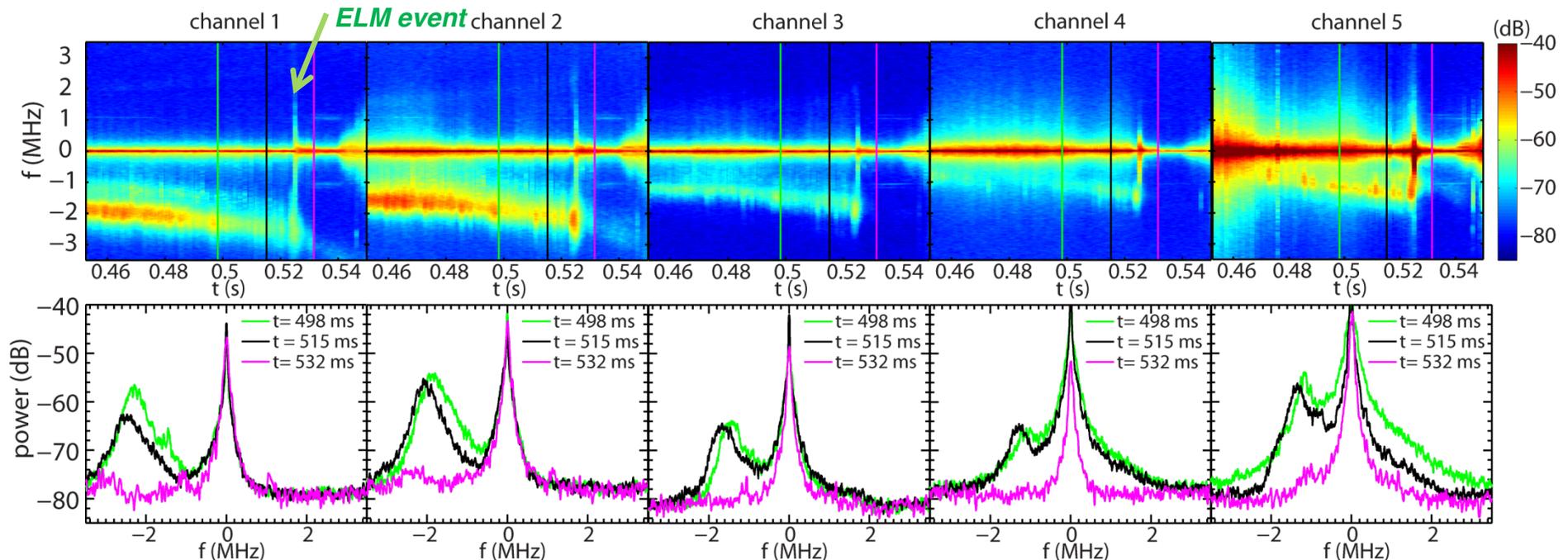
[Jenko et al. \(2001\)](#)

- $Z_{\text{eff}}$  (and  $s/q$ )\* dependence opposite to ETG expectations

$$\left( \frac{R}{L_{Te}} \right)_{\text{crit}}^{\text{ETG}} \sim \left( 1 + Z_{\text{eff}} \frac{T_e}{T_i} \right) \left( 1.3 + 1.9 \frac{s}{q} \right) (\dots)$$

\* Guttenfelder et al., *Scaling of linear microtearing stability for a high collisionality NSTX discharge*, submitted to Phys. Plasmas (Oct, 2011)

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