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## **Corrigendum: Error field correction in DIII-D Ohmic plasmas with either handedness**

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The expression for the overlap, equation (1) in the published paper (Park *et al* 2011 *Nucl. Fusion* **51** 023003), was incorrect and should be replaced by

$$\mathcal{C} \equiv \sqrt{\frac{\oint \mathrm{d}\varphi' [\oint \mathrm{d}a(\delta\vec{B}^x \cdot \hat{n}_b)(-\vartheta, \varphi' - \varphi)(\delta\vec{B}^x_d \cdot \hat{n}_b)(\vartheta, \varphi)]^2}{\oint \mathrm{d}a(\delta\vec{B}^x \cdot \hat{n}_b)^2 \oint \mathrm{d}a(\delta\vec{B}^x_d \cdot \hat{n}_b)^2}}.$$
(1)

We found that it would be worthwhile to add more descriptions. This is the root-mean-square integration over  $\varphi'$ , after the convolution integral to  $(\vartheta = 0, \varphi = \varphi')$  between the two distribution functions of magnetic field. The extra degree of freedom, only in the toroidal angle, is because the reference toroidal angle of the dominant field distribution  $\delta \vec{B}_d(\vartheta, \varphi)$  can be arbitrary due to the toroidal symmetry in a tokamak. This expression becomes simply the dot product

of the two normalized matrix vectors in the complex Fourier space, if one decomposes a magnetic field as

$$\sqrt{\mathcal{J}|\vec{\nabla}\psi|}(\delta\vec{B}^x\cdot\hat{n}_b)(\vartheta,\varphi) = \sum_{mn} \Phi_{mn} \mathrm{e}^{\mathrm{i}(m\vartheta - n\varphi)}.$$
 (2)

Note the additional factor associated with the Jacobian of the surface area,  $\mathcal{J}|\vec{\nabla}\psi|$ , to have the surface area integral  $da = \mathcal{J}|\vec{\nabla}\psi|d\vartheta \,d\varphi$  independent of coordinates. It is then straightforward to show equation (1) becomes the dot product between a given magnetic field  $\vec{\Phi} = \{(m, n)|\Phi_{mn}\}$  and the dominant magnetic field  $\vec{\Phi}_d = \{(m, n)|\Phi_{mn}\}$  as

$$C \equiv \frac{|\vec{\Phi} \cdot \vec{\Phi}_d|}{|\vec{\Phi}||\vec{\Phi}_d|}.$$
(3)