

# Abstract

Radial carbon ion density profiles are measured in NSTX plasmas using charge exchange recombination spectroscopy (CHERS). Determination of the carbon ion density requires knowledge of the neutral beam parameters and geometry, a neutral beam attenuation calculation, electron density and temperature measurements from Thomson scattering, an absolute intensity calibration, and a calculation of the effective emission rate for charge exchange. Profiles of the effective charge,  $Z_{eff}$ , are then calculated assuming carbon is the dominant impurity in the core plasma. A comparison of the expected visible bremsstrahlung emission based on the inferred  $Z_{eff}$  profiles is consistent with measured values from CHERS spectra in the core plasma; excessive emission is observed for edge sightlines.  $Z_{eff}$  profiles are typically flat for L-mode plasmas and evolve to hollow profiles during H mode in NSTX. The characteristic "ears" in the electron density during H mode seems to be the result of increased carbon ion density in the region of the edge density gradient.

# **Carbon Density and Z<sub>eff</sub> using Charge Exchange Recombination Spectroscopy**

The charge exchange recombination spectroscopy (CHERS) diagnostic observes the C VI line at 5290.5 Å. The emission due to charge exchange is given by

where  $B^{CX}$  is the measured brightness with the line integral along the intersection of the line of sight with the neutral beam.

The effective charge is defined as

Assuming c



# Z<sub>eff</sub> Profiles Using Charge Exchange Recombination Spectroscopy on NSTX R. E. Bell, T. M. Biewer, B. P. LeBlanc, V. A. Soukhanovskii

$$E^{CX} = n_c n_b \left\langle \sigma^{CX} v \right\rangle^{eff},$$

where  $n_c$  is the carbon density,  $n_b$  is the beam neutral density determined from a neutral beam attenuation calculation, and  $\langle \sigma^{CX} v \rangle^{e\!\!f\!f}$  is the effective emission rate due to charge exchange integrated over all beam components. An effective emission rate is determined using a collisional radiative model

The carbon density can be determined from the measured brightness,

$$P_C = \frac{B^{CX}}{\frac{1}{4\pi} \int n_b \left\langle \sigma^{CX} v \right\rangle^{eff} d\ell}$$

$$Z_{eff} = \frac{1}{n_e} \sum_{i} n_i Z_i^2.$$

$$Z_{eff} = 1 + \frac{n_c}{n} (Z_c^2 - Z_c) = 1 + 30 \frac{n_c}{n}$$

### **Z**<sub>eff</sub> using Visible Bremsstrahlung

Visible bremmstrahlung emission due to electron-ion collisions is a function of the electron density, electron temperature, and effective charge:

$$E^{VB} = 7.63 \times 10^{-15} \frac{n_e^2 Z_{eff} \overline{g}_{ff}}{T_e^{\frac{1}{2}} \lambda} \left[ \frac{photons}{s \text{ cm}^3 \text{ sr Å}} \right],$$

where  $n_e(cm^{-3})$  is the electron density,  $T_e(eV)$  is the electron temperature,  $\lambda(A)$  is the wavelength, and  $\overline{g}_{ff}(\lambda, T_e, Z_{eff})$  is the temperature averaged Gaunt factor.

The measured bremsstralung signal is a chord averaged measurement,

$$B_i^{VB} = \sum_i L_{ij} E_j^{VB} \quad ,$$

where  $L_{ii}$  is a length matrix and  $B_i^{VB}$  is the brightness through chord *i*. To obtain  $Z_{eff}$ profiles from a visible bremsstrahlung measurement, independent measurements of electron temperature and density and an Abel inversion are required. The plasma edge is troublesome to visible bremmstrahlung measurements: even with a region of the visible spectrum free of line radiation, other pollutant emission due to molecular emission is present in the visible.

Here we attempt to find a line-free region of the CHERS spectrum to measure the visible bremsstrahlung brightness along our sightlines. Comparing that measured value to a calculated value based on Zeff assuming carbon is the only impurity, we can get some confidence in the carbon density from charge exchange spectroscopy.

# **Z**<sub>eff</sub> and **Density Profiles** L-mode 0.230 sec 108420 0.227 sec RADIUS (cm) 5 109071 0.330 sec 5 109071 0.270 sec 0.267 sec RADIUS (cm) 0.510 sec 10906 0.517 sec

RADIUS (cm)

# **Steps to Determine** $Z_{eff}$ and $B_{VB}$

- $T_e$ ,  $n_e$  from Thomson scattering •  $n_{h}$  from beam attenuation calculation
- $< \sigma^{cx} v >^{eff}$  from integrating CX cross sections
- $B^{cx}$  from fit
- $n_c(B^{cx}, n_b, \langle \sigma^{cx}v \rangle^{eff})$  from fit
- $Z_{eff}$  assuming carbon is sole impurity
- $g_{ff}(\lambda, T_e, Z_{eff})$  computed
- $B_{VB}^{calc}$  from  $Z_{eff}$  (dashed line)
- $B_{VB}$  measured from CHERS spectra at  $\lambda_{VB}$  (solid line)





# **Fluorine Influx**

- Discrepancy between bremsstrahlung calculated using carbon density and measured bremsstrahlung for shot #109079 indicates that non-carbon emission was present
- Spectra from SPRED spectrometer show strong line emission from fluorine late in this discharge
- Fluorine density in core can be calculated based on "extra" bremsstrahlung emission using  $Z_F = 9$ .

**Z**<sub>eff</sub> Evolution



![](_page_0_Figure_51.jpeg)

![](_page_0_Figure_52.jpeg)

# 1084200.503.0<u>9</u> 0.40 2.5 0 ₽ 0.30 H-mode L-mode 0.20 100 110 120 130 140 150 Radius (em)

- Z<sub>eff</sub> typically flat during L-mode
- Carbon density profile peaks near edge after H-mode transition
- Z<sub>eff</sub> drops in core, increases at edge after Hmode transition
- Typical values  $Z_{eff} \sim 1.2-1.5$  in core,  $Z_{eff} \sim 3-$ 3.5 at edge during H-mode
- "Ear" in electron density during H mode due to carbon density peaking at edge

![](_page_0_Figure_59.jpeg)

- hollow during H mode due to accumulation of carbon near
- Particle transport barrier at edge

![](_page_0_Picture_67.jpeg)

### Summary

- Carbon density profiles measured using charge exchange recombination spectroscopy
- Carbon is the dominant impurity in NSTX
- $Z_{eff}$  profiles obtained using CHERS carbon density
- typically flat during L-mode
- plasma edge

### References

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