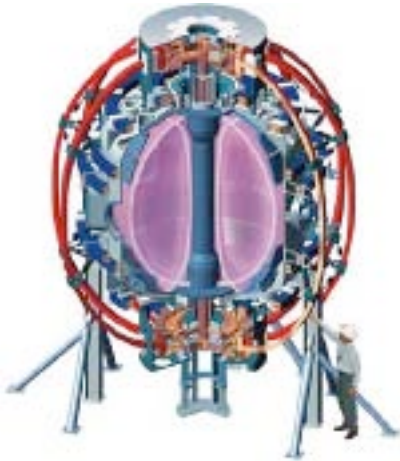


Investigation of fast-ion mode nonlinear dynamics and spatial structure in NSTX*



*by NA Crocker¹, WA Peebles¹, S. Kubota¹,
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Abstract (Revised)



Investigation of fast-ion mode nonlinear dynamics and spatial structure in NSTX*

NA Crocker, WA Peebles, SK Kubota (UCLA); ED Fredrickson, NN Gorelenkov, GJ Kramer, H Park (PPPL); WW Heidbrink (UCI); KC Lee, CW Domier, NC Luhmann Jr (UCD)

Neutral beam heated plasmas in NSTX exhibit a rich spectrum of fast-ion driven coherent modes that includes fishbones as well as toroidicity-induced and compressional Alfvén eigenmodes (TAE and CAE). These modes are of significant interest because they can induce fast-ion transport and channel fast-ion energy into the plasma. In recent experiments, the spatial structure of fishbone density perturbations has been investigated through the simultaneous application of a 288 GHz radial interferometer and three fixed-frequency microwave reflectometers operated by UCLA, three tangential far-infrared interferometers operated by UCD and an array of magnetic sensing coils external to the plasma. The coils and the UCLA diagnostics have also been utilized in a similar investigation of TAEs. The results may be compared with predictions from NOVA-K. Nonlinear three-wave interactions between fishbones, TAEs and CAEs are also studied. These interactions transfer energy in space and time and can significantly influence the effect of the modes on fast ions.

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Why study global modes excited by fast ions?

- **Fast ions important in fusion plasmas**

- produced by
 - heating techniques: neutral beams, radio frequency power
 - fusion products: alpha particles
- must be confined to heat plasma

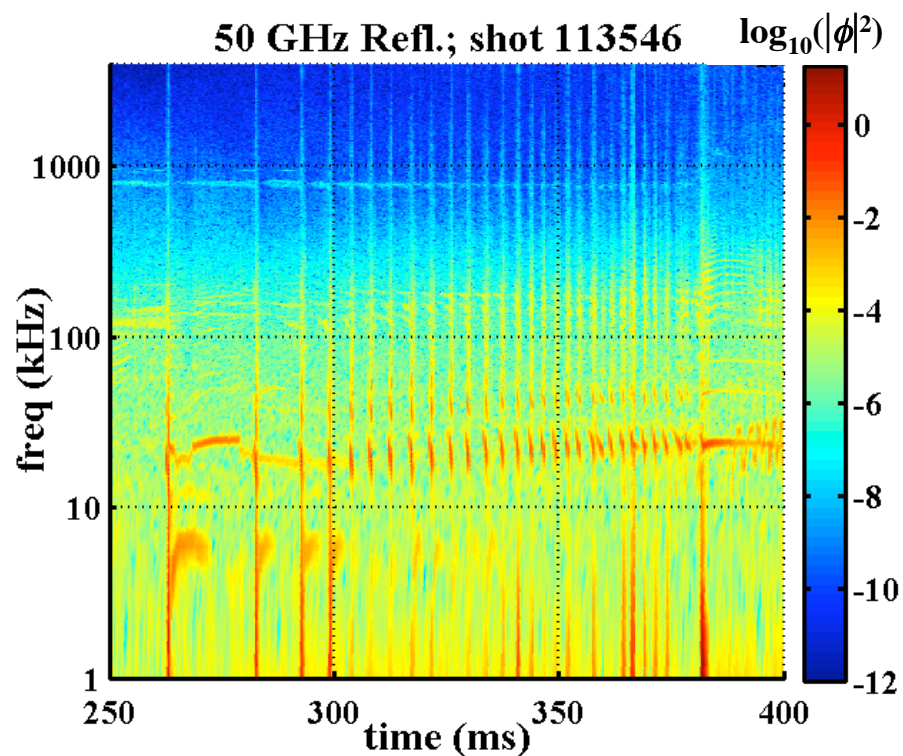
- **Fast-ion modes affect fast-ion transport**

- modes may modify orbits
 - redistribute fast ions \Rightarrow change heat deposition, force balance, etc.
 - degrade fast-ion confinement

Summary of Results

- Three wave interactions between fast-ion modes observed for first time
 - interaction across multiple scales: EPM—TAE, EPM—CAE, TAE—CAE
 - fast-ion loss events influenced
 - universal effect: wave-packet formation - lower frequency mode spatially concentrates energy of higher frequency mode
- TAEs and EPMs structure measured - observed with array of diagnostics:
 - multiple reflectometers \Rightarrow radial structure
 - radial & tangential interferometers (EPMs only for tang. interf.) \Rightarrow constrain reconstruction
 - toroidal array of Mirnov coils outside plasma \Rightarrow determine toroidal mode number
- TAE mode structure calculated by NOVA-K (linear stability code)
 - calculation uses conditions similar to experiment - some refinement necessary
 - future work - compare calculated TAE structure to measurements

Neutral beam heated plasmas in NSTX exhibit rich spectrum of fast-ion modes



- **Compressional Alfvén Eigenmodes (CAE)**
 - 0.4 to > 2 MHz
 - natural plasma resonance
- **Toroidal Alfvén Eigenmodes (TAE)**
 - 40 – 200 kHz
 - natural plasma resonance
- **Energetic Particle Modes (EPM)**
 - $\lesssim 100$ kHz
 - mode defined by fast-ion parameters
 - strong frequency chirping common

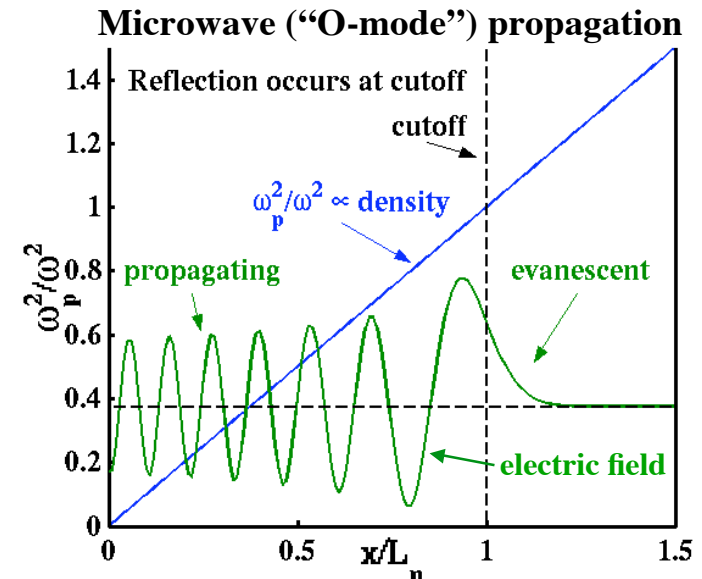
Fast-ion mode structure probed by internal density fluctuation diagnostics

- Multiple reflectometers measure local density perturbation \Rightarrow radial structure of mode (low field side)
 - interpretation of reflectometer signal for coherent modes confirmed by comparison with BES data on DIII-D.
- Radial 1mm & tangential FIR interferometers
 - survey of mode activity across entire plasma diameter
 - detection of modes localized on high field side
 - constrain reconstruction of spatial structure
- Plans to upgrade 1mm interferometer to multi-channel radially viewing polarimeter
 - allows measure of magnetic fluctuations

Reflectometers measure local density fluctuation in plasma

- Microwaves with low enough frequency ($\omega < \omega_p$) reflect from "cutoff" layer in plasma

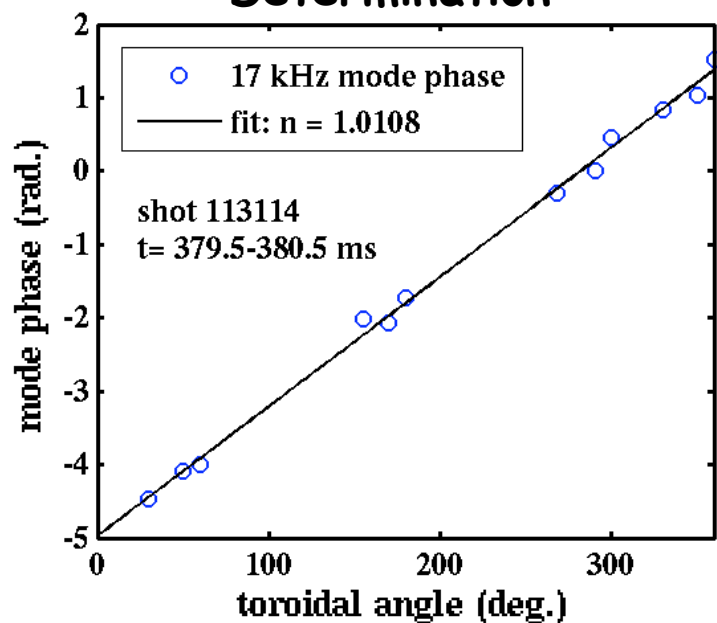
- ω_p^2 is proportional to density: $\omega_p^2 = e^2 n_0 / \epsilon_0 m_e$
- Dispersion relation for "ordinary mode" ("O-mode") microwaves: $\omega^2 = \omega_p^2 + c^2 k^2$
- $k \rightarrow 0$ as $\omega \rightarrow \omega_p$
- Microwaves reflect at "cutoff" surface, where $\omega = \omega_p$, $k = 0$



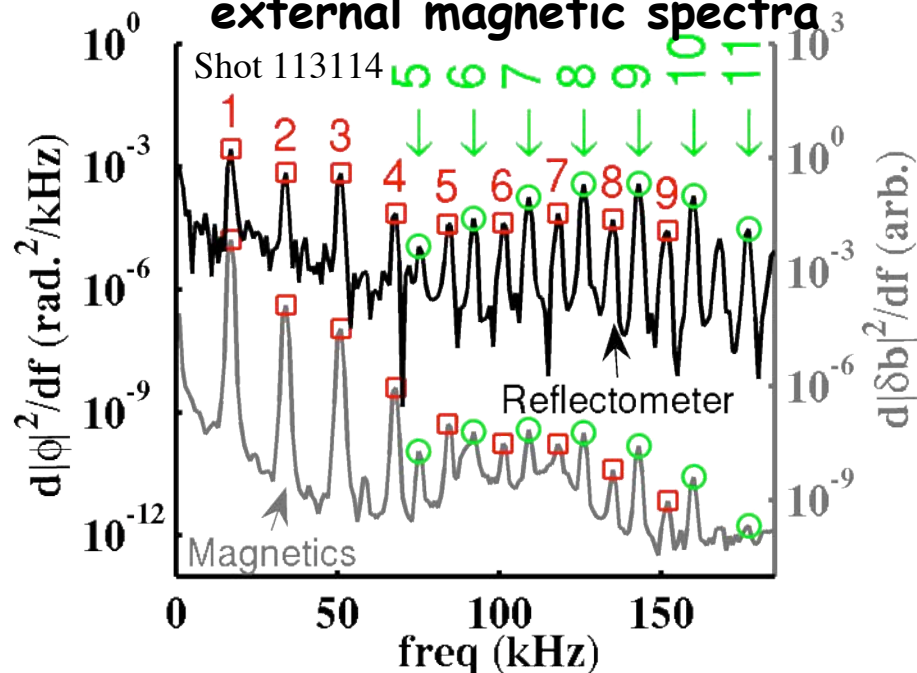
- Microwaves launched into plasma.
Relative phase of reflected and launched waves determined
- Wave propagation controlled by density \Rightarrow phase fluctuations proportional to density fluctuations (for large scale modes):
 $\delta n/n_0 \sim \delta \phi / (2k_{\text{vac}} L_n)$

Toroidal mode numbers of coherent modes may be determined from Mirnov array

Example of Mode Number Determination



50 GHz reflectometer and external magnetic spectra



- **Toroidal array of Mirnov coils outside plasma**
 - modes visible in magnetic spectrum
 - mode phase varies with coil position \Rightarrow measure toroidal mode number
- **Many peaks in refl. spectrum also seen in magnetic spectrum**
 - if peaks in both spectra, mode number of refl. perturbation known

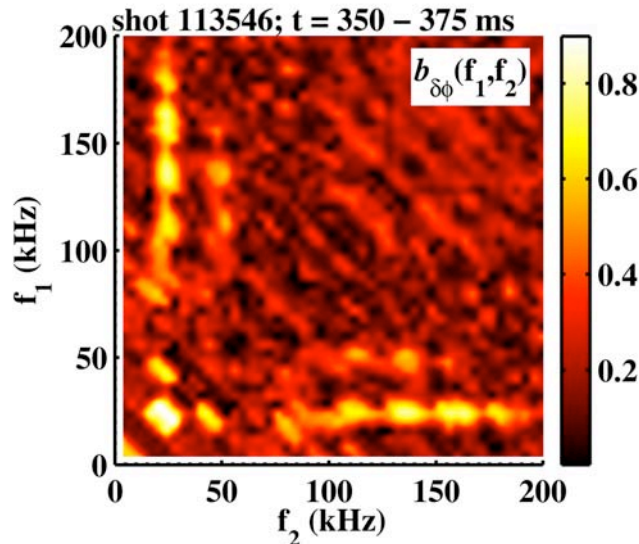
Three-wave interactions influence fast-ion loss

- **EPMs, TAEs active during fast-ion loss events:**

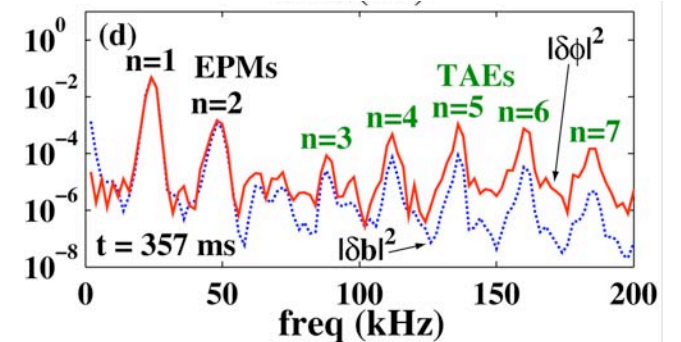
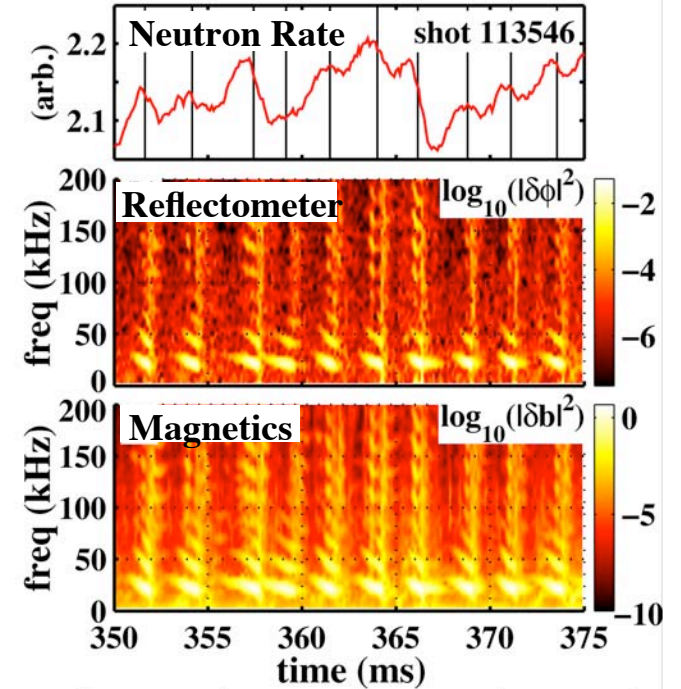
- EPM: Harmonic, $f \sim 24$ kHz, 48 kHz, $n = 1, 2$
- TAEs: $f \sim 80 - 200$ kHz, $n = 3 - 7$
 - uniformly spaced in f and n : $\Delta f \sim 24$ kHz, $\Delta n = 1$

- **EPMs interact with pairs of TAEs:**

- neighboring TAEs satisfy matching requirements:
 $\Delta f_{\text{TAE}} = f_{\text{EPM}}$ and $\Delta n_{\text{TAE}} = n_{\text{EPM}}$,
 so for TAE pair (f, n) and (f', n) :
 $f' = f + \Delta f_{\text{TAE}} = f + f_{\text{EPM}}$ and $n' = n + \Delta n_{\text{TAE}} = n + n_{\text{EPM}}$
- interacting modes show high bicoherence (bottom left)



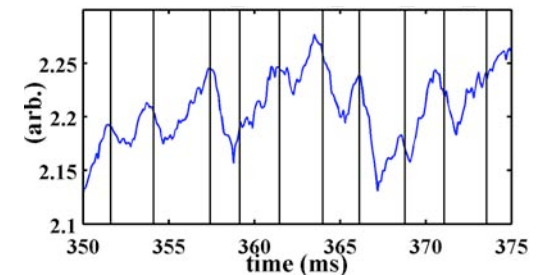
$$b_{\delta\phi}(f_1, f_2) = \frac{|\langle \delta\phi(f_1) \delta\phi(f_2) \delta\phi^*(f_1 + f_2) \rangle|}{(\langle |\delta\phi(f_1) \delta\phi(f_2)|^2 \rangle \langle |\delta\phi(f_1 + f_2)|^2 \rangle)^{1/2}}$$



Three-wave interactions couple disparate scales (TAEs and EPMs to CAEs)

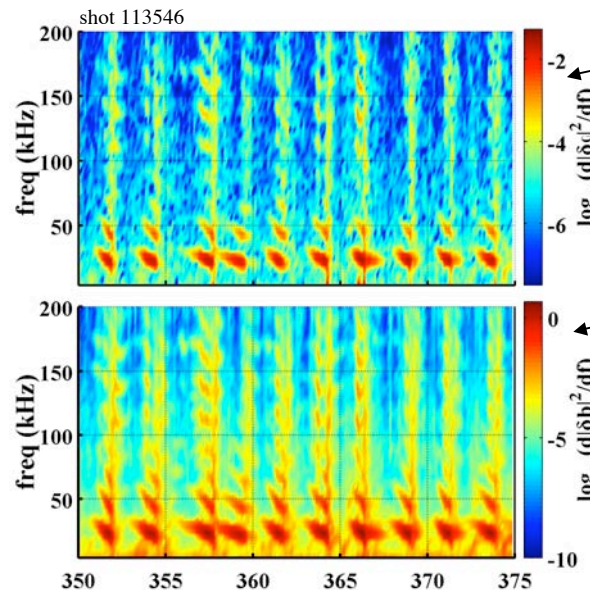
- CAE spectrum broadens during fast-ion loss events (drops in neutron rate) - sidebands appear
- bicoherence measurement indicates three-wave interactions cause broadening

Neutron Rate
(with fluctuation bursts marked)



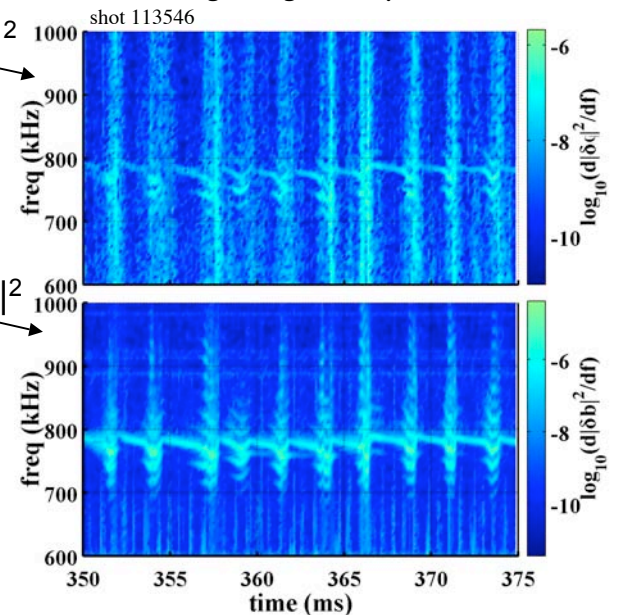
TAEs and EPMs

50 GHz reflectometer phase
and edge magnetic spectra

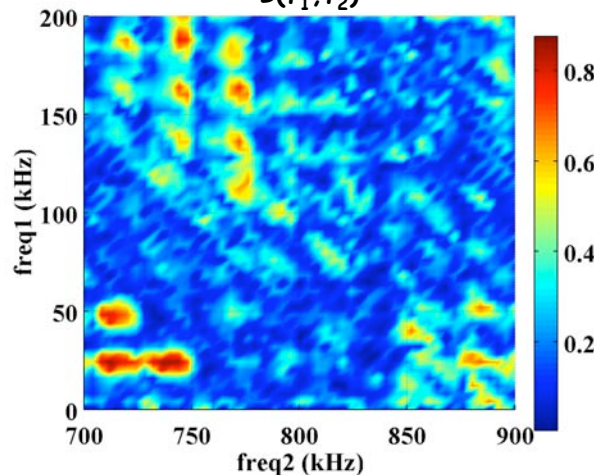


CAEs

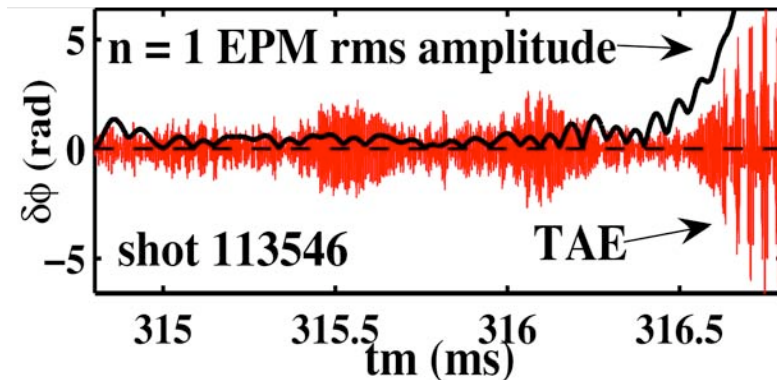
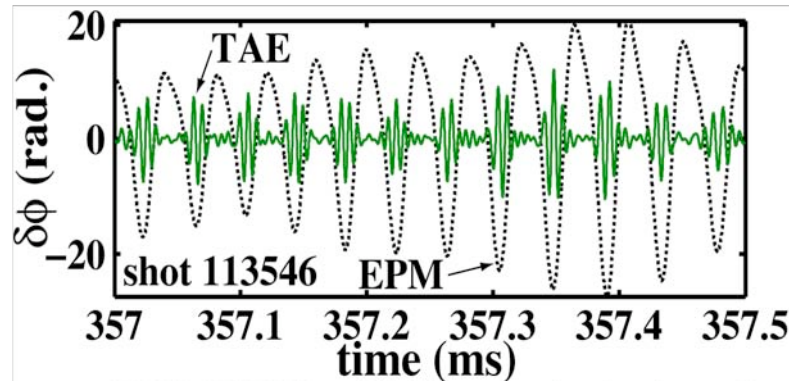
50 GHz reflectometer phase
and edge magnetic spectra



Bicoherence of δb
 $B(f_1, f_2)$



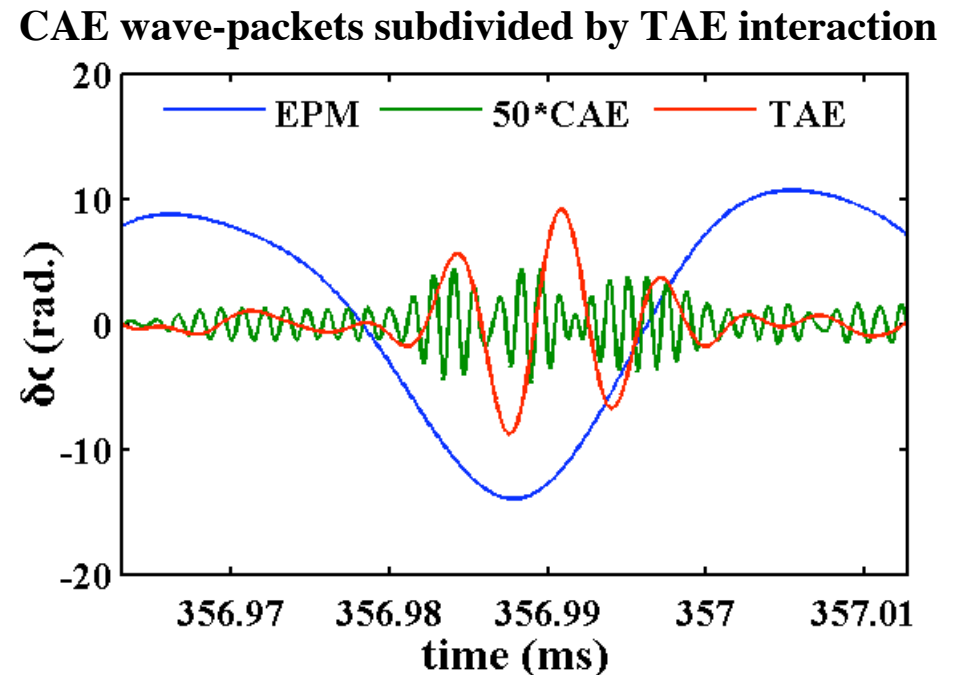
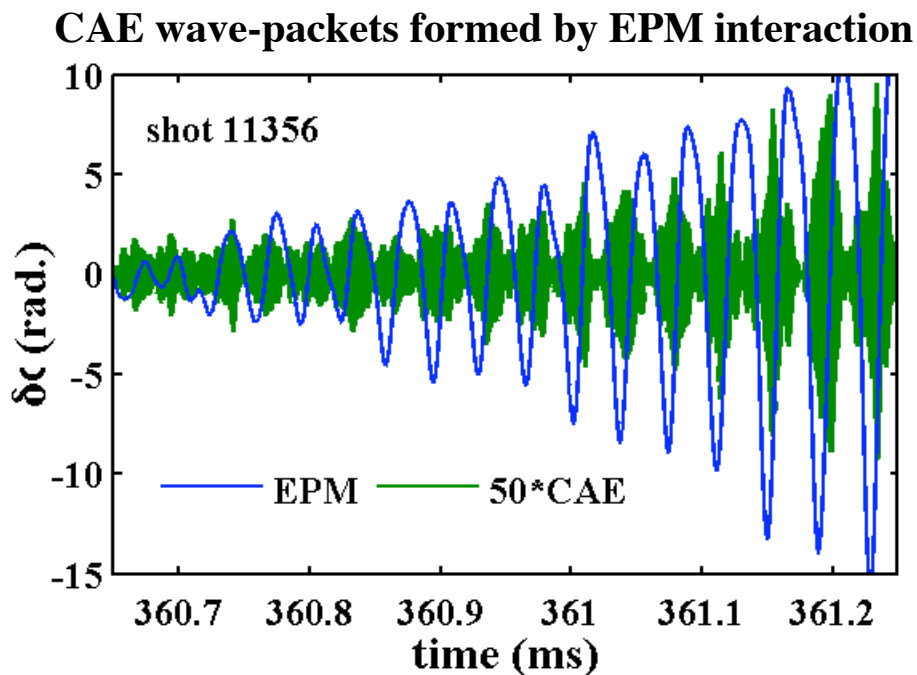
Three-wave interactions with EPMs spatially concentrates TAE energy



- **Interaction phase locks TAEs to form coherent structure:**
 - interaction with EPM TAEs imposes uniform frequency and wave number separation
 - uniform spacing \Rightarrow well defined group velocity = EPM phase velocity
 - TAEs form coherent (long-lived) structure that propagates in lock-step with EPM
- **Coherent TAE structure is toroidally localized “wave-packet” (top left)**
 - reflectometer phase band-pass filtered to extract TAE and EPM components
 - TAE component shows wave-packet: seen as pulse every time it passes reflectometer
 - pulse always occurs in same phase of EPM
- **TAE energy “spatially concentrated” ONLY when EPM active (top right)**
 - TAEs frequently active without EPMs
 - Wave-packet forms only when EPM strong

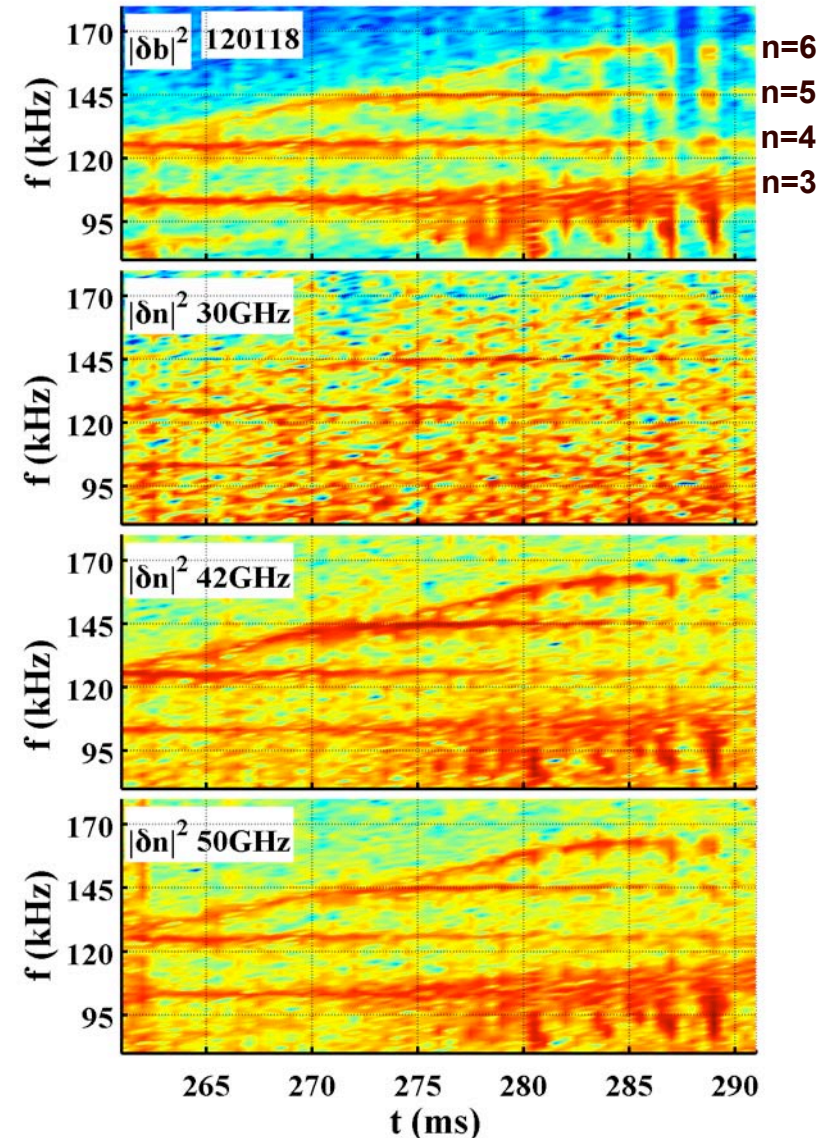
Three-wave interactions spatially concentrate CAE energy

- Interaction with EPMs forms CAEs into wave-packet (bottom left) - analogous to EPM—TAE interaction
- Interaction with TAEs subdivides CAE wave-packets into smaller packets (bottom right)
- Conjecture: high frequency wave-packets form because fast-ion population (i.e. free energy source) localized toroidally by low-frequency mode.



TAE spatial structure investigated

- TAEs measurements available from:
 - external toroidal Mirnov array (top right)
 - three fixed-frequency reflectometers (bottom right)
 - radial chord 1mm interferometer (not shown)
- TAE measurements can be exploited in several ways:
 - compare with NOVA-K — test usefulness to predict structure
 - understand effect on fast ions — compare with fast-ion measurements (NPA, SSNPA, sFLIP, neutrons, etc.)
 - learn to exploit diagnostic capabilities of TAEs — generally, AEs can burst or persist and chirp \Rightarrow behavior sensitive to fast-ion and plasma properties?



TAE structure evolves substantially over lifetime

• Typical TAE investigated

- frequency sweeps from $f \sim 120$ to 145 kHz during $t \sim 263 - t \sim 271$ ms
- frequency stable after up-sweep
- $n = 5$ (from external Mirnov array)

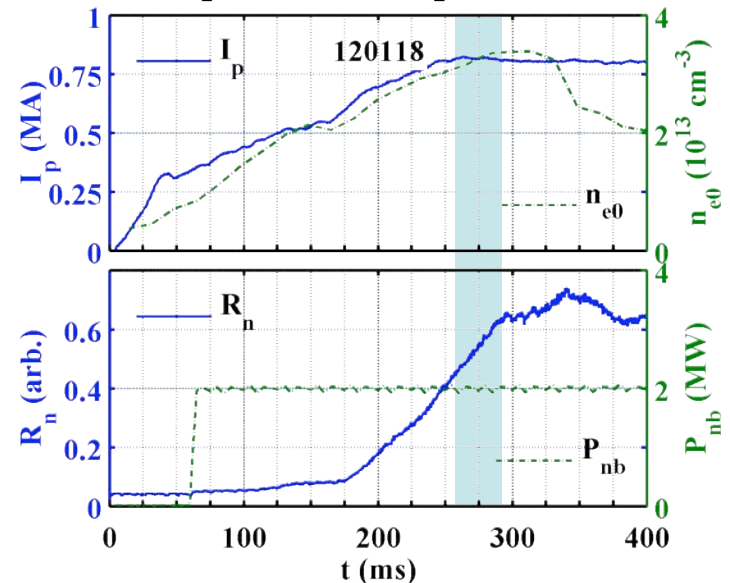
• $\Delta|\tilde{n}|/|\tilde{n}| \sim 300\%$ in 20 ms at $R \sim 122$ cm

- possible causes of change:
 - radial mode structure evolves significantly (e.g. mode peak shifts radially)
 - mode amplitude evolves
- Future work: use interferometer to distinguish between causes

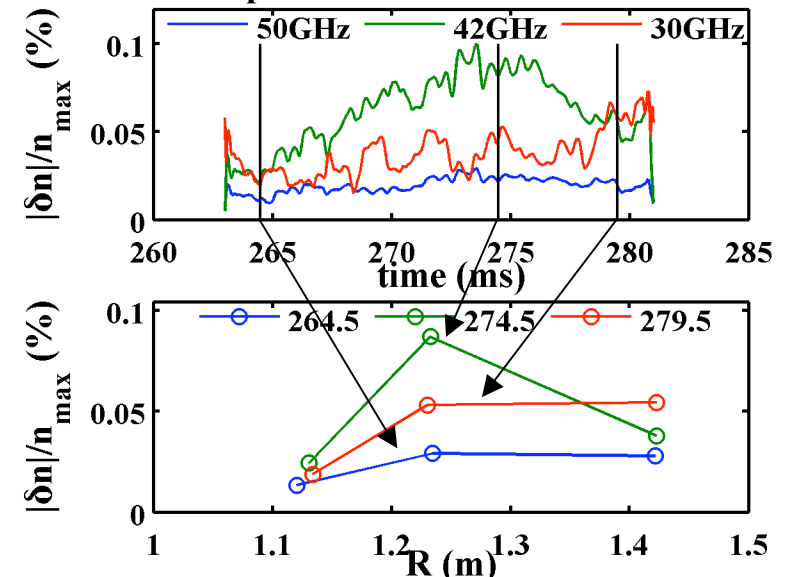
• Question: what causes rapid evolution?

- too rapid to be caused by equilibrium change?
- controlled by changes in fast-ion population? must compare with fast-ion diagnostics
- why is slope of f same for multiple modes? (see previous slide)

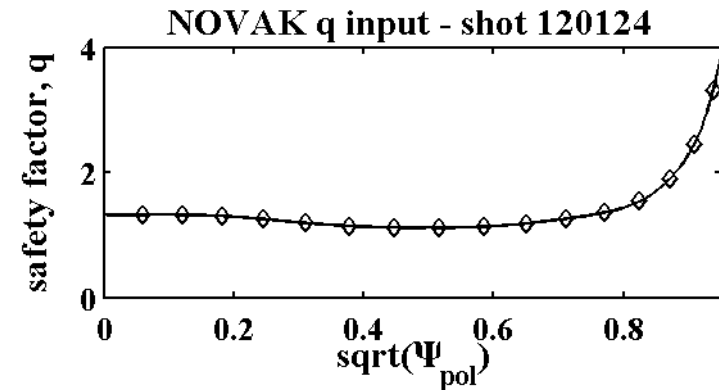
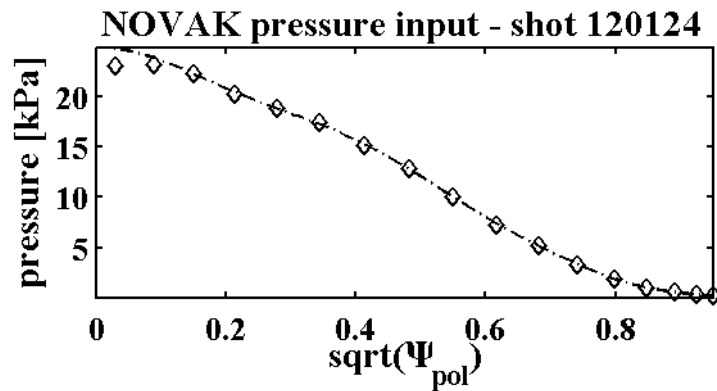
Plasma equilibrium temporal evolution



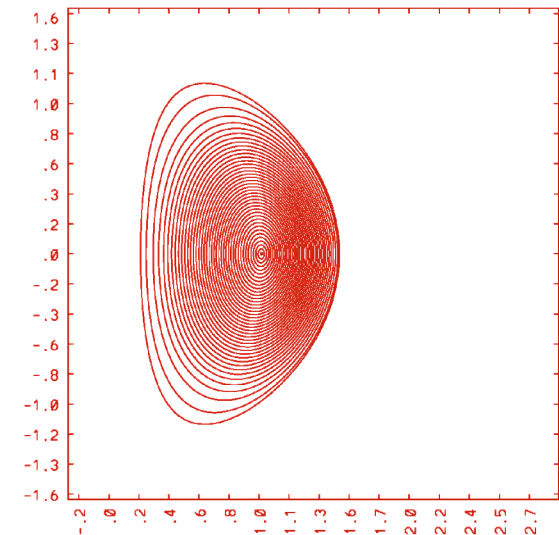
Mode amplitude - shot 120118 - $n = 5$ TAE



NOVA-K solves linear ideal MHD stability equation for eigenmodes



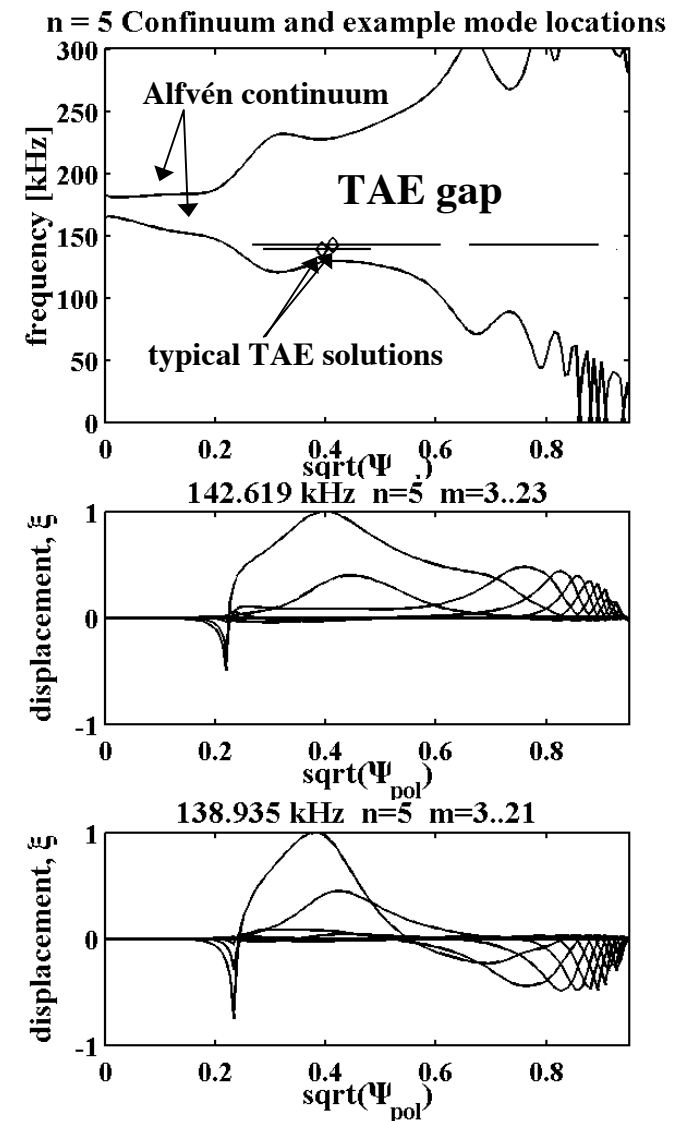
NOVAK input geometry - shot 120124



- **Solves in “perturbative” limit**
 - fast ions affect growth rate, not mode structure and frequency
- **Input from experiment:**
 - equilibrium geometry
 - pressure profile
 - q profiles
- **Solution assumes no rotation**
 - must account for rotation in experiment – typically, mode considered Doppler-shifted by rotation at mode peak

NOVA-K solutions include toroidicity-induced Alfvén eigenmodes (TAE)

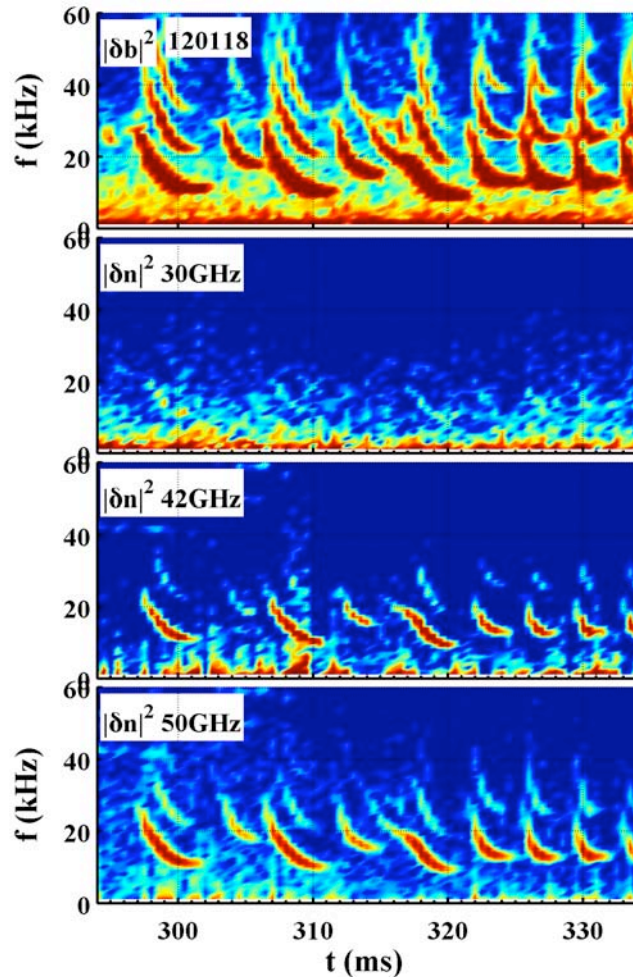
- TAEs are shear Alfvén waves
 - pure toroidal Fourier modes
 - multiple poloidal harmonics
- TAE frequencies lie in “TAE gap” in Alfvén continuum - heavily damped if intersect continuum
- Typical solutions shown on right
 - Calculated with “acoustic filtering” for speed. Must recalculate without. Structure may change substantially, but not frequency.
- Future work: Compare measurements to solutions with similar frequencies
 - frequency match may be inexact due to errors in plasmas equilibrium used by NOVA-K



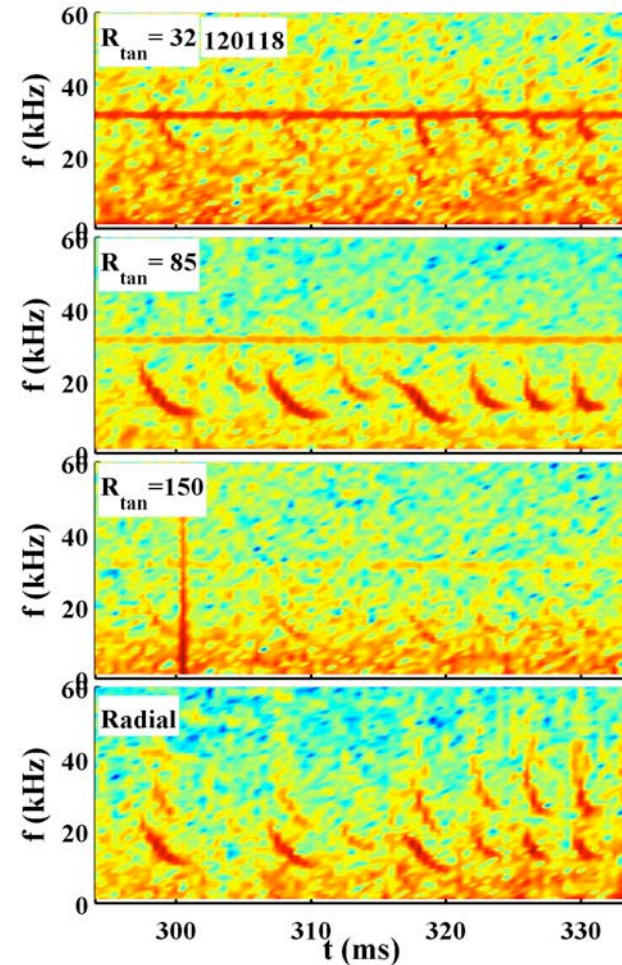
EPM spatial structure investigated

- EPM measurements available from:

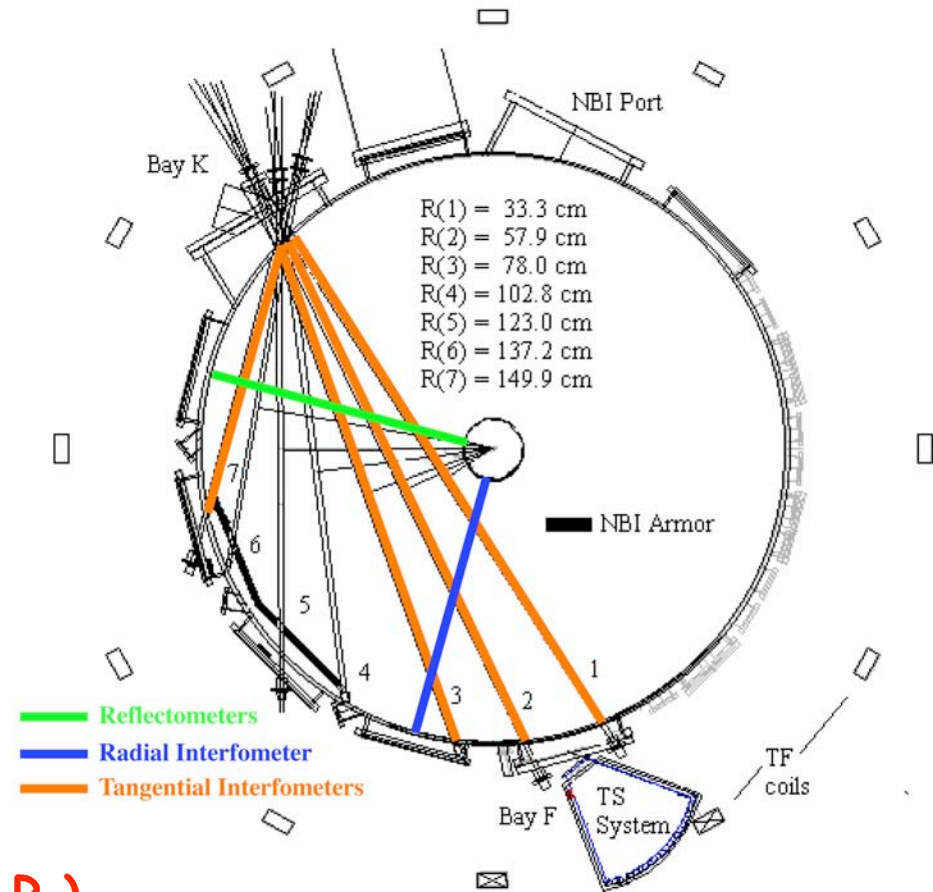
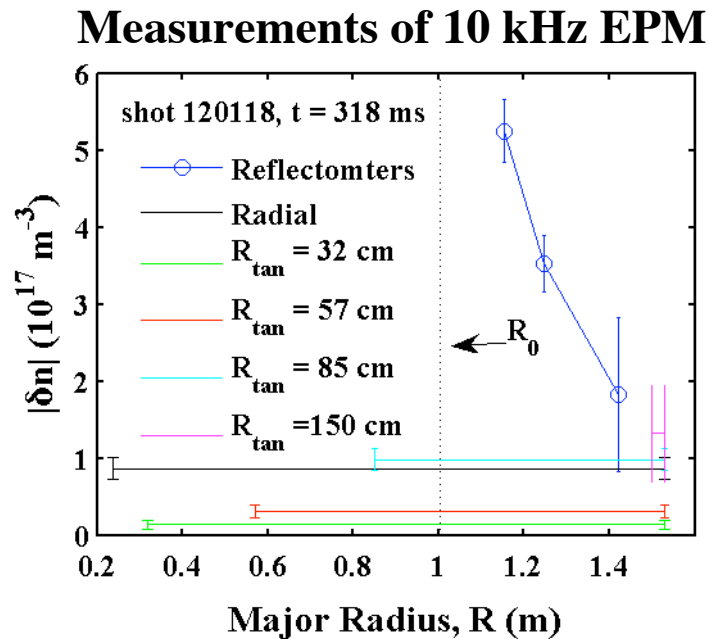
- external toroidal Mirnov array (top left)
- three fixed-frequency reflectometers (bottom 3 left)



- radial chord 1mm interferometer (bottom right) and tangential FIR interferometers (top 3 right)
- complementary data available from USXR chord arrays (not shown)



EPM measurements indicate centrally peaked structure



- Reflectometer measurements of 10 kHz EPM indicate centrally peaked structure (i.e. peak near R_0)
- Future work: use interferometer measurements to constrain reconstruction of EPM
 - choose basis for reconstruction (e.g. splines)
 - consider relative phase of signals
 - consider simplifying assumptions (e.g. assume mode is antisymmetric flux function)

Conclusions

- **Three wave interactions between fast-ion modes observed for first time**
 - interaction across multiple scales: EPM—TAE, EPM—CAE, TAE—CAE
 - fast-ion loss events influenced
 - **universal effect: wave-packet formation** - lower frequency mode spatially concentrates energy of higher frequency mode
- **TAEs and EPMs structure measured - observed with array of diagnostics:**
 - multiple reflectometers \Rightarrow radial structure
 - radial & tangential interferometers (**EPMs only for tang. interf.**) \Rightarrow constrain reconstruction
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- **TAE mode structure calculated by NOVA-K (linear stability code)**
 - calculation uses conditions similar to experiment - some refinement necessary
 - future work - compare calculated TAE structure to measurements

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