## Dependence of SOL widths on plasma parameters in NSTX

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## Dependence of SOL widths on nebar

## Edge density profile is 'fixed' during the H -mode density ramping period



- $\mathrm{T}_{\mathrm{e}, \text { sep }}$ estimated from 2 point model (i.e. from power balance)
- Re-align TS profiles as a function of R-Rsep with separatrix location determined by $T_{e, \text { sep }}$
- $\mathrm{n}_{\mathrm{e}, \text { sep }}$ is fixed irrespective of nebar

This enables us to ensemble TS profiles at different time slices for identical H -mode shots to construct a profile with higher spatial resolution

## Reliability of ensembled TS profiles



- $X^{2}$ stays low up to $T_{e, \text { sep }}=45 \mathrm{eV}\left(T_{e, \text { sep }}=35 \mathrm{eV}\right.$ from 2-point model $)$
- For $\mathrm{T}_{\mathrm{e}, \text { sep }}>45 \mathrm{eV}$, ensembled profile becomes more scattered
$\rightarrow \mathrm{T}_{\mathrm{e}, \text { sep }}=35 \mathrm{eV}$ appears to be in right ball park


## Density scan in ELM-free H-mode: $\lambda_{\text {ne }}$ stays constant



- nebar continuously rises during the H -mode, by a factor of $\sim 2$, with $\mathrm{n}_{\mathrm{e}, \text { sep }}$ fixed
- $\lambda_{\text {ne }}$ stays constant at $\sim 1.1 \mathrm{~cm}$


## Density scan in ELM-free H-mode: $\lambda_{\text {Te }}$ stays constant




- nebar continuously rises during the H -mode, by a factor of $\sim 2$, with $\mathrm{n}_{\mathrm{e}, \text { sep }}$ fixed
- $\lambda_{\text {Te }}$ stays constant at $\sim 0.9 \mathrm{~cm}$


## Density scan in ELM-free H-mode: $\lambda_{\mathrm{q}}$ stays constant



- nebar continuously rises during the H -mode, by a factor of $\sim 2$, with $\mathrm{n}_{\mathrm{e}, \text { sep }}$ fixed
- $\lambda_{\mathrm{q}}$ stays constant at $0.4-0.5 \mathrm{~cm}$


## Density scan result for Type-V ELMy H-mode: $\lambda_{q}$ stays constant



- nebar increase by $\sim 40 \% \rightarrow \lambda_{q}$ stays at $0.8-0.9 \mathrm{~cm}$
- $\lambda_{T e}$ and $\lambda_{\text {ne }}$ also stay constant irrespective of nebar


## Dependence of SOL widths on $I_{\mathrm{p}}$ and power

## $\lambda_{\mathrm{Te}}$ from TS measurement decreases with increasing Ip while $\lambda_{\text {ne }}$ stays constant, in Type-V ELMy H-mode



$$
\text { Ip: 1MA } \rightarrow \text { 800kA }
$$

- $\lambda_{\text {Te }}: 0.55 \mathrm{~cm} \rightarrow 0.75 \mathrm{~cm}$
- $\lambda_{\text {ne }}$ : staying at $\sim 0.65 \mathrm{~cm}$


## $\lambda_{\mathrm{Te}}$ from probe measurement decreases with increasing Ip while $\lambda_{\text {ne }}$ stays constant, in Type-V ELMy H-mode




Ip: 1MA $\rightarrow 0.8 \mathrm{MA}$
$\lambda_{\text {Te }}$ increases by a factor of $\sim 2$

- $\lambda_{\text {ne }}$ remains unchanged


## $\lambda_{\text {jsat }}$ and $\lambda_{q}$ decrease with increasing Ip in Type-V ELMy H-mode


$R$-Rsep for fast probe $J_{\text {sat }}$ data is from equilibrium reconstruction


Ip: 1MA $\rightarrow 700 \mathrm{kA}$

- $\lambda_{\text {jsat }}: 0.4 \mathrm{~cm} \rightarrow 0.9 \mathrm{~cm}$
- $\lambda_{\mathrm{q}}: 0.5 \mathrm{~cm} \rightarrow 1.1 \mathrm{~cm}$


## Power scan result: $\lambda_{\mathrm{Te}}$ and $\lambda_{\mathrm{q}}$ decrease with increasing power while $\lambda$ ne stays constant, in Type-V ELMy H-mode



$\mathrm{P}_{\mathrm{NBI}}: 1 \mathrm{MW} \rightarrow 2 \mathrm{MW}$

- $\lambda_{\mathrm{Te}}: 0.6 \mathrm{~cm} \rightarrow 0.44 \mathrm{~cm}$
- $\lambda_{\text {ne }}$ : staying at $\sim 0.65 \mathrm{~cm}$
- $\lambda_{\mathrm{q}}: 0.69 \mathrm{~cm} \rightarrow 0.54 \mathrm{~cm}$


## Comparison of observed dependences with H-mode SOL width scaling laws from conventional tokamaks

$$
\begin{array}{ll}
\lambda_{T e} \propto n_{e}^{0.92 \pm 0.18} I_{p}^{-1.79 \pm 0.27}\left(P_{t o t}-P_{r a d}\right)^{-0.63 \pm 0.09} & \text { from ASDEX-U1 } \\
\lambda_{n e} \propto n_{e}^{1.11 \pm 0.13} I_{p}^{-2.25 \pm 0.16} & \text { from ASDEX-U1 } \\
\lambda_{q} \propto B_{t}^{-0.93} q_{95}^{0.41} P_{S O L}^{-0.48} n_{e, u}^{0.15} & \text { from JET }{ }^{2}
\end{array}
$$

- $\lambda_{T e}, \lambda_{\text {ne }}$, and $\lambda_{q}$ dependence on $I_{p}$ and power: consistent with scaling laws
- $\lambda_{\mathrm{q}}$ dependence on density: consistent with scaling law
- $\lambda_{\mathrm{Te}}$ and $\lambda_{\mathrm{ne}}$ dependence on density: different trend from scaling laws
${ }^{1}$ K. McCormick, J. Nuclear Material 266-269 (1999) 99
${ }^{2}$ W.Fundamenski, Nucl. Fusion 44 (2004) 20


## Role of ELMs in determining SOL widths

## $\lambda_{\mathrm{Te}}$ is strongly affected by ELMs and turbulent blobs



- Probe measurement is continuously affected by ELMs and blobs
$\rightarrow$ measured Te shows high scatter
$\rightarrow$ Te SOL width broadens
- Probe I-V data with ELM affected portions removed
$\rightarrow$ re-process probe data
$\rightarrow$ Te SOL width becomes narrower
- TS measurement is instantaneous
$\rightarrow$ misses many ELM filaments in the near SOL
$\rightarrow$ effectively represents inter-ELM profile with narrower $\lambda_{\text {Te }}$
$\rightarrow$ Probe $\lambda_{\text {Te }}$ (with ELMs and blobs removed) similar to TS $\lambda_{\text {Te }}$
$\mathrm{T}_{\mathrm{e}}$ profile is highly affected by ELMs and turbulent blobs


## $\lambda_{\mathrm{ne}}$ is little affected by ELMs and turbulent blobs



- $\lambda_{\text {ne }}$ from probe is only a little broader than $\lambda_{\text {ne }}$ from TS
- Probe data 'without ELMs and bobs' produces only a little narrower $\lambda_{\text {ne }}$, compared with $\lambda_{\mathrm{ne}}$ 'with ELMs and blobs'
- Change in Te affects density only to a limited extent because of stronger contribution of jsat (i.e., $n_{e} \propto I_{\text {sat }}^{+} / \sqrt{T_{e}}$ )
'Time-average' density profile is not sensitive to ELMs and turbulent blobs


## ELMs may be responsible for broad Te and heat flux profiles

$$
\left.\begin{array}{l}
\text { 1. Characteristics of Type-V ELMs }{ }^{1}: \\
\mathrm{n}=3-4, \mathrm{~L}_{\Phi} \sim 0.15 \mathrm{~m}, \mathrm{~L}_{\mathrm{r}} \sim 0.1 \mathrm{~m}, \mathrm{~L}_{\mathrm{z}} \sim 0.1 \mathrm{~m}, \\
\mathrm{v}_{\Phi} \sim 8.6 \mathrm{~km} / \mathrm{s}, \mathrm{v}_{\mathrm{r}} \sim 0.2 \mathrm{~km} / \mathrm{s} \\
\text { 2. Outer circumference of plasma } \\
\text { surface at midplane }=9 \mathrm{~m}
\end{array}\right\}
$$

ELM toroidal coverage at outer midplane $=(0.45-0.6) / 9=5-7 \%$

- TS measurement is made at a specific toroidal location instantaneously $\rightarrow$ Probability of a specific TS measurement to detect ELM is $5-7 \%$
$\rightarrow$ Ensemble averaged TS Te profile effectively represents inter-ELM profile
- Probe continuously measures IV curves with $\mathrm{t}_{\text {sweep }}=0.25 \mathrm{~ms}$ and is hit by toroidally rotating ELM filaments, while moving across SOL
$\rightarrow$ Probe Te profile represents an ELM-averaged one (broad profile)
- IR camera measures heat flux profile for $33 \mathrm{~ms} \rightarrow$ reflects multiple ELM stripes at the divertor target (broader profile than without ELMs)

[^0]
## TS measurement for ELM-free H -mode provides evidence of the role of ELM filaments in determining $\lambda_{T \mathrm{e}} / \lambda_{\mathrm{q}}$ ratio



1. $\lambda_{\mathrm{Te}} / \lambda_{\mathrm{q}}=\frac{7}{2}\left(\frac{T_{e}-T_{e 1}}{T_{e}-C q_{1} T_{e}^{-5 / 2}}\right)^{1}$ is expected if e-conduction is dominant, in light of flat $\mathrm{T}_{\mathrm{e}}$ and heat flux profiles in the far SOL (i.e., $\lambda_{T \mathrm{e}} / \lambda_{\mathrm{a}} \sim 2$ )
2. Experimental results:

- Type-V ELM H-mode: $\lambda_{\mathrm{Te}} / \lambda_{\mathrm{q}} \sim 2$ (Probe vs IR)
- Type-V ELM H-mode: $\lambda_{\mathrm{Te}} / \lambda_{\mathrm{a}} \sim 1$ (TS vs IR)
- ELM-free H-mode: $\lambda_{\mathrm{Te}} / \lambda_{\mathrm{q}} \sim 2(\mathrm{TS}$ vs IR)

More evidence is to be investigated in the next campaign:

- Take probe measurement for ELM-free H-mode and compare with IR
- Take fast ( $\sim 20 \mathrm{kHz}$ ) IR camera measurement to avoid effect of ELM stripes
${ }^{1}$ J-W. Ahn, submitted to Phys. Plasmas (2008)


## Summary and conclusions

1. Dependences of $\lambda_{\mathrm{T}}, \lambda_{\mathrm{ne}}, \lambda_{\mathrm{q}}$, and $\lambda_{\text {jsat }}$ on $\mathrm{I}_{\mathrm{p}}$, nebar, and input power in H -mode plasmas were identified negative dependence on $I_{p}$ and power

## no dependence on nebar

$\lambda_{\mathrm{q}}$ change appears to be driven primarily by $\lambda_{\mathrm{Te}}$ change
2. $\lambda_{\text {ne }}$ remains largely unchanged by any parameter scan, due to the stiff edge density gradient in H -mode
3. Comparison with H -mode scaling laws from conventional tokamaks
$\lambda_{\mathrm{Te}}, \lambda_{\mathrm{ne}}$, and $\lambda_{\mathrm{q}}$ dependence on $\mathrm{I}_{\mathrm{p}}$ and power is consistent with scaling laws
$\lambda_{\mathrm{q}}$ dependence on density also consistent with the scaling law $\lambda_{\text {Te }}$ and $\lambda_{\text {ne }}$ dependence on density shows different trend
4. ELM filaments appear to broaden Te and heat flux SOL widths Probe and ensembled TS measurements represent different aspect of type-V ELMy H-mode profiles, i.e. ELM-averaged and inter-ELM profiles, respectively

## Future Work

1. More data points and confidence

SOL width scaling laws for $\lambda_{\mathrm{q}}, \lambda_{\mathrm{Te}}$ and $\lambda_{\mathrm{ne}}$
Comparison with theoretical models eg, SOLPS results and analytic $\lambda_{q}$ models
Extrapolation to future machines
2. More detailed investigation of parallel SOL transport

Roles of parallel convection to account for remaining discrepancy in Near SOL, using the Mach probe in 2009
Probe measurement in ELM-free H-mode to double-check the role of e-conduction in parallel transport
3. Fast IR camera measurement to separate ELM effects
4. Investigate relation of SOL widths with edge turbulence characteristics, eg, Probe $\mathrm{I}_{\text {sat }}$ data, GPI data, Edge reflectometry data Investigation of perpendicular SOL transport

## Sign up for pre-print!


[^0]:    ${ }^{1}$ R. Maingi, et al., Phys. Plasmas 13, 092510 (2006)

