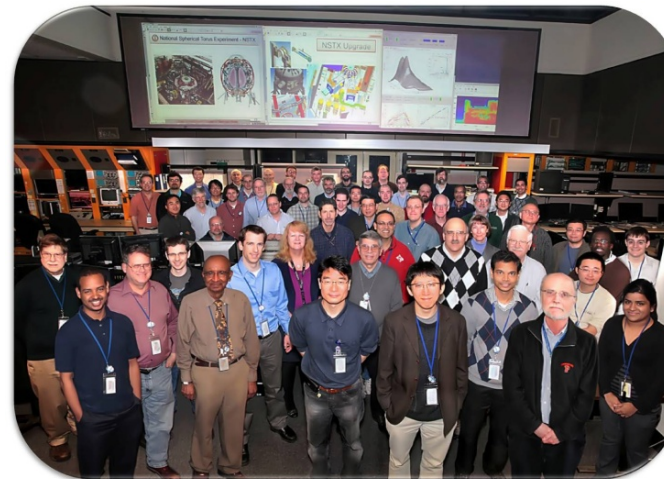
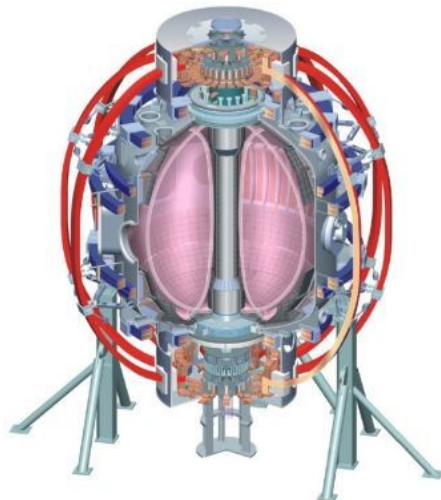


L-H power threshold scaling with magnetic geometry on NSTX and the role of ion orbit loss

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and the NSTX Research Team

**53rd APS DPP Meeting
Salt Lake City, UT
November 16, 2011**



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Outline

- X-transport: suppression of neoclassical ion orbit loss is a constraint on the edge E_r
 - Constraint on E_r scales with T_i and magnetic geometry, especially R_X
 - $E_r \times B$ shear predicted to play a role in L-H transition
- NSTX results: lend strong support for the X-transport theory
 - Critical edge T_i , and thus P_{LH} , increase as R_X moves inward
 - Increase in T_i nearly matches increase in calculated critical ion loss energy
 - Critical T_i mostly independent on divertor pumping and CS fueling
 - But large variation in $P_{LH} \rightarrow$ these things impact coupling between core heating and edge T_i
- X-transport should be included in L-H transition models
 - “Hidden variables” of L-H transition qualitatively consistent with X-transport physics

Connection between E_r and L-H transition

- Steady-state E_r maintains ambipolar transport

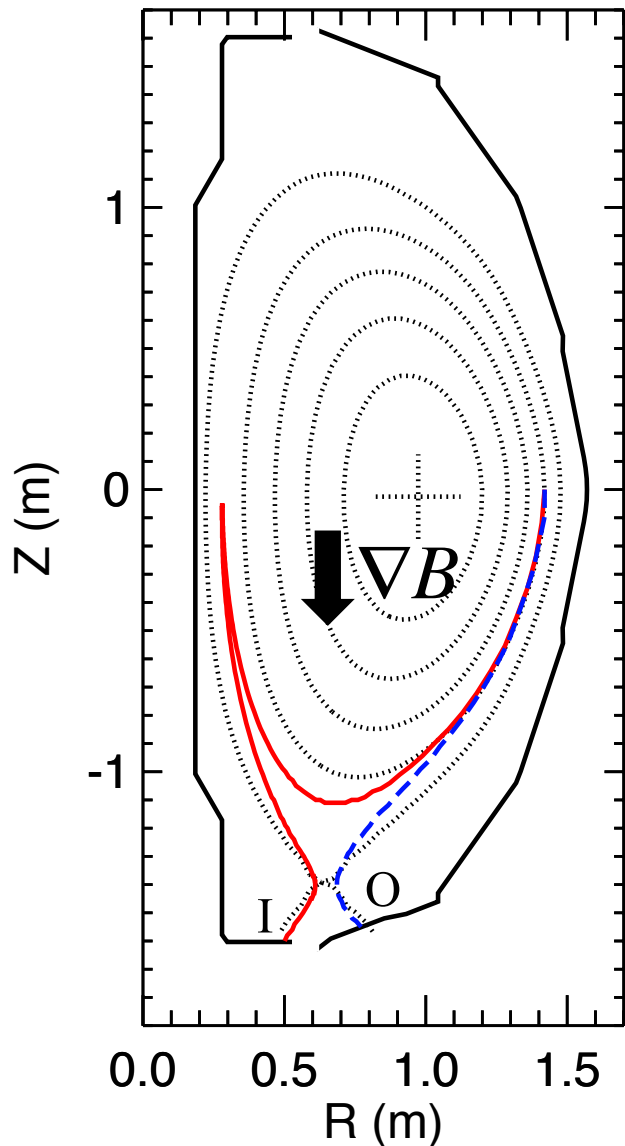
$$\varepsilon \frac{\partial}{\partial t} E_r = e \left[\Gamma_e(E_r) - \Gamma_i(E_r) \right] = 0$$

- Self-consistent E_r calculation must consider all non-ambipolar transport processes
 - X-transport: significant non-ambipolar process near plasma edge
- Edge $E_r \times B$ shear thought to play a role in L-H transition
 - Transport suppression mechanism still under investigation
 - Possible: decorrelation of turbulent eddies by $E_r \times B$ flow shear ^[1]
 - Possible: interaction of mean and turbulent driven $E_r \times B$ flow shear ^[2]
 - May be a critical $E_r \times B$ shear for triggering L-H transition
 - $E_r \times B$ shear \sim minimum in E_r well

[1] K.H. Burrell, Phys. Plasmas, **4** (1997) 1499-5185

[2] G.D. Conway, et. al. Phys. Rev. Lett. **106** (2011) 065001

X-transport: suppression of non-ambipolar transport of ions on neoclassical orbits constrains the edge E_r



- X-point amplifies grad-B drift
 - X-point: low B_θ , slows poloidal transit
 - Non-ambipolar: ion drift \gg electron drift
- Lowest energy loss orbits:
 - Start at outboard midplane
 - Bounce at inboard midplane
 - Lost to inner divertor leg in favorable grad-B
- Negative E_r acts to confine ions
 - Constraint on E_r : must be negative enough to nearly suppress non-ambipolar ion loss

C.S. Chang, S. Ku, H. Weitzner, *Phys. Plasmas*, **9** (2002)

Ion velocity loss hole has a critical energy near ion thermal energy in edge region when $E_r = 0$

- Single particle guiding-center orbit tracing with $E_r = 0$, no collisions

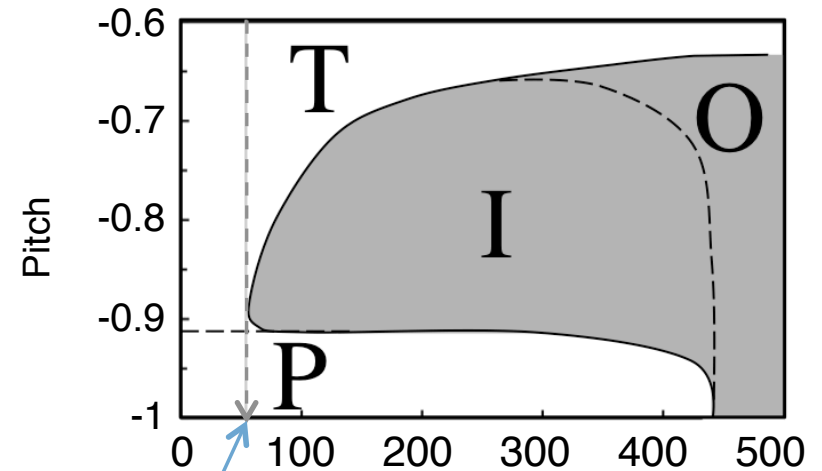
- White: Confined orbits
 - T: Trapped, P: Passing
- Gray: Unconfined orbit
 - I: Strike inner div
 - O: Strike outer div

- K_{crit} within Maxwellian T_i in edge
 - X-transport important only in edge pedestal region
 - Negative E_r pushes K_{crit} curve to higher energies

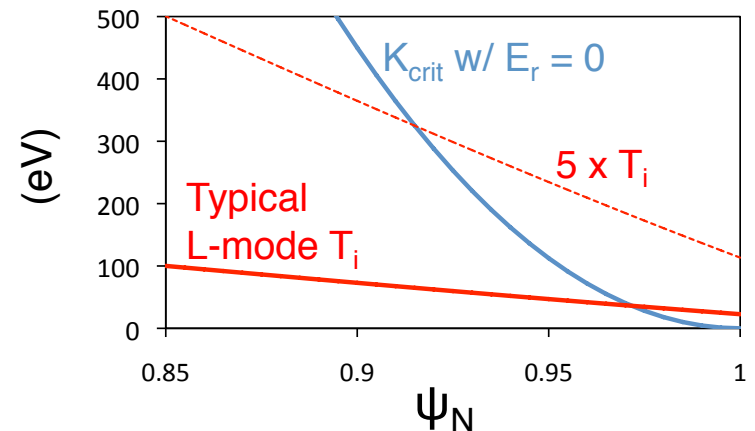
Guiding-center orbit tracing code:

S. Ku, H. Baek, C. S. Chang, *Phys. Plasmas* **11** (2004)

$\psi_N = 0.96$ (~1 cm from separatrix) at midplane



Critical ion energy for loss:
 $K_{crit} = 71\text{eV}$



Analytical model for ion orbit loss with $E_r = 0$ illustrates impact of plasma parameters on K_{crit}

- K_{crit} : Critical energy for collisionless ion loss with $E_r = 0$
 - B, B_X, B_B : Magnetic field at launch point, bounce point & X-point
 - $\Delta\psi = \psi_X - \psi_{launch}$, $\Delta r = R_{out} - R_{launch}$ on midplane

$$K_{crit} = \frac{1}{8} \left(\frac{e}{\pi c} \right)^2 \frac{Z^2}{M_i} \Delta\psi^2 \left[\frac{1}{B} \sqrt{1 - \frac{B}{B_B}} + \frac{1}{B_X} \sqrt{1 - \frac{B_X}{B_B}} \right]^{-2}$$

A.V. Chankin, G.M. McCracken,
Nucl. Fusion **33** (1993)

$$\sim \frac{Z^2}{M_i} (\Delta r)^2 \left(\frac{I_P}{C_P} \right)^2 \left[\sqrt{1 - \frac{R_{in}}{R_{out}}} + \frac{R_X}{R_{out}} \sqrt{1 - \frac{R_{in}}{R_{out}}} \right]^{-2}$$

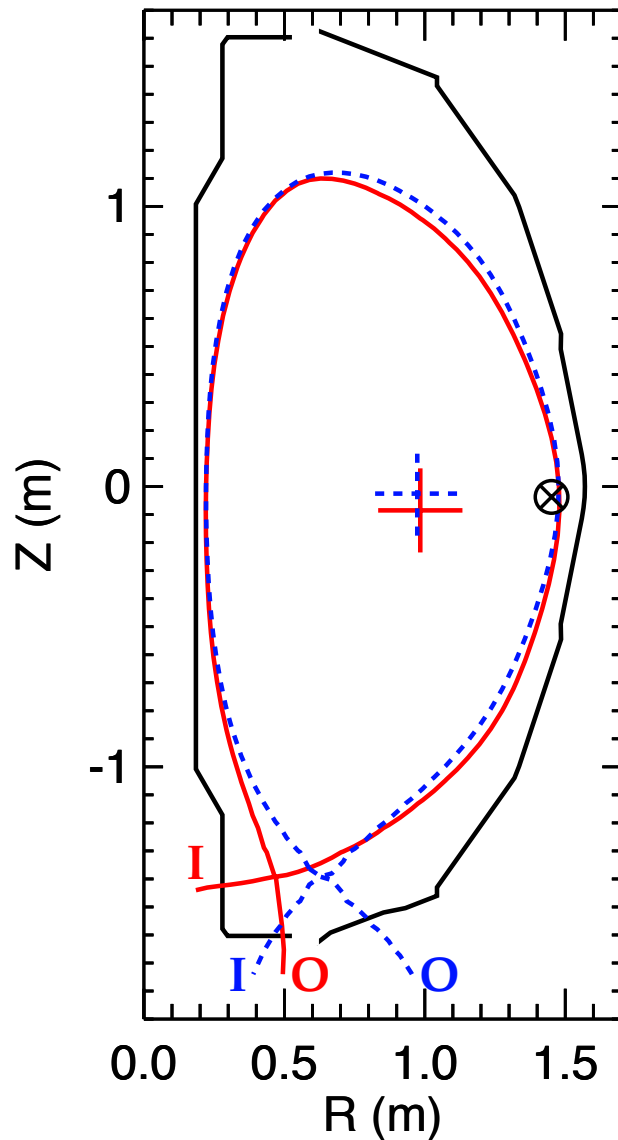
Launch at outboard midplane, $E_r = 0$
(Large aspect ratio approximation)

D: 0.5
H: 1.0
He: 1.0

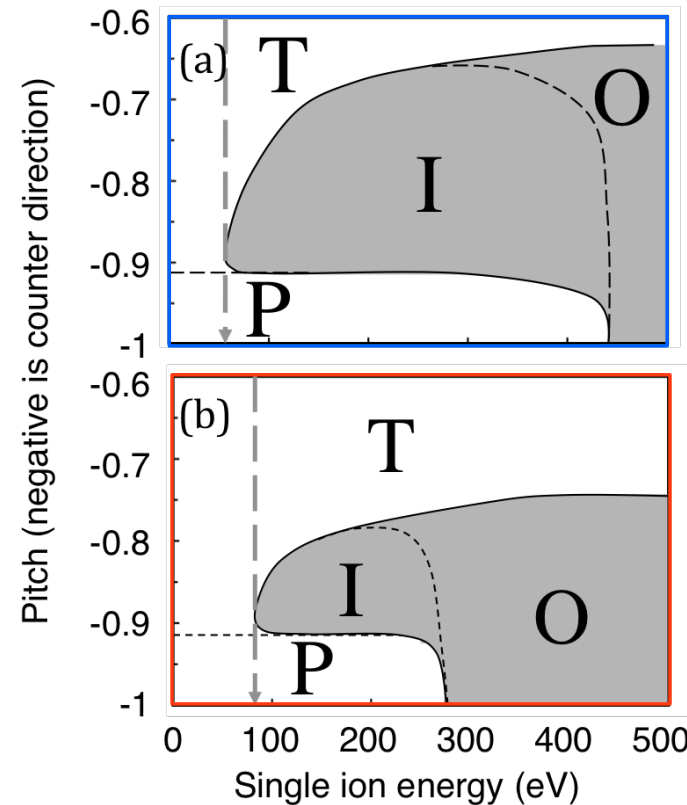
$K_{crit} \sim 2.5$ times larger as R_X goes from R_0 to R_{in}
 K_{crit} decreases with aspect ratio, provided $B_{\phi,out} \gg B_{\theta,out}$

K_{crit} increases with larger I_p and smaller plasma circumference (C_p)

Critical energy for ion loss increases as R_x decreases



$\psi_N = 0.96$ (~ 1 cm from separatrix) at midplane



Large- R_x
(Low δ)
 $K_{i,crit} = 71$ eV

Small- R_x
(High δ)
 $K_{i,crit} = 96$ eV

- Critical energy increases $\sim 35\%$ as R_x moves in: $0.64 \rightarrow 0.47$ m

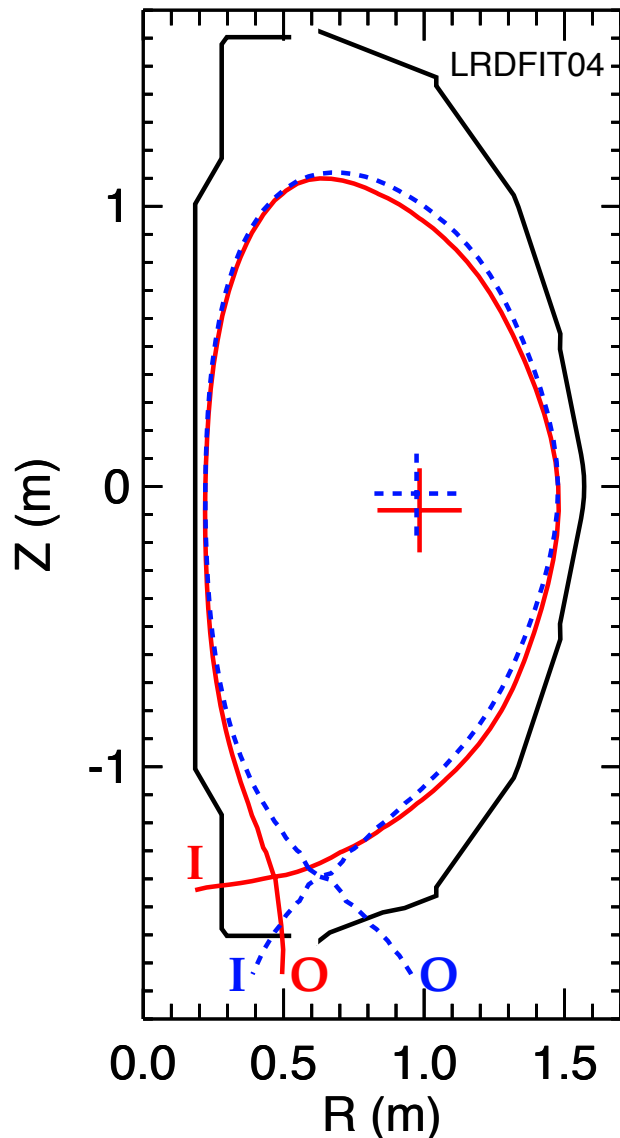
Full self-consistent E_r solution must consider collisions

- Low collisionality regime (ion loss time * $\nu_{ii} \ll 1$)
 - Ions scattered into loss hole are lost on collisionless orbits
 - Increase collisionality: more ions scattered into loss hole → more negative E_r
- High collisionality regime
 - Ions scattered into loss hole do not complete collisionless orbit
 - Increase collisionality: fewer collisionless ion loss orbits → more positive E_r
- Ion loss time impacts edge E_r
 - Grad-B drift direction: longer ion loss orbit time in unfavorable
 - X-point region: large null region (snowflake) increases ion loss time

Connection between X-transport theory and L-H transition

- Ion loss hole makes a significant contribution to E_r at edge
 - Contribution scales with T_i (Neoclassical scales with $\text{grad-}T_i$)
- For a given T_i , if K_{crit} is lowered, E_r becomes more negative
- More negative E_r leads to a deeper E_r well
 - Larger $dE_r/dr \rightarrow$ increased shearing rate
- If L-H transition at critical dE_r/dr , then transition at critical T_i
 - Critical T_i depends on plasma geometry and collisionality
 - Expect edge T_i and heating power to be connected $\rightarrow P_{\text{th}}$

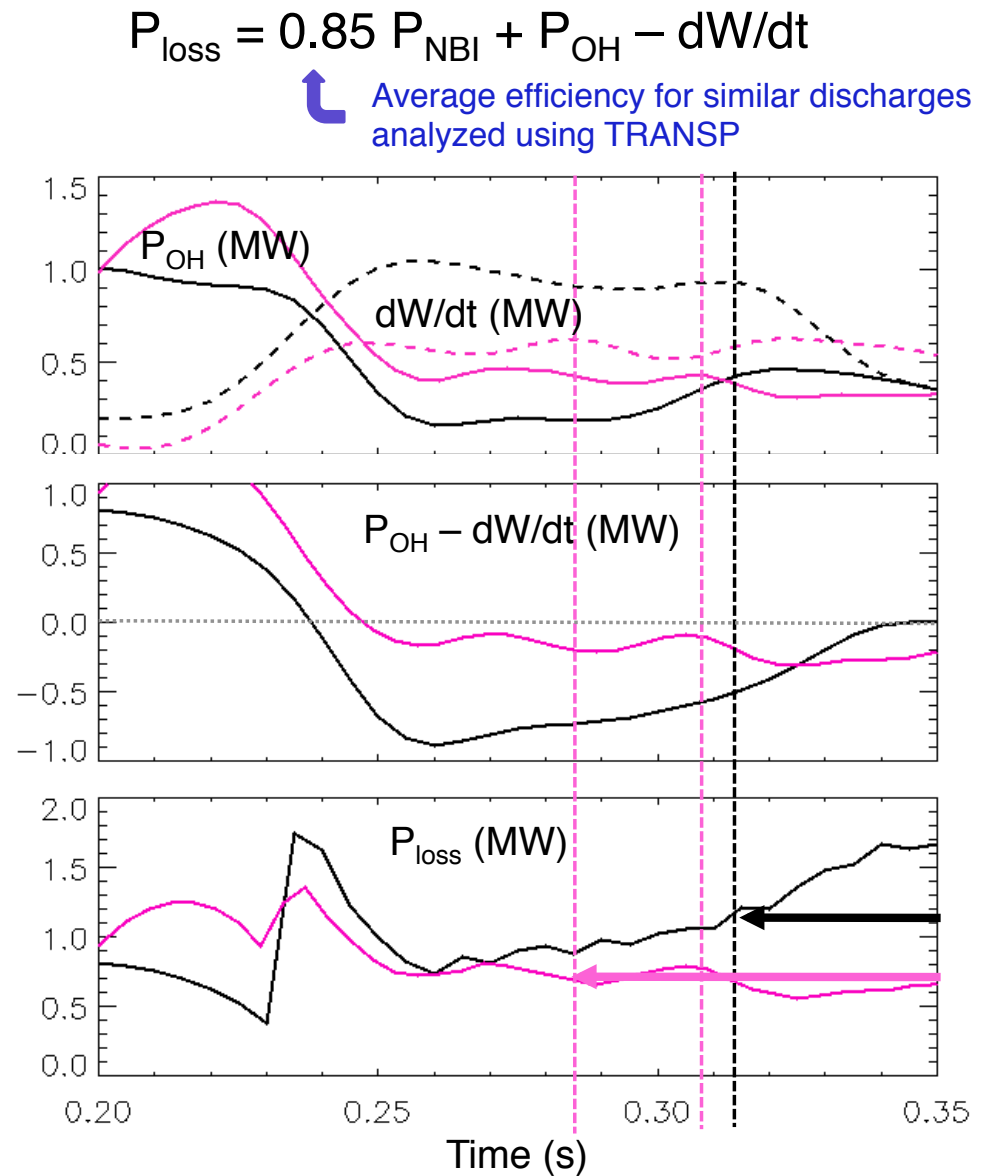
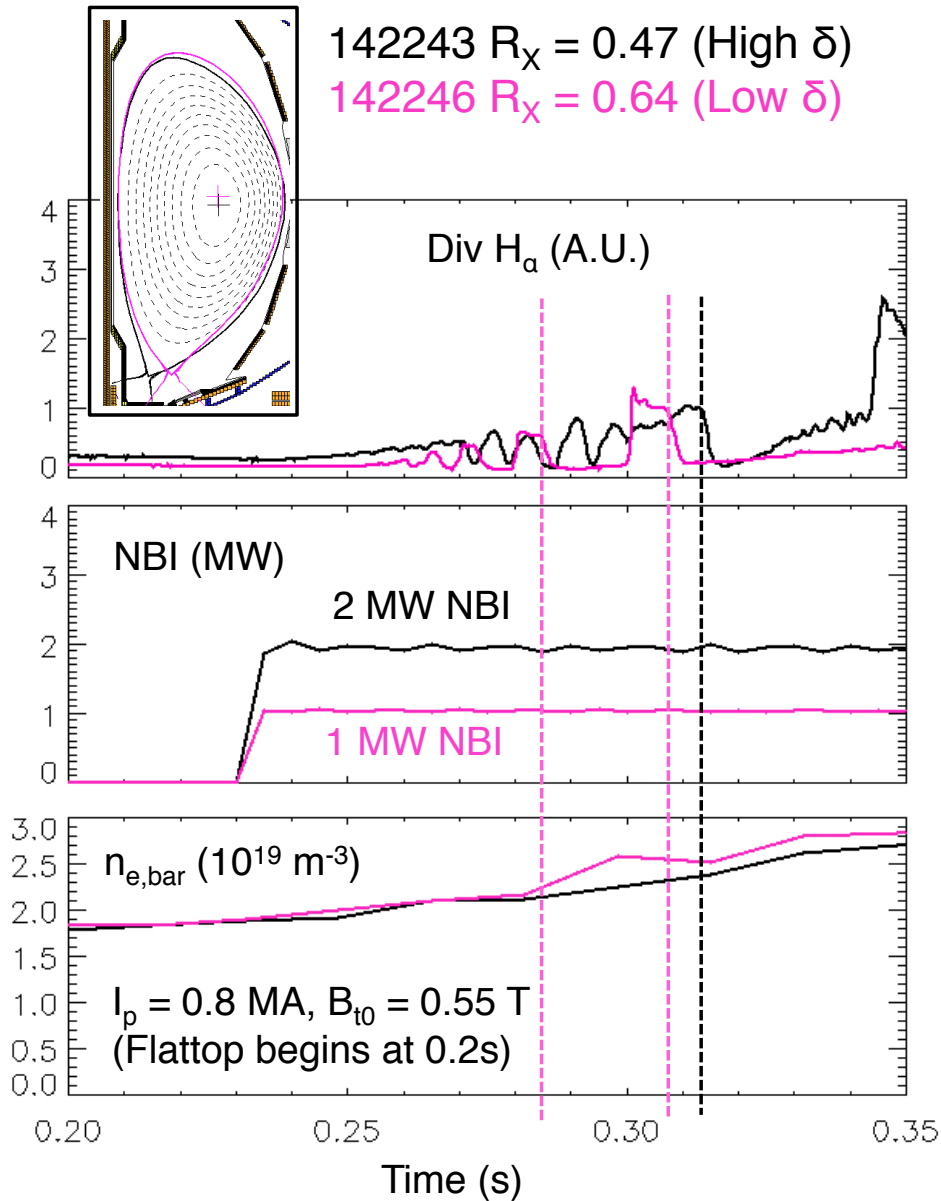
X-transport predictions motivated a dedicated P_{LH} versus R_X experiment on NSTX



- Goal: quantify P_{LH} vs R_X on NSTX
 - Match R_{IN} , R_{OUT}
 - Nearly match X-point height, surface area, B_{T0} , B_{OUT} , density
 - Reproduce shapes under different pumping & fueling conditions
 - Shot-to-shot change in NBI < 300 kW
- Delay L-H to > 40 ms after NBI turn-on to reduce error in P_{LH}
 - P_{OH} – dW/dt slowly varying

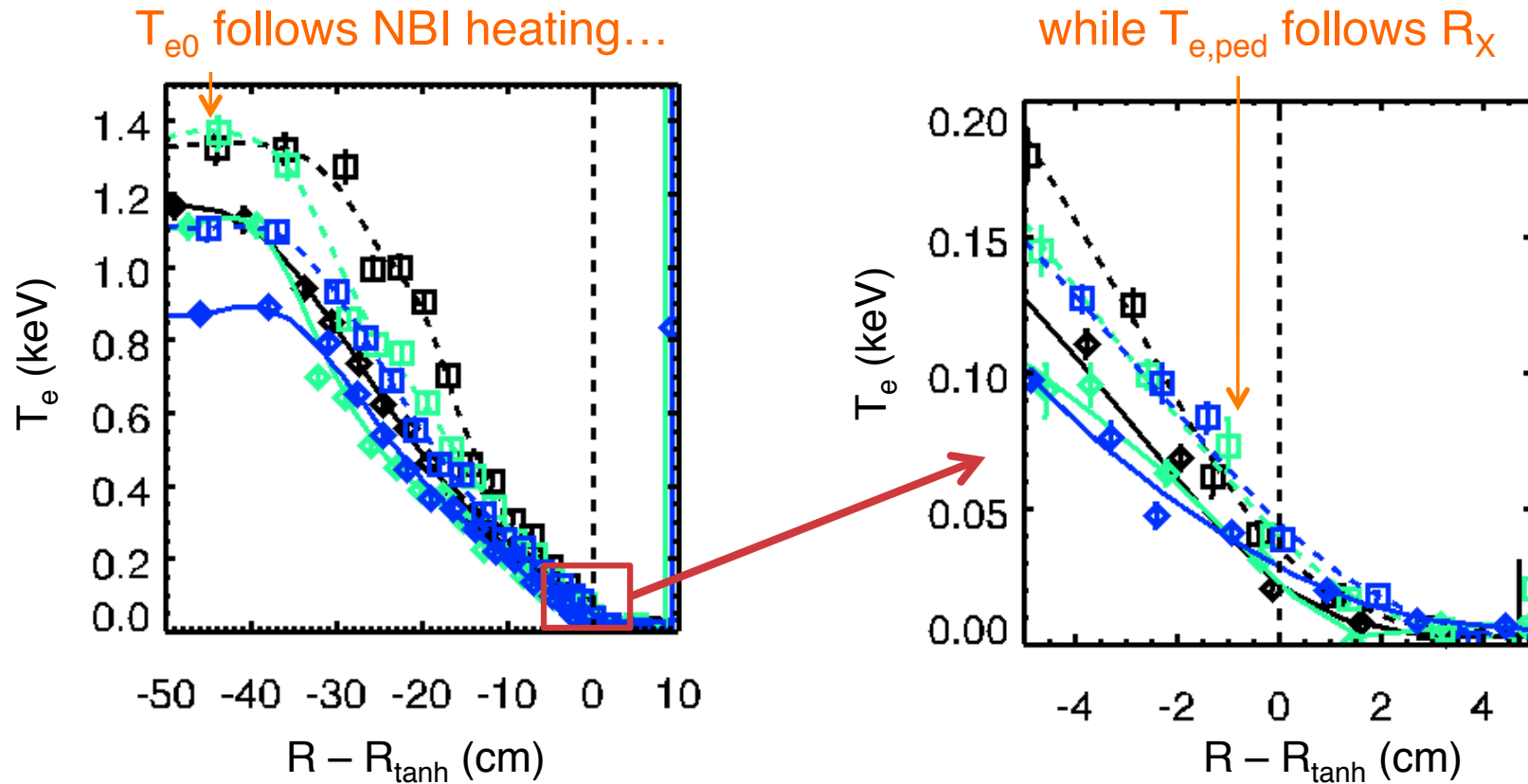
First results reported in: R. Maingi, et. al. *Nucl. Fusion* **50** (2010)

LH transitions occur when $P_{OH} - dW/dt$ is slowly varying



Slowly varying profiles provide good constraint to L-mode electron profiles right before L-H transition

Last T_e profile before LH transition
(X-axis normalized to density tanh symmetry point)



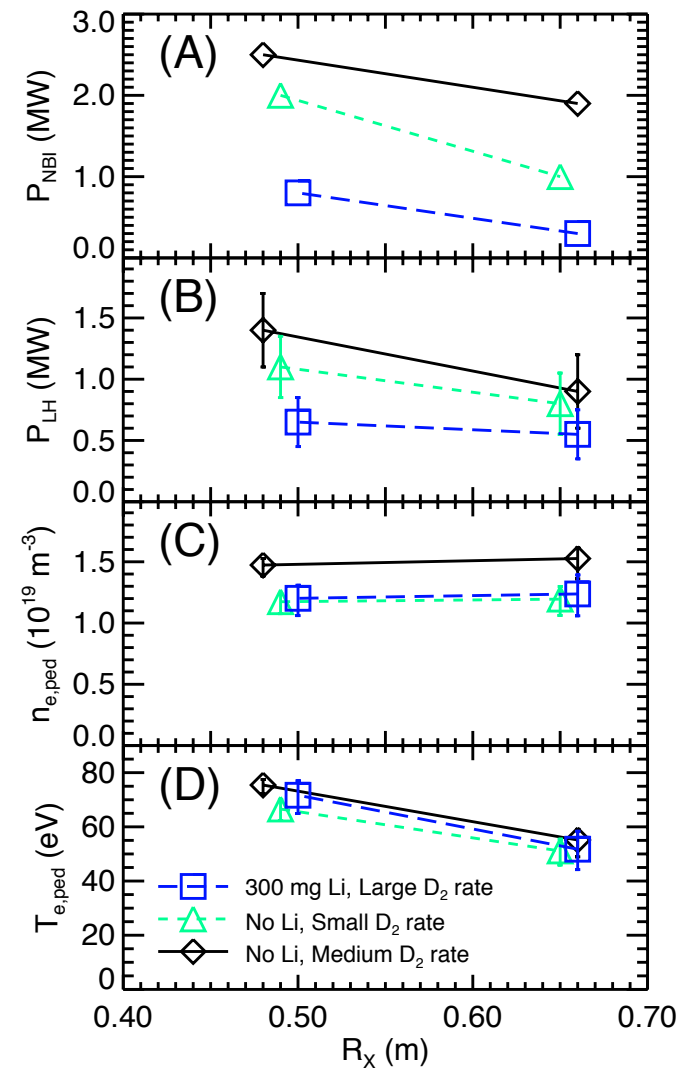
300 mg Lithium, Large fueling
Small R_x : dotted (0.6 MW)
Large R_x : solid (0.3 MW)

0 mg Lithium, Small fueling
Small R_x : dotted (2.0 MW)
Large R_x : solid (1.0 MW)

0 mg Lithium, Medium fueling
Small R_x : dotted (2.5 MW)
Large R_x : solid (2.0 MW)

Edge temperature prior to LH transition larger for small- R_x than large- R_x over wide range of heating and pumping

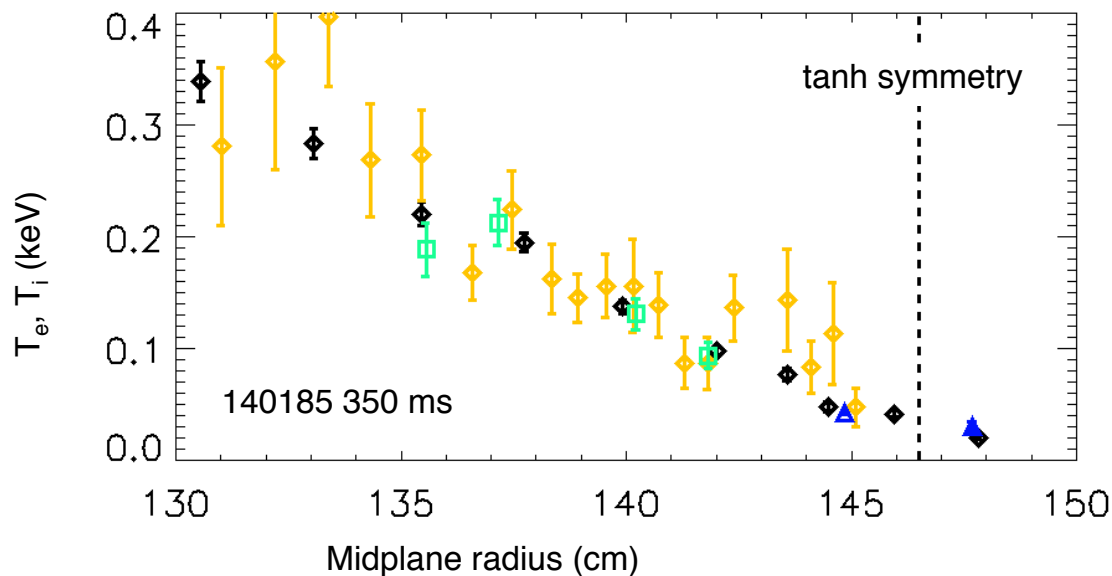
- Three pumping & fueling conditions
 - Divertor pumping on NSTX varied using inter-shot Lithium deposition rates
- P_{LH} increases with...
 - Less divertor pumping
 - Higher $n_{e,ped}$ and/or HFS fueling rate
 - Smaller R_x
- Pedestal T_e near L-H time ...
 - Nearly independent of pumping, fueling density, NBI power
 - Increases ~40% as R_x 0.64 \rightarrow 0.37 m
 - Unfortunately, T_i not available for all shots



D.J. Battaglia, et. al., PRL, *submitted*

T_e can be used as a proxy for T_i in the L-mode edge of NSTX

- Infer $T_i \sim T_e$ in L-mode edge
 - Thermal equilibration time < electron thermal confinement
 - T_i is the critical parameter in X-transport theory
- Thus, $T_{i,LH}$ increases $\sim 40\%$ as $R_X: 0.64 \rightarrow 0.47$
 - Leads to 20 – 60% larger P_{LH}



Example of L-mode T_e, T_i close to L-H transition

T_e (MPTS)
 T_i (CVI CX Spectroscopy)
 T_i (CVI intrinsic emission inverted)
 T_i (CIII intrinsic emission)

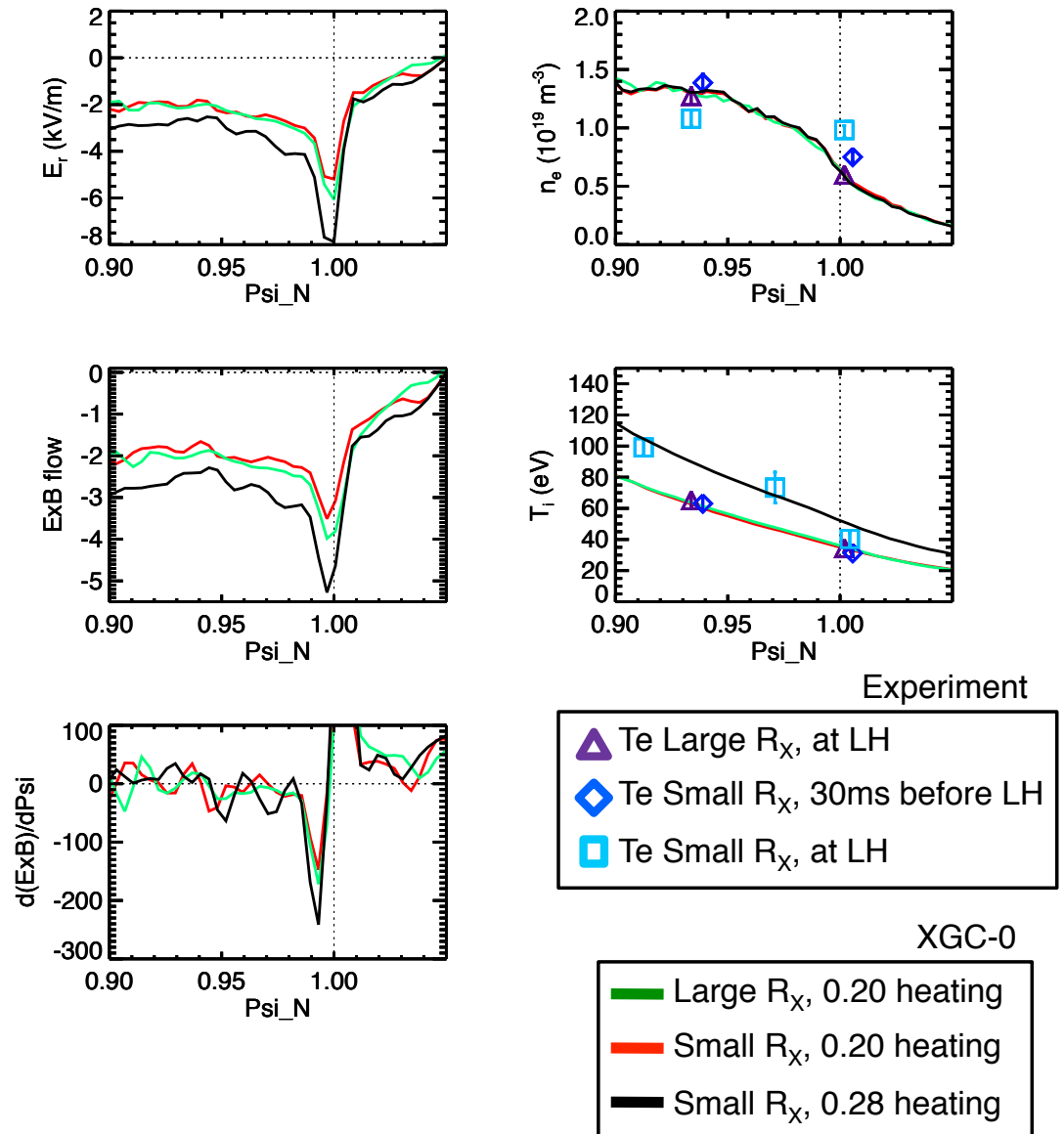
Experimental data is consistent with X-transport model

- Increase in T_i matches increase in K_0
 - As R_x : 0.64 \rightarrow 0.47 m ...
 - K_0 increases 35% (guiding center calculation)
 - Edge T_i increases 40% (derived from measurements)
 - Full self-consistent E_r calculations needed to confirm quantitative agreement
- P_{LH} dependence on pumping and fueling
 - T_i similar for all three conditions ... but P_{LH} varies considerably
 - Recycling cools edge \rightarrow need larger core heating to reach critical T_i
- Does change in strike points with R_x impact div. recycling?
 - Most likely and it will have an effect on P_{LH}
 - But it should not effect $T_{i,LH}$ (since it appears insensitive to large changes in divertor recycling)

Preliminary self-consistent calculations of E_r using XGC-0

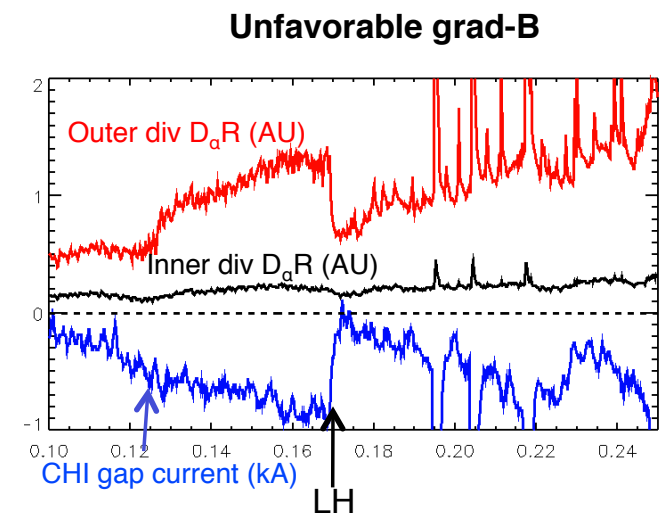
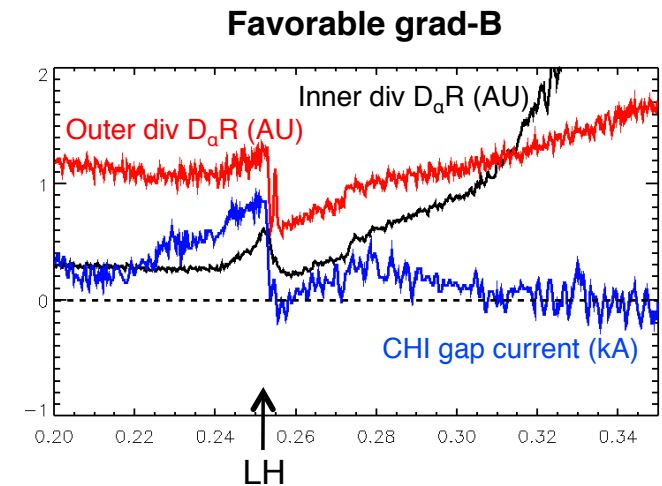
- XGC-0 simulations
 - 5D ion & electron + neutrals kinetic model
 - Divertor sheath and recycling models
- Tailor model parameters to match L-mode profiles
 - Inboard heating, outboard neutral fueling, anomalous transport rate
- E_r well deeper for large- R_X shape with matched profiles

Preliminary results. No electron physics included.



Measurements may suggest there is finite ion orbit loss to the divertor

- NSTX: D_α increases prior to L-H
 - Favorable: biggest change in inboard div
 - Unfavorable: biggest change in outboard div
- NSTX: Divertor current prior to L-H
 - Direction of current follows grad-B direction
- Are ions lost on neoclassical orbits?
 - Signals consistent with ions lost on neoclassical orbits to the divertor
 - Large ion loss not expected in steady state, but may be possible in a dynamic solution
- Use XGC-0 to explore L-H dynamics and impact on SOL / divertor



Acknowledgements

This work was funded by the US Department of Energy under Contract Numbers DE-AC02-09CH11466 and DE-AC05-00OR22725.

Reprints
