

Density fluctuation measurement by FIRETIP for the Enhanced Pedestal H-mode on NSTX

K. C. Lee¹, C. W. Domier¹, N. C. Luhmann, Jr¹,
R. Kaita², R. Maingi³
and the NSTX Research Team

UC DAVIS
UNIVERSITY OF CALIFORNIA

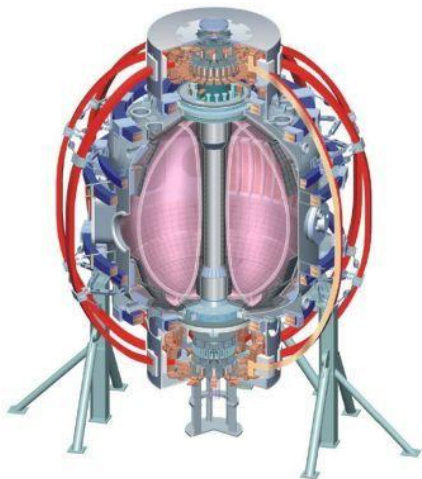
Columbia U
CompX
General Atomics
FIU
INL
Johns Hopkins U
LANL
LLNL
Lodestar
MIT
Nova Photonics
New York U
ORNL
PPPL
Princeton U
Purdue U
SNL
Think Tank, Inc.
UC Davis
UC Irvine
UCLA
UCSD
U Colorado
U Illinois
U Maryland
U Rochester
U Washington
U Wisconsin

¹ University of California at Davis,

² Princeton Plasma Physics Laboratory, ³ Oak Ridge National Laboratory

53th APS DPP meeting
Salt Lake City, Utah. Nov. 14-18 2011

Culham Sci Ctr
U St. Andrews
York U
Chubu U
Fukui U
Hiroshima U
Hyogo U
Kyoto U
Kyushu U
Kyushu Tokai U
NIFS
Niigata U
U Tokyo
JAEA
Hebrew U
Ioffe Inst
RRC Kurchatov Inst
TRINITY
NFRI
KAIST
POSTECH
ASIPP
ENEA, Frascati
CEA, Cadarache
IPP, Jülich
IPP, Garching
ASCR, Czech Rep

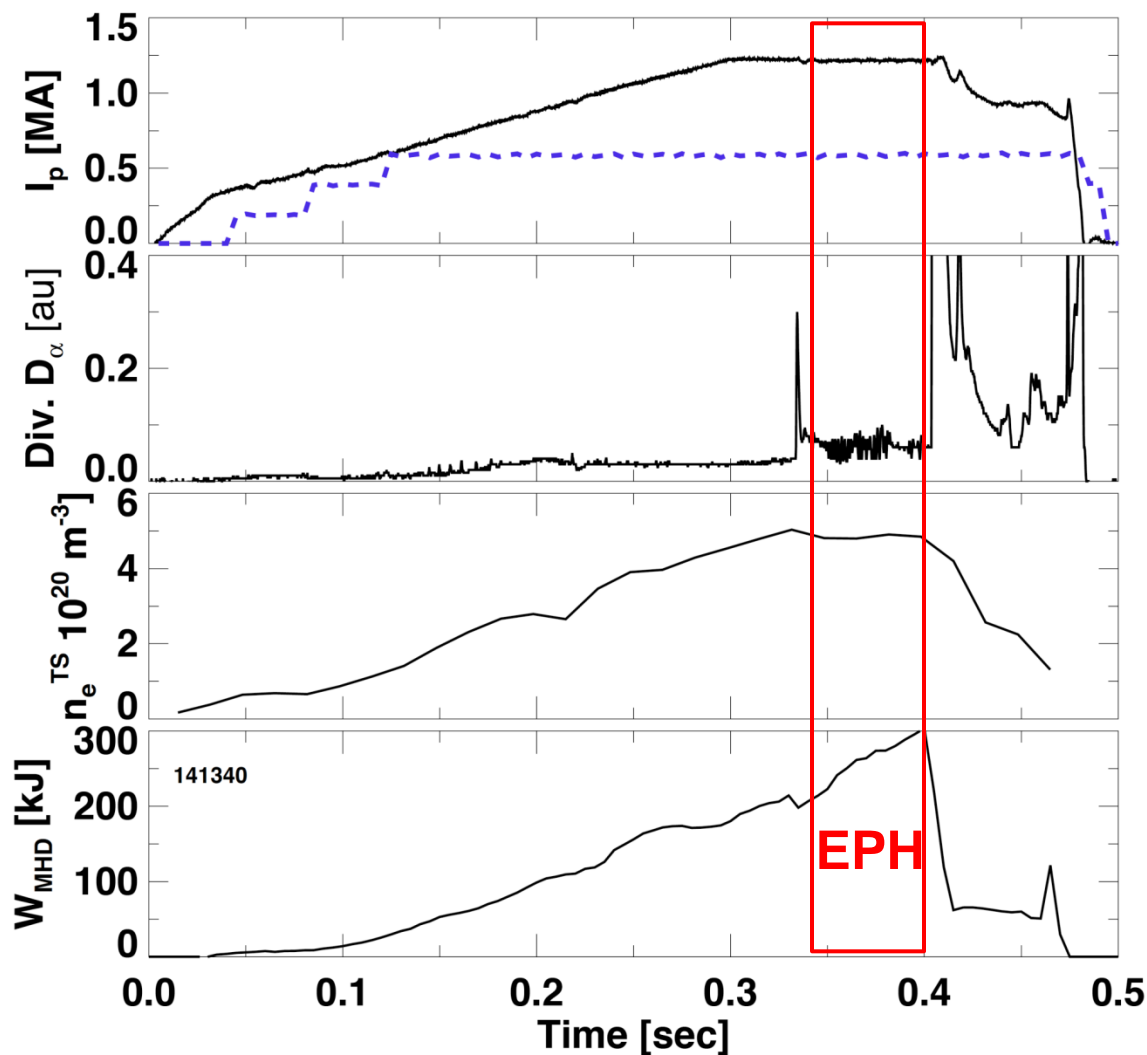


Abstract

The multi-channel Far Infrared Tangential Interferometry/Polarimetry (FIReTIP) system has been used to measure changes in the electron density fluctuation spectrum for the Enhanced Pedestal H-mode (EPH-mode) on the National Spherical Torus Experiment (NSTX). Data shows dramatic density fluctuation suppression as the EPH-mode is triggered, similar in nature to the turbulence reduction present at the conventional L\H transition. Coherent fluctuations are observed by FIReTIP during the EPH-mode with frequencies greater than 10 kHz. Density fluctuation measurements from FIReTIP edge channels with different tangency radii during the EPH-mode are compared with L-mode and H-mode cases, and are presented together with a discussion of a possible EPH-mode triggering mechanism based on the gyro-center shift (GCS) theory.

**This work is supported by U.S. Department of Energy Grant Nos. DE-FG02-99ER54518 and DE-AC02-09CH11466.*

Enhance Pedestal H-mode (EP H-mode) on NSTX

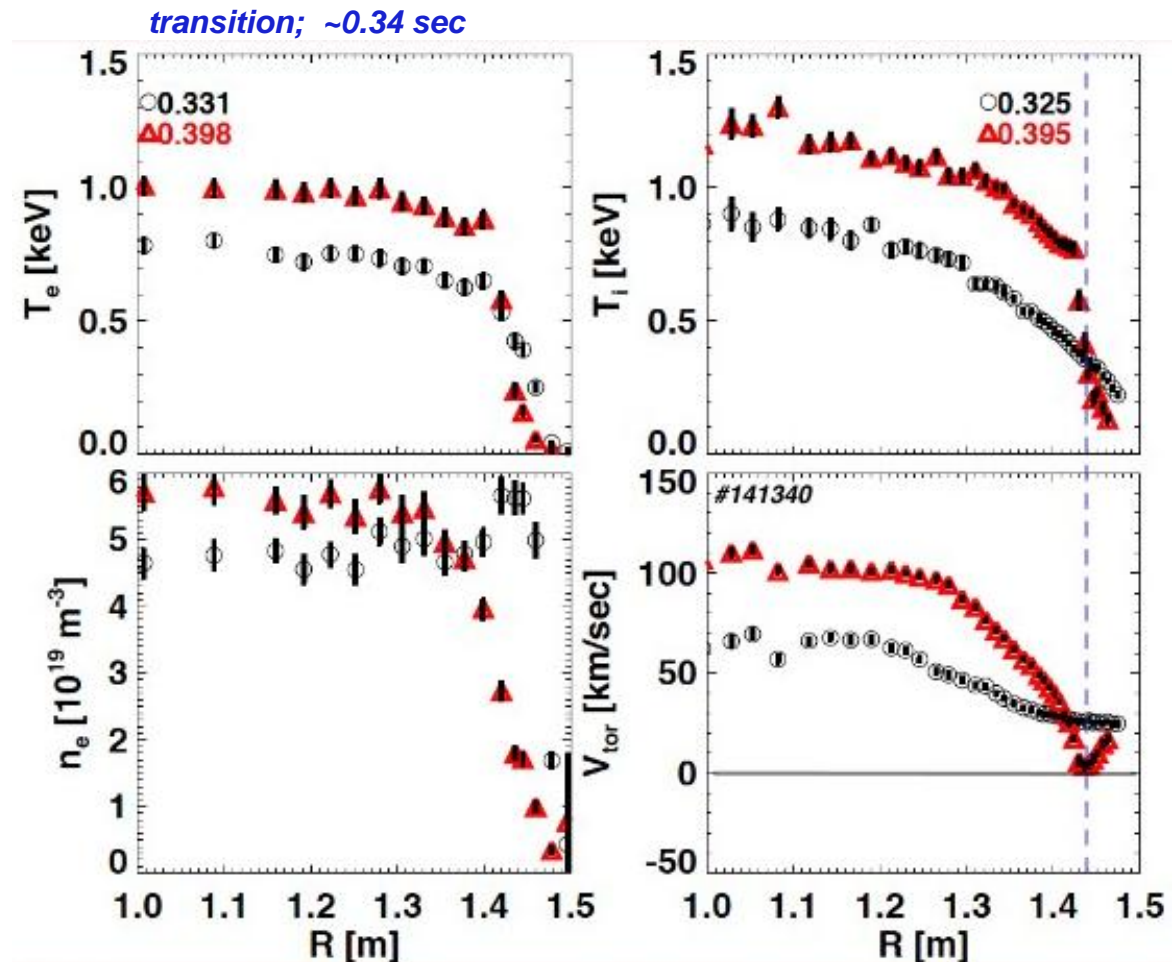


- EP H-mode triggered by natural ELM
- energy confinement increases 50%
- transition after large ELM, either natural or externally triggered by 3D fields

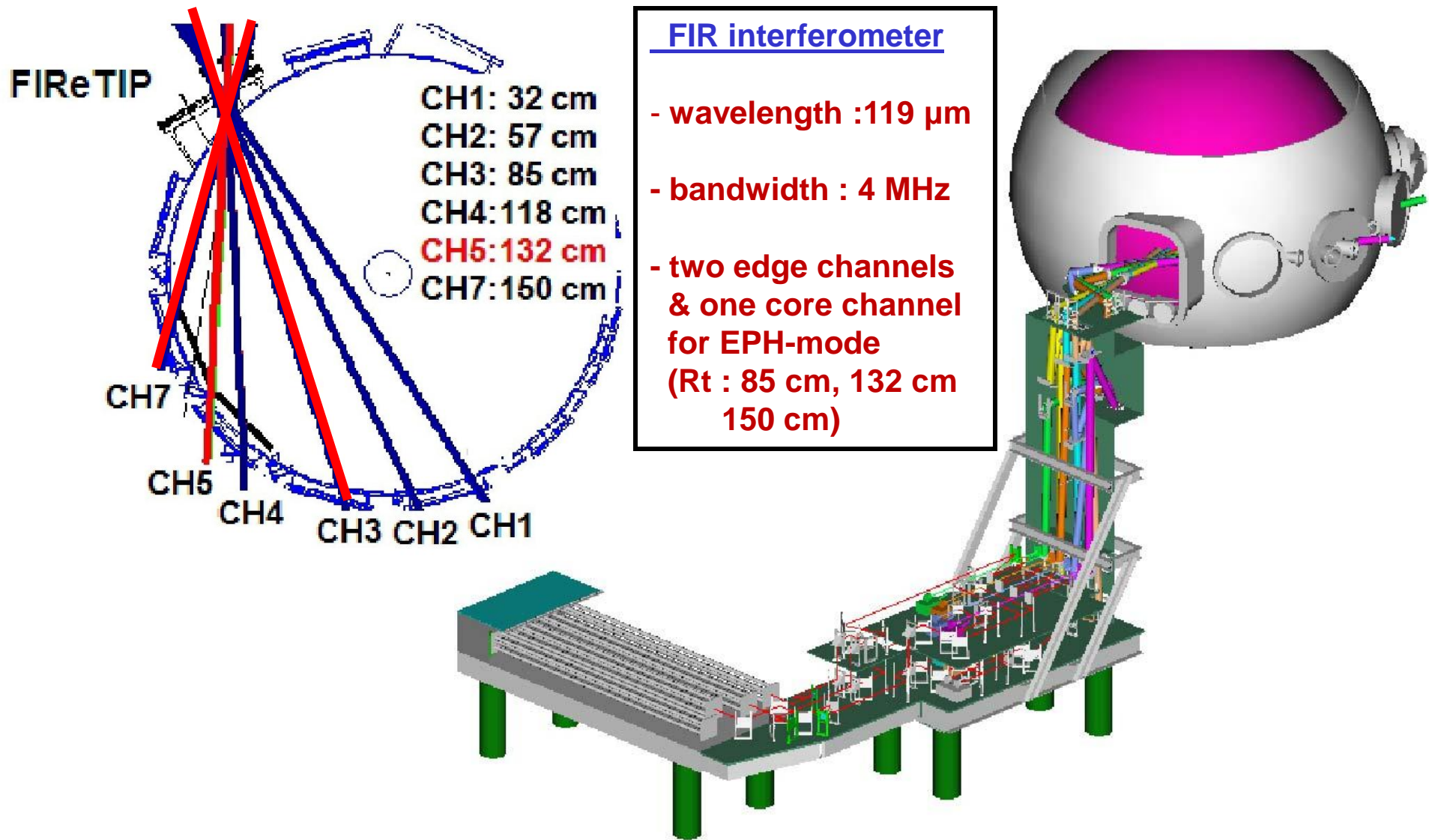
[R. Maingi, *et al.*, PRL, **105** 135004, 2010]

Profile changes before and after the EP H-mode transition

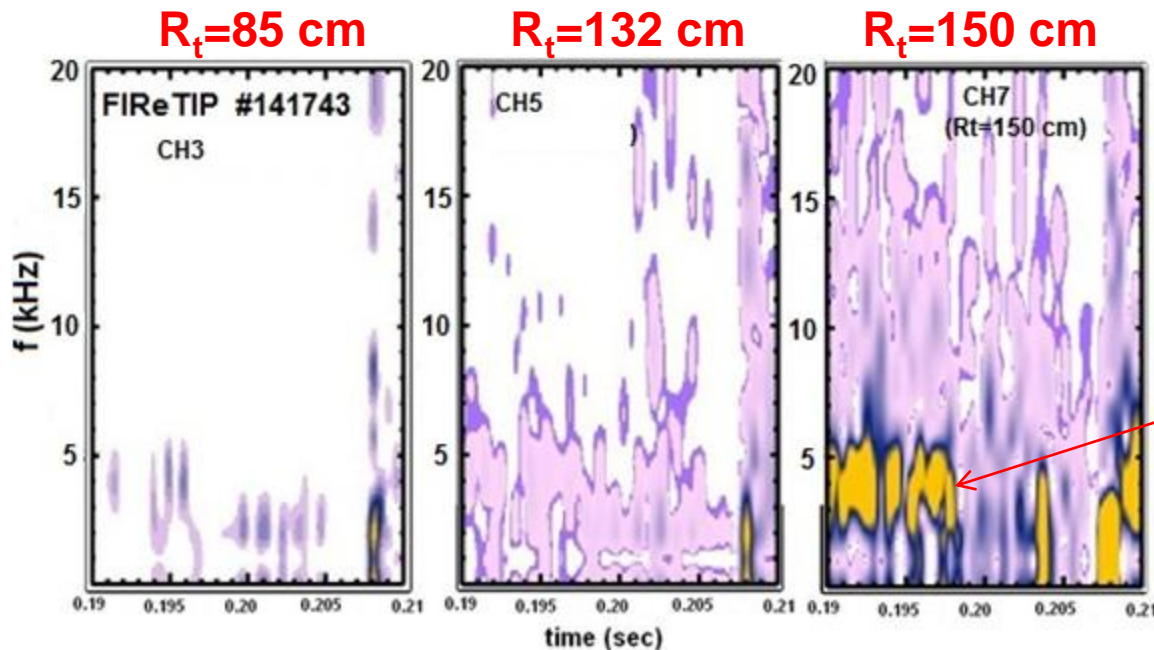
- pedestal temperature doubles after transition
- transport barrier grows inward from the edge
- sharp pedestal correlated with toroidal velocity locked to zero near $q=3$ surface



FIReTIP measurement of n_e fluctuations on NSTX

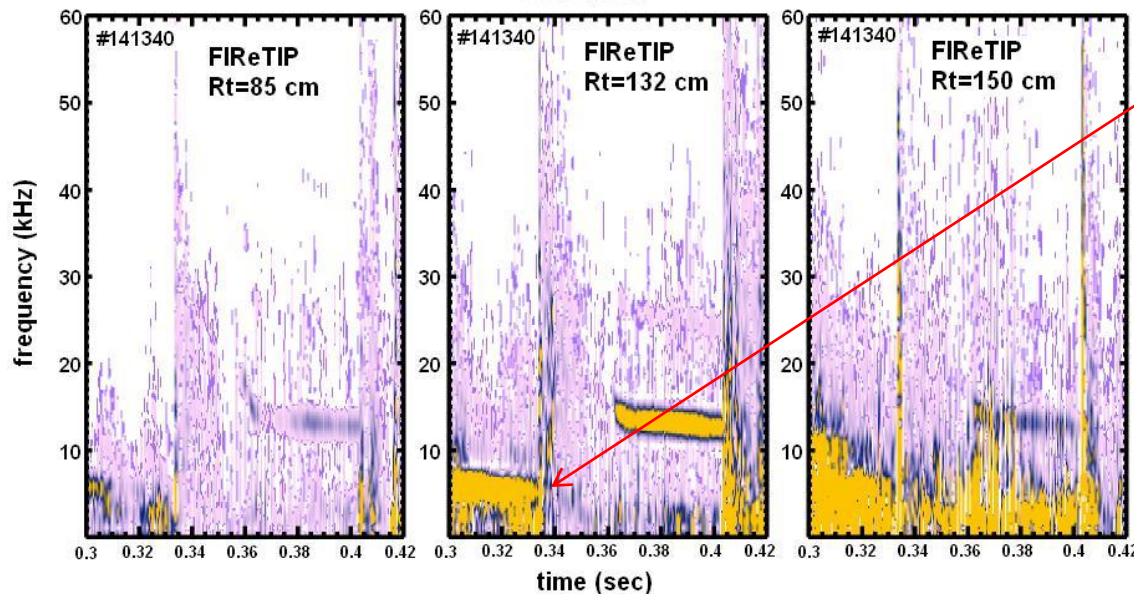


FIReTIP showed ∂n_e suppression on EPH-mode



(FIReTIP) n_e fluctuation spectrum H-mode versus EP H-mode

LH Transition

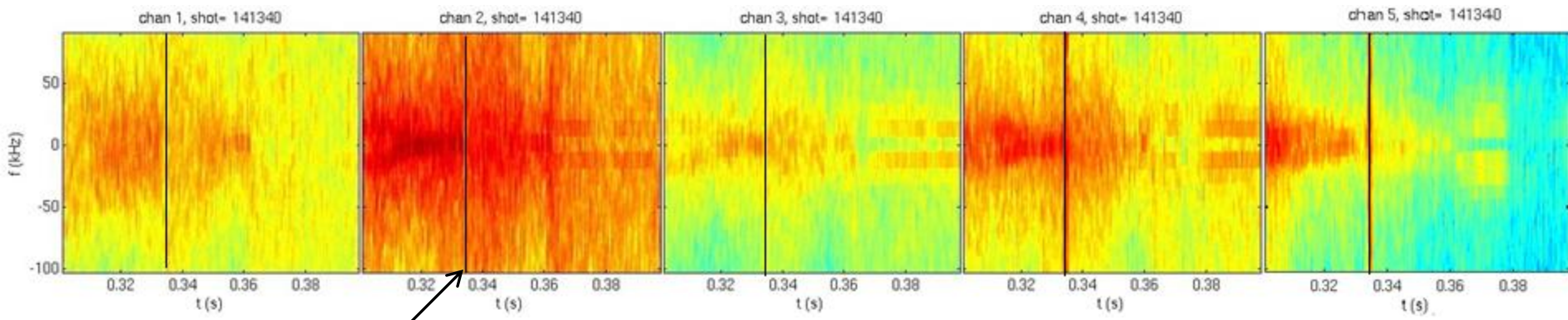


EP H-mode Transition

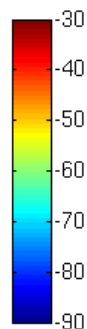
► both reduced fluctuations but at different location (magnetic fluctuations by Mirnov coil showed increase on EPH-mode?)

Intensity: magenta < dark blue < yellow

High-k measurement also showed fluctuation reduction



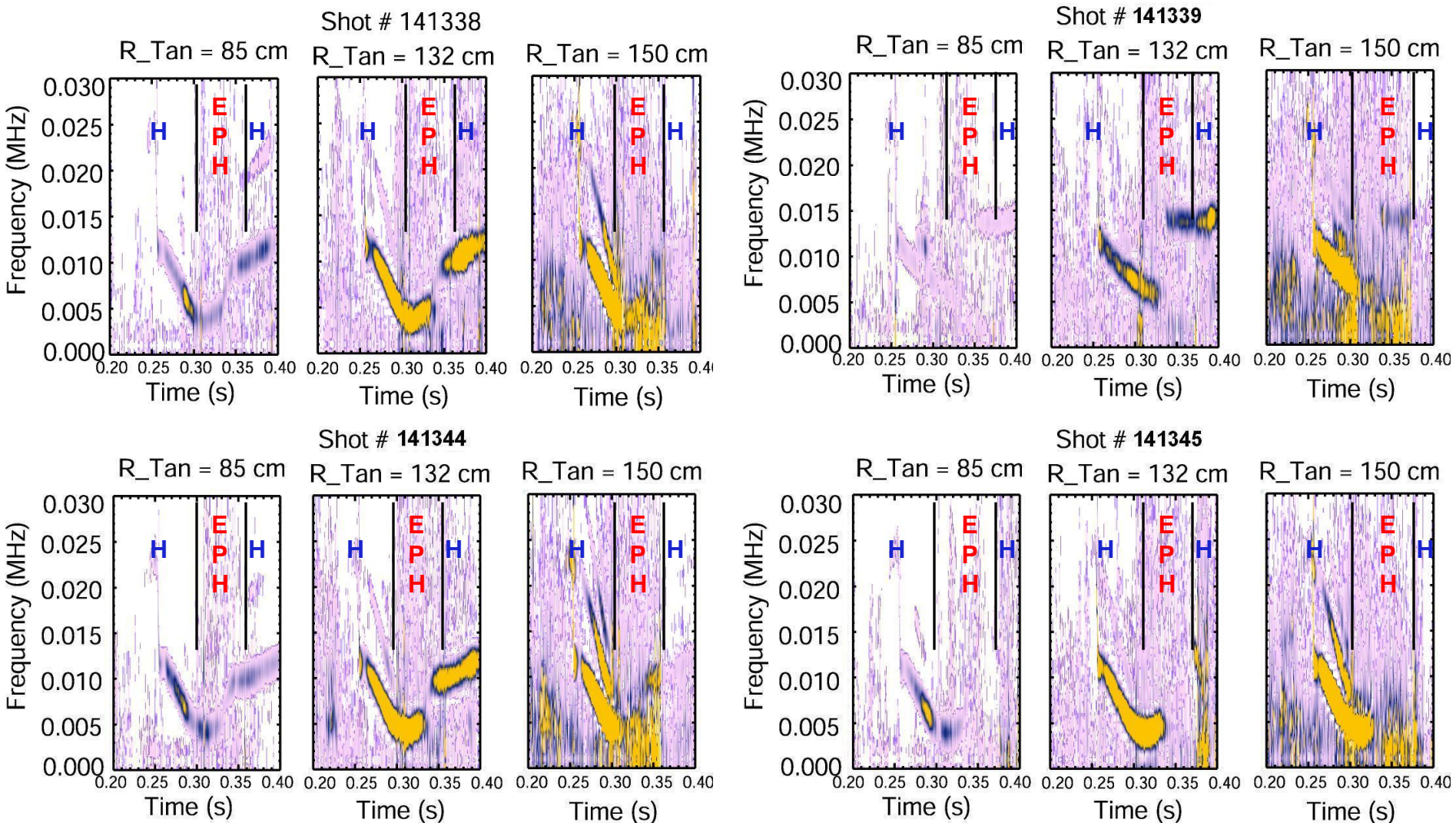
EPH-mode transition



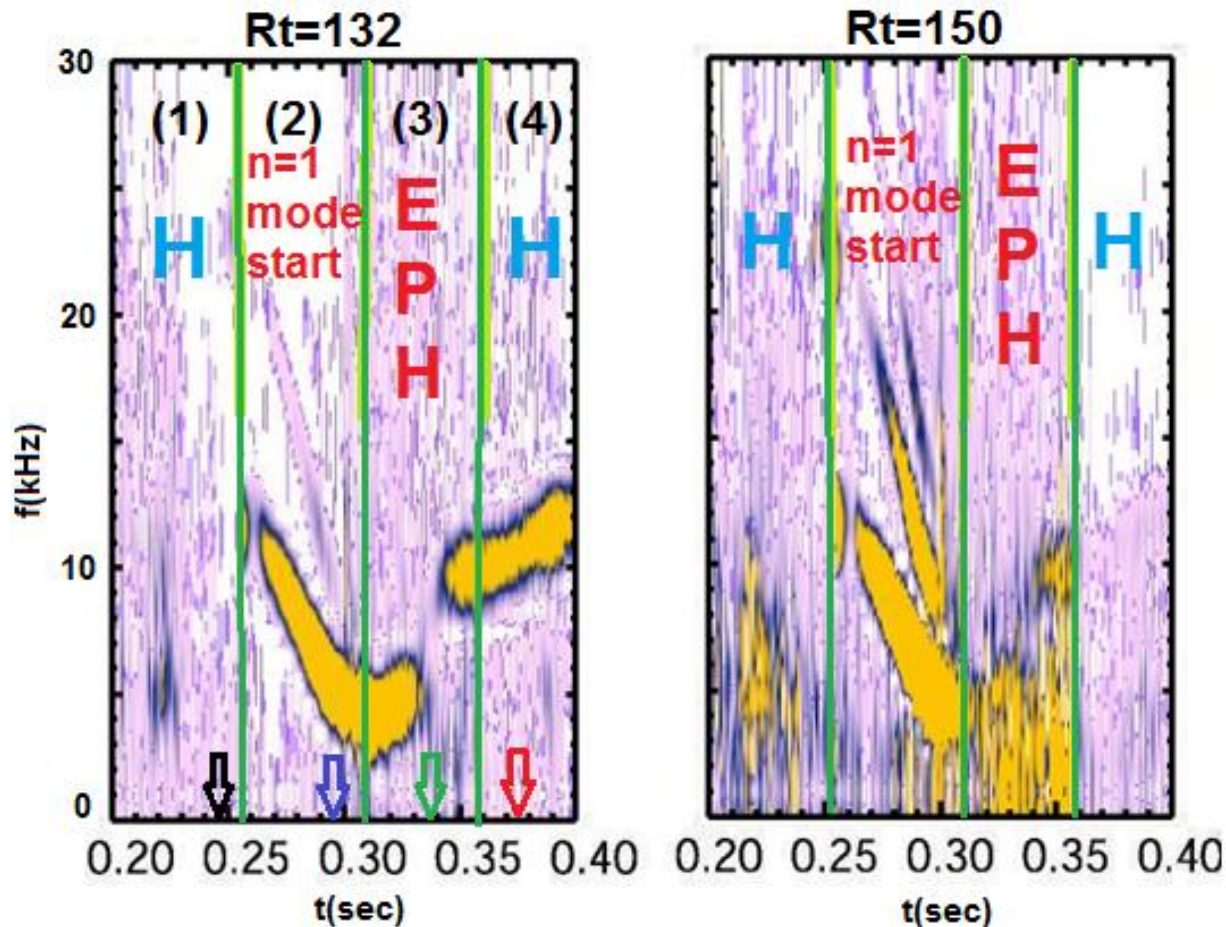
- ▶ measured for $R_T=143-144$ cm
- ▶ most intensive Ch2 showed clear reduction at EPH transition
- ▶ frequency range is similar to FReTIP measurement

both FReTIP and High-k showed fluctuation reduction on EPH-mode transition at few cm inner location than normal H-mode transport barrier

n_e fluctuation spectra by FReTIP showed $n=1$ MHD (5~15 kHz) on EPH-mode



Characteristics of $n=1$ MHD measured by FIRETIP



there are 4 stages:

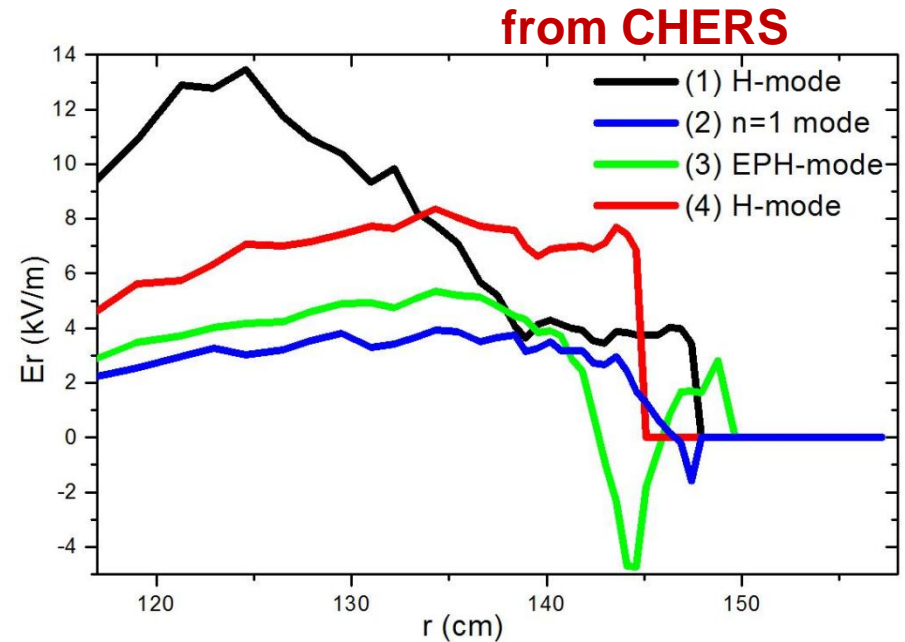
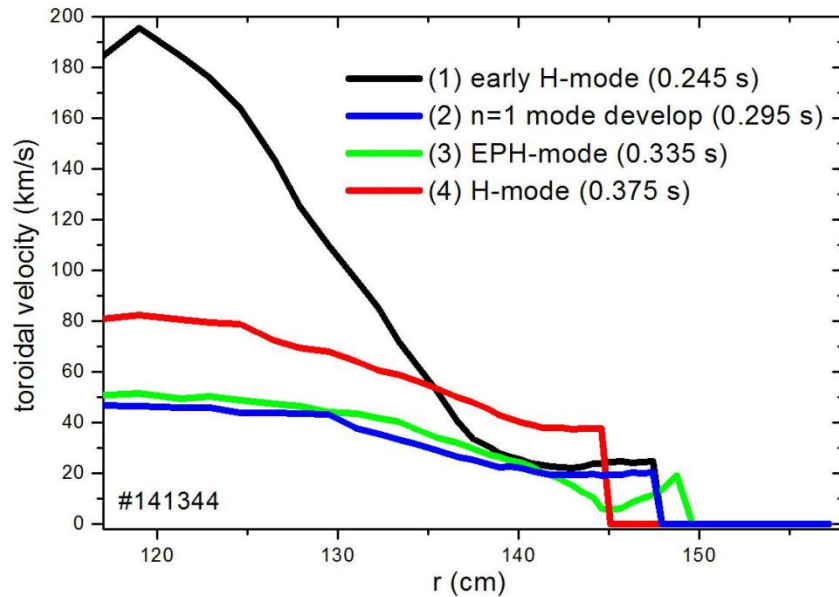
- (1) early H-mode
- (2) $n=1$ mode start
- (3) EPH-mode by ELM
- (4) H-mode by ELM

▶ $n=1$ mode frequency decreases with toroidal rotation reduction in stage (2)

▶ on EPH-mode start by ELM, $n=1$ mode frequency begins to increase

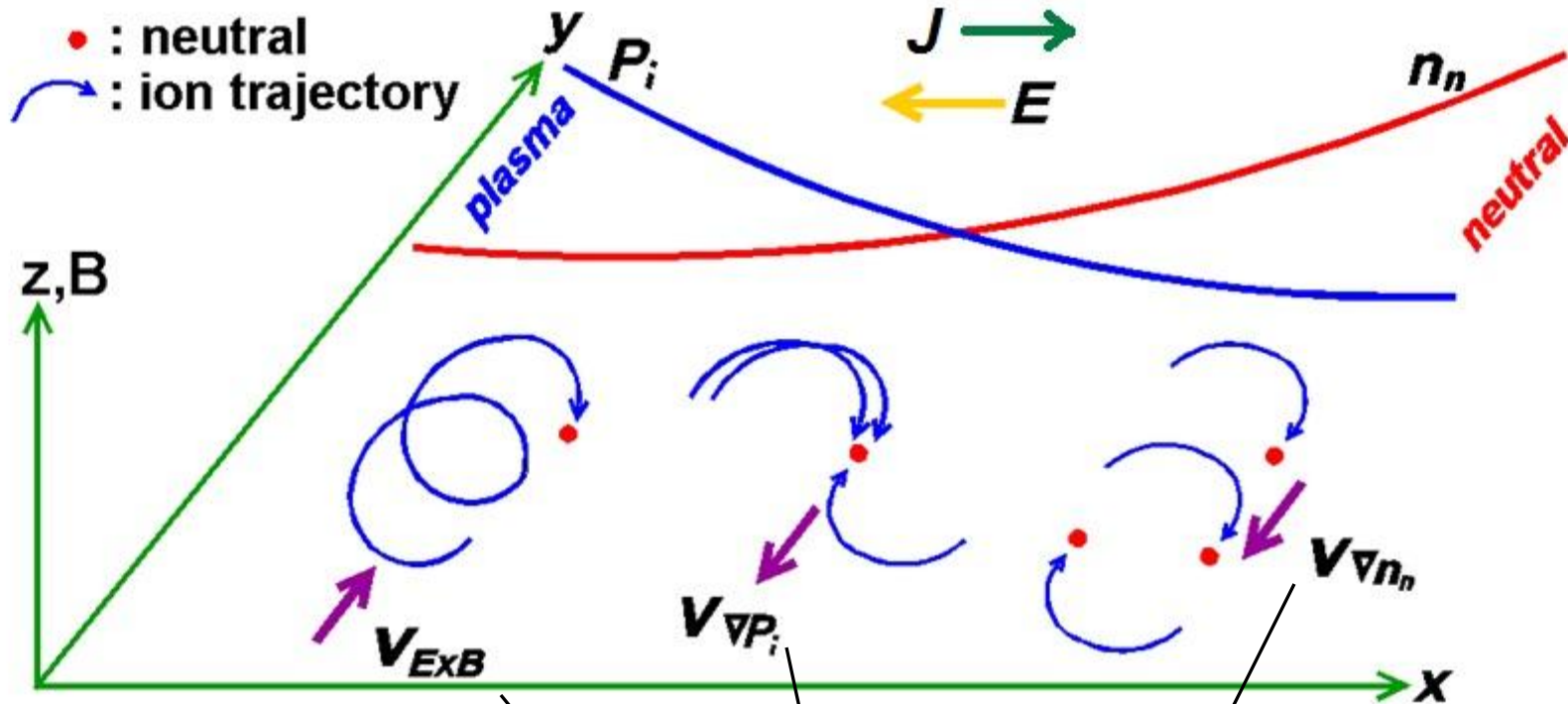
- ▶ there is always $n=1$ mode in stage (2) but behavior of $n=1$ mode is not consistent in stage (3) or (4)

V_t and E_r profile changes on EPH-mode development



- ▶ in stage (2), while n=1 mode develops, toroidal rotation decreases, and there is additional change in E_r for boundary
- ▶ by ELM, V_t and E_r profiles changed to have inside points of dip; stage (3)
- ▶ In EPH-mode, there are inflection points for E_r and V_t ($\nabla^2 E_r = 0$, $\nabla^2 V_t = 0$)
- ▶ at second ELM, EPH-mode transforms back to H-mode with inner boundary than earlier H-mode.

Theory of gyro-center shift (GCS)



$J \times B$ = ion momentum loss rate due to collision with neutrals

$$J_r^{GCS} = en_i \frac{r_{Li}}{\lambda_{i-n}} \left(\frac{E}{B} - \frac{1}{eB} \frac{\nabla P_i}{n_i} + \frac{kT_i}{eB} \frac{\nabla n_n}{n_n} \right)$$

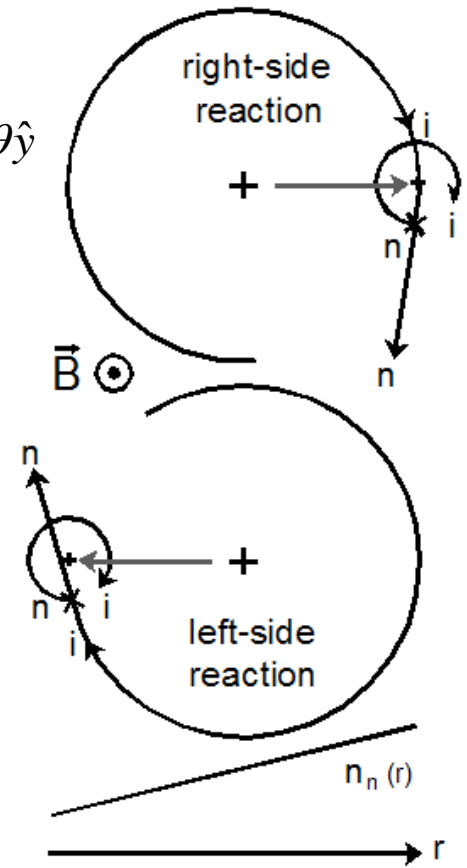
return current term

driving terms

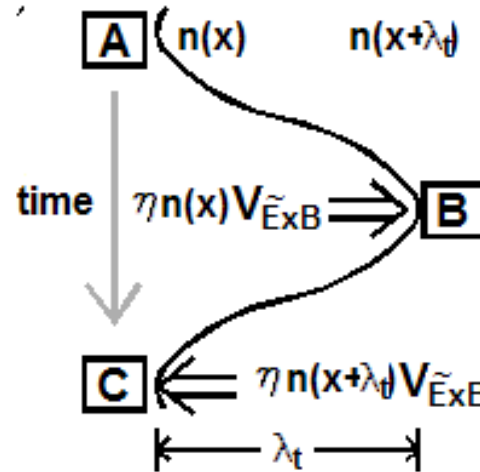
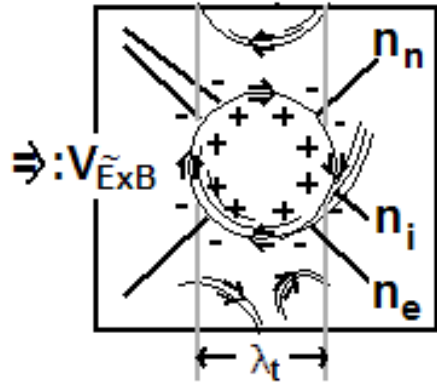
- When ions collide with neutrals, they lose their momentum; ex) charge exchange. y-components of ion velocity in above picture ($V_{\nabla p_i}$, $V_{\nabla n_n}$) generate a current in x-component (radial in tokamak) by $\mathbf{J} \times \mathbf{B}$ = momentum loss. $V_{\nabla p_i}$ is ion diamagnetic drift and $V_{\nabla n_n}$ is effective ion velocity comes from the neutral density gradient calculated by;

$$v_{av} = \frac{\sigma v_i \oint v n_n(x) d\theta}{\sigma v_i \oint n_n(x) d\theta} = \frac{1}{2} r_L v_{\perp} \frac{1}{n_n} \frac{\partial n_n}{\partial r}, \quad v = v_{\perp} \sin \theta \hat{r} + v_{\perp} \cos \theta \hat{y}$$

- These two terms ($V_{\nabla p_i}$, $V_{\nabla n_n}$) drive gyro-center shift current (J_r^{GCS}) and forms electric field as the source of return current.
- The theory of gyro-center shift explains the origin of the radial electric field on the boundary of tokamaks. [Lee, POP 2006]



Turbulence induced diffusion by GCS theory [Lee, PPCF 09']



$$\Gamma = \underbrace{\partial n}_{\eta \lambda_t \nabla n} \cdot \underbrace{\tilde{v}}_{\frac{1}{\pi} \frac{\tilde{E}}{B}} \quad \rightarrow \quad D = \frac{\eta}{\pi} \frac{\tilde{E} \lambda_t}{B} \quad \rightarrow \quad \boxed{D = \frac{2}{\pi} \eta^2 \frac{kT_e}{eB}}$$

$\tilde{E} \lambda_t \approx 2\eta \frac{kT_e}{e} \left(\frac{e\tilde{\phi}_t}{kT_e} \approx \frac{\tilde{n}_e}{n_e} : \text{Boltzmann relation, } \tilde{\phi}_t \approx \tilde{E} \frac{\lambda_t}{2} \right)$

▶ turbulence induced ion and electron diffusion : $\eta \lambda_t \nabla n$

▶ turbulence induced charge diffusion : $-\eta \lambda_t \nabla \rho$

▶ ion and electron move toward boundary => diffusion

▶ charge (ρ) moves toward core => dilution current =>

saturation by J^{GCS}

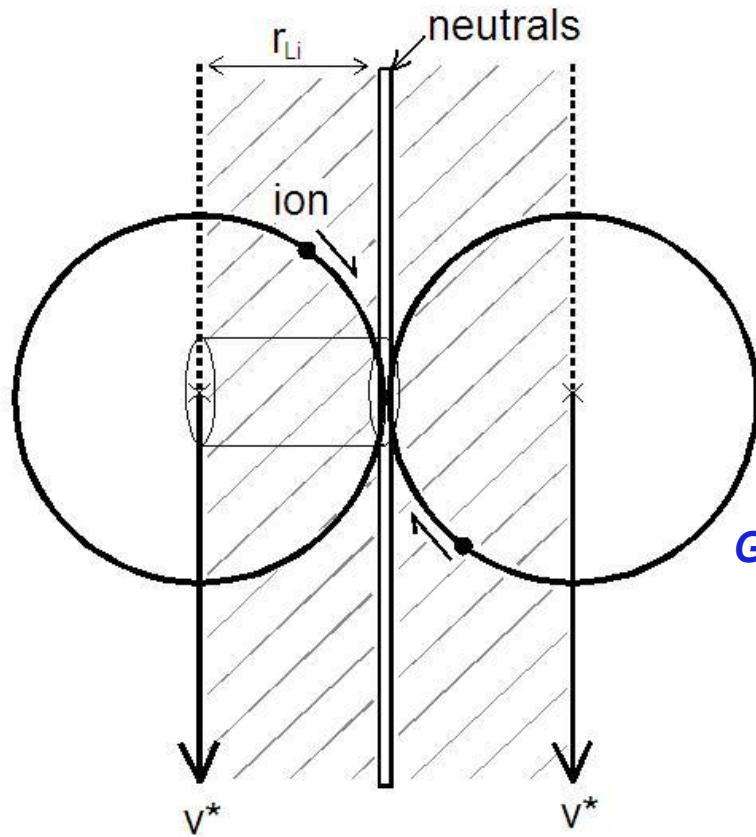
Reynolds number of ion-neutral friction

Turbulence

- ▶ ions and electrons move toward boundary => diffusion
- ▶ charge (ρ) moves toward core => dilution current => **saturation condition**

$$J_r^{GCS} = en_i \frac{r_{Li}}{\lambda_{i-n}} \left(\frac{E}{B} - \frac{1}{eB} \frac{\nabla P_i}{n_i} + \frac{kT_i}{eB} \frac{\nabla n_n}{n_n} \right)$$

v^*



inertia force ↓

$$Re \equiv \frac{n_i m_i v^{*2} / r_{Li}}{n_i m_i \nu_{i-n} v^*} = \frac{eB}{kT_i} \lambda_{i-n} v^*$$

viscosity force ↑

(saturation condition : $J_r^{GCS} = D \nabla \rho$)

$$Re = \frac{2}{\pi} \eta^2 \frac{B}{m_i n_i (\sigma_{i-n} n_n)^2 v_{\perp}} \nabla \rho$$

Gauss's law in slab geometry ↓

$$\nabla E = \frac{\rho}{2\epsilon_0}$$

$$Re = \frac{4\epsilon_0}{\pi} \eta^2 \frac{B}{m_i n_i (\sigma_{i-n} n_n)^2 v_{\perp}} \nabla^2 E$$

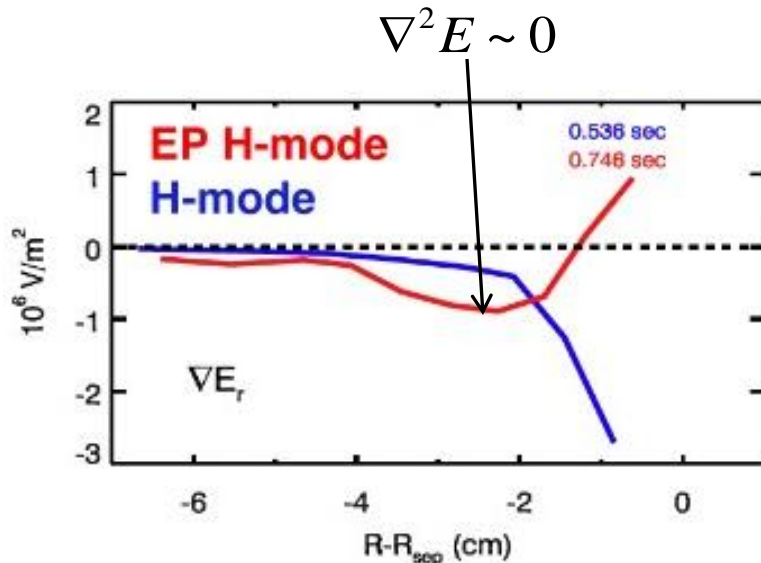
Reynolds number for EPH-mode

$$\text{Re} = \frac{4\epsilon_0}{\pi} \eta^2 \frac{B}{m_i n_i (\sigma_{i-n} n_n)^2 v_{\perp}} \nabla^2 E$$



If $\nabla^2 E \sim 0$, Re can be smaller than Re^* which determine turbulence

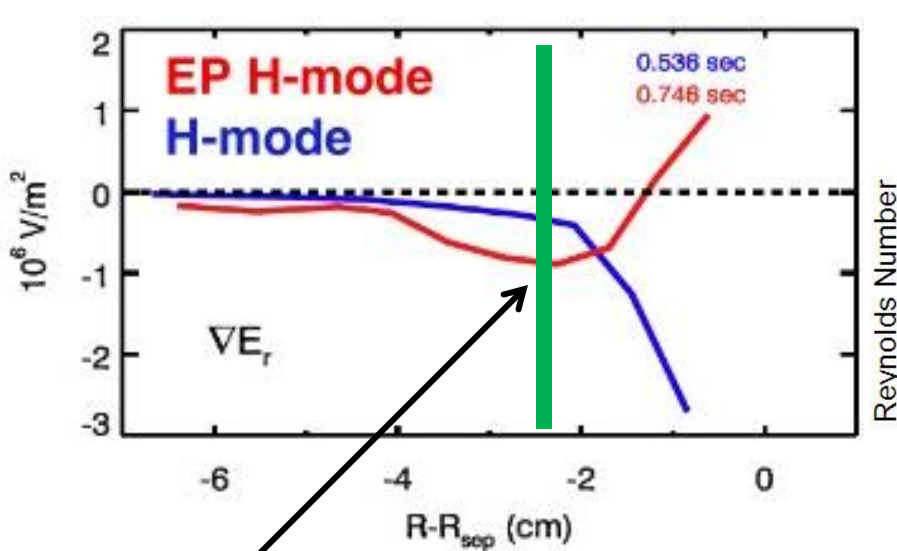
$\text{Re}^* \equiv$ critical Reynolds number (2000~3000), $\text{Re} < \text{Re}^*$ (laminar, H-mode)



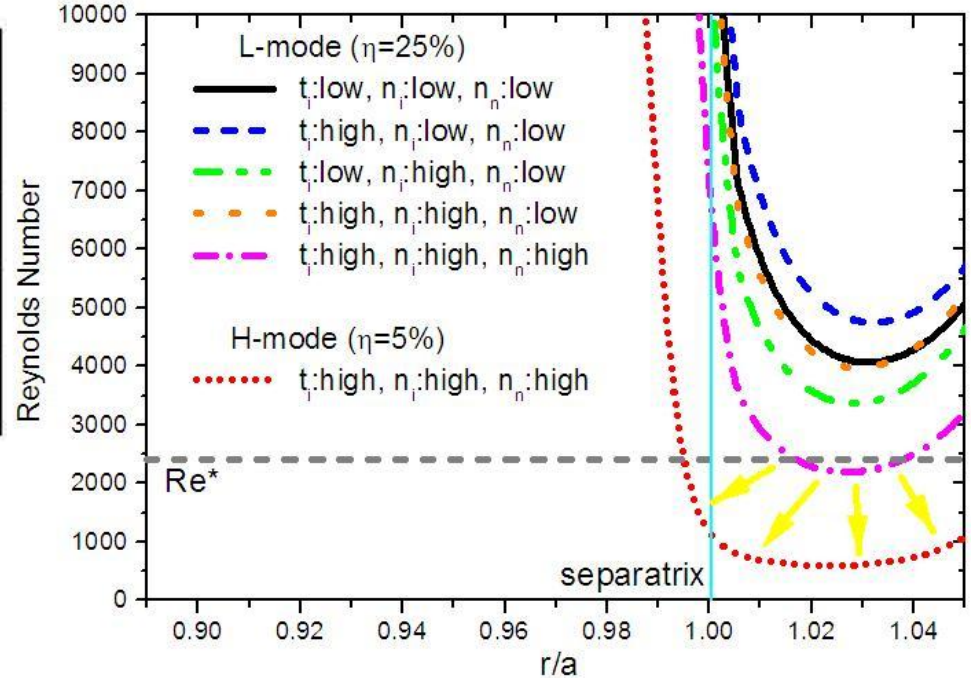
possible scenario of EPH-mode transition

- ELM** → change of ion velocity profile
- change of E_r profile
- making $\nabla^2 E \sim 0$ where inner part of H-mode transport barrier
- $\text{Re} < \text{Re}^*$ (critical Reynolds number)
- turbulence suppression
- confinement increase

Re for H-mode and EP H-mode transitions



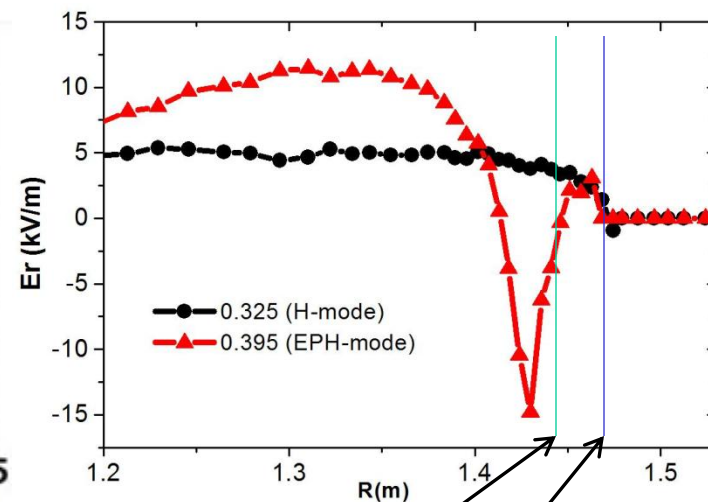
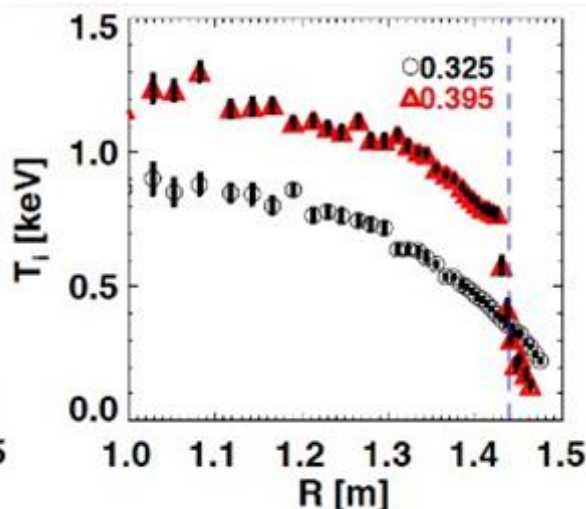
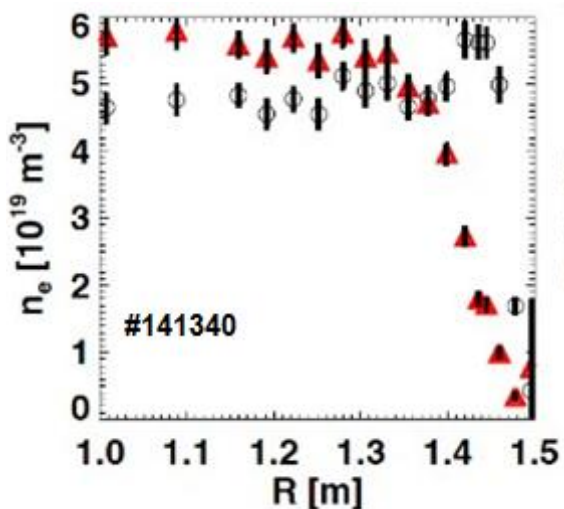
$\nabla^2 E \rightarrow 0$ (Inflection point)



$$Re = \frac{4\epsilon_0}{\pi} \eta^2 \frac{B}{m_i n_i (\sigma_{i-n} n_n)^2 v_{\perp}} \nabla^2 E$$

- ▶ for normal H-mode, $|\nabla^2 E| > 0$ and n_n determines Re profile
- ▶ for EPH-mode, $\nabla^2 E \rightarrow 0$ and $Re < Re^*$, and n_n can be lower than H-mode

Analysis of EPH-mode based on GCS theory



transport barrier for EPH-mode
transport barrier for early H-mode

- ▶ location of EPH-mode transport barrier is close to $\nabla^2 E \rightarrow 0$ area
- ▶ transport barrier further inside of separatrix \rightarrow lower neutral density
- ▶ lower neutral density \rightarrow lower fueling & cooling \rightarrow low n_e & high T_i

Summary

- ▶ EP H-mode is triggered by ELM and shows 50% increase of energy confinement on NSTX
- ▶ density fluctuation measurement by FReTIP and high-k scattering system showed reduction at EP H-mode transition

- ▶ FReTIP data showed n=1 mode development as pre-stage of EPH-mode transition.
- ▶ toroidal ion velocity and E_r profiles of EP H-mode changes to have $\nabla^2 E \rightarrow 0$ area inside of separatrix.

- ▶ Reynolds number of ion-neutral collision from gyro-center shift theory is proportional to the second gradient of E_r ($Re \propto \nabla^2 E_r$)
- ▶ location of EPH-mode transport barrier is close to $\nabla^2 E \rightarrow 0$ area

- ▶ possible mechanism of EP H-mode transition :
ELM $\Rightarrow \nabla^2 E \rightarrow 0 \Rightarrow Re < Re^* \Rightarrow$ turbulence suppression

* Please e-mail to kcllee@pppl.gov for reprint and further questions