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Comparison of Pedestal Stability on the Power Scrape-Off Layer Width in NSTX

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Focused efforts have confirmed the importance of plasma current scaling for both pedestal gradients and divertor heat flux widths

- Results of coordinated U.S. experiments to understand the scaling of:
 - the divertor power flux width, λ_q and
 - the H-mode pedestal with engineering parameters
- λ_q was found to scale ~ B_{pol}^{mid} under attached divertor conditions
- Pedestal height was found to be well described by Peeling-Ballooning modes
- Given the common dependence on Ip on both λ_q and the pedestal, a correlation is possible



In NSTX, lithium wall conditioning has also been shown to have a strong impact on λ_q

- 0 mg (boronized) and 150 mg of Li yield similar heat fluxes (inter-ELM averaged)
- λ_q contracts with addition of 150 mg of Li
 - Likely due to the elimination of small Type V that can be ubiquitous in boronized conditions [Maingi2005]
- With sufficient Li (300 mg), heat flux is also reduced
 - Measurements made with DBIR camera to account for surface emissivity effects
 - λ_q contracts further still



TK Gray, IAEA-FEC ...

There are numerous definitions for the SOL width, λ_{q}



- Integral: $\lambda_{q, int}$
- Numerical integration of radial heat flux profile
- Assumes axisymetric heat deposition profile

$$\lambda_{q,int} = \frac{\int \left(q(s) - q_{\rm BG}\right) ds}{q_{\rm max}}$$

- Diffusive-Gaussian (D-G) Model[Eich2012]
 - Simple semi-empirical model
 - Assumes λ_q is an exponential in the SOL before entering the divertor

$$q(\bar{s}) = q_0 \exp\left(-\frac{\bar{s}}{\lambda_q f_{exp}}\right)$$

- The exponential "diffuses" into the private flux region as it enters the divertor
- No mechanism for this diffusion is put forth
 - But is clearly observed in all divertor footprints
- 5 free parameters which require nonlinear least squares fitting to determine

 $q(\bar{s}) = \frac{1}{2}q_0 \exp\left(\left(\frac{S}{2\lambda_q f_{exp}}\right)^2 - \frac{\bar{s}}{\lambda_q f_{exp}}\right) \operatorname{erfc}\left(\frac{S}{2\lambda_q f_{exp}} - \frac{\bar{s}}{S}\right) + q_{BG}$

s = radial coordinate s₀ = Strike Point Location

 $\overline{S} = S_{q_0} = S_{ak}$ heat flux $q_{BG} = background heat flux$ $\lambda_q (\lambda_{q, Eich}) = e-folding length of q in SOL$ S = Gaussian diffusion parameter

$\lambda_{q, Eich}$ contracts with I_p

- Diffusive-Gaussian analysis used to analyze heat flux profile
 - $\lambda_{q, \ \text{Eich}} \sim \lambda_{q, \ \text{int}}$
- λ_{q, Eich} is similar between 0 and 150 mg data
 - Perhaps less scatter in the data
 - This suggests that λ_q is similar at the outer midplane (OMP) for 0 and 150 mg data
 - But their must be more diffusion in the 0 mg data to achieve larger $\lambda_{q, int}$'s
- 300 mg data is still contracted under this analysis



At high I_p , for all Li deposition amounts, $1 \le \lambda_q \le 3.5$ mm

Could Peeling-Ballooning (P-B) or Kinetic Ballooning Modes (KBMs) set λ_q?

- Pedestal profiles are characterized using modified TANH functions
- The peak pedestal pressure is found to scale as ~ Ip²
- The pedestal pressure height is set by Peeling-Ballooning physics
- Useful to describe the pedestal with a normalized pressure gradient:

$$\alpha = -Rq^2 \frac{d\beta}{dr}$$



Lithium wall conditioning has a strong effect on the electron density (and therefore pressure) pedestal



Increasing amounts of Li have a nearly continuous effect on the pedestal



- With increasing I_p, the peak of $-\nabla_{\perp}P_e$ increases
- $abla_{\perp}P_e$ near the seperatrix remains largely unchanged at -20 kPa/m with I_p
- According to Peeling-Ballooning theory, the onset of ELMs occur when a critical value of the normalized pressure gradient occurs, acrit
 - Possible that P-B sets the edge transport into the SOL even before the critical value is reached



Preliminary analysis shows only a weak correlation between λ_q and the α parameter

- Comparison of λ_q's from from a range of discharges
 - high δ , LSN
 - $P_{NBI} = 4 MW$
 - $0.7 \le I_p \le 1.2$ MA
 - 150 and 300 mg of
 Li
- Only a weak correlation is found with α_{MHD} evaluated at ψ_N = 0.85 or 0.9



Conclusions

- λ_q has been found to scale strongly with I_p and predischarge Li wall conditioning
 - I_p dependence varies with amount of pre-discharge Li wall conditioning
 - At high $I_p,\,\lambda_q\sim 2$ 4 mm regardless of amount of Li deposited
- Likewise, the electron pressure pedestal varies with I_p and Li
- However, only a weak correlation between λ_q and α has been found at $\psi_N=0.85$ or 0.9
 - No correlation has been found with α evaluated closed to the seperatrix

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NSTX Divertor Diagnostics

- 2, 30 Hz IR cameras
 - 1 viewing the lower divertor, 1 the upper
 - sensitive to 6 13 μm
- 1 Fast IR camera (≤ 16 kHz)
 - viewing the lower divertor
 - sensitive to 6 13 μm
 - equipped with dual band optics
- 2 Fast Phantom cameras (≤ 100 kHz)
 - each viewing nearly the entire lower divertor
 - bandpass filtered
- 2 1D CCD array cameras
 - measuring Dα and Li II emission
- 3 channels of lower divertor bolometry
- Flush mounted divertor Langmuir probes



The Dual-Band IR Camera Allows Measurement of Divertor Surface Temperature with Variable Surface Emissivity

- The addition of lithium complicates the measurement of divertor surface temperature, T_{surf}
 - Li emissivity (clean) ~ 0.1
 - Graphite ~ 0.8
 - Lithium and Carbon are eroded and redeposited constantly through out the discharge
- Using a Santa Barbra Focal Plane Camera
 - 128x128 pixels
 - frame rate ≤ 16 kHz (typical operation at 1.6 kHz with dual band optics)
- 2 different IR wavelength bands are imaged simultaneously[McLean 2012]
 - MWIR: 7 10 μm
 - LWIR: 10 13 μm
- The ratio of the 2 bands yields T_{surf}(r,t)
 - Assumes the surface emissivity is similar across both wavelength bands
 - Not a bad assumption for a diffuse, grey body emitter
- Once T_{surf} is known, heat flux can be calculated
 - 2D finite difference calculation (THEODOR)

