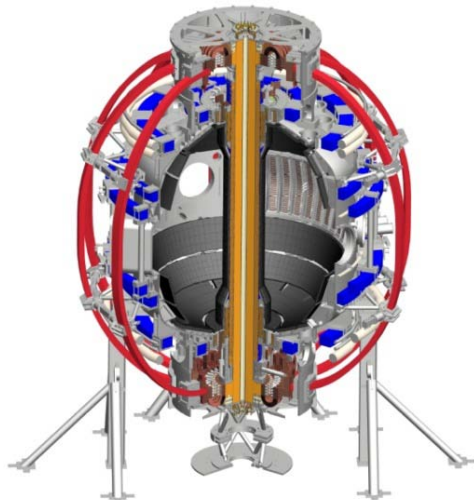


M3D-K Simulations of Wave-Particle Interaction with Realistic Fast Ion Distribution

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Lodestar
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Nova Photonics
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54th APS-DPP Meeting
Providence, RI, Oct. 29–Nov. 2 2012



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Initialization with Realistic Fast Ion Distributions could Improve Accuracy of Hybrid simulations

- Fluid/kinetic hybrid simulation codes such as M3D-K typically use a simple analytic form to specify the initial fast ion distribution.
→ This could limit the simulation accuracy when the analytic distribution is a poor fit to the actual distribution.
- An interface between the output of tokamak fast ion Monte Carlo code TRANSP/NUBEAM and M3D-K has been developed.
→ More realistic fast ion distributions from NUBEAM can be directly used as M3D-K inputs to improve simulation accuracy.
- M3D-K linear simulations of non-resonant kink mode and TAE mode on NSTX have been performed with analytic and realistic fast ion distributions to validate the new method and check the possible benefits.

M3D-K: Global Kinetic/MHD Hybrid Code

$$\rho \frac{d\mathbf{v}}{dt} + \rho(\mathbf{v}_i^* \cdot \nabla) \mathbf{v}_\perp = -\nabla \cdot \mathbf{P} - \nabla \cdot \mathbf{P}_h + \mathbf{J} \times \mathbf{B} - \mathbf{b}\mathbf{b} \cdot \nabla \cdot \Pi_i$$

$$\mathbf{P}_h = P_\perp \mathbf{I} + (P_\parallel - P_\perp) \mathbf{b}\mathbf{b}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}_i) = 0$$

$$\rho \left[\frac{\partial \mathbf{v}}{\partial t} + \mathbf{v} \cdot \nabla \mathbf{v} + (\mathbf{v}_i^* \cdot \nabla) \mathbf{v}_\perp \right] = -\nabla \mathbf{p} + \mathbf{J} \times \mathbf{B} + \mu \nabla^2 \mathbf{v}$$

$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \eta \mathbf{J} - \frac{\nabla_\parallel p_e}{ne}$$

$$\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$$

$$\mathbf{J} = \nabla \times \mathbf{B}$$

$$\frac{\partial p}{\partial t} + \mathbf{v} \cdot \nabla p = -\gamma p \nabla \cdot \mathbf{v} + \nabla \cdot n \chi_\perp \nabla \left(\frac{p}{\rho} \right) - \mathbf{v}_i^* \cdot \nabla p - \gamma p \nabla \cdot \mathbf{v}_i^* + \frac{\mathbf{J} \cdot \nabla p_e}{ne} + \gamma p_e \mathbf{J} \cdot \nabla \left(\frac{1}{ne} \right)$$

$$\frac{\partial p_e}{\partial t} + \mathbf{v} \cdot \nabla p_e = -\gamma p_e \nabla \cdot \mathbf{v} + \nabla \cdot n \chi_{\perp e} \nabla \left(\frac{p_e}{\rho} \right) + \frac{\mathbf{J}_\parallel \cdot \nabla p_e}{ne} - \gamma p_e \nabla \cdot \left(\mathbf{v}_e^* - \frac{\mathbf{J}_\parallel}{ne} \right)$$

$$\text{where } \mathbf{v}_e^* \equiv -\frac{\mathbf{B} \times \nabla p_e}{neB^2}, \quad \mathbf{v} \equiv \mathbf{v}_i - \mathbf{v}_i^* = \mathbf{v}_e - \mathbf{v}_e^* + \frac{\mathbf{J}_\parallel}{ne}, \quad \mathbf{v}_i^* \equiv \mathbf{v}_e^* + \frac{\mathbf{J}_\perp}{ne}$$

*Fu et al. Phys. Plasmas 2006

➤ Extended MHD equations
(Single fluid, **Two-fluid (but not used in this work)**)

➤ Hybrid (**Kinetic fast ions**)

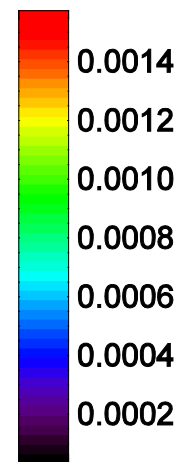
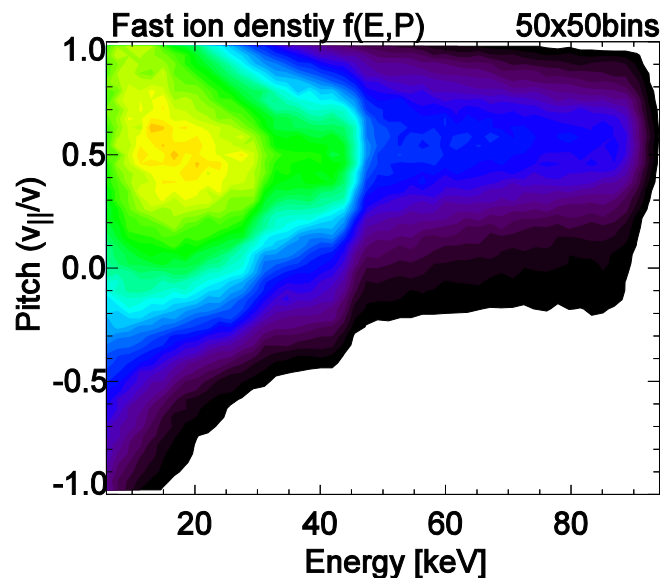
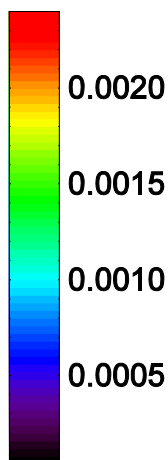
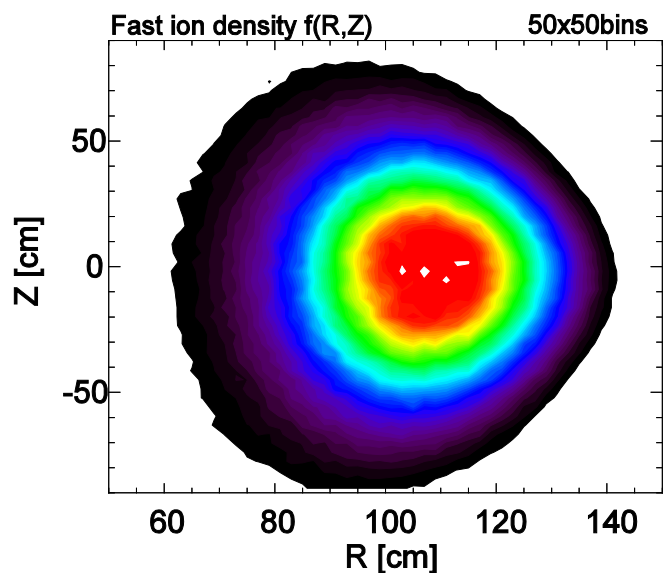
➤ Kinetic fast ions couple back to fluid model through pressure tensor, which is calculated by gyrokinetic or drift-kinetic equation

Main Features of M3D-K

- Use realistic geometry
- Use experimental parameters and profiles
- Global
- Linear and nonlinear
- Non-perturbative fast ion effects (i.e., can model EPM)
- Numerical equilibrium calculation with plasma rotation and anisotropic fast ion pressure
- Interface with TRANSP/NUBEAM fast ion distribution is developed and tested.

Numerical Fast Ion Distribution is Calculated by Monte Carlo Modeling NUBEAM

- NUBEAM module in TRANSP: a Monte Carlo code for 4D ($R, Z, \lambda=v_{||}/v, E$) time dependent simulation of fast ion transport in tokamak devices.
- Assume fast ions behave classically; could add anomalous radial diffusion operator based on user provided diffusivities
- Includes guiding center drift orbiting, collisional, and atomic physics effects



Fast Ion Distribution Needed to be Transformed from (R,Z,λ,E) to Constants-of-Motion Space

- M3D-K code needs fast ion distribution $f(R,Z,\lambda,E)$ expressed in terms of constants of the motion (P_ϕ, μ, E) and the function must be smooth enough to allow derivatives to be taken with E and toroidal angular momentum P_ϕ
- In addition to P_ϕ, μ, E , the sign of $v_{||}$ is also needed to resolve orbit type.
- NUBEAM calculated fast ion distribution $f(R,Z,\lambda,E)$ needs to be transformed to $f(P_\phi, \mu, E)$ coordinates.

$$\int f(x, y, z, vx, vy, vz) d^3\mathbf{x} d^3\mathbf{v} \quad \text{where}$$

$$= \int f(R, Z, \lambda, E) R \sqrt{E} dR dZ d\lambda dE$$

$$= \sum_{v_{sign}} \int f_{v_{sign}}(P_\phi, \mu, E) \mathfrak{T}_{v_{sign}}(P_\phi, \mu, E) dP_\phi d\mu dE$$

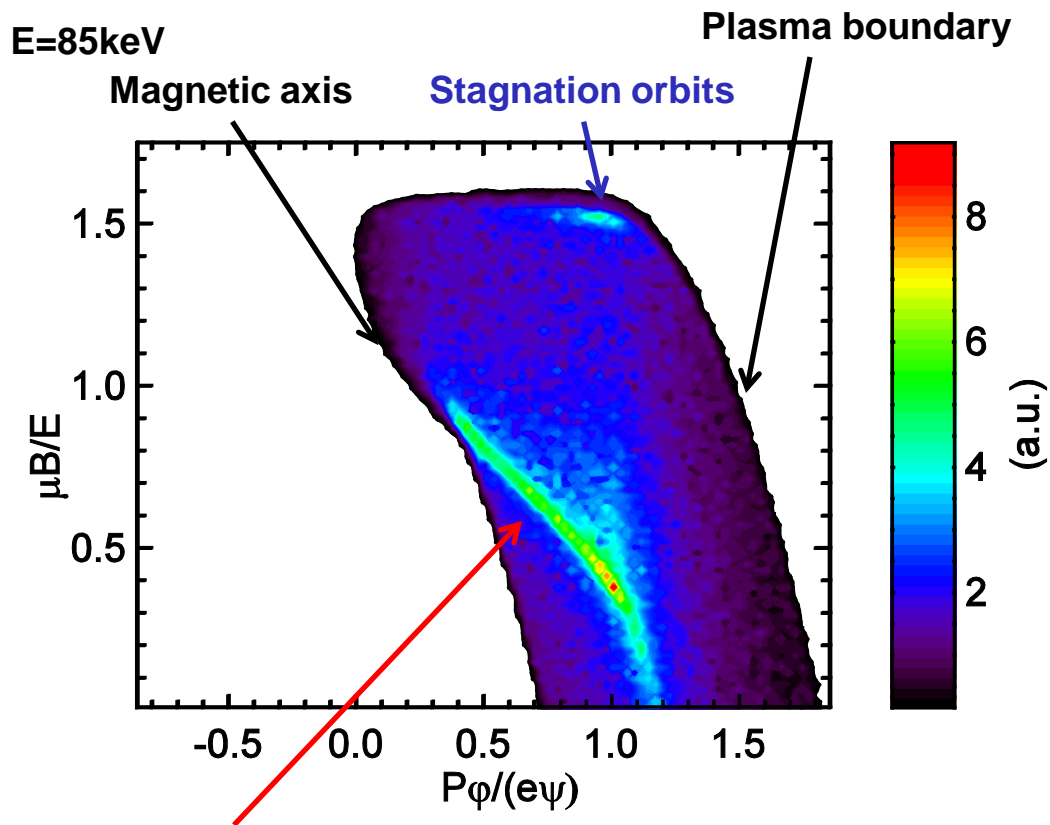
$$P_\phi = (eA_\phi + mv_\phi)R \approx e\psi(R, Z) + mv_{||}R \frac{B_\phi(R, Z)}{B(R, Z)}$$

$$\mu = \frac{\frac{1}{2}mv_\perp^2}{B} = \frac{[1 - (v_{||}/v)^2]E}{B(R, Z)} \quad \lambda = v_{||}/v$$

- Use TRANSP “Plasma State” utility to get ψ and B at each particle location (R, Z) , and then calculate P_ϕ, μ

Jacobian of the Transform Strongly Depends on Particle Orbit Topology

counter-current Jacobian of the transform
from $(R, Z, v_{\perp}^2, v_{\parallel})$ to (P_{φ}, μ, E)



Trapped/passing boundary

➤ Sharp change of Jacobian appears in the boundary of different particle orbit topology because

$$\mathfrak{J}(P_{\varphi}, \mu, E) = \frac{4\pi^2 \tau_{bounce}}{qm^2}$$

*J. Egedal Nucl. Fusion 2005

Jacobian of the Transform is Calculated by Monte Carlo Method

$$\int f(x, y, z, vx, vy, vz) d^3\mathbf{x} d^3\mathbf{v} = \int f(R, Z, \lambda, E) R \sqrt{E} dR dZ d\lambda dE = \sum_{v_{sign}} \int f_{v_{sign}}(P_\varphi, \mu, E) \mathfrak{J}_{v_{sign}}(P_\varphi, \mu, E) dP_\varphi d\mu dE$$

➤ An exact analytic $\mathfrak{J}_{v_{sign}}(P_\varphi, \mu, E)$ cannot be calculated because it requires integration over phase space orbits.

➤ A Monte Carlo approach is used to calculate Jacobian

➤ Procedure:

(1) Generate a large population of ~50 million particles uniformly in phase space (R, Z, λ, E) , within the same range as in the NUBEAM distribution

(2) Transform coordinates to constants-of-motion space

(3) Bin the particles in (P_φ, μ, E) space for each $v_{||}$ sign.

(4) The particle number in each (P_φ, μ, E) bin proportional to $\mathfrak{J}_{v_{sign}}(P_\varphi, \mu, E)$

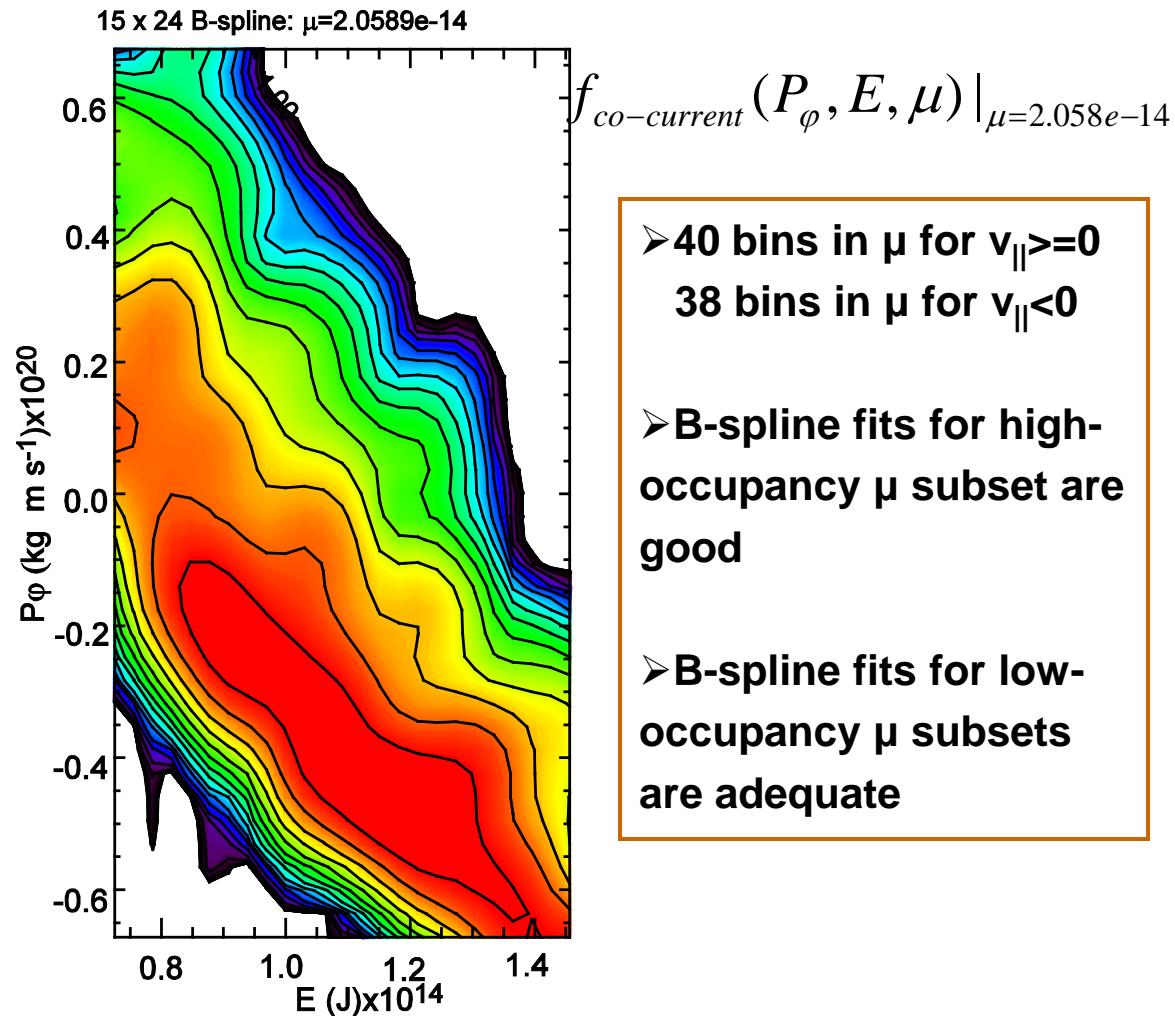
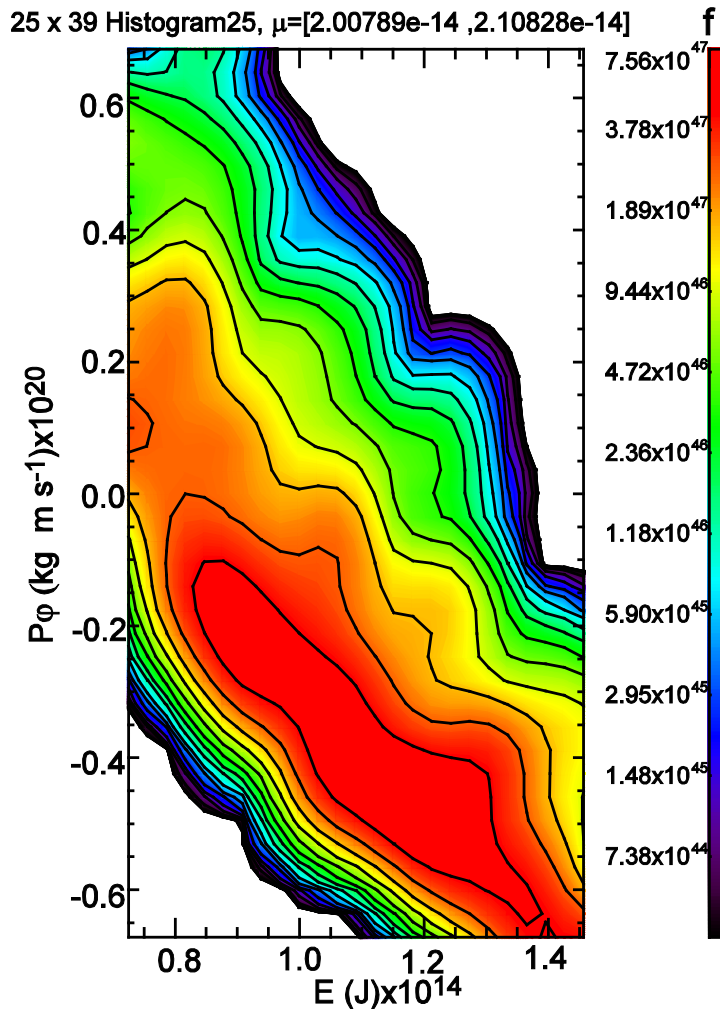
Constructing Fast Ion Density Function in Constants-of-Motion Space

- Divide NUBEAM output particle data into two subsets based on v_{\parallel} sign
- For each subset, sort particles by μ and divide them into several subpopulations of equal width in μ .
- Divide each subpopulation into a number of bins in the P_{ϕ} and E directions. The bin width in P_{ϕ} and E direction are the same for each subpopulation.
- Multiply by the numerically calculated Jacobian.
- Apply Gaussian smoothing in both P_{ϕ} and E directions.
- Fit 2D cubic B-spline to the smoothed data using GSL routines, with uniform knots and a number of coefficient s in each direction approximately $5/8$ the number of bins.
- Store the spline coefficients in a file, which can be used to construct fast ion density function in constants-of-motion space and perform quick spline and derivative interpolations at arbitrary location.
- **Improved from the work of Breslau et al. Sherwood Meeting 2011*

Good Agreement between Raw Fast Ion Density Function and Spline Fit in Constants-of-motion Space

Raw data from NUBEAM

Reconstructed from cubic B-spline fit

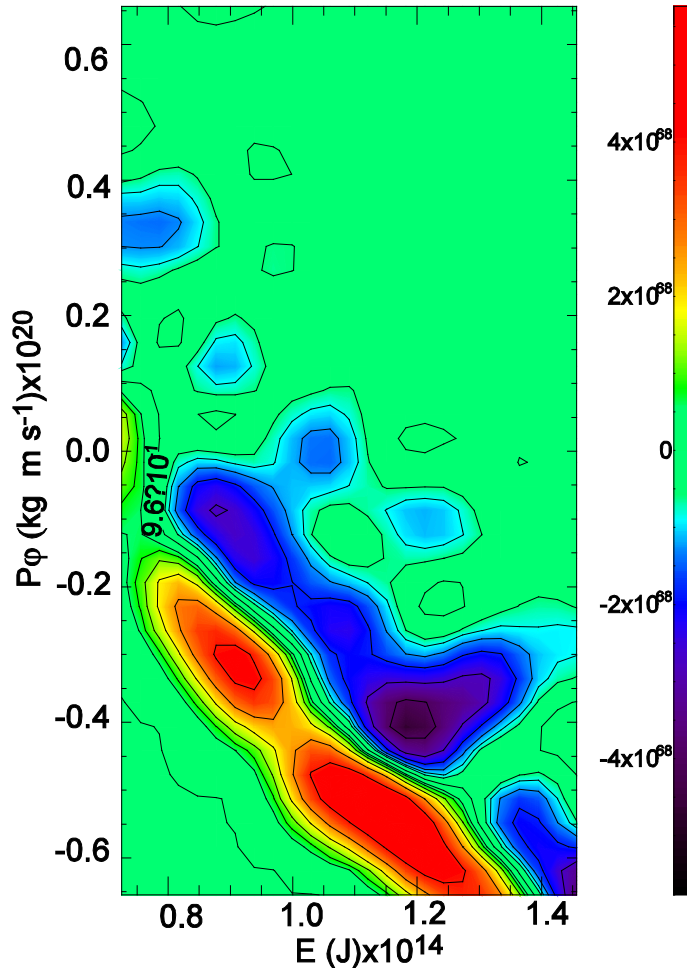


- 40 bins in μ for $v_{||} \geq 0$
- 38 bins in μ for $v_{||} < 0$
- B-spline fits for high-occupancy μ subset are good
- B-spline fits for low-occupancy μ subsets are adequate

Gradients are Smooth, Match well with NUBEAM Raw Data

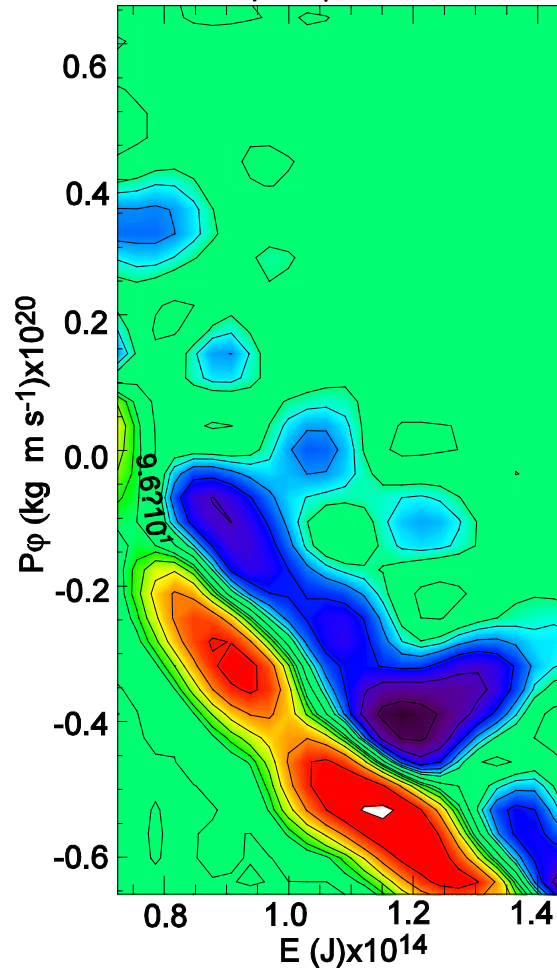
Raw data from NUBEAM

25 x 39 Histogram25, $\mu=[2.00789\text{e-}14, 2.10828\text{e-}14]$



Reconstructed from cubic B-spline fit

15 x 24 B-spline: $\mu=2.0589\text{e-}14$



$$\frac{\partial f_{co-current}(P_\phi, E, \mu)}{\partial P_\phi}$$

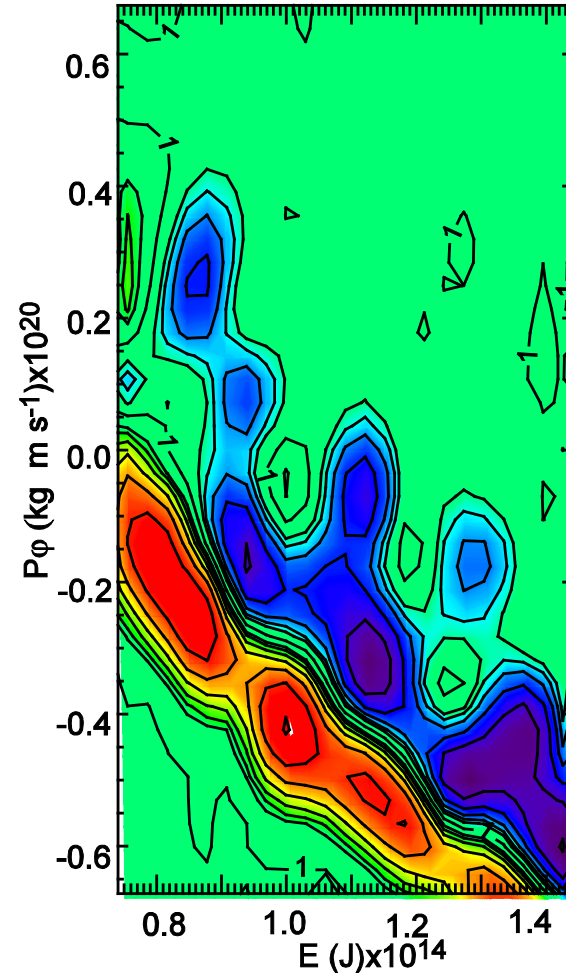
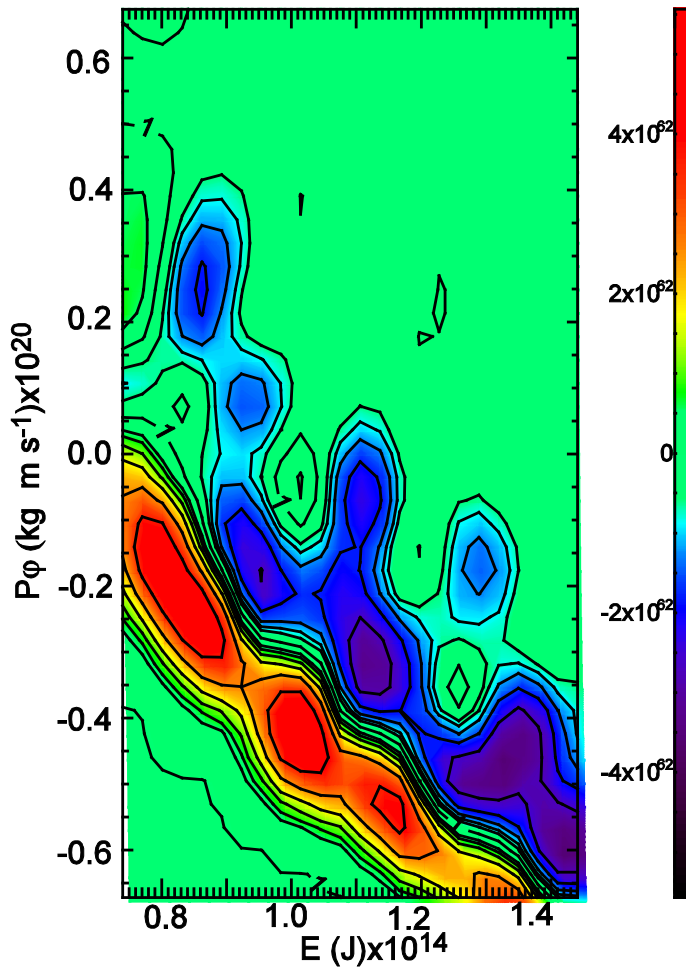
Gradients are Smooth, Match well with NUBEAM Raw Data

Raw data from NUBEAM

Reconstructed from cubic B-spline fit

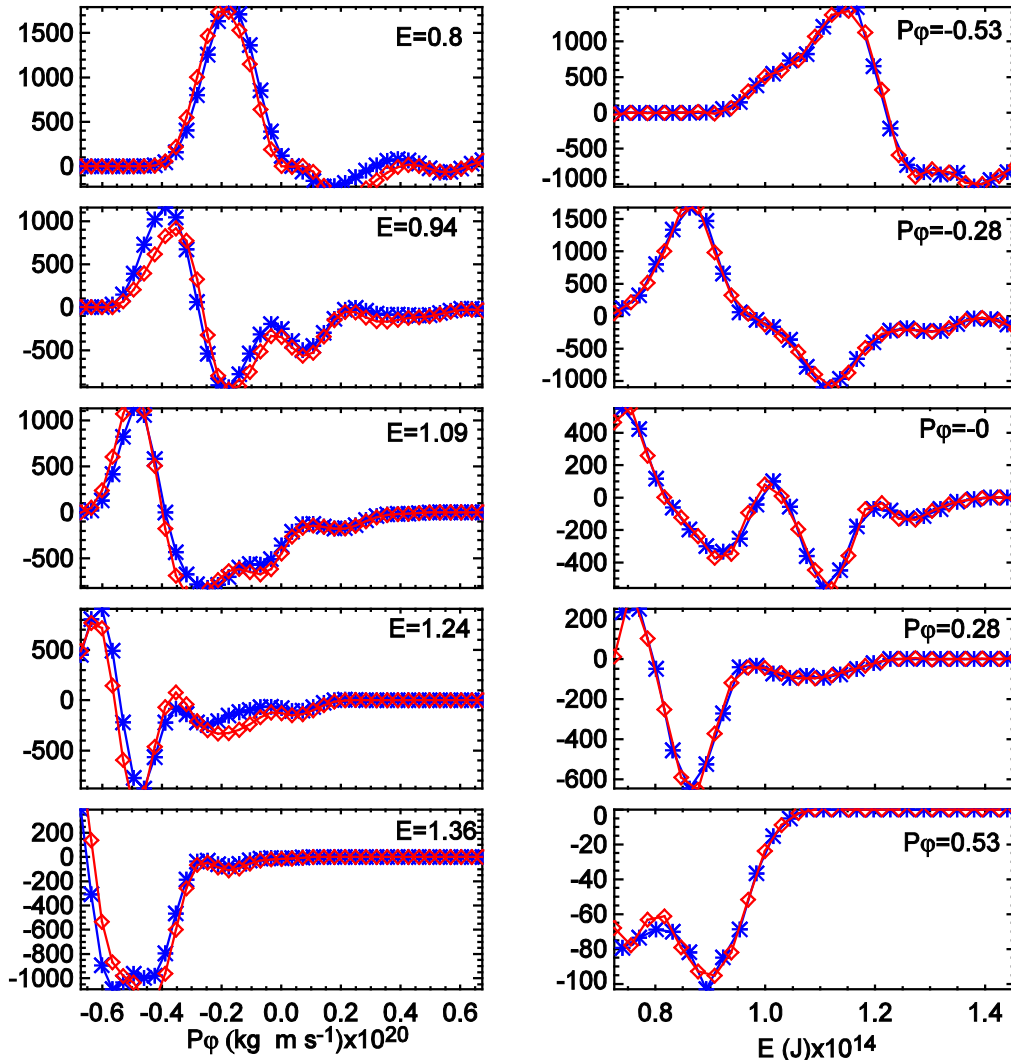
25 x 39 Histogram25, $\mu=[2.00789\text{e-}14, 2.10828\text{e-}14]$

15 x 24 B-spline: $\mu=2.0589\text{e-}14$



$$\frac{\partial f_{co-current}(P_\phi, E, \mu)}{\partial E}$$

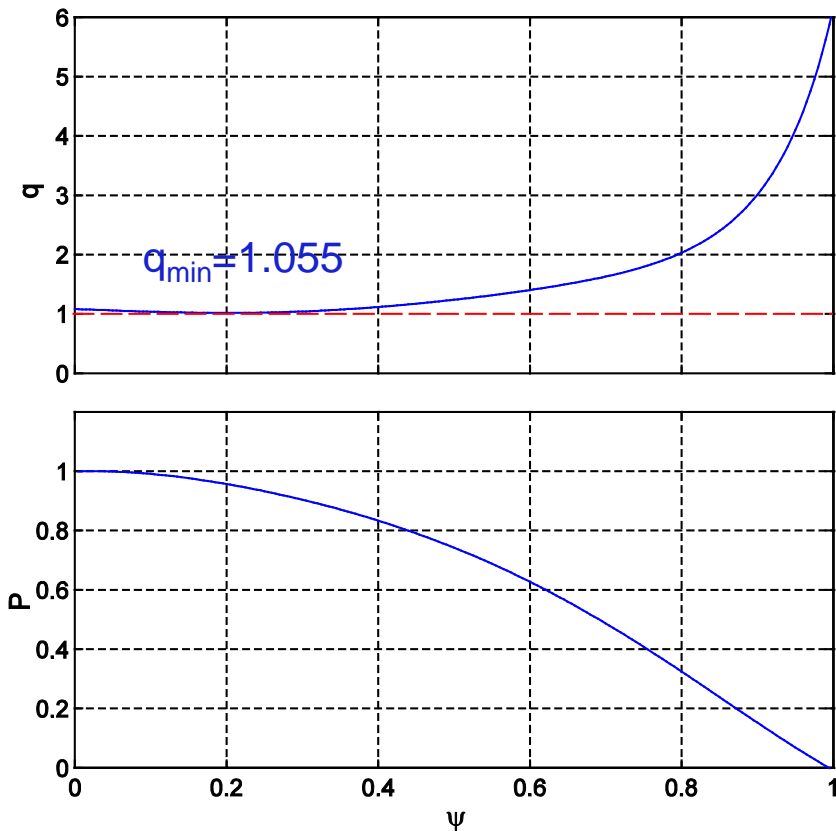
Gradients are Smooth, Match well with Raw Data (cont'd)



Blue: raw data from NUBEAM
 Red: reconstructed from cubic B-spline fit

$$\frac{\partial f_{co-current}(P_\phi, E, \mu)}{\partial E}$$

M3D-K Linear Simulations of Non-Resonant Kink Mode on NSTX



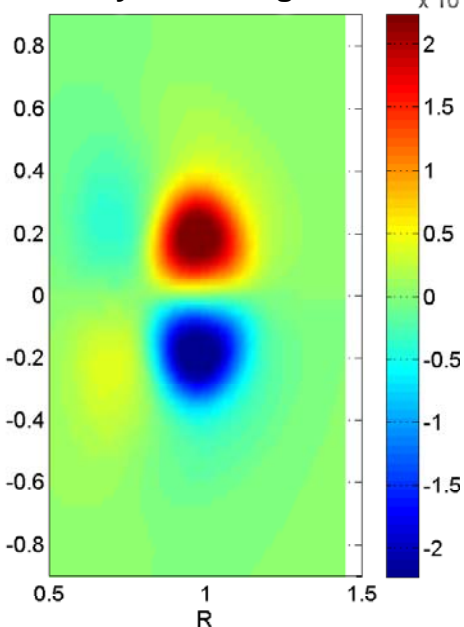
- Both experiments and simulations show that (1,1) non-resonant kink mode is unstable in NSTX when the shear is low and central safety factor is close but greater than one.
- The mode is a pressure-driven mode. It can saturate nonlinearly and persist, and can develop $m=2$ magnetic islands.
- The mode could severely impact performance: flatten rotation, redistribute fast ions, trigger neoclassical tearing modes.

*S. Gerhardt *Nucl. Fusion* 2011; J. Breslau *Nucl. Fusion* 2011;
F. Wang *Sherwood Meeting* 2012

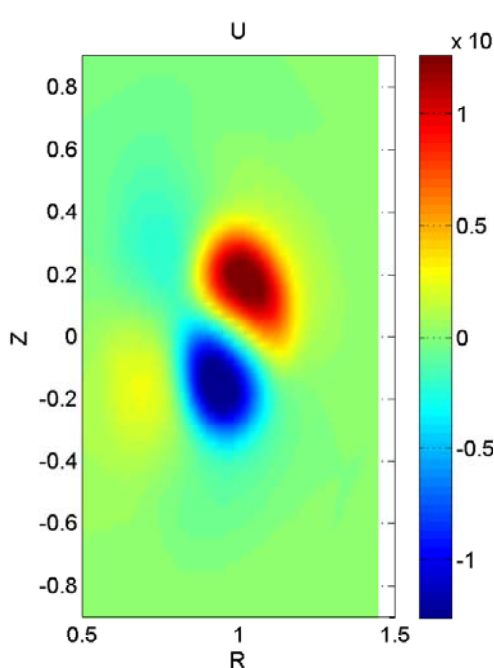
- NSTX #124379 at $t=0.635$ s equilibrium profiles
- $R=0.86$ m, $a=0.6$ m, $B_t(0)=0.44$ T, $n_e(0)=9.3 \times 10^{13}$ cm $^{-3}$
- $q_{\min}=1.055$, $\beta_{\text{thermal}}(0)=24\%$, $\beta_{\text{fi}}(0)=5\%$

Reasonable M3D-K Simulations Results Obtained with both Analytic and NUBEAM Fast Ion Distributions

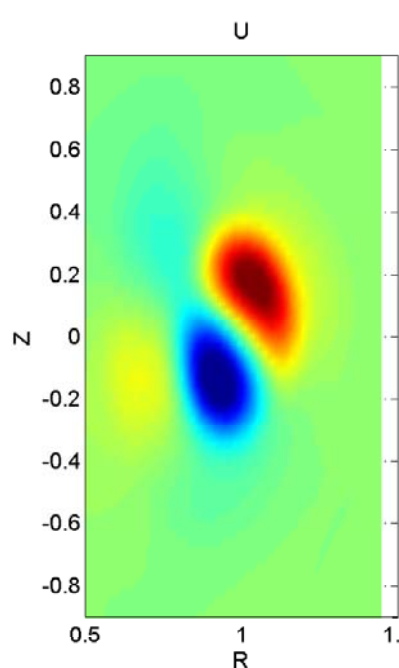
Velocity streaming function U



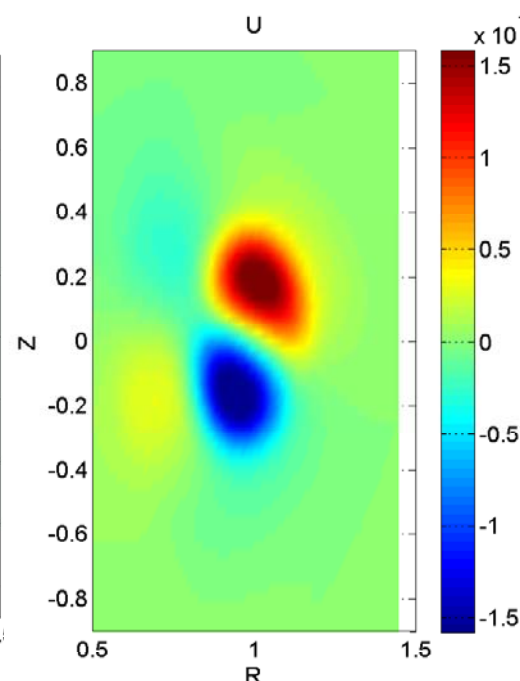
Without fast ions
 $\gamma/\omega_A = 3.04\%$



Analytic isotropic
fast ions
 $\gamma/\omega_A = 2.56\%$



Analytic anisotropic
fast ions
 $\gamma/\omega_A = 2.80\%$

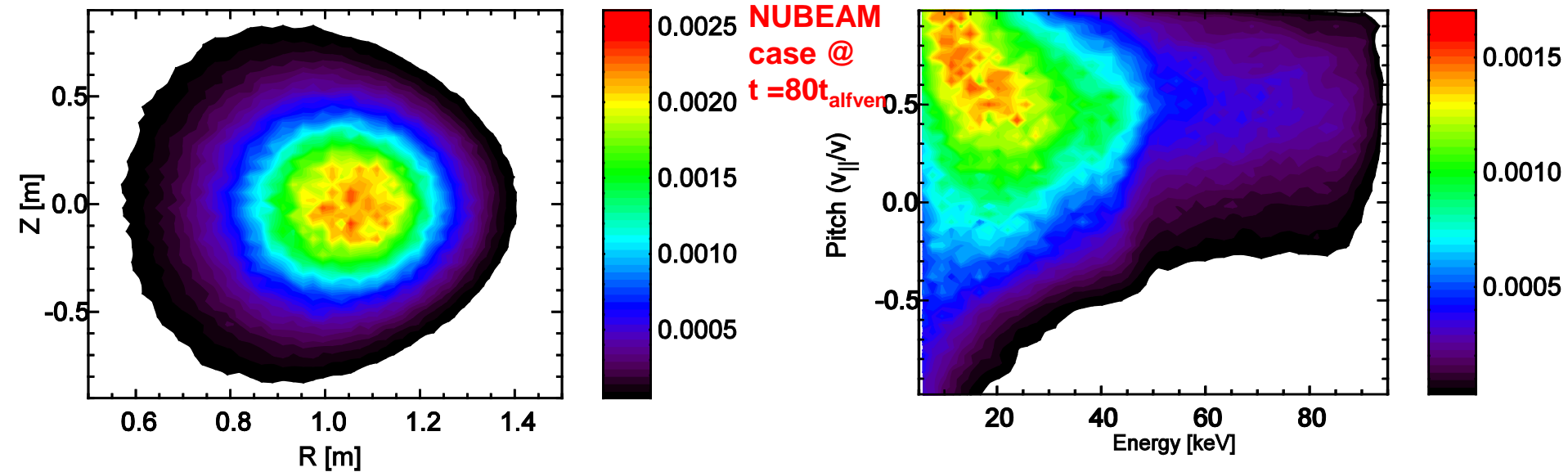
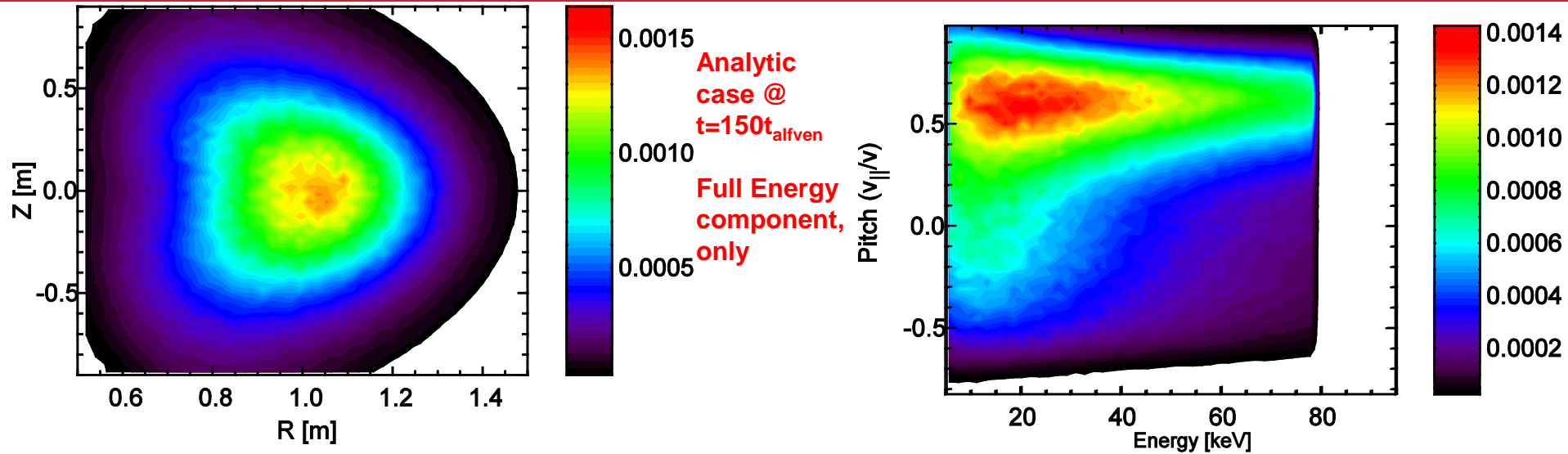


NUBEAM fast ions
 $\gamma/\omega_A = 2.70\%$

$$f = \frac{cH(v_0 - v)}{v^3 + v_c^3} \exp(-\psi / \Delta\psi) \exp[-(\Lambda - \Lambda_0)^2 / \Delta\Lambda^2], \Lambda = \frac{\mu B}{E}$$

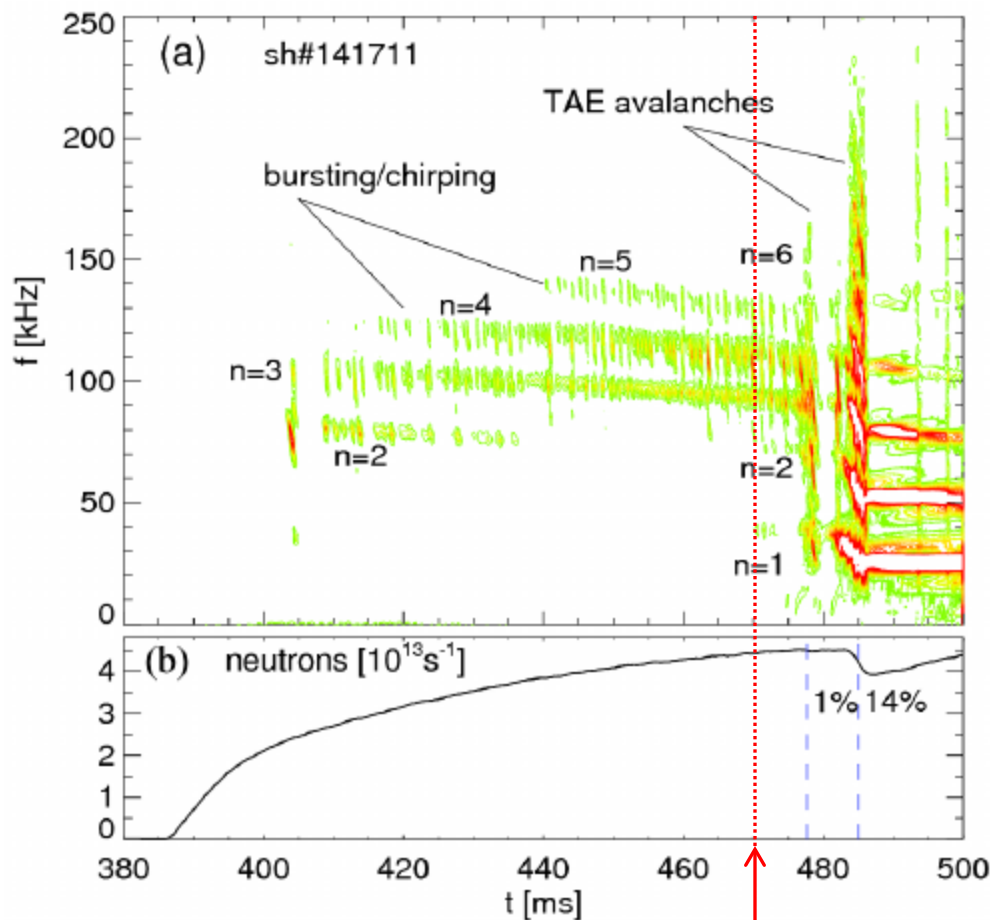
- $\beta_{\text{thermal}}(0) = 24\%$ fixed in all cases, $\beta_{\text{fi}}(0) / \beta_{\text{tot}}(0) = 5\%$ for the cases with fast ions;
- Fast ions weakly stabilize non-resonant kink mode, and induce mode rotation;
- Different fast ion distributions can change mode growth rate and mode frequency.

Initial Analytic & NUBEAM Fast Ion Distributions are Different, and they remain Unchanged during Linear Simulation

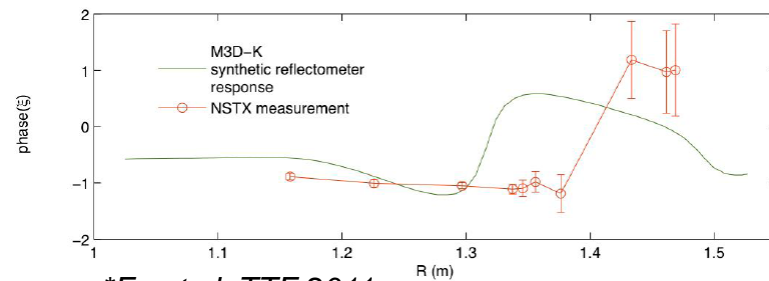
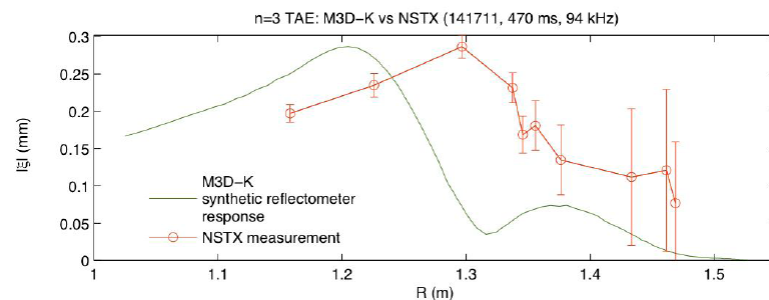


Similar to original NUBEAM distribution in slide 5

Previous M3D-K Linear Simulations Qualitatively Agree with Experimental Measurements



- Equilibrium and plasma profiles at $t=470$ ms are used as M3D-K inputs.
- Previous M3D-K linear simulation results (Fu TTF 2011) are similar to the measurements with respect to mode frequency, core-localized radial structure, and phase shift, but there is a significant radial shift between M3D-K simulations and measurements from multi-channel reflectometer.

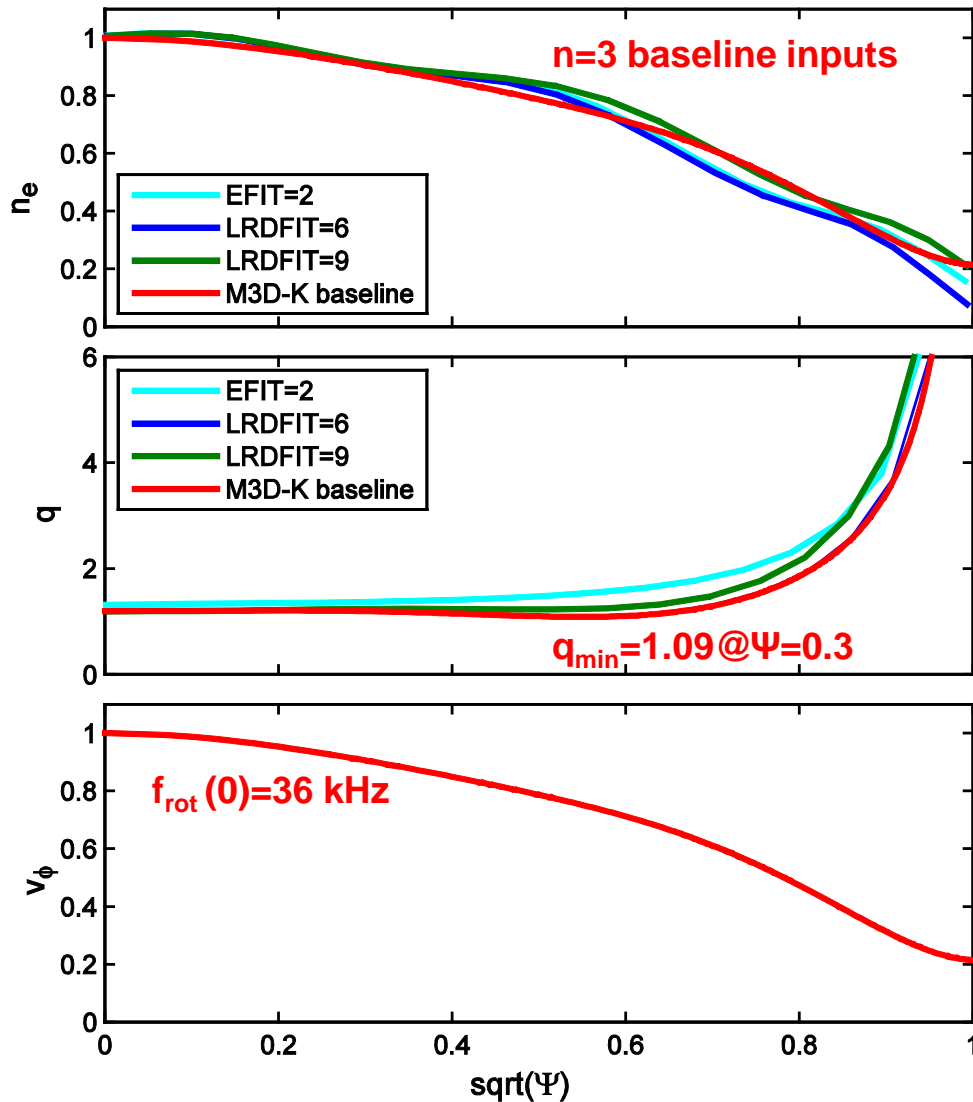


*Fu et al. TTF 2011

*at Courtesy of M. Podesta et al

$t=470$ ms

Parameters and Profiles for M3D-K Simulation of n=3 TAE on NSTX



➤ NSTX parameters ($B=0.55\text{T}$, $R=0.85\text{m}$, $a=0.67\text{m}$) and equilibrium profiles at 470 ms of shot 141711 are used as M3D-K inputs

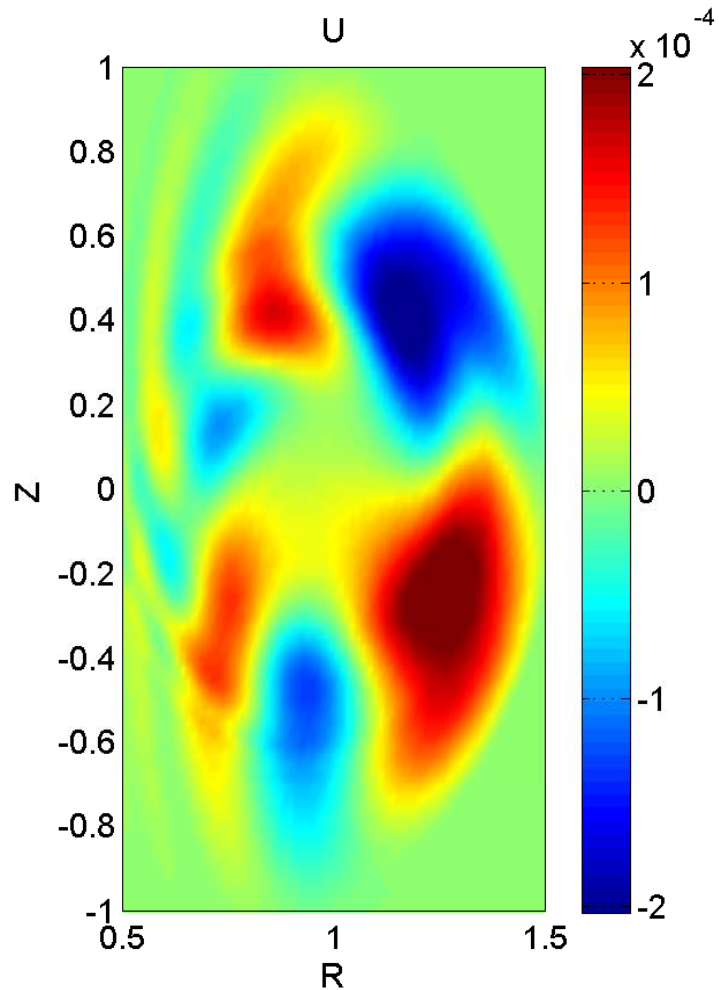
- $n_e(0) = 4.4 \times 10^{13} \text{ cm}^{-3}$
- $T_e(0) = 1.4 \text{ keV}$
- $T_i(0) = 1.3 \text{ keV}$

➤ Self-consistent equilibrium with plasma rotation and fast ion pressure included

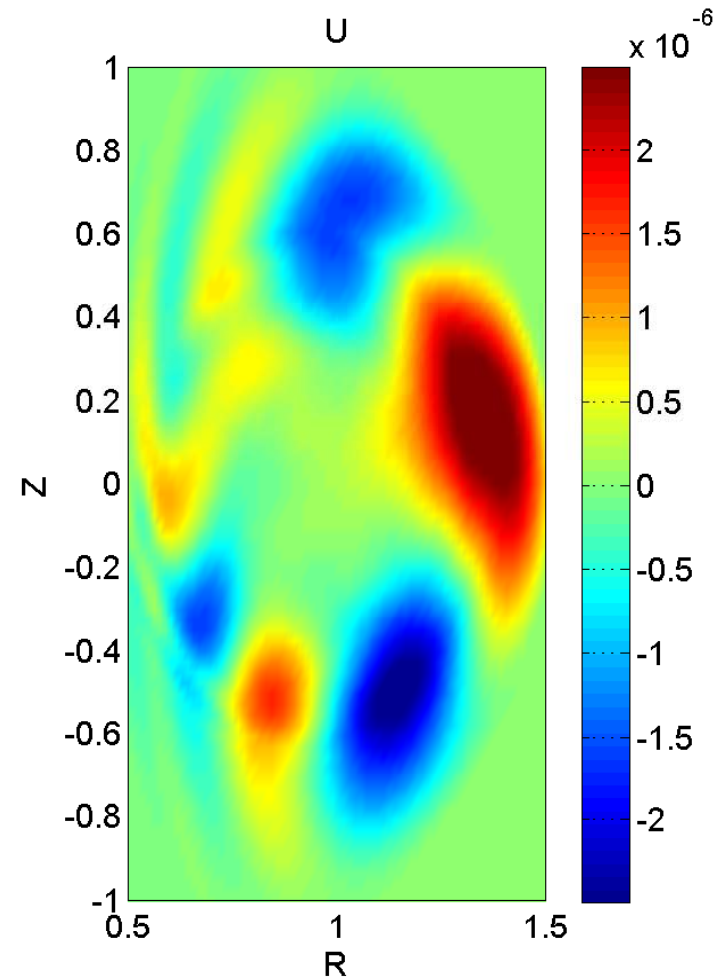
- $\beta_{\text{tot}}(0) = 24\%$
- $\beta_{\text{fi}}(0) / \beta_{\text{tot}}(0) = 0.67$

$P_{\text{fi_perpendicular}} \neq P_{\text{fi_parallel}}$

Similar Mode Structure Obtained with Analytic and NUBEAM Fast Ion Distributions

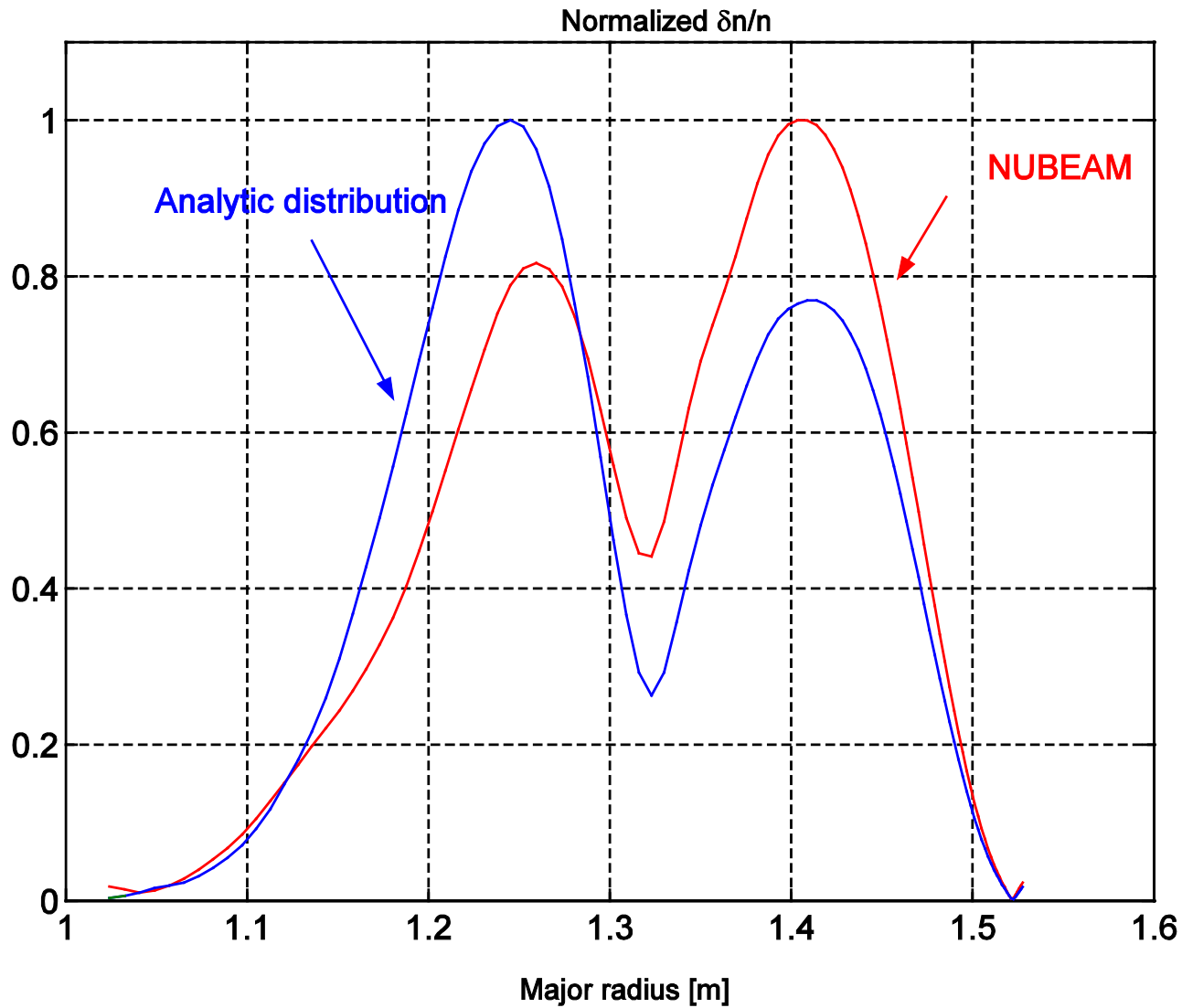


Analytic distribution, pitch(0)=0.64
 $\gamma/\omega_A = 6.60\%$, $f=113\text{kHz}$



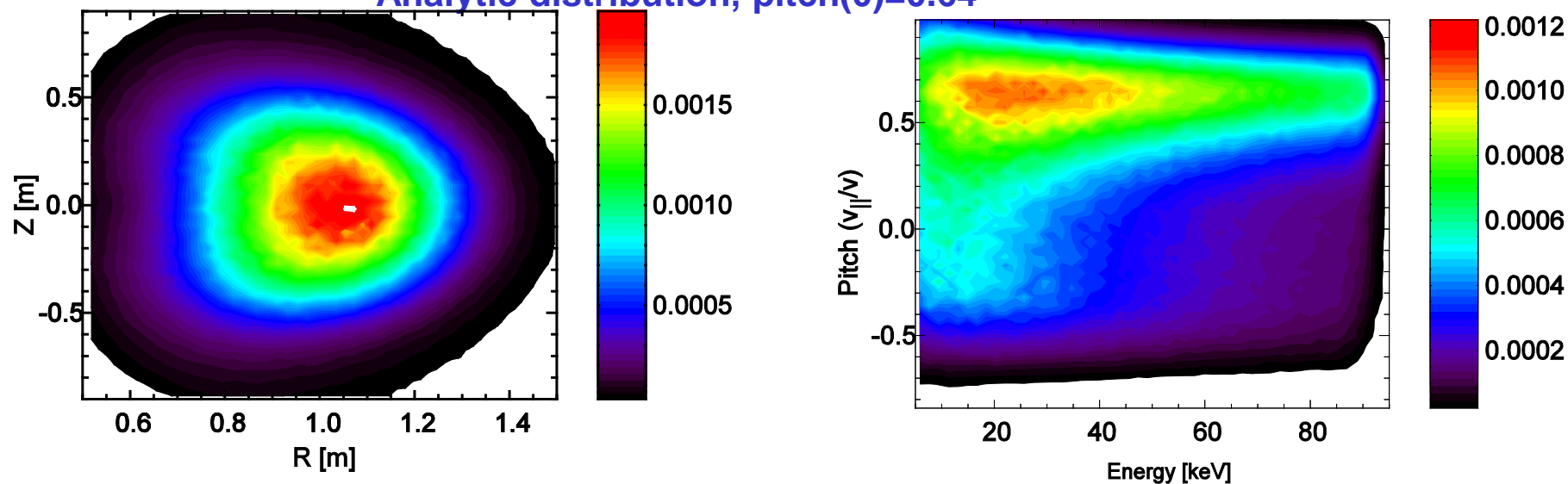
NUBEAM distribution
 $\gamma/\omega_A = 4.7\%$, $f=110\text{kHz}$

Detailed Mode Structure Affected by Initial Fast Ion Distribution

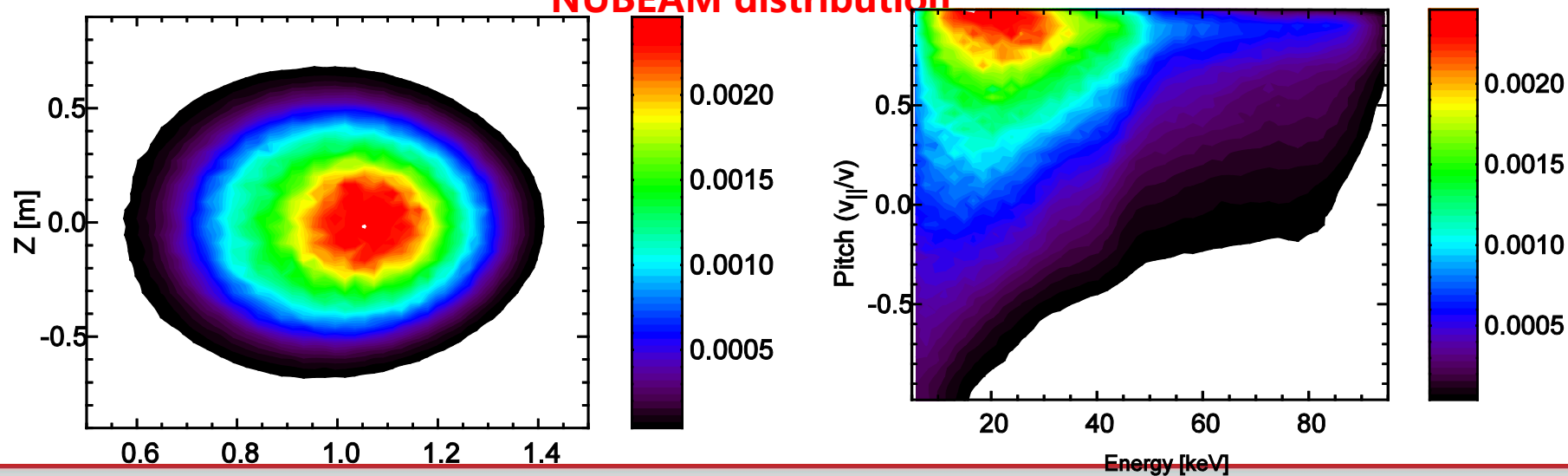


Comparison of Analytic and NUBEAM Fast Ion Distributions

Analytic distribution, pitch(0)=0.64



NUBEAM distribution



Conclusions and Future Plan

- **An interface between NUBEAM output and M3D-K has been developed and tested. More realistic NUBEAM fast ion distribution now can be used as M3D-K inputs to improve simulation accuracy.**
- **M3D-K linear simulation of non-resonant kink mode on NSTX show that fast ions can weakly stabilize the mode. Different fast ion distributions (analytic distribution and NUBEAM fast ion distribution) can affect mode growth rate and frequency.**
- **Preliminary simulation results of $n=3$ TAE mode with NUBEAM fast ion distribution are similar to previous results with analytic fast ion distribution. But detailed mode structure and frequency are affected by the input NUBEAM fast ion distribution.**
- **Future work**
 - **Improve the accuracy of Jacobian calculation near the confined/lost boundary, trapped and passing boundary**
 - **Convergence study and plasma parameter scan for $n=3$ TAE case**
 - **Compare mode structure with reflectometer measurements and NOVA-K results**