Characteristics of Neoclassical Toroidal Viscosity in NSTX and KSTAR for Rotation Control and the Evaluation of Plasma Response

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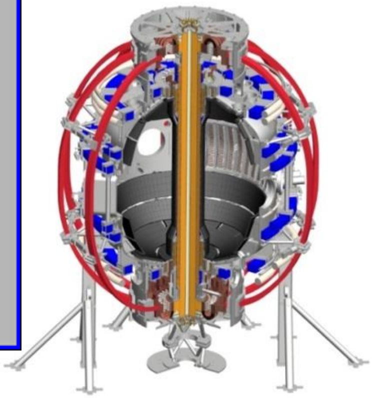

56th Annual APS DPP Meeting

October 28th, 2014

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V1.3p

The physical characteristics of NTV investigated in tokamaks for rotation control and the evaluation of plasma response

Motivation

- Low magnitude ($\delta B/B_0 \sim O(10^{-3})$) 3D magnetic fields are used favorably used in tokamaks (e.g. ELM suppression, MHD mode control)
- 3D fields of this magnitude can produce neoclassical toroidal viscosity (NTV), which can:
  - Alter plasma rotation
  - Significantly reduce fusion gain, Q, by increased alpha particle transport ($\delta B/B_0 \sim O(10^{-4})$)
- Therefore, it is important to understand NTV in tokamaks, backed by accurate ($\sim O(1)$) quantitative modeling

K.C. Shaing, et al., Nucl. Fusion **54** (2014) 033012

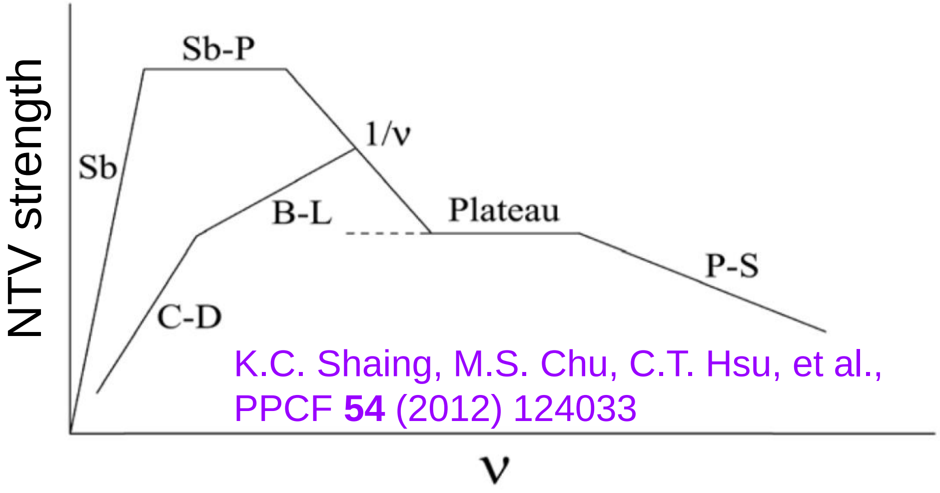
K.C. Shaing, et al., IAEA FEC 2014 Paper TH/P1-11

Outline

- NTV physical characteristics; comparison of theory to experiment
- NTV assessment of plasma response
- Application of NTV to plasma rotation control for NSTX-U

Neoclassical Toroidal Viscosity (NTV) studied through the application of 3D fields in NSTX and KSTAR

- Theory: NTV strength varies with plasma collisionality ν, δB^2, rotation



K.C. Shaing, M.S. Chu, C.T. Hsu, et al., PPCF **54** (2012) 124033

NTV force in “1/ν” collisionality regime

$$\left\langle \hat{e}_t \bullet \vec{\nabla} \bullet \overleftrightarrow{\Pi} \right\rangle_{(1/\nu)} = B_t R \left\langle \frac{1}{B_t} \right\rangle \left\langle \frac{1}{R^2} \right\rangle \frac{\lambda_{1i} p_i}{\pi^{3/2} \nu_i} \varepsilon^{3/2} (\omega_\phi - \omega_{NC}) I_\lambda$$

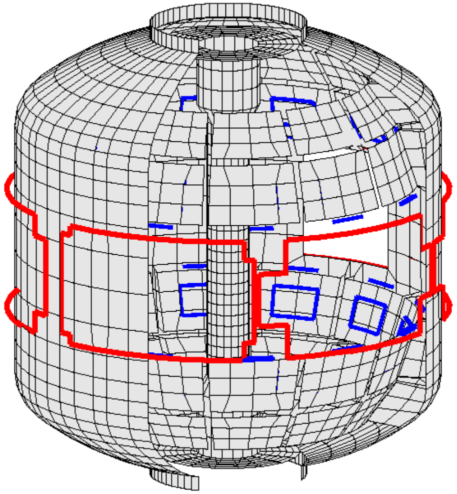
δB^2

$T_i^{5/2}$

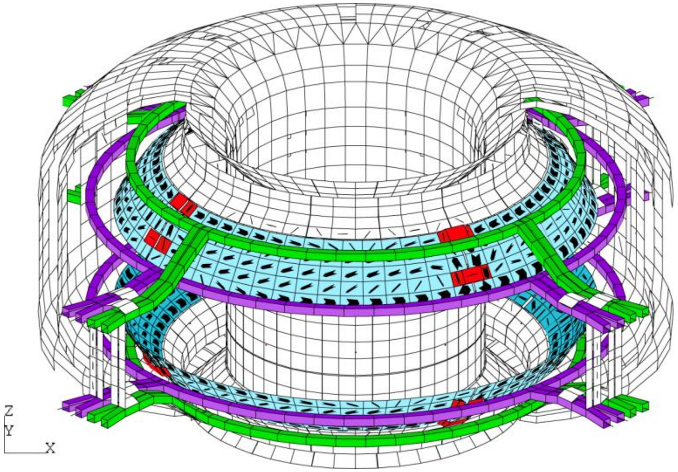
plasma rotation

K.C. Shaing, et al., PPCF **51** (2009) 035004

NSTX 3D coils



KSTAR 3D coils

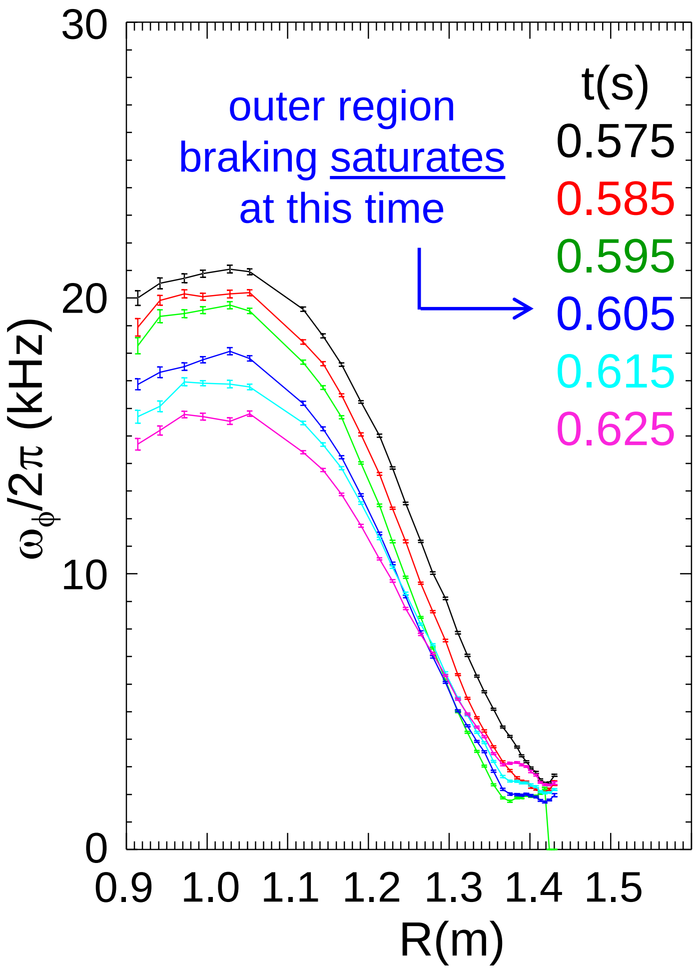


NTV physical characteristics are generally favorable for rotation control

- Non-resonant NTV characteristics (e.g. in NSTX and KSTAR)
  - Experimentally, NTV torque, T_{NTV}, is radially extended, with a relatively smooth profile
  - NTV changes continuously as the applied 3D field is increased
  - Can alter the ω_ϕ profile without mode locking
  - T_{NTV} is not simply an integrated torque applied at the plasma boundary, but a radial profile – e.g. ω_ϕ shear can be changed
    - potential for mode control
- Questions remain
  - e.g. Is there hysteresis when ω_ϕ is altered by NTV?

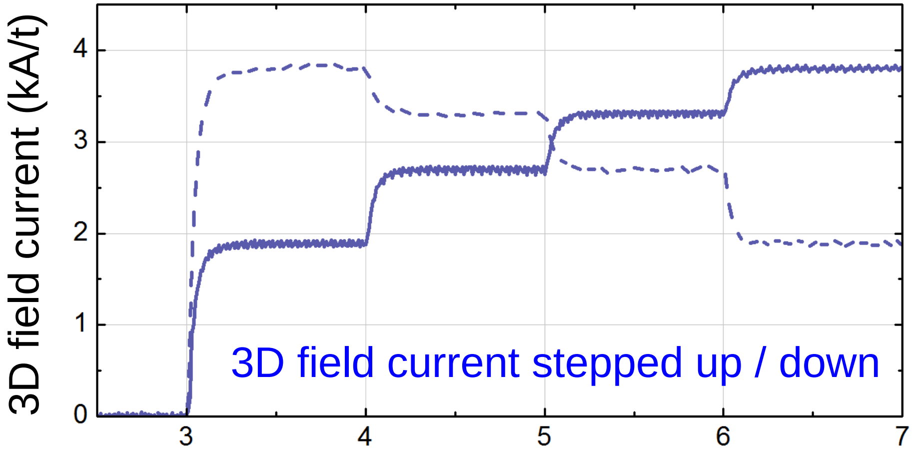
Suggested in: Y. Liang, et al., NF **50** (2010) 025013

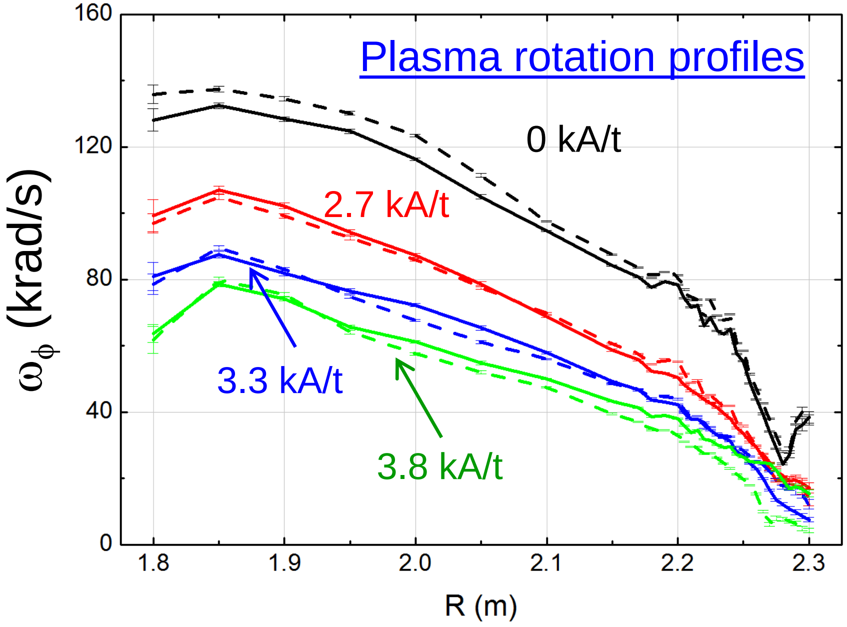
ω_ϕ alteration by n = 2 applied field configuration in NSTX



KSTAR experiments show essentially no hysteresis in steady-state ω_ϕ profile vs. applied 3D field strength

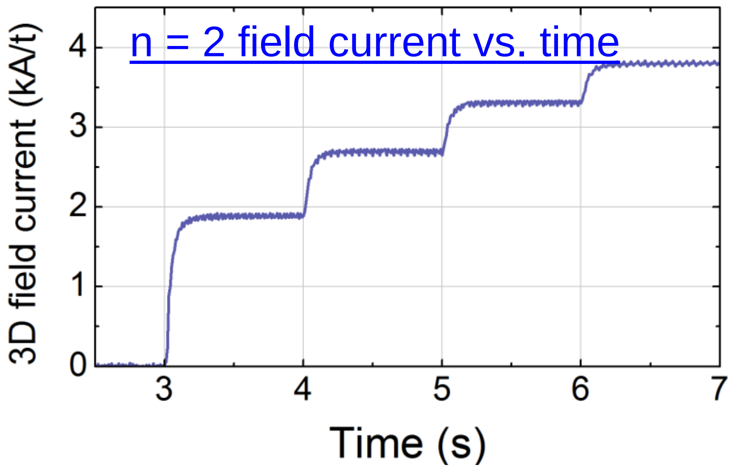
KSTAR non-resonant (“n = 2”) NTV experiments

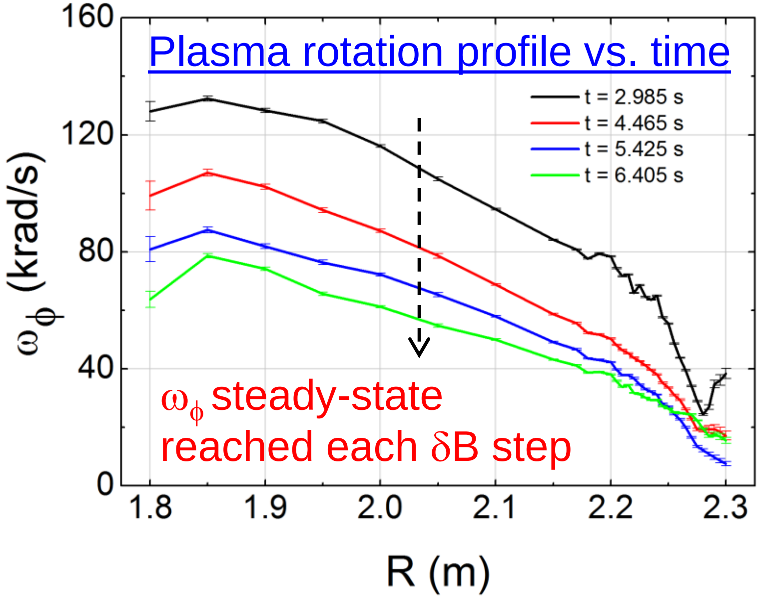




- Steady-state rotation profile reached at each 3D field level is generally independent of the starting point of ω_ϕ
- Absence of hysteresis further confirmed in very recent experiments with 6 steps in 3D field current

KSTAR experiments show Neoclassical Toroidal Viscosity varies as δB^2, and $T_i^{2.27}$, expected by theory





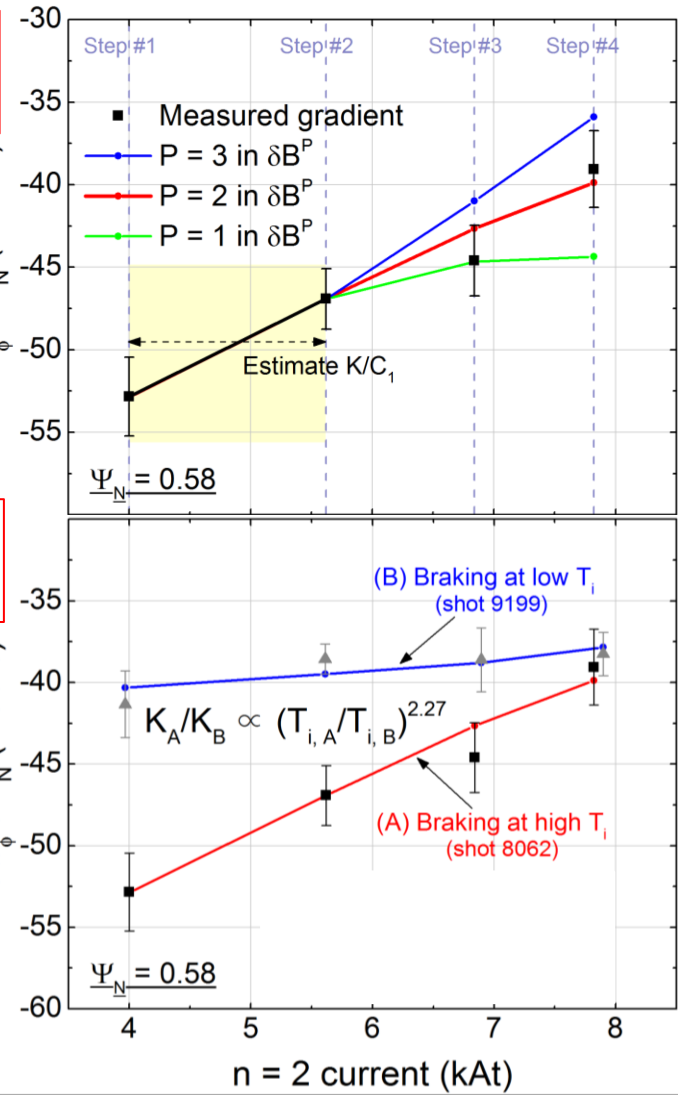
ω_ϕ steady-state reached each δB step

- ❑ T_i scaling confirms NSTX results

S.A. Sabbagh, et al., Nucl. Fusion 50 (2010) 025020

Best fit:
$\propto \delta B^2$

Best fit:
$\propto T_i^{2.27}$



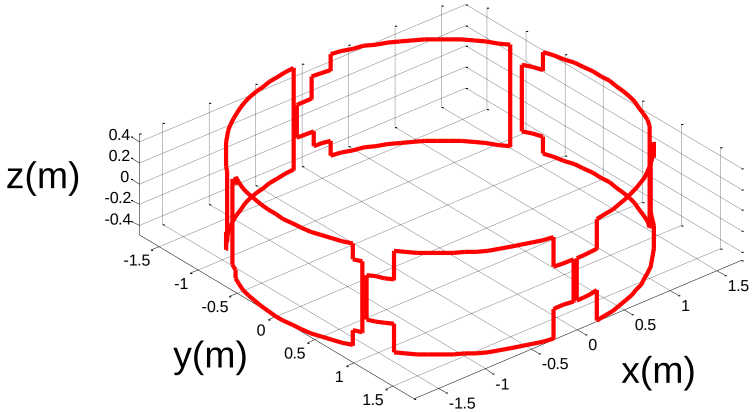
Y.S Park, et al., CP8.03 (Mon. PM)

3D field perturbation experiments measure the T_{NTV} profile in NSTX

- ❑ Apply/change 3D field on a timescale significantly faster than momentum diffusion time, τ_m
  - ❑ Analysis before/after 3D field application isolates T_{NTV} in the momentum diffusion equation; $-dL/dt = T_{NTV}$ (other parameters ~ constant)
- ❑ *dL/dt* measured experimentally and compared to theoretically computed T_{NTV} on this timescale
  - ❑ Important, as *dL/dt* profile changes significantly on timescales > τ_m, (diffuses radially, broadens, leads to significant error compared to T_{NTV})
- ❑ Focus on non-resonant applied 3D field configurations to avoid driving strong MHD modes

Theoretical NTV torque density profiles, T_{NTV} are computed for NSTX using theory applicable to all collisionality regimes

Non-axisymmetric coils fully modelled in 3D



- NTV analysis of NSTX – data interfaced to NTVTOK

  (Y. Sun, Liang, Shaing, et al., NF **51** (2011) 053015)

  - Use Shaing's "connected NTV model", covers all ν, superbanana plateau regimes

    (K.C. Shaing, Sabbagh, Chu, NF **50** (2010) 025022)

  - Full 3D coil specification and δB spectrum, ion and electron components computed

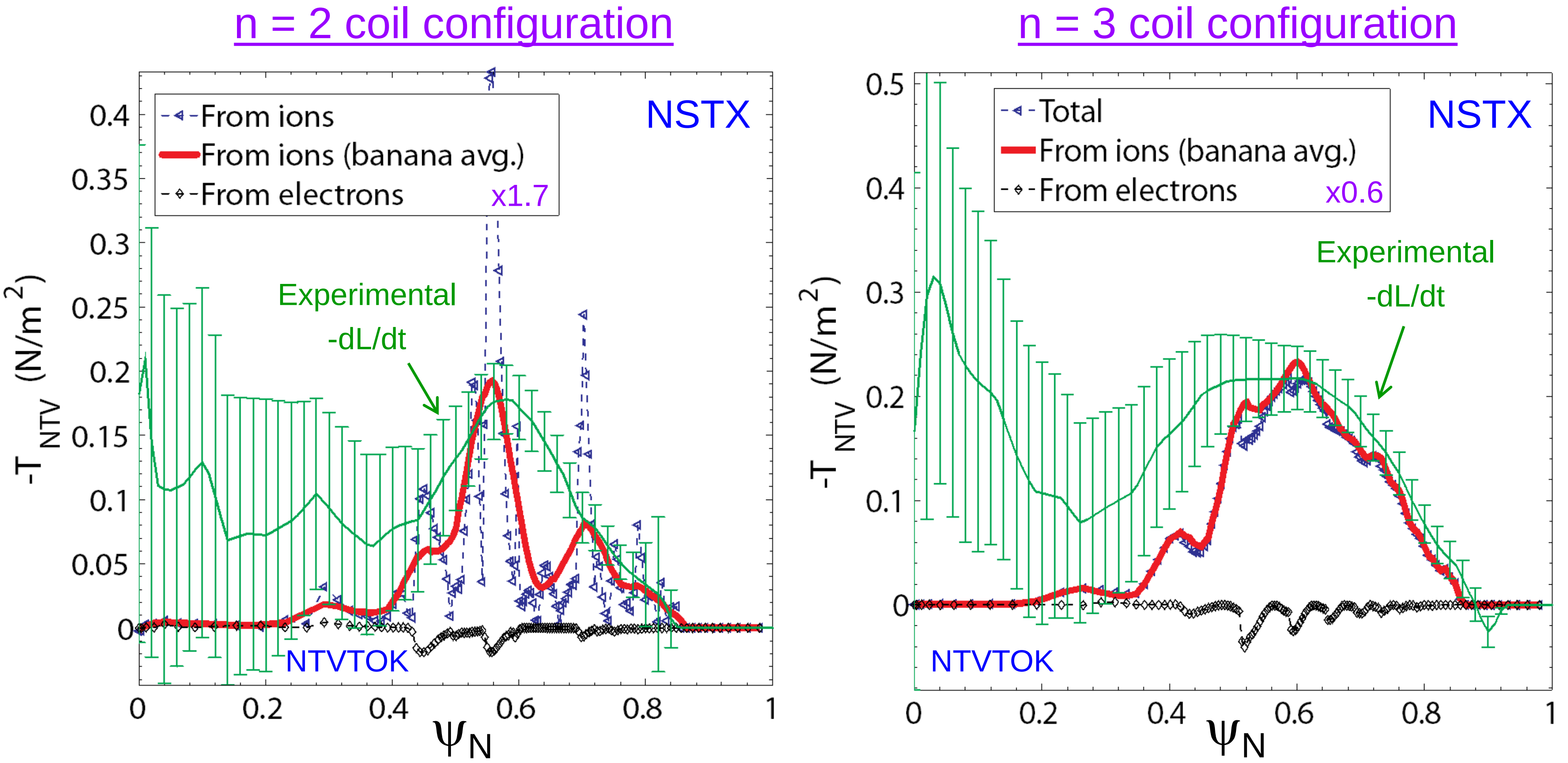
3D field definition

$$\delta B = \vec{b} \bullet \left(\vec{B} / B\right) + \left(\vec{\xi} \bullet \nabla B\right)$$

plasma displacement

- General considerations
  - In tokamaks, ξ not measured in detail, can lead to large error
  - "Fully-penetrated field constraint" used to define ξ $\left(\vec{B}_{2D} \bullet \nabla \vec{\xi} = \vec{b}\right)$
  - Computed $|\xi|$ ~ 0.3 cm << $\varepsilon^{0.5}\rho_i$, therefore, ion banana width-averaging is used for ion channel
    - Can explain why strong resonant peaks in NTV profile are not observed in experiment

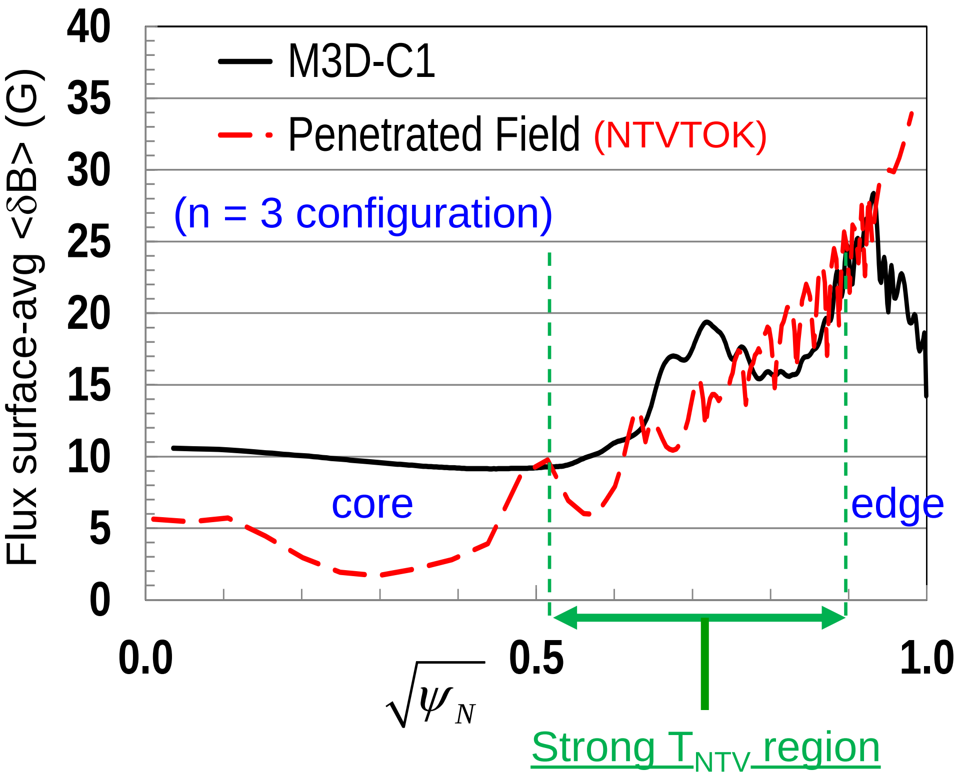
Measured NTV torque density profiles quantitatively compare well to computed T_{NTV} using fully-penetrated 3D field



- T_{NTV} (theory) scaled to match *peak* value of measured *-dL/dt*
  - Scale factor ($(dL/dt)/T_{NTV}$) = 1.7 and 0.6 for cases shown above – O(1) agreement
  - O(1) agreement using “fully-penetrated 3D field” indicates that plasma response is *not* strongly amplified from this “vacuum field assumption” ($T_{NTV} \sim \delta B^2$)

Plasma response from fully-penetrated 3D field used in NTV experimental analysis matches M3D-C^1 single fluid model

Surface-averaged δB from fully penetrated model vs. M3D-C^1 single fluid model



- NTV experimental data is a strong quantitative constraint on plasma response of δB
  - $T_{NTV} \propto \delta B^2$,
- Good agreement between NTVTOK / M3D-C^1 single fluid models in strong NTV region
  - M3D-C^1 core <δB> larger than NTVTOK
    - Core mode in M3D-C^1
  - M3D-C^1 edge <δB> smaller
    - Experimental T_{NTV} too small in this region to constraint δB

Non-resonant NTV and NBI used as actuators in state-space rotation feedback controller designed for NSTX-U

- Momentum force balance – ω_ϕ decomposed into Bessel function states

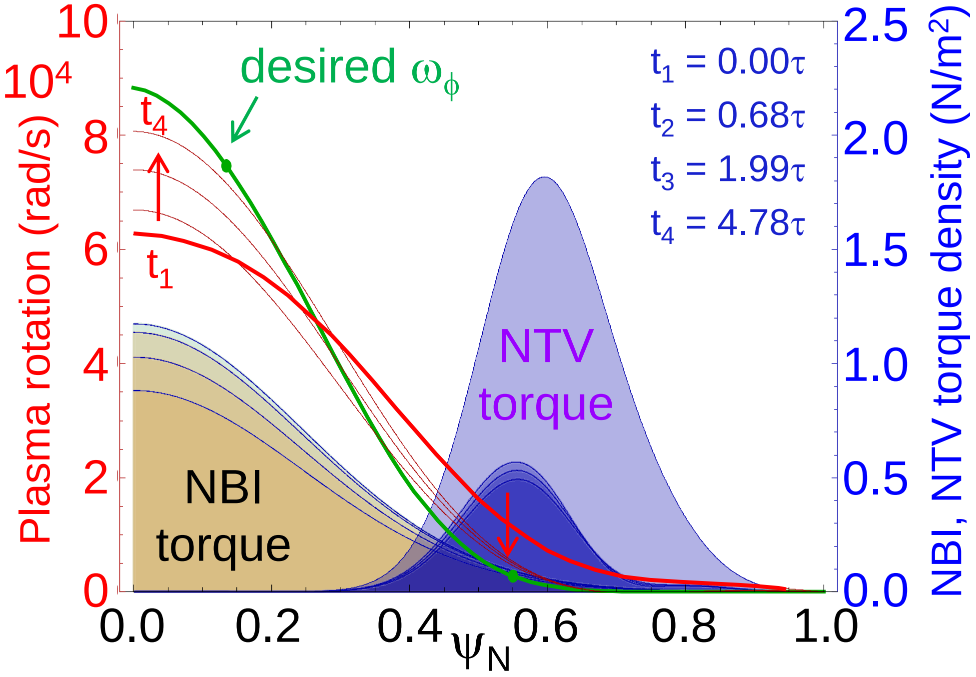
$$\sum_i n_i m_i \left\langle R^2 \right\rangle \frac{\partial \omega}{\partial t} = \left(\frac{\partial V}{\partial \rho} \right)^{-1} \frac{\partial}{\partial \rho} \left[\frac{\partial V}{\partial \rho} \sum_i n_i m_i \chi_\phi \left\langle \left(R \nabla \rho \right)^2 \right\rangle \frac{\partial \omega}{\partial \rho} \right] + T_{NBI} + T_{NTV}$$

- NTV torque:

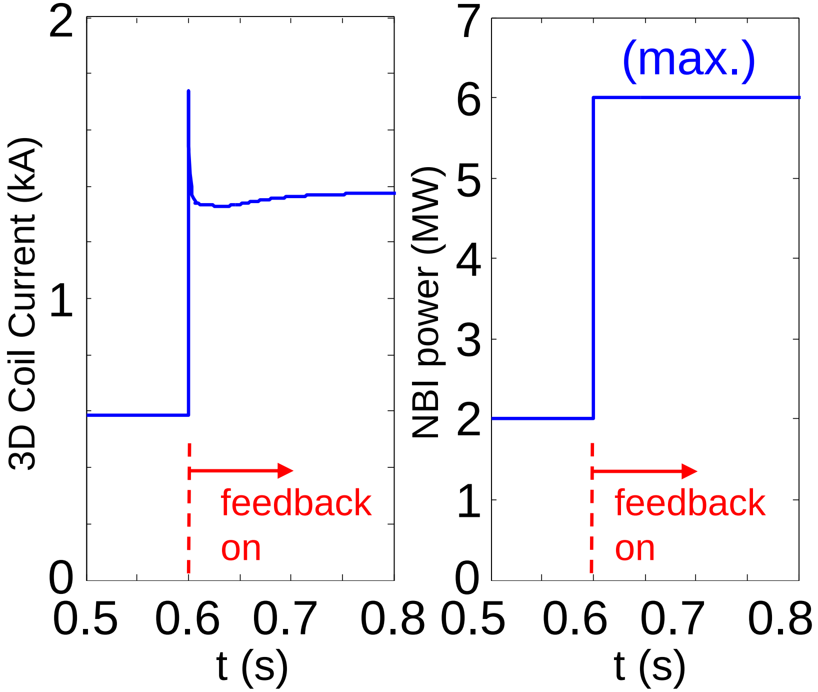
$$T_{NTV} \propto K \times f\left(n_{e,i}{}^{K1} T_{e,i}{}^{K2} \right) g\left(\delta B(\rho) \right) \left[I_{coil}{}^{2} \omega \right]$$

(non-linear)

Rotation evolution and NBI and NTV torque profiles



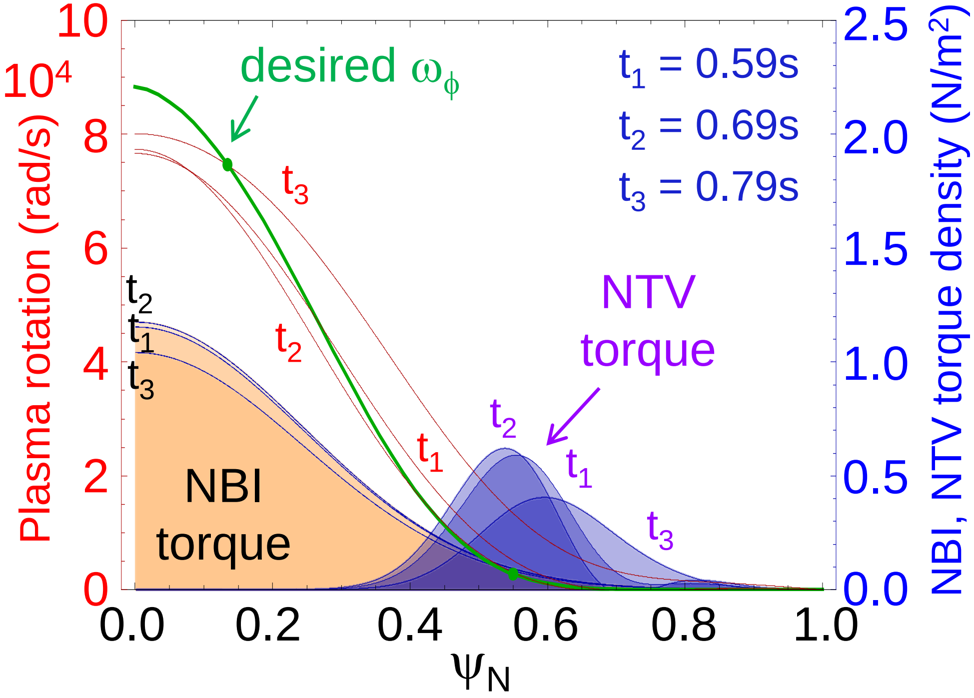
3D coil current and NBI power (actuators)



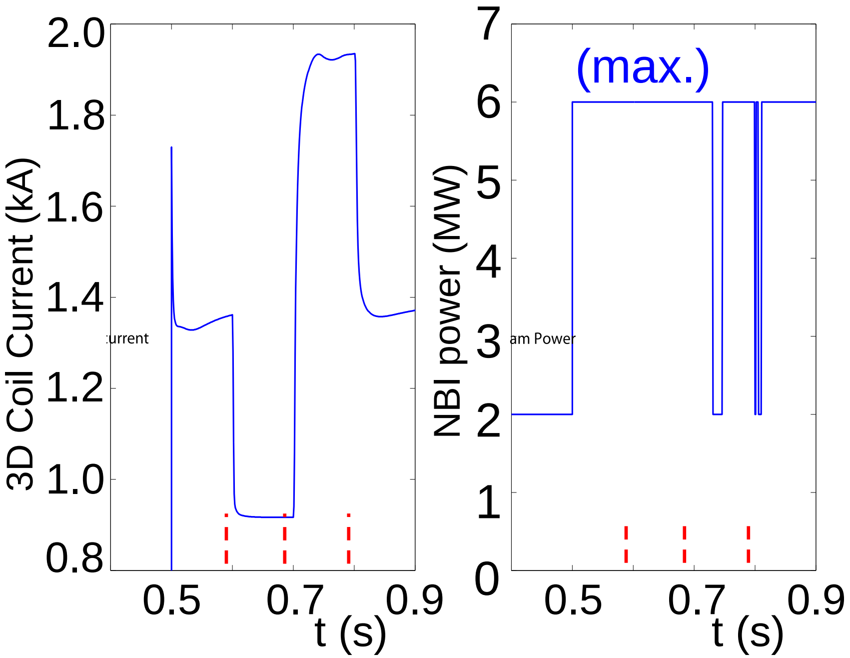
I. Goumiri, et al., PP8.053 (Wed. PM)

When T_i is included in NTV rotation controller model, 3D field current and NBI power can compensate for T_i variations

Rotation evolution and NBI and NTV torque profiles

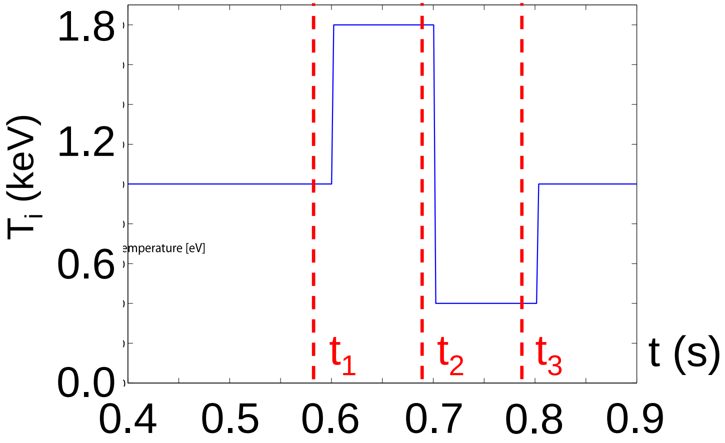


3D coil current and NBI power (actuators)



$$T_{NTV} \propto K \times f\left(n_{e,i}^{\;K1} T_i^{\,K2}\right) g\left(\delta B(\rho)\right)\left[I_{coil}^{\;2}\,\omega\right]$$

K1 = 0, K2 = 2.5



- NTV torque profile model for feedback dependent on ion temperature

Physical characteristics of NTV are investigated in tokamaks for rotation control and the evaluation of plasma response

❑ NTV characteristics / comparison to theory

- ❑ Non-resonant NTV torque is radially extended, relatively smooth profile
- ❑ KSTAR shows $T_{NTV} \propto (\delta B_{3D})^2$; $T_{NTV} \propto T_i^{2.27}$; no hysteresis on the rotation profile (key for control), confirms NSTX
- ❑ Measured T_{NTV} profile in NSTX quantitatively compares well between experiment and Shaing's "connected NTV theory"

K.C. Shaing, et al., NF **50** (2010) 025022)

❑ Plasma response

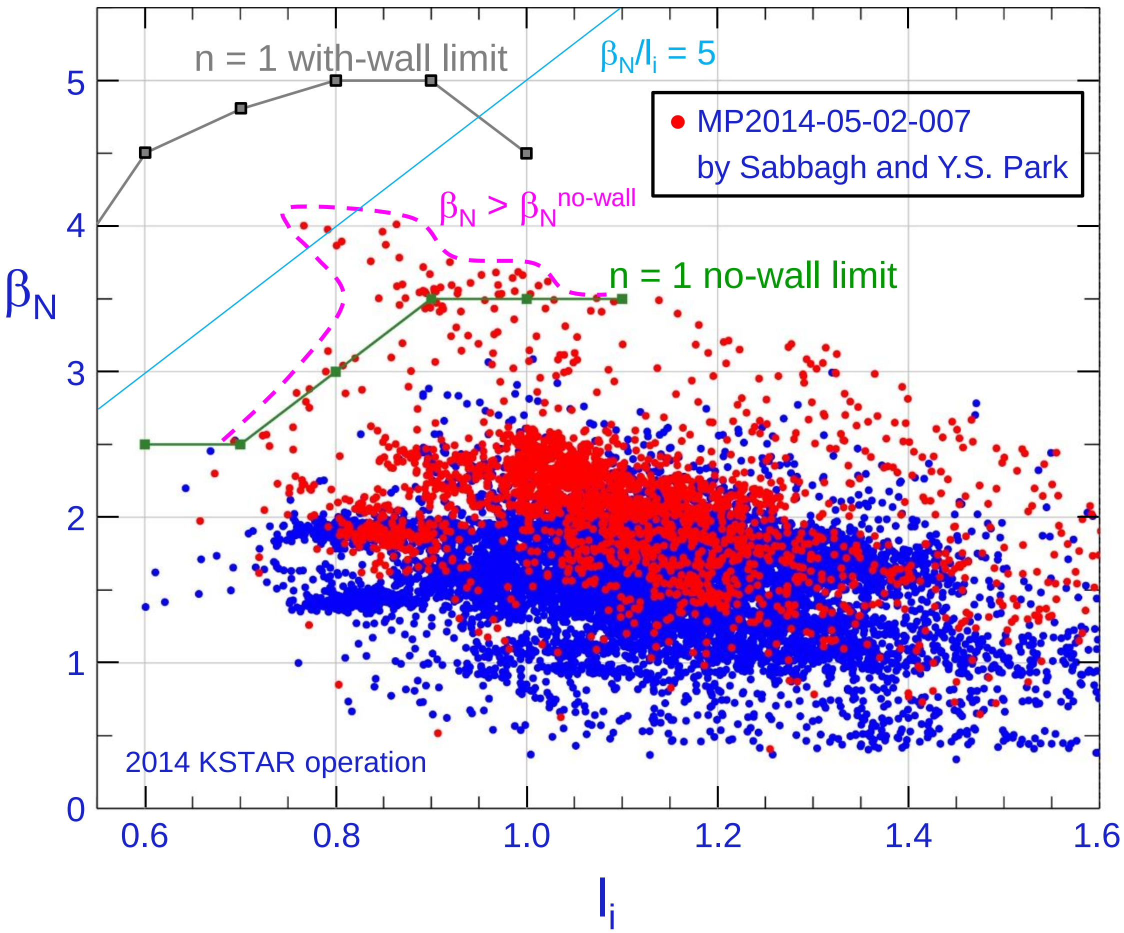
- ❑ Non-resonant T_{NTV} profile in NSTX quantitatively consistent with "fully-penetrated field" assumption without amplification
- ❑ Flux surface-averaged 3D field profile from M3D-C^1 single fluid model consistent with field used for quantitative NTV agreement in experiment

❑ Rotation control

- ❑ Model-based, rotation controller using NTV and NBI designed/tested for NSTX-U

Backup slides

Very recently, high beta plasmas transiently reached β_N = 4 in 2014 campaign


KSTAR operating space containing ~11,500 equilibria

Y.S Park, et al., IAEA FEC 2014 paper EX/P8-05 (Fri. PM)

S.W. Yoon, et al., IAEA FEC 2014 paper OV/3-4 (Tues. AM)

- Values obtained using fully converged KSTAR EFIT reconstructions
- High values reached transiently at lowered B_t
  - B_T in range 0.9 - 1.2 T
  - β_N up to 4 with l_i ~ 0.8 for duration longer than τ_E ~60 ms in these discharges
  - β_N/l_i = 5 is ~ 40% over the computed n = 1 ideal MHD no-wall limit
- Adding newly available 3rd neutral beam source may further increase the operating performance in the ongoing device campaign

Non-resonant Neoclassical Toroidal Viscosity (NTV) physics will be used for the first time in rotation feedback control

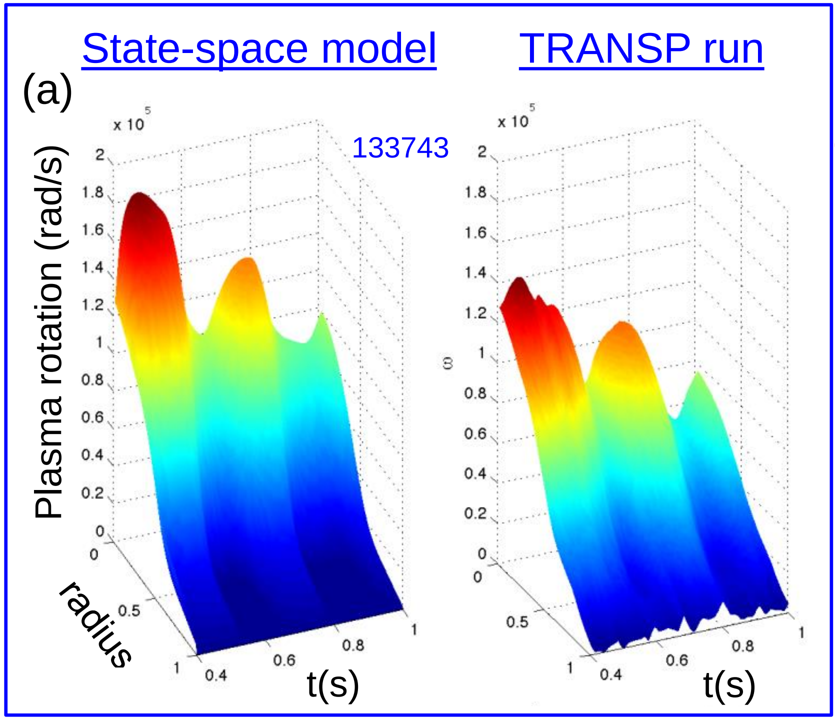
- Momentum force balance – ω_ϕ decomposed into Bessel function states

$$\sum_i n_i m_i \left\langle R^2 \right\rangle \frac{\partial \omega}{\partial t} = \left(\frac{\partial V}{\partial \rho}\right)^{-1} \frac{\partial}{\partial \rho}\left[\frac{\partial V}{\partial \rho} \sum_i n_i m_i \chi_\phi \left\langle \left(R \nabla \rho\right)^2 \right\rangle \frac{\partial \omega}{\partial \rho}\right] + T_{NBI} + T_{NTV}$$

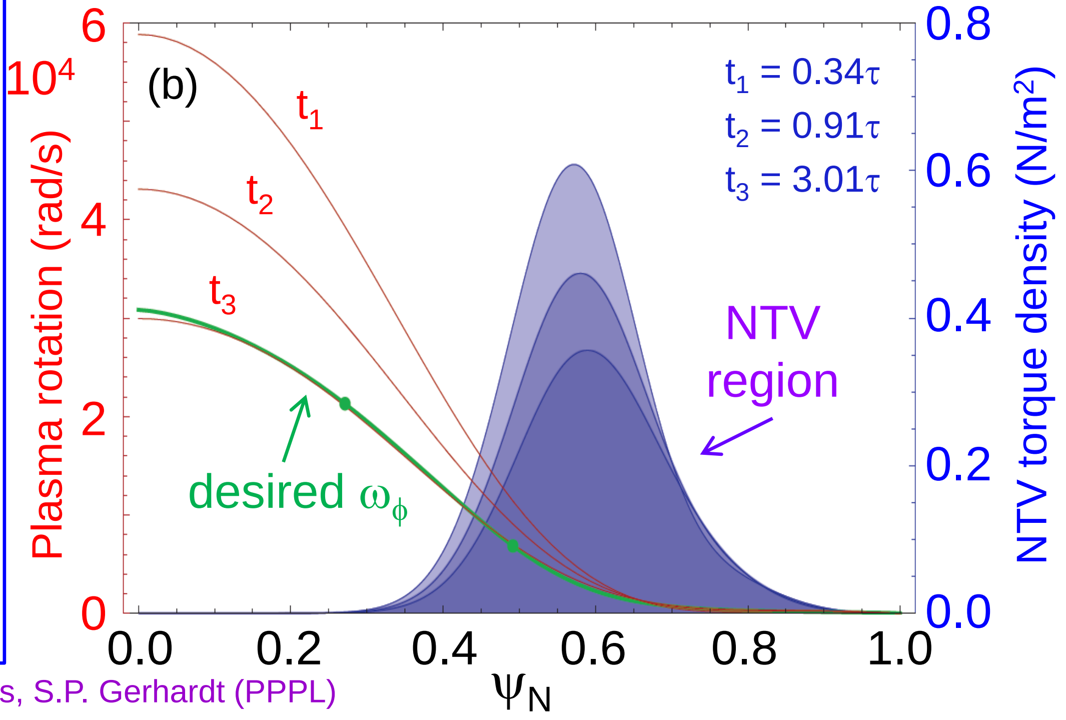
- NTV torque:

$$T_{NTV} \propto K \times f\left(n_{e,i}{}^{K1} T_{e,i}{}^{K2}\right) g\left(\delta B(\rho)\right) \left[I_{coil}{}^{2} \omega\right]$$

(non-linear)

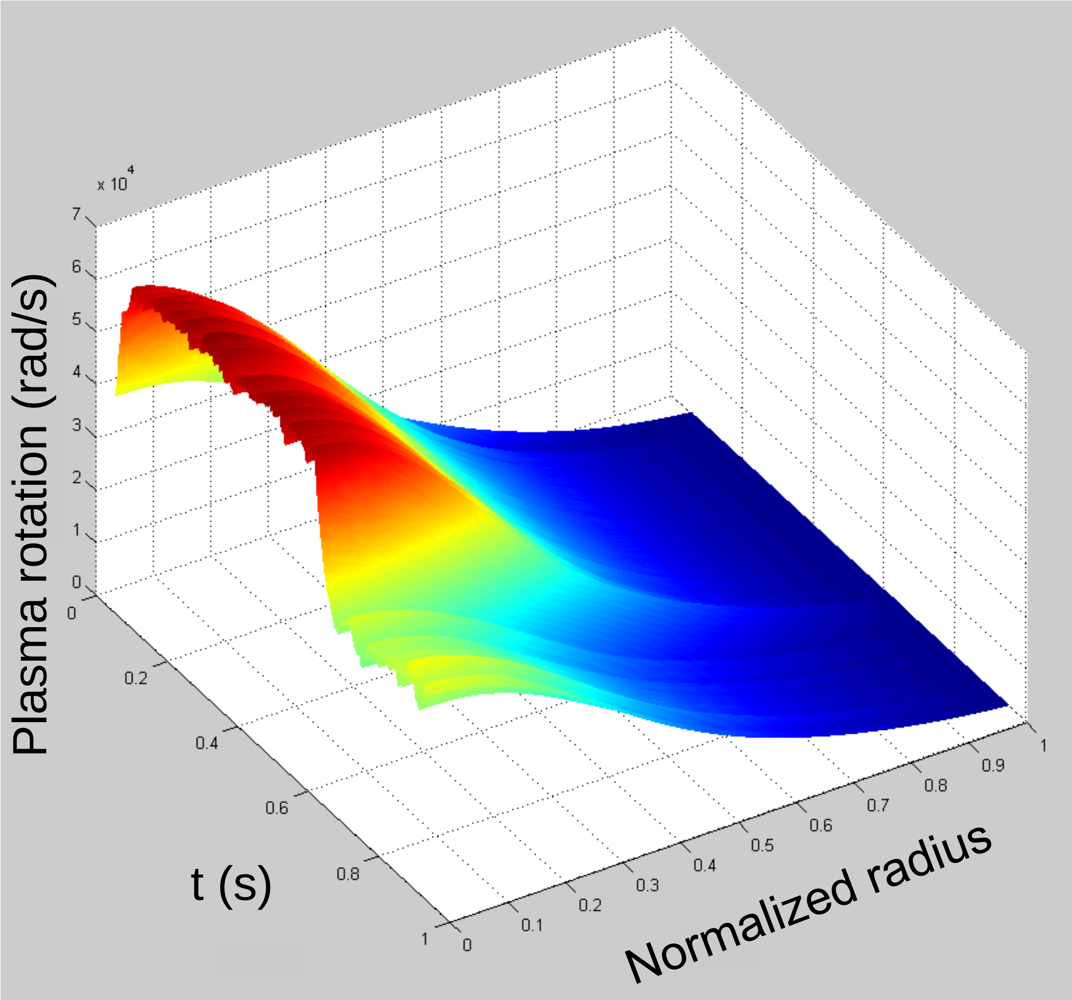

State-space model TRANSP run

Feedback using NTV: "n=3" $\delta B(\rho)$ spectrum



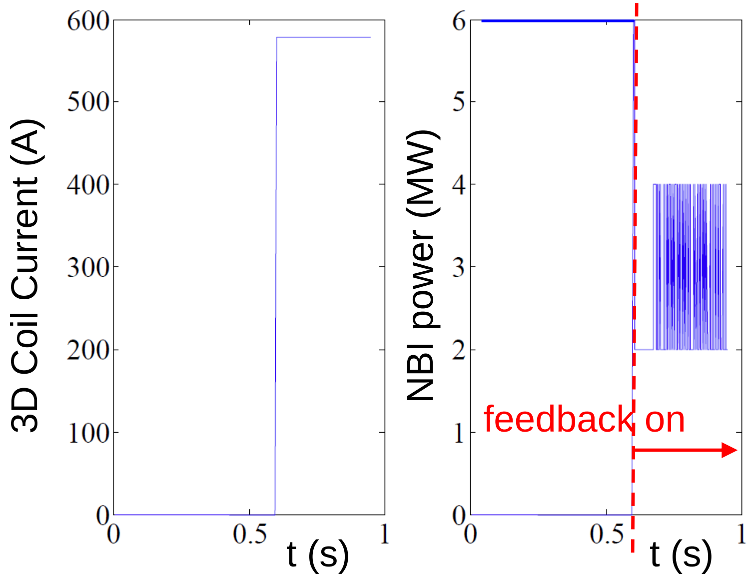
I. Goumiri (PU), S.A. Sabbagh (Columbia U.), D.A. Gates, S.P. Gerhardt (PPPL)

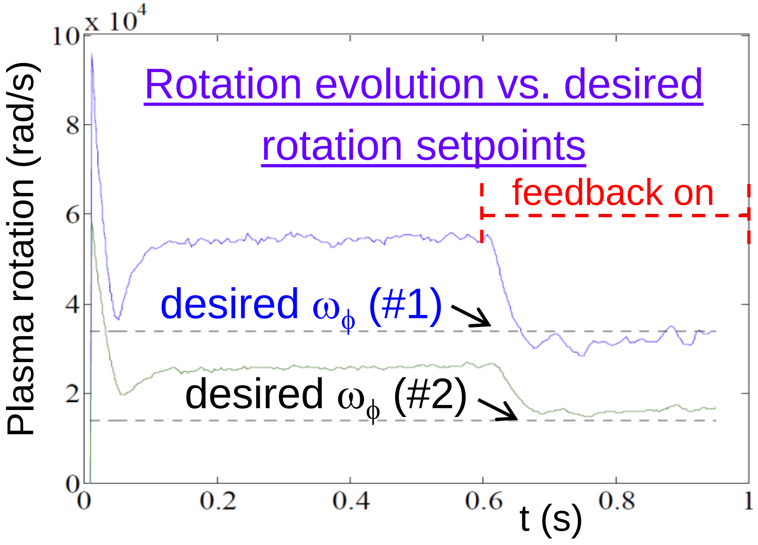
Plasma rotation control has been demonstrated for the first time with TRANSP using NBI and NTV actuators



This case uses pre-programmed 3D coil current and NBI feedback

3D coil current and NBI power (actuators)





Several ordered publications by K.C. Shaing, et al. led to the “Combined” NTV Formulation

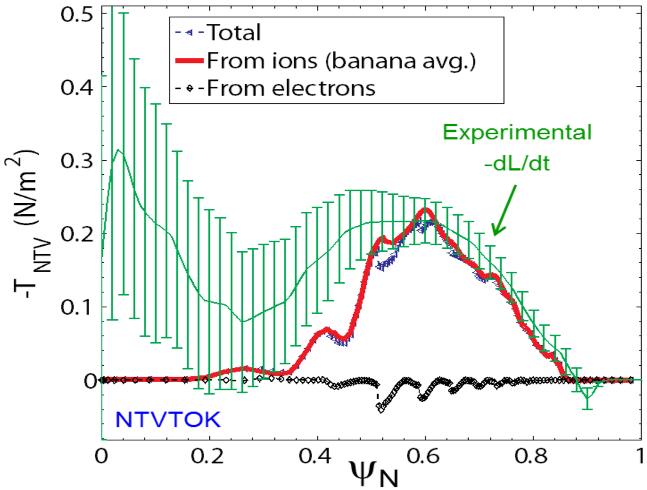
❑ Publications (chronological order)	❑ Topic
1) K.C. Shaing, S.P. Hirschman, and J.D. Callen, Phys. Fluids **29** (1986) 521.	➢ Plateau transport
2) K.C. Shaing, Phys. Rev. Lett., **87** (2001) 245003.	➢ Island NTV
3) K.C. Shaing, Phys. Plasmas **10** (2003) 1443.	➢ Collisional, $1/\nu$ regimes
4) K.C. Shaing, Phys. Plasmas **13** (2006) 052505.	➢ Banana, $1/\nu$ regimes
5) K.C. Shaing, S. A. Sabbagh, and M. Peng, Phys. Plasmas **14** (2007) 024501.	➢ Multiple trapping
6) K.C. Shaing, S. A. Sabbagh, M.S. Chu, et al., Phys. Plasmas **15** (2008) 082505.	➢ Orbit squeezing
7) K.C. Shaing, P. Cahyna, M. Becoulet, et al., Phys. Plasmas **15** (2008) 082506.	➢ Coll. b’dary layer, $\nu^{0.5}$
8) K.C. Shaing, S. A. Sabbagh, and M. S. Chu, PPCF **51** (2009) 035004.	➢ Low ν regimes
9) K.C. Shaing, S. A. Sabbagh, and M. S. Chu, PPCF **51** (2009) 035009.	➢ Superbanana plateau
10) K.C. Shaing, S. A. Sabbagh, and M. S. Chu, PPCF **51** (2009) 055003.	➢ Superbanana regime
11) K.C. Shaing, M. S. Chu, and S. A. Sabbagh, PPCF **51** (2009) 075015.	➢ Bounce/transit/drift res.
12) K.C. Shaing, M. S. Chu, and S. A. Sabbagh, PPCF **52** (2010) 025005.	➢ $J_{\text{bootstrap}}$ w/resonances
13) K.C. Shaing, S. A. Sabbagh, and M. S. Chu, Nucl. Fusion **50** (2010) 025022.	➢ Combined NTV formula
14) K.C. Shaing, J. Seol, Y.W. Sun, et al., Nucl. Fusion **50** (2010) 125008.	➢ ∇B drift in CBL analysis
15) K.C. Shaing, M. S. Chu, and S. A. Sabbagh, Nucl. Fusion **50** (2010) 125012.	➢ Flux/force gen. coords.
16) K.C. Shaing, T.H. Tsai, M.S. Chu, et al., Nucl. Fusion **51** (2011) 073043.	➢ SBP regime refinement
17) K.C. Shaing, M.S. Chu, C.T. Hsu, et al., PPCF **54** (2012) 124033.	➢ NTV brief overview

GO3.04: Physical Characteristics of Neoclassical Toroidal Viscosity in Tokamaks for Rotation Control and the Evaluation of Plasma Response

Highlights

- ❑ Experimental NTV characteristics
  - ❑ NTV experiments on NSTX and KSTAR
  - ❑ NTV torque T_{NTV} from applied 3D field is a radially extended, relatively smooth profile
  - ❑ Perturbation experiments measure T_{NTV} profile
- ❑ Aspects of NTV for rotation control
  - ❑ Varies as $\delta \boldsymbol{B}^2$; $T_{NTV} \propto T_i^{5/2}$ in primary collisionality regime for large tokamaks
  - ❑ No hysteresis on the rotation profile when altered by non-resonant NTV is key for control
  - ❑ Rotation controller using NTV and NBI tested for NSTX-U; model-based design saves power
- ❑ NTV analysis to assess plasma response
  - ❑ Non-resonant NTV quantitatively consistent with fully-penetrated field assumption
  - ❑ Surface-averaged 3D field profile from M3D-C^1 single fluid model consistent with field used for quantitative NTV agreement in experiment

Perturbation experiments measure NTV torque profile and compare to theory



Rotation controller using NTV and NBI