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The Potential Contribution of RF Sheaths to Field-Aligned SOL Losses on NSTX

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Significant HHFW losses in SOL can perhaps be understood by analyzing Langmuir probes

Midplane camera image



RF affects Langmuirprobe floating potential



Only probes underneath spiral are affected

Up to 2 MW/m² heat flux measured under spiral

RF rectification is a candidate mechanism to explain the change in Langmuir probe signals

- RF electric field drives enhanced electron current to wall
 - Due to nonlinear IV characteristic for sheath
- Probe floating potential shifts to more negative value

$$\exp\left[\frac{\Delta V_f}{T_e}\right] = I_0 \left(\frac{V_{RF}}{T_e}\right)$$

- RF rectification blamed for ICRF-induced impurity contamination: enhanced ion-bombardment and erosion
- This talk focuses on a far-field sheaths
 - Formed away from antenna due to propagating waves

[A. Boschi and F. Magistrelli, Il Nuovo Cimento 29 (1963) 487.] [J. C. Hosea, PhD thesis, Stanford Univ., 1966]

Fitting probe characteristics with RF rectification yields reasonable fits around V_{fl} and yields V_{RF} = 43.7 V



- Assume same I_{sat} and T_e values for both probes
 - Only change V_{fl}
- Resulting fits are reasonable for both characteristics especially in vicinity of floating potential

Similar results found for RF-only discharges $V_{RF} = 64 \text{ V}$ for $P_{RF} = 1.1 \text{MW}$ and $V_{RF} = 33 \text{ V}$ for $P_{RF} = 0.55 \text{ MW}$



- RF-only shots; no NBI
 - 141836, P_{RF} = 1.1 MW
 - 141830, P_{RF} = 0.55 MW
 - t = 0.294 s
- V_{RF} increases with P_{RF} from 0.55 MW to 1.1 MW
- V_{RF} is not ~ $\sqrt{P_{RF}}$ but SOL plasma conditions are different as well
- Assuming RF rectification gives V_{RF} ~ 30 – 70 V

V_{RF} enhances sheath transmission factor (heat flux to surface) Indicates RF rectification contributes substantially to spiral

- Heat flux to probe biased at V is : $q_{surface} = \gamma(V) * j_{sat} * T_e$
 - $\gamma(V)$ is the sheath transmission factor
- V_{fl-RF} and $V_{fl-noRF}$ floating potentials relative to ground (tile voltage)

$$\begin{split} \gamma_{noRF} = -\frac{V_f}{T_e} + \frac{V_{fl-noRF}}{T_e} + 2.5 \frac{T_i}{T_e} + \frac{2}{1-\sigma_e} \exp\left[-\frac{V_{fl-noRF}}{T_e}\right] \\ \text{[P.C. Stangeby Plasma Physics Series (2000)]} \end{split}$$

- Add RF potential $V_{RF} \sin(\omega t)$ and average over RF period

$$\gamma_{RF} = -\frac{V_f}{T_e} + \frac{V_{fl-noRF}}{T_e} + 2.5 \frac{T_i}{T_e} + \frac{2}{1-\sigma_e} \exp\left[-\frac{V_{fl-RF}}{T_e}\right]$$

• For the RF + NBI case (141899):

 No RF
 γ_{noRF} = 7.12
 q_{noRF} = 0.103 MW/m²

 With RF
 γ_{RF} = 14.43
 q_{RF} = 0.209 MW/m²

• In rough agreement with IR measurements in the vicinity of the vessel gap

Increase of gamma should be even larger if measured at more intense portion of spiral



Langmuir probe location

Cylindrical cold-plasma model



- Model retains finite electron mass
- Still being validated
 - Poynting flux computation agrees with antenna loading calculation to within 1%

Computed V_{RF} is of same order as probe analysis



Conclusion: RF rectification may be contributing substantially to fast-wave losses in the SOL

- If shift in probe characteristic is RF rectification, then $V_{\rm RF}$ ~ 30 70 V
 - Better accuracy expected with fast-acquisition probes using coaxial feeds and broadened sweep range
- Cold-plasma model gives V_{RF} on the same order (~ 20 30 V)
 - Given simplifications involved in model, this is encouraging
 - Will perform similar calculations with AORSA results
- Sheath transmission factors predicted to be enhanced (x2) with RF
 - Better IR coverage in NSTX-U
- Emerging self-consistent picture that RF rectification could be playing a substantial role in fast-wave losses in the SOL

V_{RF} enhances sheath transmission factor (heat flux to surface) Indicates RF rectification contributes substantially to spiral

Heat flux to probe biased at voltage V is : $q_{surface} = \gamma(V) * j_{sat} * T_e$ $\gamma(V) = -\frac{V_f}{T_e} + \frac{\Delta V}{T_e} + 2.5 \frac{T_i}{T_e} + \frac{2}{1 - \sigma_e} \exp\left[-\frac{\Delta V}{T_e}\right]$

 $\Delta V=V-V_f$, with V_f relative to V_{pl} [P.C. Stangeby, Plasma Physics Series (2000)]

- $V_{\text{fl-RF}}$ and $V_{\text{fl-noRF}}$ floating potentials relative to ground (tile voltage) with and without RF
- Add RF potential ΔV = V_{fl-0} + V_{RF} sin(ωt) and average over the RF period

$$\begin{split} \gamma_{noRF} &= -\frac{V_f}{T_e} + \frac{V_{fl-noRF}}{T_e} + 2.5\frac{T_i}{T_e} + \frac{2}{1-\sigma_e}\exp\left[-\frac{V_{fl-noRF}}{T_e}\right]\\ \gamma_{RF} &= -\frac{V_f}{T_e} + \frac{V_{fl-noRF}}{T_e} + 2.5\frac{T_i}{T_e} + \frac{2}{1-\sigma_e}\exp\left[-\frac{V_{fl-RF}}{T_e}\right] \end{split}$$

- Energy flux to probe at $V_{pr} = 0$ (tile V) is enhanced for RF relative to no RF
- For the RF + NBI case (141899): $\gamma_{0-RF=0} = 7.12$ $q_{0-RF=0} = 0.103 \text{ MW/m}^2$ $\gamma_{0-RF} = 14.43$ $q_{0-RF} = 0.209 \text{ MW/m}^2$

Levels of heat flux to the probe compare well with estimates from IR measurements of q_{surface}



 Q is not measured at Bay B – rough estimates of Q_{bay B} from IR measurements at Bay I are made in vicinity of the probe radial location

• Background ~ 0.3 MW/m²; with RF $\Delta Q \sim 0.3$ at peak

- From Bay G measurement (top): background ~ 0.2; with RF Δ Q ~ 0.1
- Thus, heat flux measurements to probe are comparable with these results within the uncertainties associated with the probe and IR measurements