

# Experimental Study of Density Gradient Stabilization Effects on High-k Turbulence in NSTX

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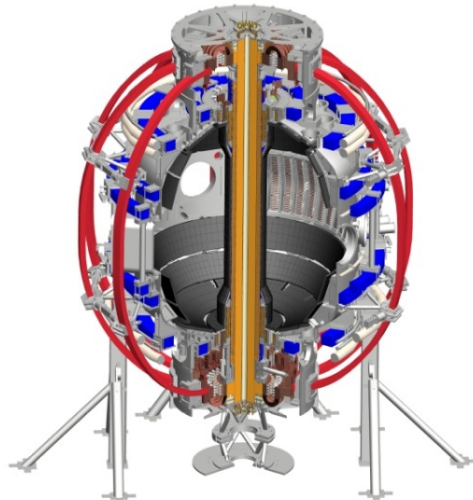
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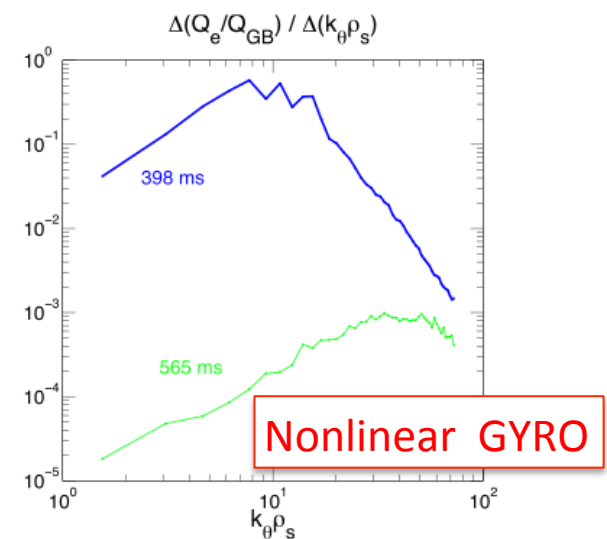
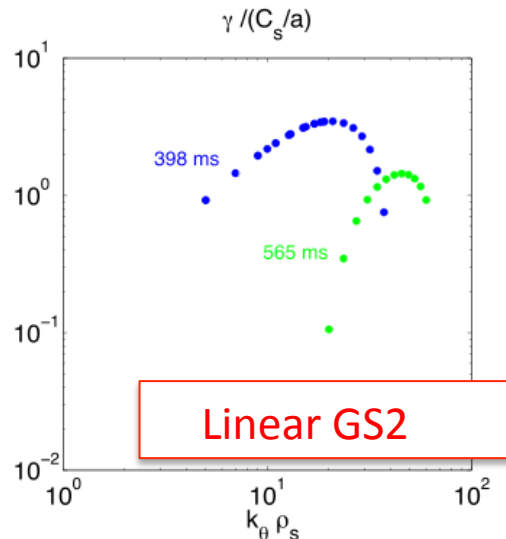
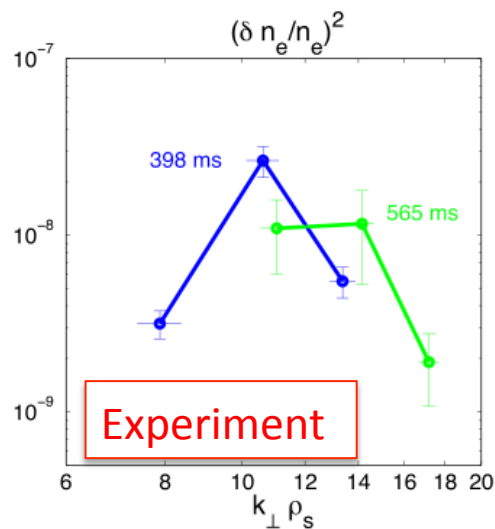


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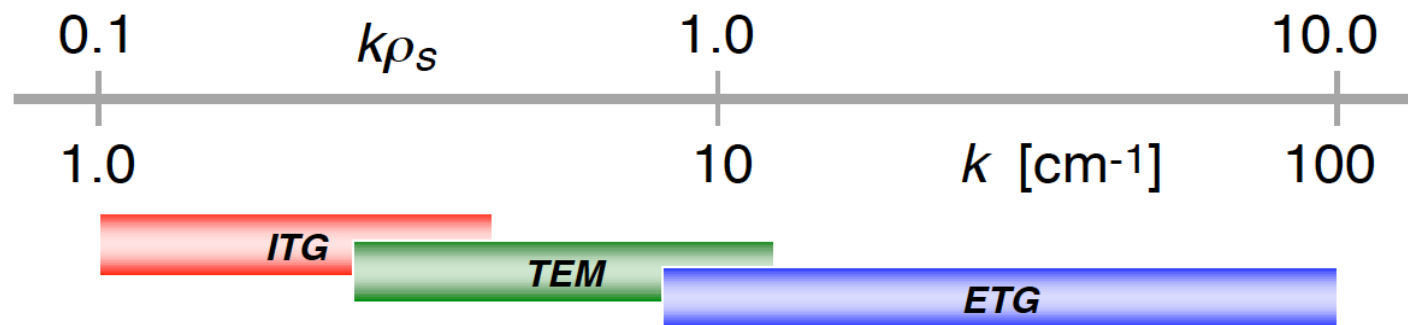
# Linear and Nonlinear Gyrokinetic Simulations are Compared with Measured Trends in High-k Turbulence at NSTX

- $(R/L_{Te}^{\text{exp}}) - (R/L_{Te})_{\text{crit}}$  determines linear threshold for instability and is correlated with the presence of observed fluctuations.
- Increasing  $R/L_{ne}$  shifts high-k fluctuations to higher  $k$  values (stabilizing) and decreases real frequency  $\omega_r$ , consistent with Doppler subtracted plasma frame frequency of detected fluctuations.
- Electron density gradient increases the ETG nonlinear threshold, consistent with experimental observations of reduced fluctuation amplitude, and reduces electron heat flux and stiffness.



# Electron Thermal Transport is Anomalous in All NSTX Confinement Regimes

- NSTX H-mode plasmas exhibit ion thermal transport close to neoclassical levels due to low- $k$  turbulence suppression by ExB shear [*cf. Kaye NF 2007*]. **Electron thermal transport is always found anomalous.**
- **ETG** turbulence is a candidate for anomalous electron thermal transport in some NSTX and NSTX-U operating regimes.
- A **microwave scattering diagnostic** is used at NSTX to measure electron-scale density fluctuations indicative of **high- $k$  turbulence** ( $k_{\perp}\rho_s > 1$ ).
- Linear and nonlinear gyrokinetic simulations are used to study high- $k$  turbulence and electron thermal transport in an NBI-heated H-mode plasma.
- We observe a **stabilizing effect of the electron density gradient** on experimental high- $k$  fluctuation levels and predicted electron heat flux (article accepted to PoP 2015).



From Mazzucato PPPL presentation

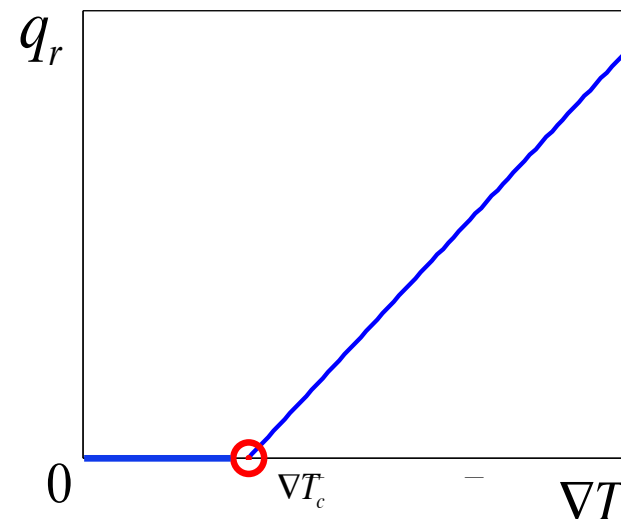
# Critical Gradient and Critical ETG Formula

- Normalized gradient of quantity X

$$R / L_X = R |\nabla X / X|$$

- Critical gradient

$$q_r^{turb} = \underbrace{\chi_{GB} f(\hat{s}, q, \nabla n_e, \dots)}_{\text{Nonlinear dependence}} \underbrace{\left( R / L_{T_e} - (R / L_{T_e})_c \right)}_{\text{Linear threshold}}$$



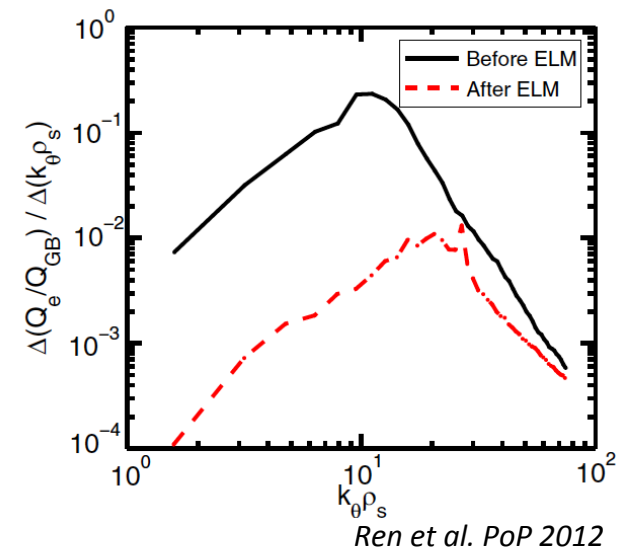
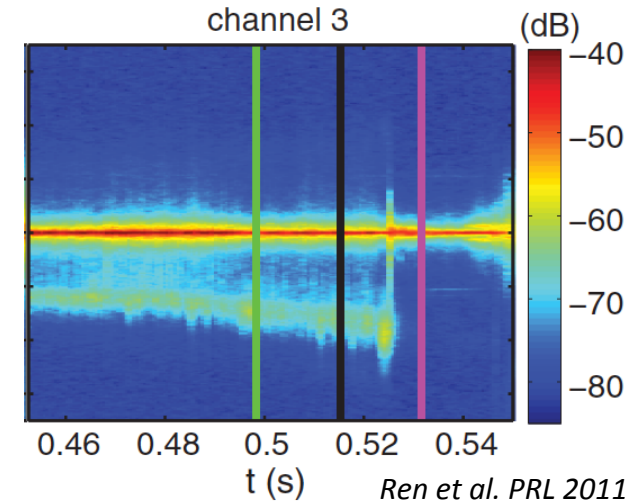
- Jenko critical temperature gradient [*cf. Jenko Phys. Plasmas 2001*].

$$(R / L_{T_e})_{crit} = \max \left\{ \begin{array}{l} 0.8R / L_{ne} \\ (1 + \tau)(1.33 + 1.91\hat{s} / q)(1 - 1.5\varepsilon)(1 + 0.3\varepsilon d\kappa / d\varepsilon) \end{array} \right. \quad \text{with} \quad \tau = Z_{eff} T_e / T_i$$

Applicability: low- $\beta$ , positive  $\hat{s}$  and large aspect ratio with local Miller equilibrium (Miller *et al* PoP 1998).

# Previous Work Suggested Density Gradient Stabilized Turbulence after ELM Event

- First direct experimental demonstration of density gradient stabilization of  $e^-$ -scale turbulence (*Ren et al. PRL 2011*). Shot 140620.
  - **ELM event** at  $t \sim 525$  ms  $\rightarrow$  change in density gradient.
  - Stabilization of lower- $k$   $e^-$ -scale fluctuations ( $k_{\perp} \rho_s < 10$ ).
- Nonlinear gyrokinetic simulations show the effect of density gradient on transport (*Ren et al. PoP 2012*). Shot 140620.
- New work presents a detailed analysis of the linear and nonlinear stabilizing effect of density gradient on ETG fluctuations *focusing on changes in critical gradient and stiffness* during controlled current ramp down experiment (Shot 141767. Article accepted to PoP 2015).

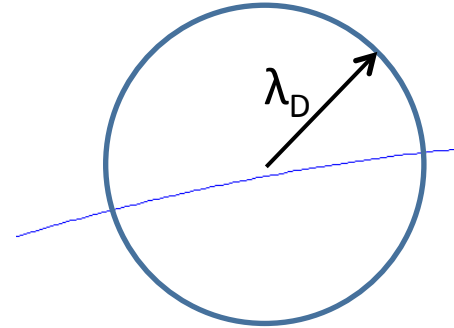


# Experimental Set-Up

# Collective Scattering is Used to Measure High-k Turbulence

- Collective/coherent scattering

$$k\lambda_D \leq 1$$

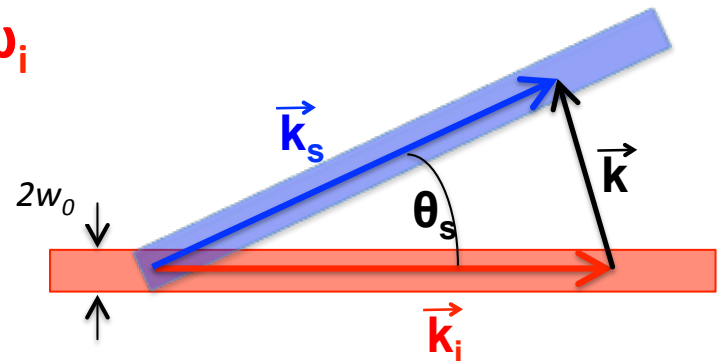


- Scattered power density  $\frac{d^2 P}{d\Omega d\nu} = P_i r_e^2 L_z |\Pi \cdot \hat{e}|^2 \frac{|\tilde{n}_e(k, \omega)|^2}{VT}$
- Three wave-coupling** between incident beam  $(\mathbf{k}_i, \omega_i)$  and plasma  $(\mathbf{k}, \omega)$

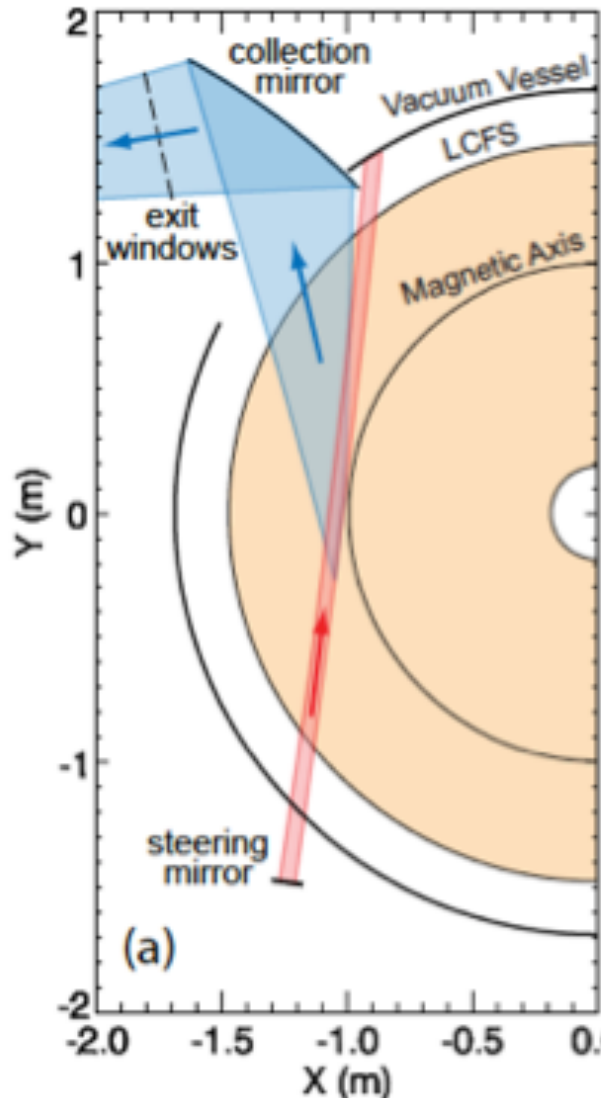
$$\vec{\mathbf{k}}_s = \vec{\mathbf{k}} + \vec{\mathbf{k}}_i \quad \omega_s = \omega + \omega_i$$

- $\omega_i, \omega_s \gg \omega$  imposes Bragg condition

$$k = 2k_i \sin(\theta_s/2)$$



# High-k Microwave Scattering Diagnostic at NSTX

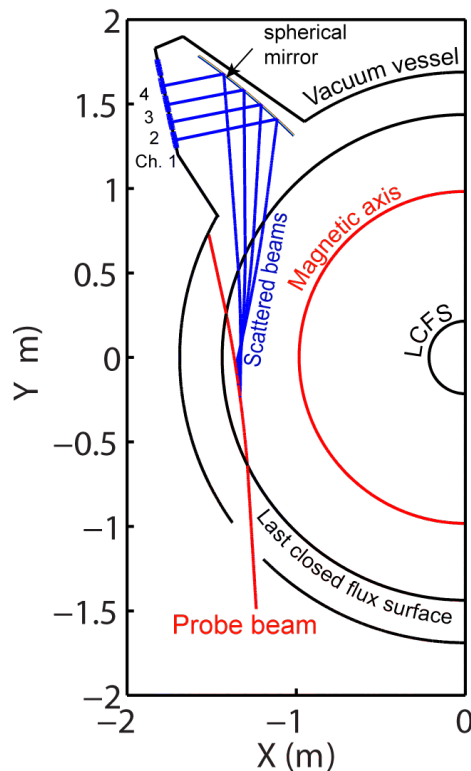


- Gaussian Probe beam: 15 mW, 280 GHz,  $\lambda_i \sim 1.07$  mm,  $a = 3$  cm ( $1/e^2$  radius).
- Propagation close to midplane  $\Rightarrow k_r$  spectrum.
- 5 detection channels  $\Rightarrow$  range  $k_r \sim 5$ -30  $\text{cm}^{-1}$  (*high-k*).
- Wavenumber resolution  $\Delta k = \pm 0.7$   $\text{cm}^{-1}$ .
- Radial coverage:  $R = 106$ -144 cm.
- Radial resolution:  $\Delta R = \pm 2$  cm (unique feature).

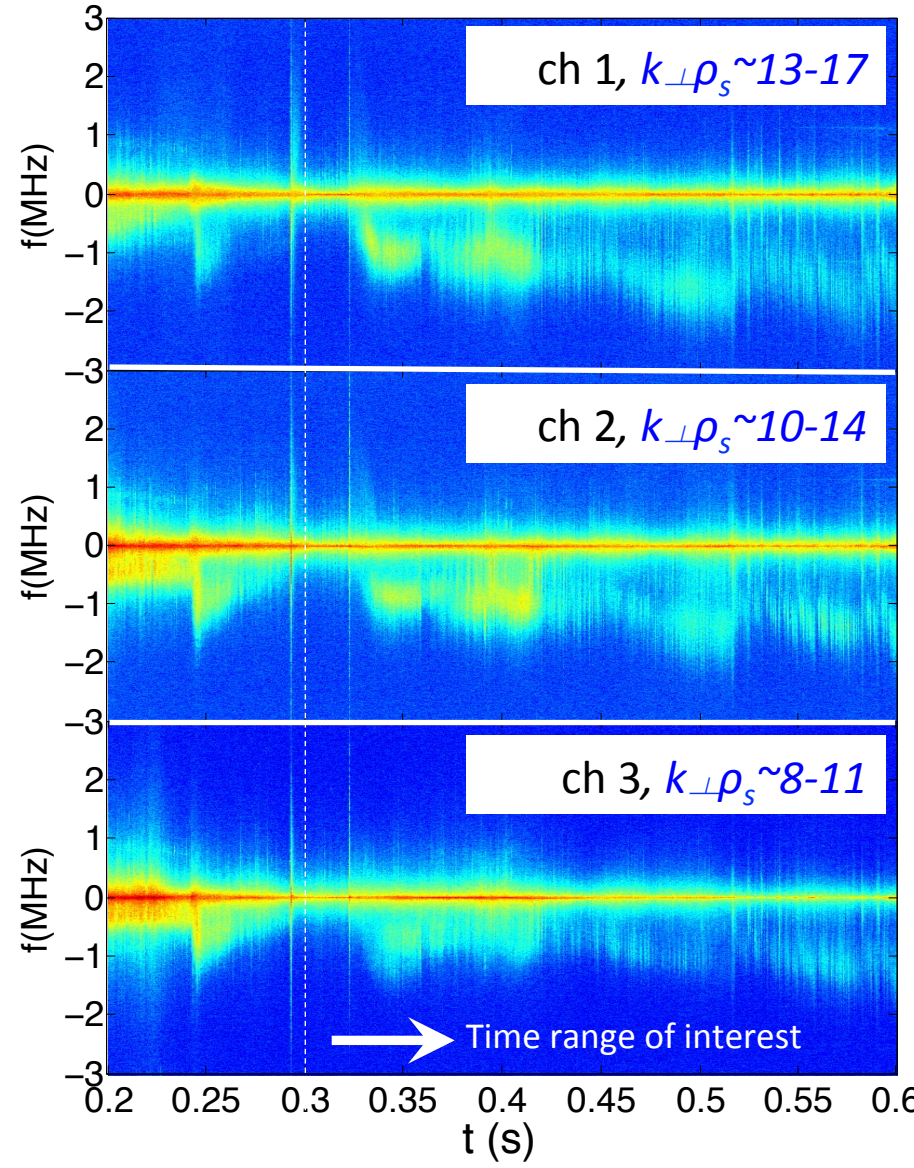


# Each Channel of the NSTX High-k Scattering System Detects a Fluctuation Wavenumber $k$

- Channel 1 detects highest  $k_{\perp}$  and  $k_t$ , Doppler shift is greatest ( $f_D = k_t v_t / 2\pi$ ).
- High peak at  $f \sim 0$  corresponds to stray radiation.
- Scattering region  $R \sim 135\text{-}136$  cm,  $r/a \sim 0.7$ . (major radius 0.85 m, minor radius 0.68 m).

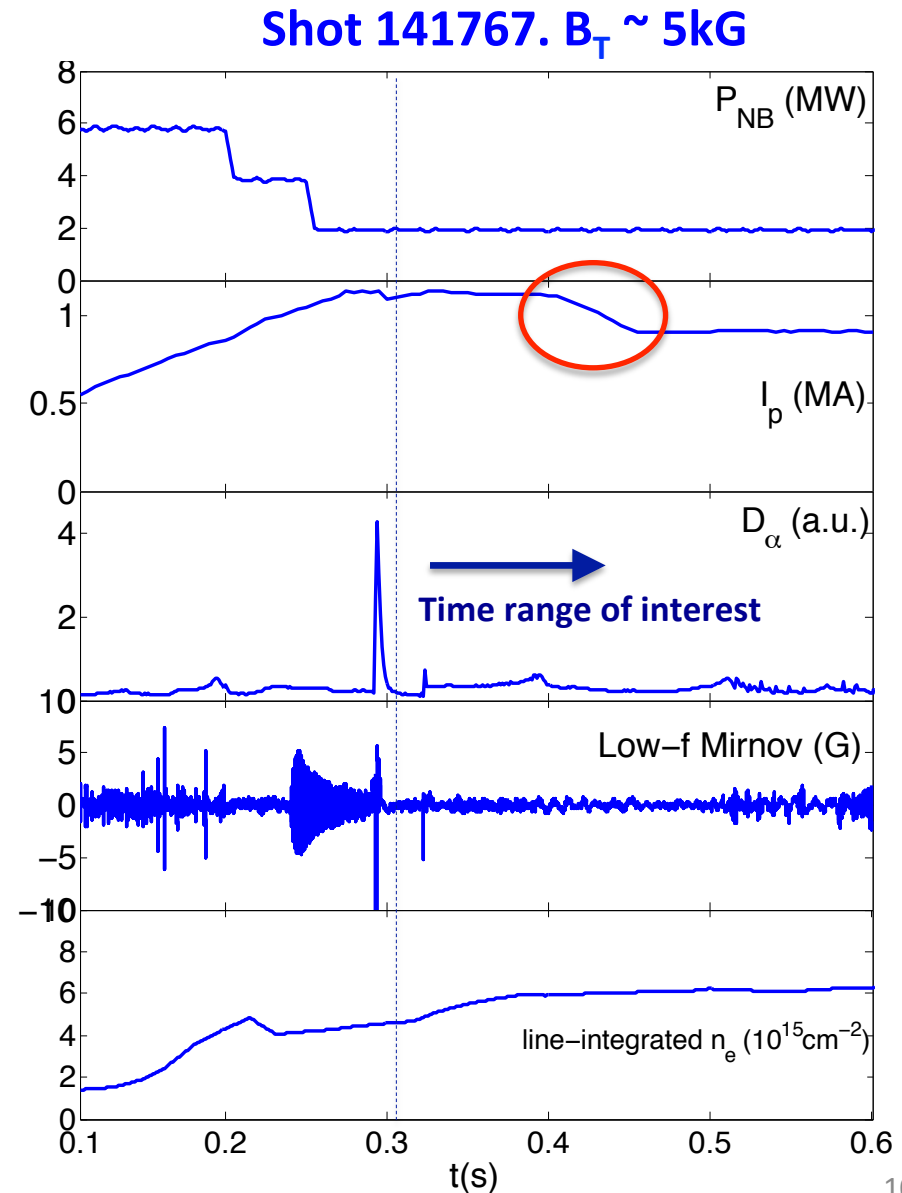


Shot 141767

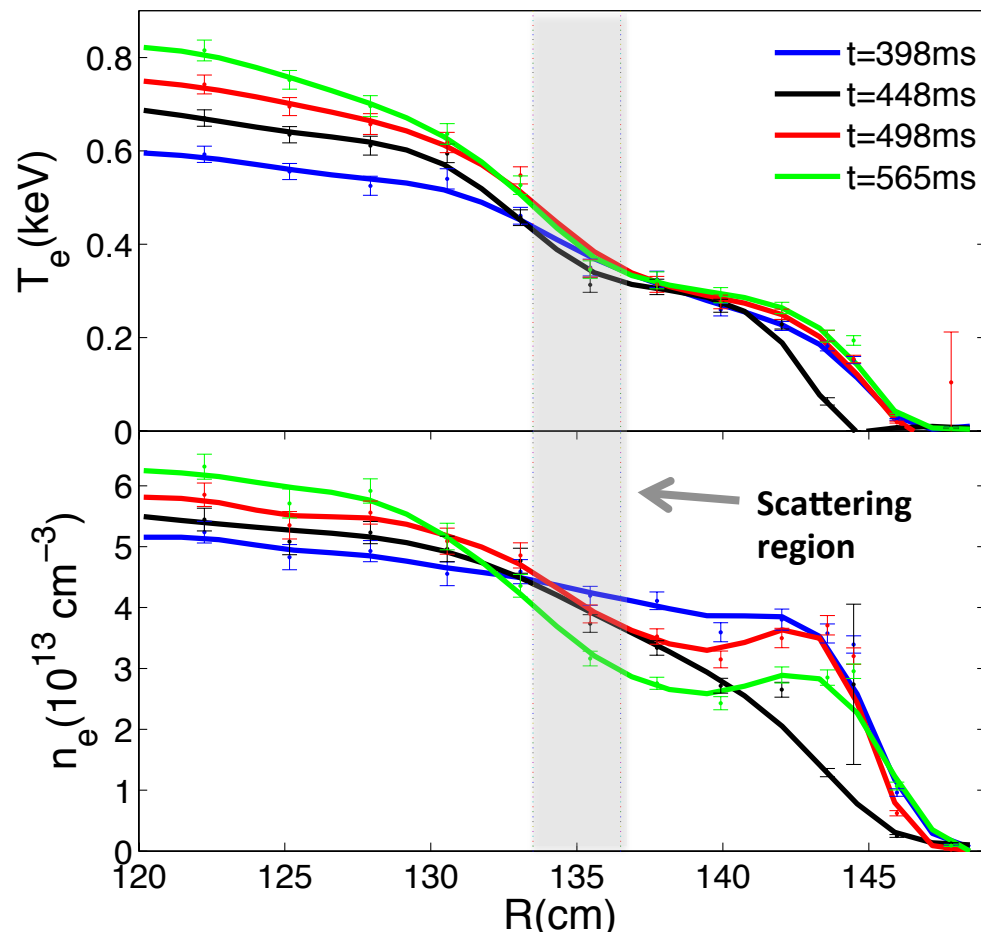
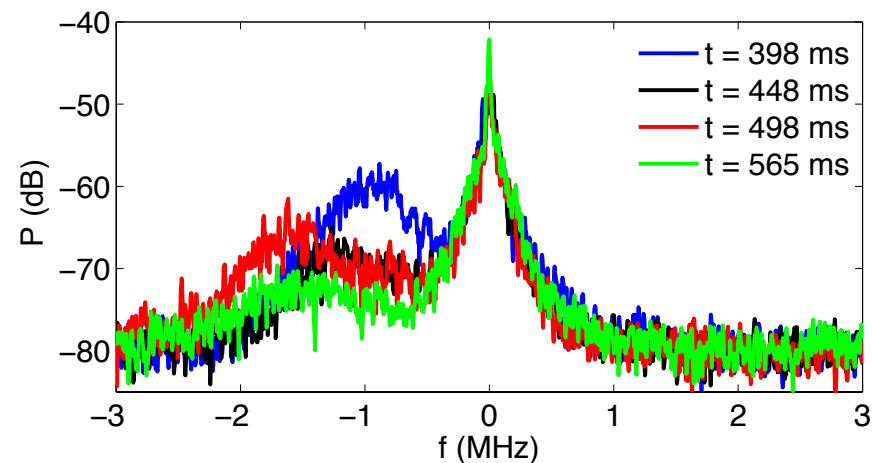
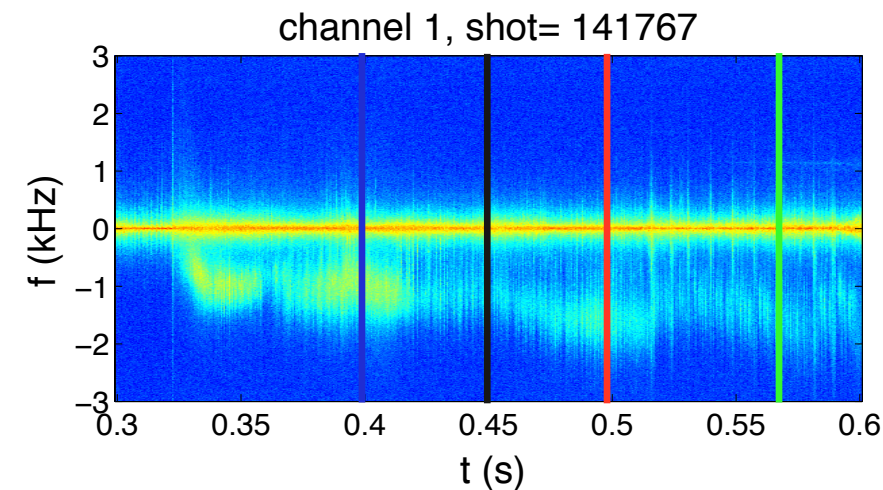


# A Set of NBI-heated H-mode Plasmas is Used to Study High-k Turbulence during Current Ramp-down

- **NBI heated**, HHFW heating is absent during the run.
- **Controlled Current ramp down** between  $t = 400$  ms and  $t = 450$  ms (from LRDFIT).
- Time range of interest is  $t > \sim 300$  ms, covering current ramp-down phase, and after ELM event at  $t \sim 290$  ms.
- **MHD activity is quiet during time range of interest.** (*cf.* low-f Mirnov signal).
- Line integrated density is fairly constant during the time range of interest.



# Observed High-k Fluctuations Correlate to Local Electron Density Gradient

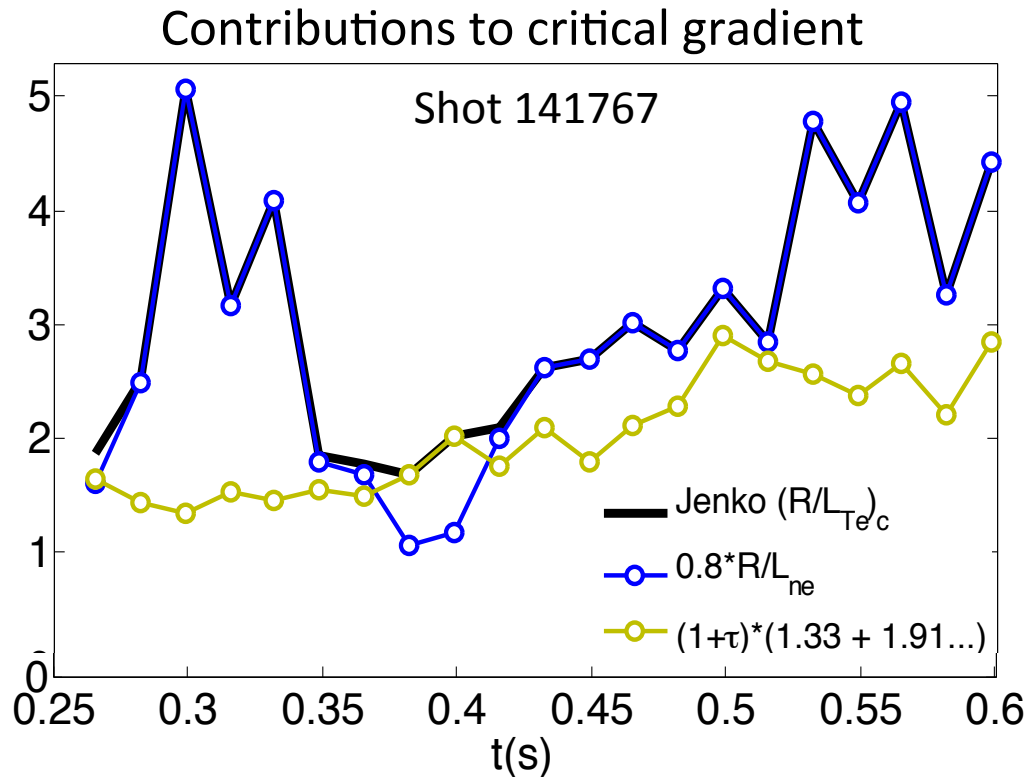


$\nabla T_e \rightarrow$  *Drives* ETG  
 $\nabla n_e \rightarrow$  *Stabilizes* ETG

} Two competing effects:  $\nabla n_e$  is dominant effect.

# Comparisons with Linear Critical Gradient Formula (Jenko)

# Electron Density Gradient Can *Linearly* Stabilize ETG Turbulence by Increasing Critical Gradient



- Jenko critical gradient is a maximum of a  $R/L_{ne}$  term and an  $s/q$  term.

$$(R/L_{Te})_{crit} = \max \begin{cases} 0.8R/L_{ne} \\ (1+\tau)(1.33+1.91\hat{s}/q)(1-1.5\epsilon)(1+0.3\epsilon d\kappa/d\epsilon) \end{cases} \quad \text{with} \quad \tau = Z_{eff}T_e/T_i$$

- Higher values of  $R/L_{ne}$  raise the critical gradient for ETG (possibly above the experimental gradient value). This *should* have a **stabilizing** effect on turbulence.

# High-k Density Fluctuations are Observed during ETG Unstable Time Periods

- Total scattered power (integrated in *freq*).

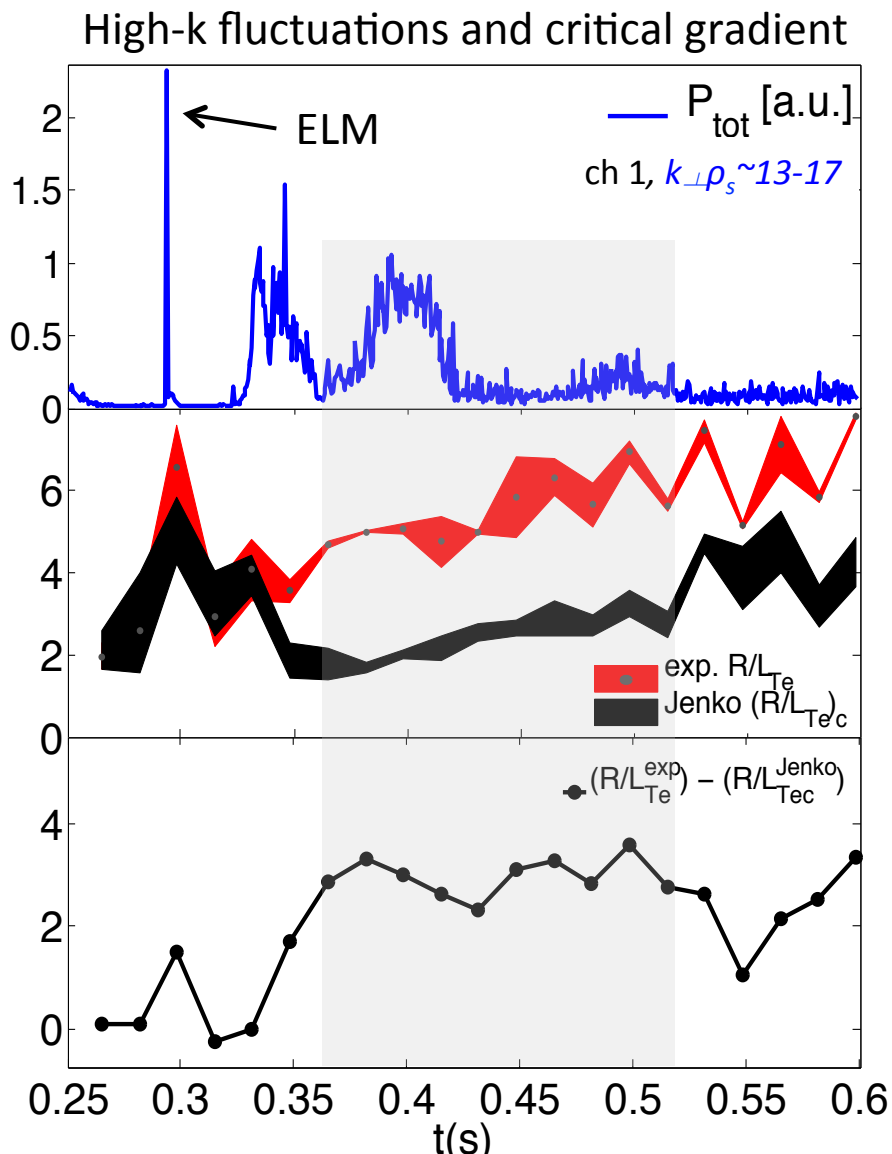
$$P_{tot} \propto (\delta n_e / n_e)^2$$

- $(R/L_{Te}^{exp}) - (R/L_{Te})_{crit}$  determines linear threshold for instability.

- $t < 320 \text{ ms}$   $(R/L_{Te}^{exp}) \sim (R/L_{Te})_{crit}$   
 → ETG marginally stable, no fluctuations.

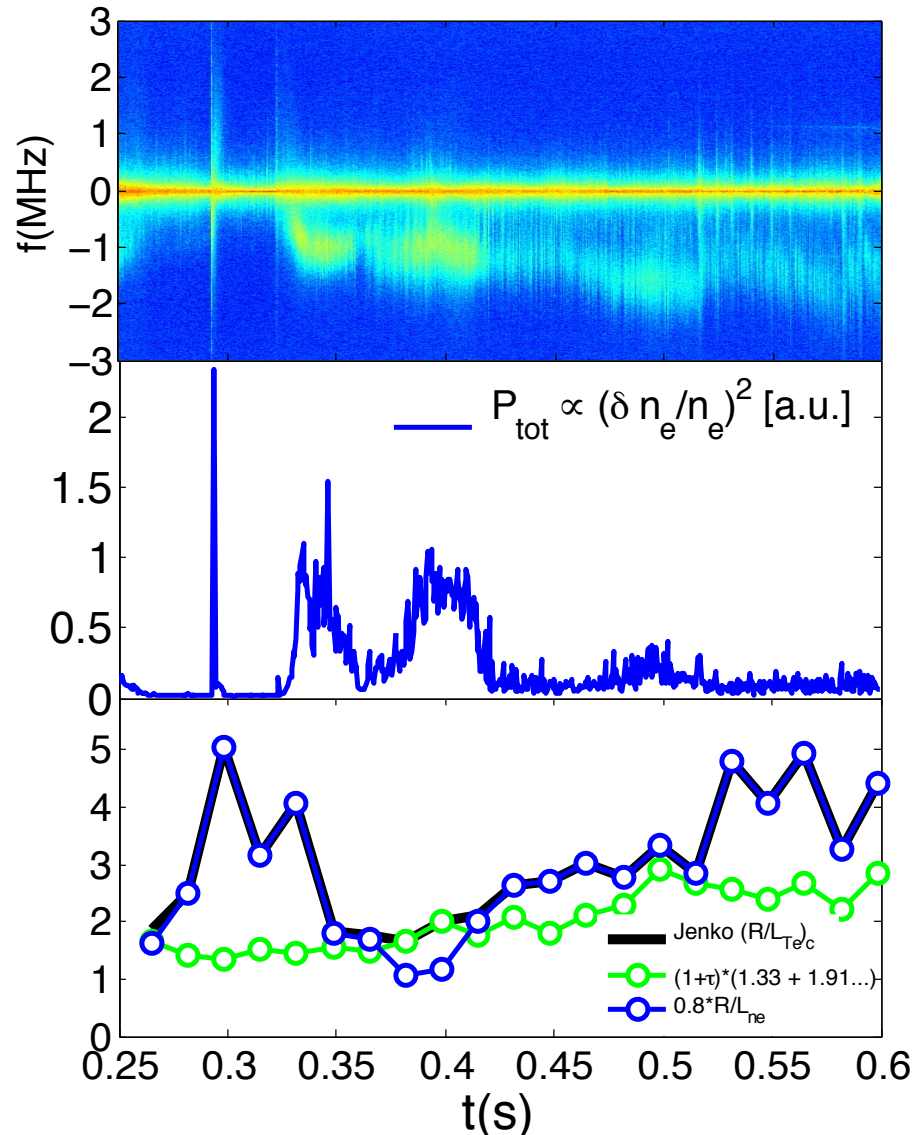
- $t > 320 \text{ ms}$   $(R/L_{Te}^{exp}) > (R/L_{Te})_{crit}$   
 → fluctuations develop.

- $360 \text{ ms} < t < \sim 520 \text{ ms}$  (gray shading)  
 Similar  $(R/L_{Te}^{exp}) - (R/L_{Te})_{crit}$  produces **VERY** different  $P_{tot}$ . Nonlinear evolution of turbulence motivates the use of nonlinear gyro-kinetic simulations (future work).



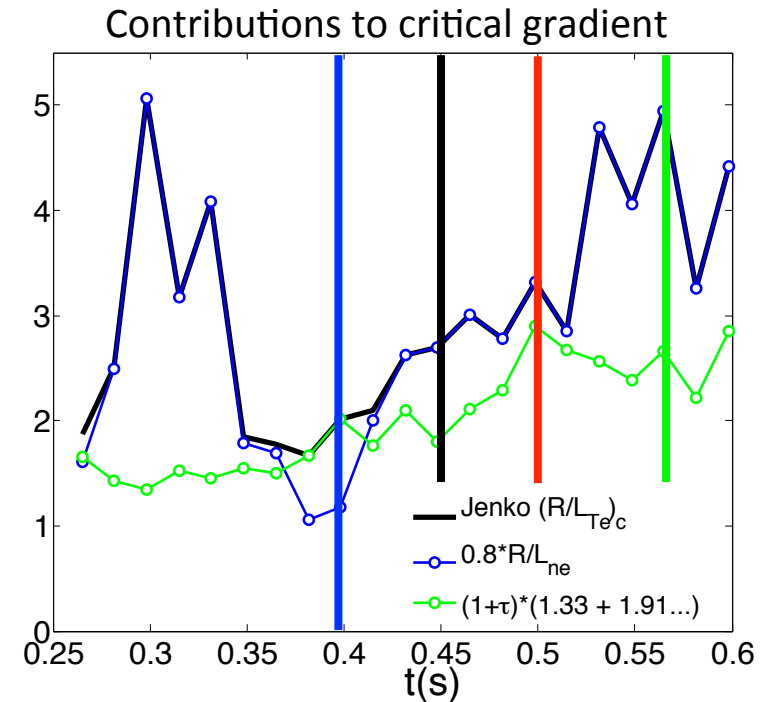
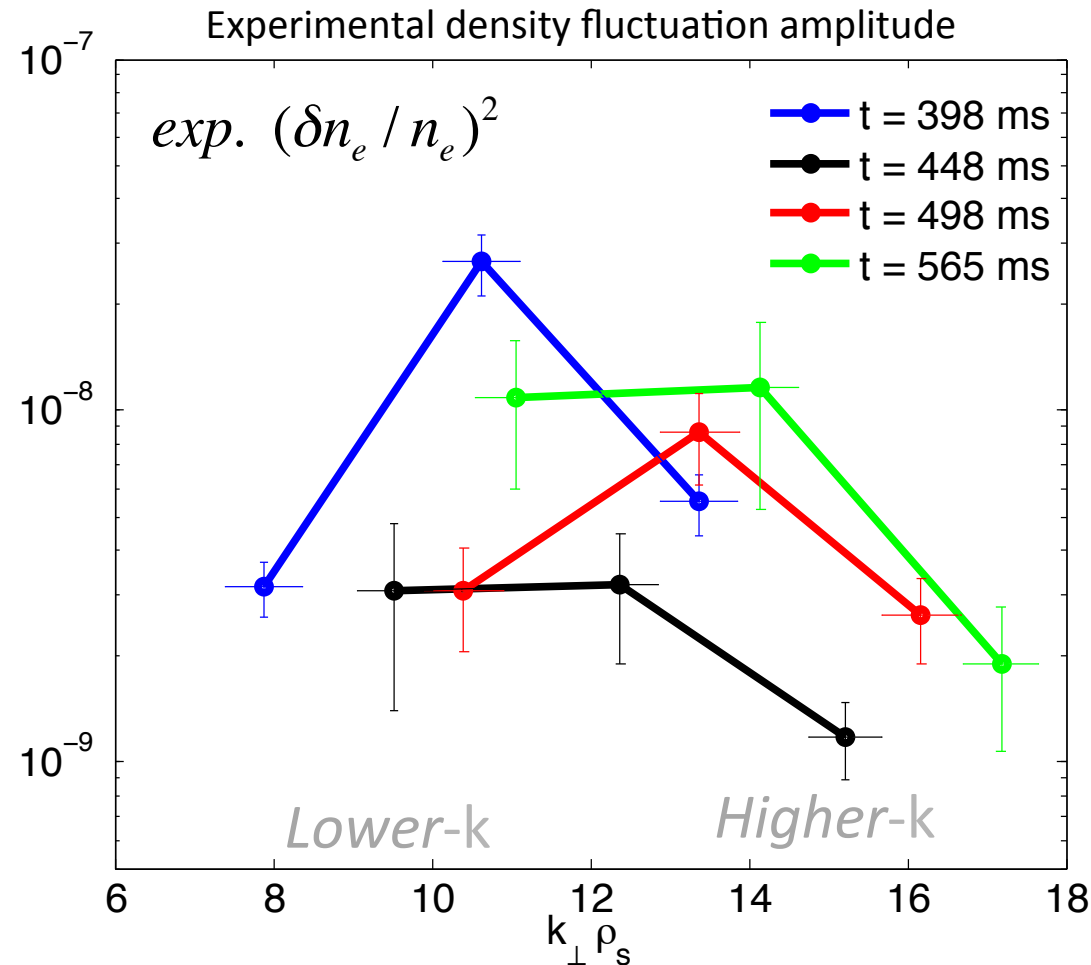
# High-k Density Fluctuations are Linearly Stabilized by Density Gradient through the Critical Gradient

shot 141767, channel 1



- As  $R/L_{ne}$  increases, it dominates in Jenko's formula  $(R/L_{Te})_{crit}$  ( $t < 340$  ms,  $t > 410$  ms &  $t > 515$  ms). Fluctuations decrease during that time.
- Previous to  $t \sim 320$  ms ETG is marginally stable with respect to Jenko critical gradient. No fluctuations are observed.
- $R/L_{ne}$  is a *linear stabilizing* mechanism when it dominates the Jenko critical gradient.
- Electron density gradient stabilization of ETG turbulence first observed by Y. Ren *et al*, Phys. Rev. Letters 2011.

# Wavenumber Spectrum of Fluctuation Amplitude and Electron Density Gradient



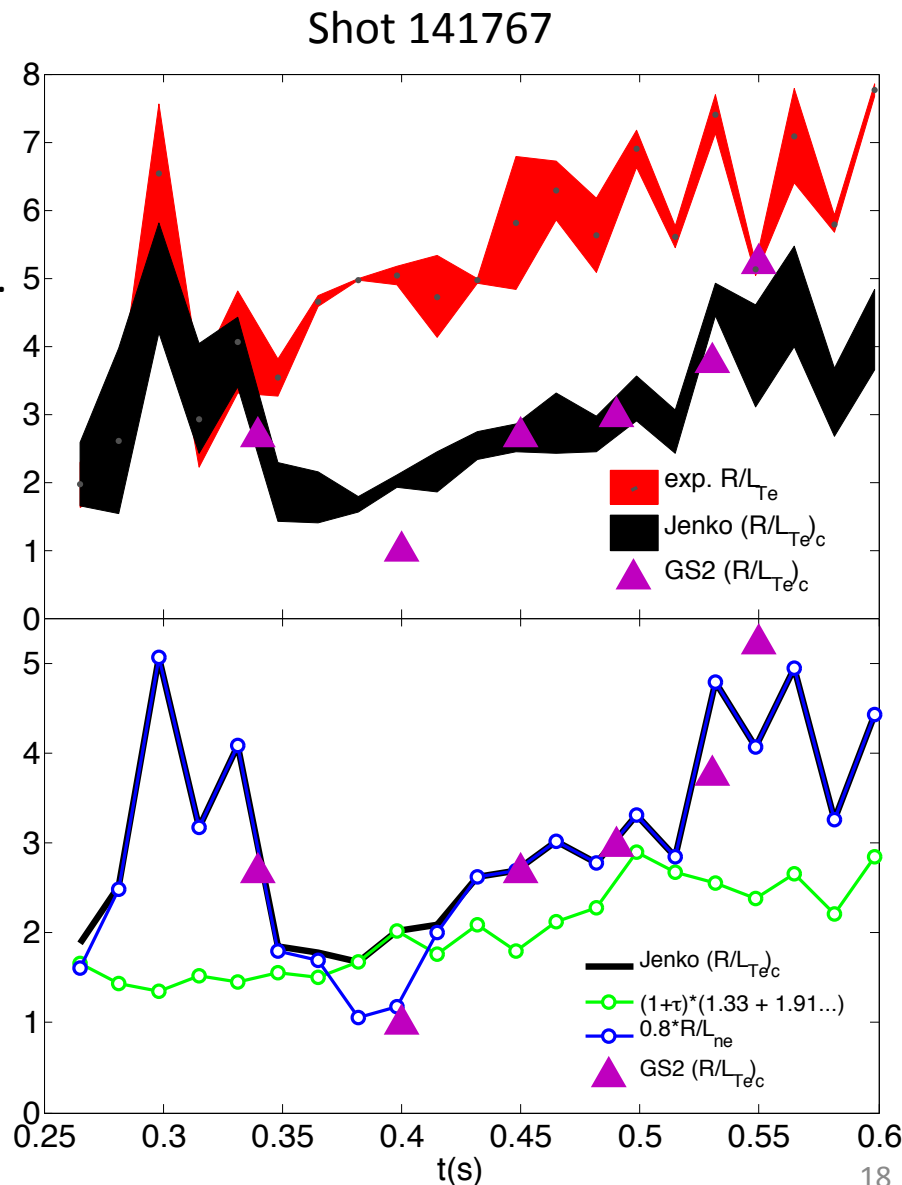
- Lower-k ( $k_{\perp} \rho_s < 10$ )  $(\delta n_e / n_e)^2$  decreases for  $398 < t < 498$  ms.
- After  $t \sim 448$  ms, higher k ( $k_{\perp} \rho_s \sim 12-16$ ) fluctuation levels increase. During that time,  $R/L_{ne}$  increases.



# Comparisons between Experiment and Linear GS2 simulations

# Critical Gradient Computed with GS2 Linear Runs Agrees with Jenko's Critical Gradient

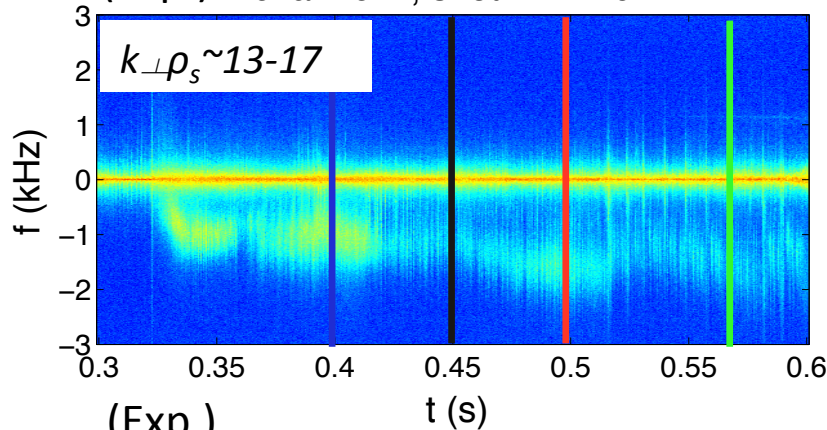
- Regime of validity of  $(R/L_{Te})_{crit}$  :
  - low- $\beta$
  - positive  $\hat{s} > 0.2$
  - large aspect ratio
 } not NSTX parameters.
- $(R/L_{Te})_{crit}$  is explicitly calculated with GS2.
- Good agreement between **GS2**  $(R/L_{Te})_{crit}$  calculations and **Jenko**  $(R/L_{Te})_{crit}$ :
  - Jenko's critical ETG formula is assumed valid in these NSTX plasmas.
- $GS2 (R/L_{Te})_{crit}$  seems to follow  $R/L_{ne}$ .



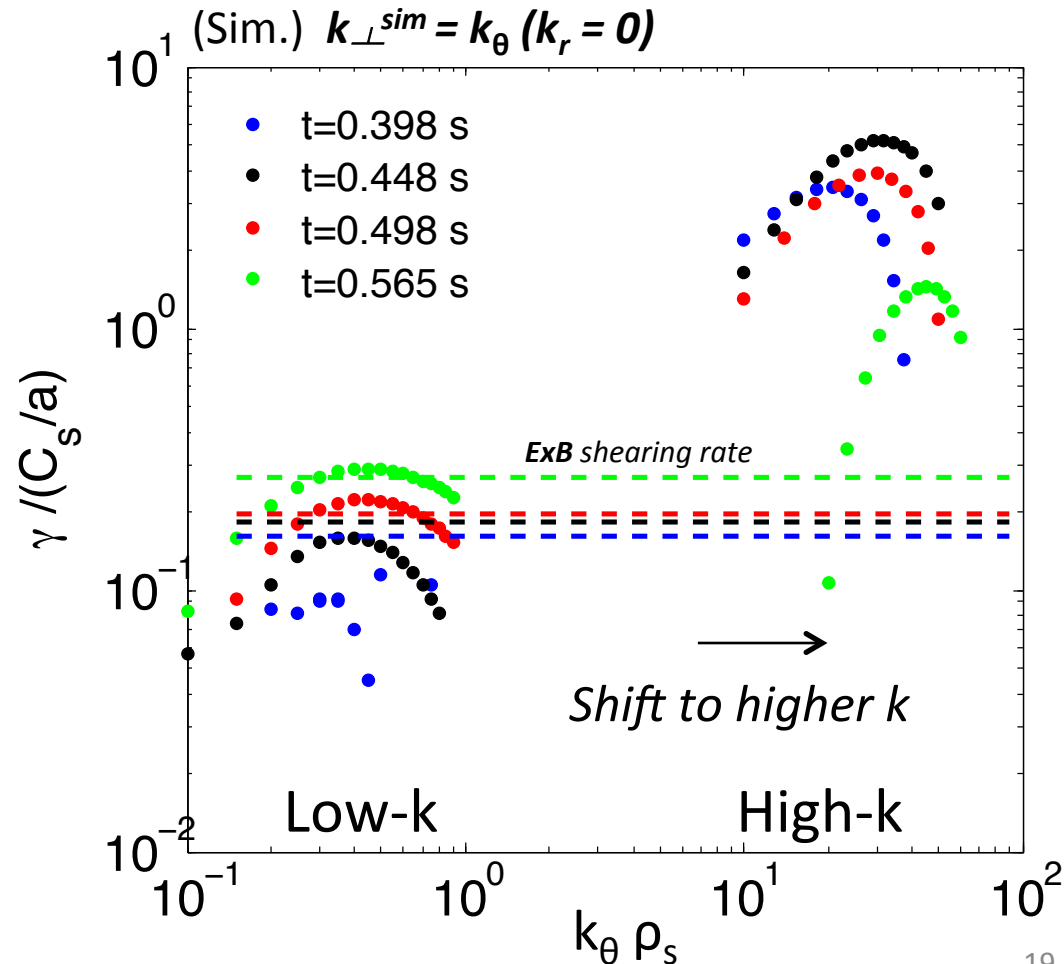
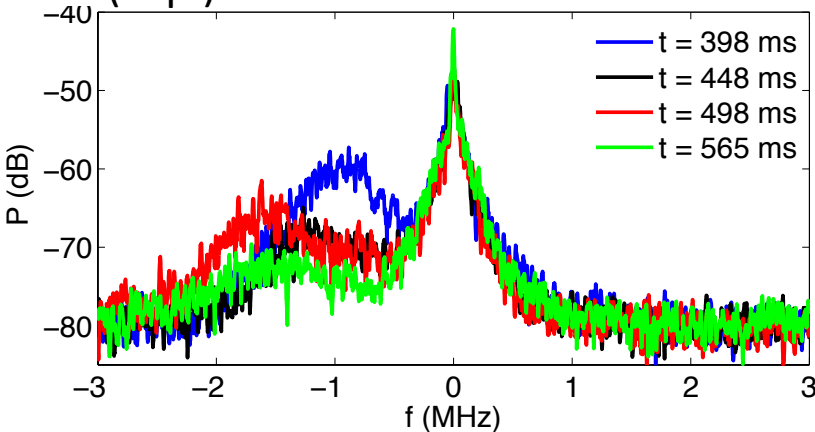
# GS2 Linear Simulations Show the Wavenumbers at Maximum Growth Rate Shift to Higher $k$ in Time

- Linear simulations compute most unstable mode ( $k_r = 0$ ). Experimental  $k$  is found to be linearly stable.
- Low- $k$  linear growth rates ( $k_\theta \rho_s \leq 1$ ) are comparable to **ExB shearing rate** levels (Waltz, Miller PoP 1999).
- High- $k$  wavenumbers corresponding to maximum linear growth rate shift towards higher- $k$ .
- Observed fluctuations decrease as  $k_b \rho_s (\gamma_{\max})$  increases.

(Exp.) channel 1, shot= 141767

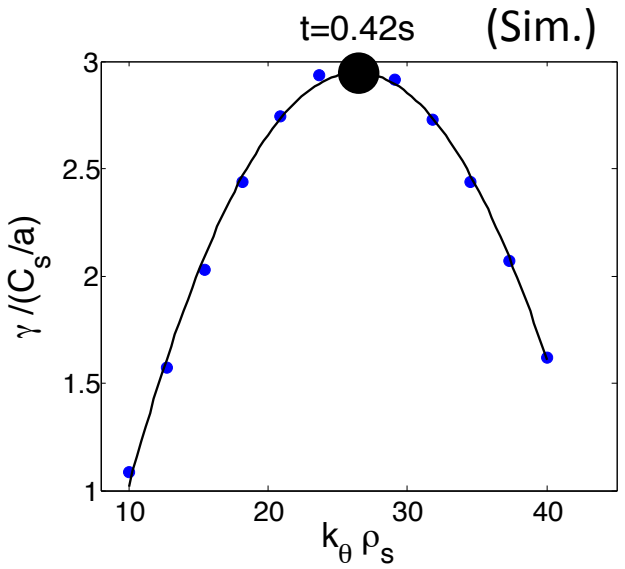


(Exp.)



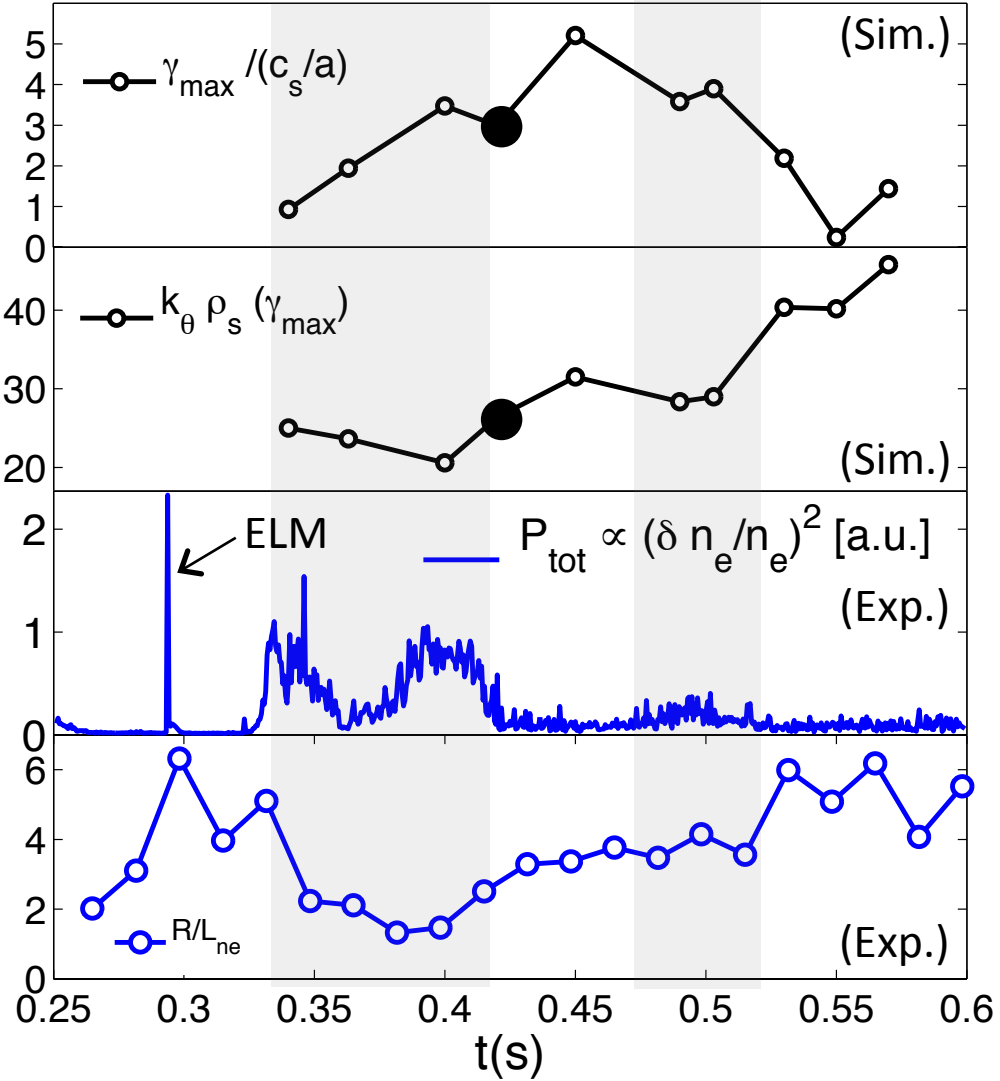
# Wavenumber at Maximum Linear Growth Rate Correlates to Electron Density Gradient and Fluctuation Amplitude

- Linear growth rates are calculated at each time: determine  $\gamma_{max}/(c_s/a)$  and  $k_{\theta}\rho_s(\gamma_{max})$  (black dot).



- $\gamma_{max}/(c_s/a)$  not correlated with and  $P_{tot}$  or  $R/L_{ne}$ .

$k_{\theta}\rho_s(\gamma_{max})$  correlates to total scattered power  $P_{tot}$  and  $R/L_{ne}$  at the scattering location (cf. evolution within time panels).



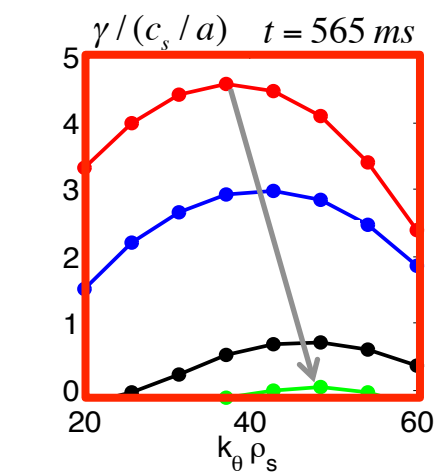
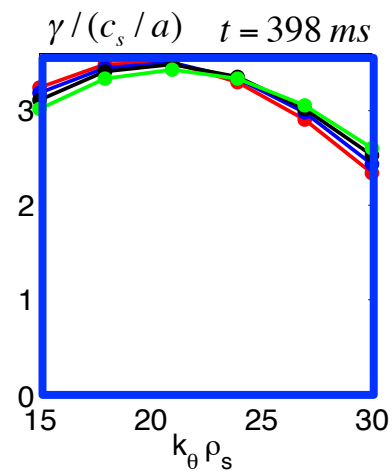
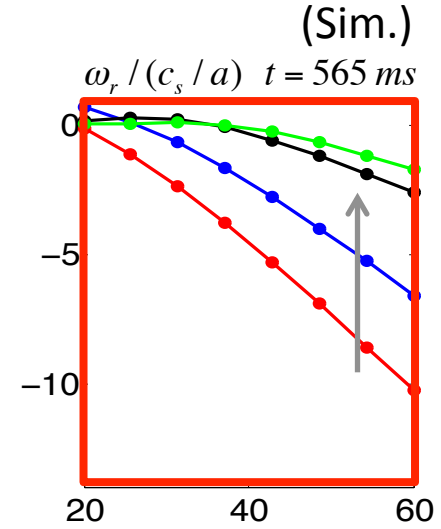
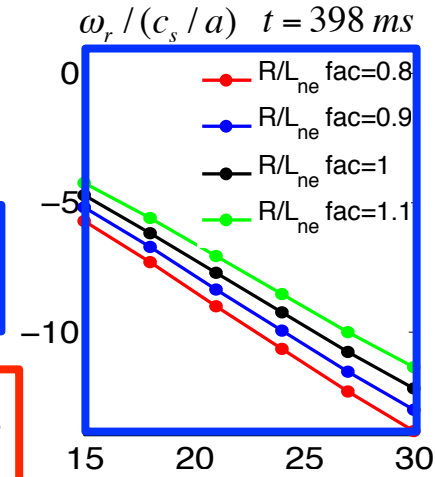
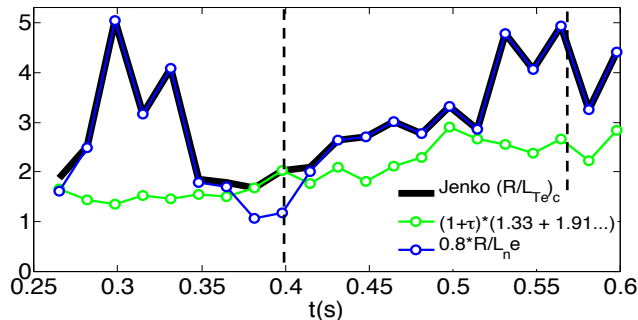
# Scan with GS2 is Performed to Confirm Effect of Electron Density Gradient on High-k Turbulence

- Real frequency  $\omega_r$  and linear growth rate  $\gamma$  are sensitive when  $0.8 \cdot R/L_{ne}$  dominates Jenko's critical gradient for ETG.

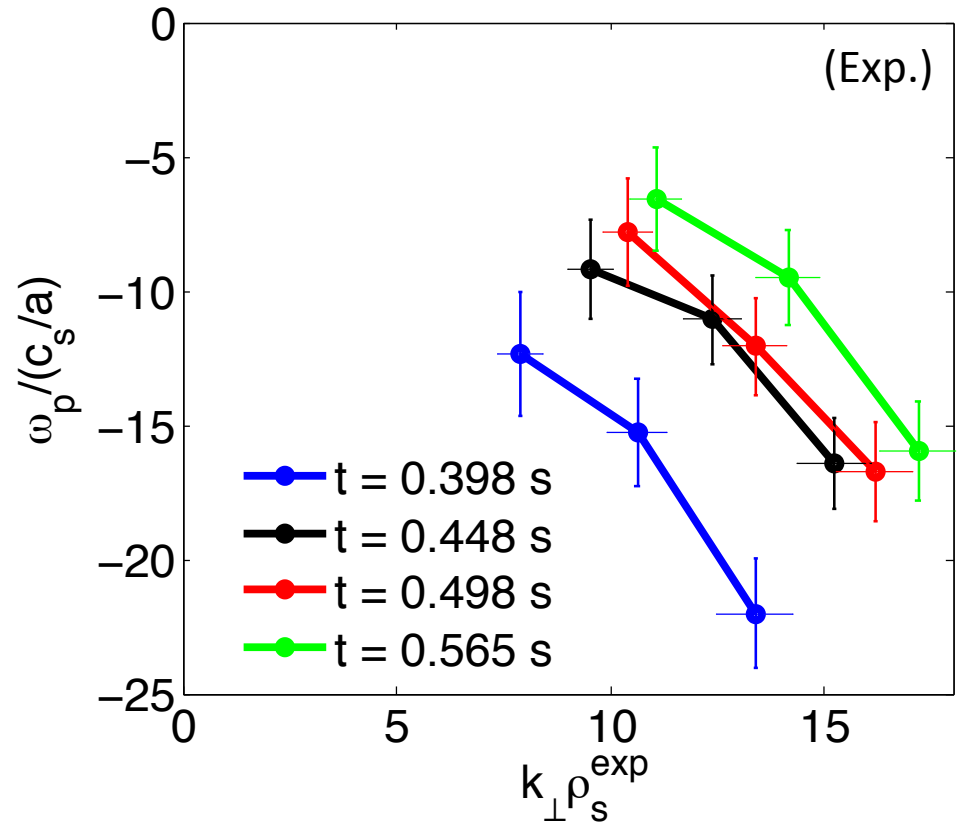
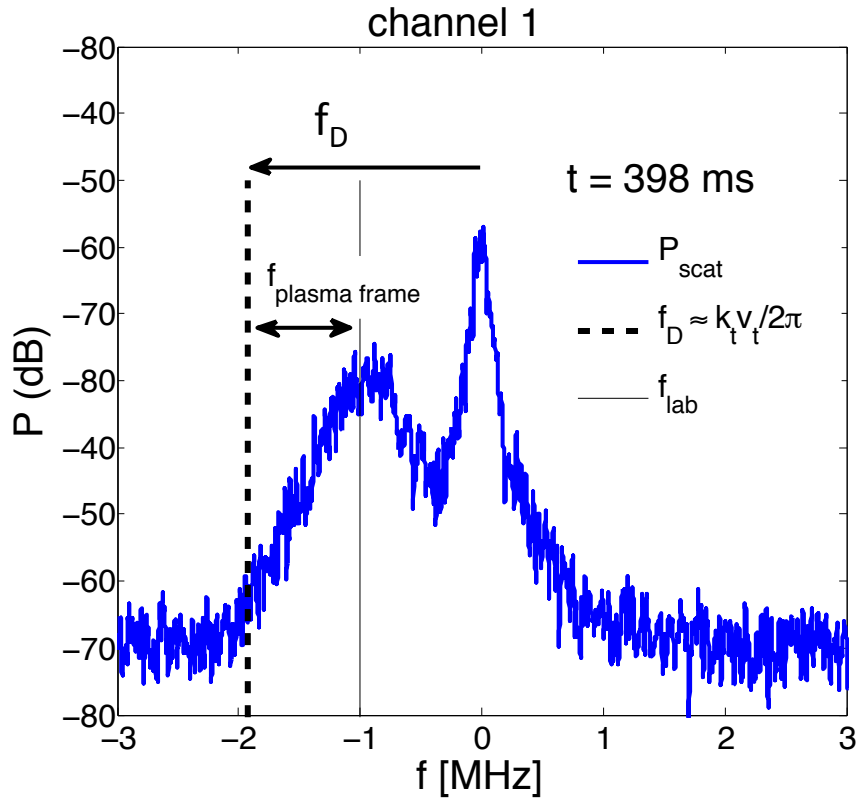
$t = 398 \text{ ms}$ ,  $0.8 \cdot R/L_{ne}$  term not dominant  
 $\rightarrow \omega_r$  and  $\gamma$  insensitive to  $R/L_{ne}$ .

At  $t = 565 \text{ ms}$   $R/L_{ne}$  dominates, and  $R/L_{ne}$  decreases  $\gamma$  and shifts  $k_{\theta} \rho_s (\gamma_{max})$  to higher-k  
 $\rightarrow$  stabilizing effect.

$|\omega_r|$  decreases with  $R/L_{ne}$ .



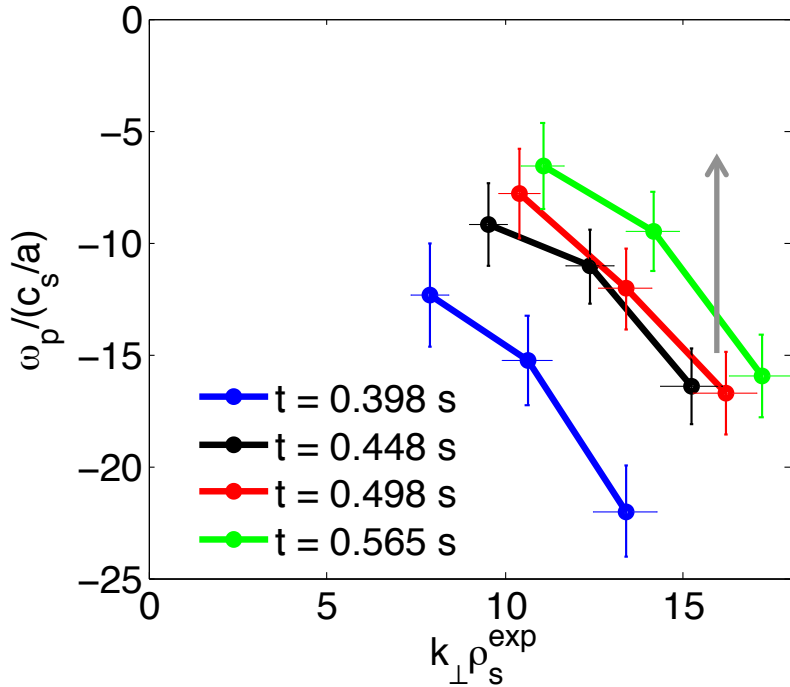
# Experimental Real Frequency of High-k Turbulence is Calculated by Subtracting Doppler Shifted Frequency



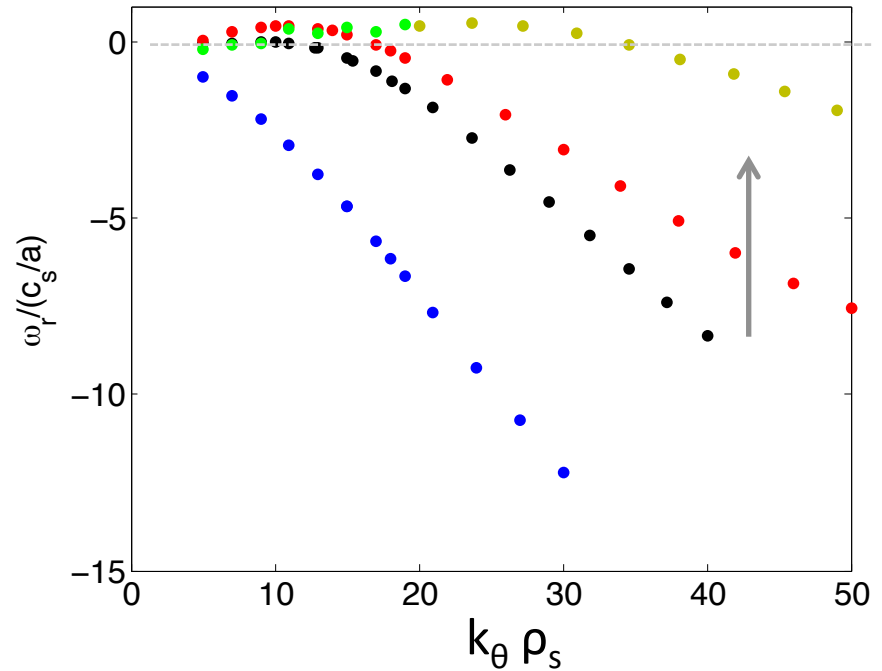
- Lab frame frequencies detected  $f_{\text{lab}}$  are Doppler shifted from plasma frame frequencies by  $f_D = k_t v_t / 2\pi$ , and  $\omega_p / 2\pi = f_{\text{lab}} - f_D$ .
- Obtain  $k_t$  from ray tracing calculations,  $v_t$  from CHERS measurement and TRANSP calculations.

# Experimental Real Frequency and GS2 Real Frequency Exhibit Similar Behavior

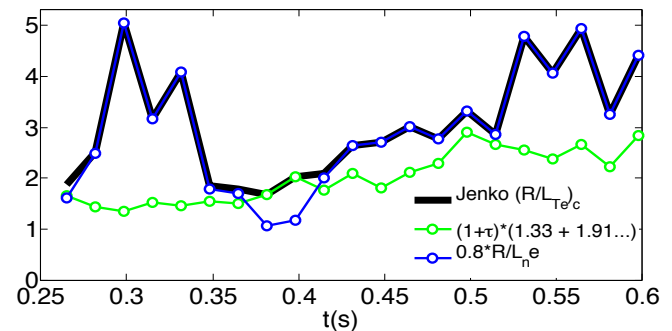
Exp:  $k_{\perp}^{exp} = \sqrt{k_r^2 + k_b^2}$  ,  $k_r/k_b \gg 1$



Sim:  $k_{\perp}^{sim} = k_{\theta}$  ( $k_r = 0$ ) (Most unstable mode)

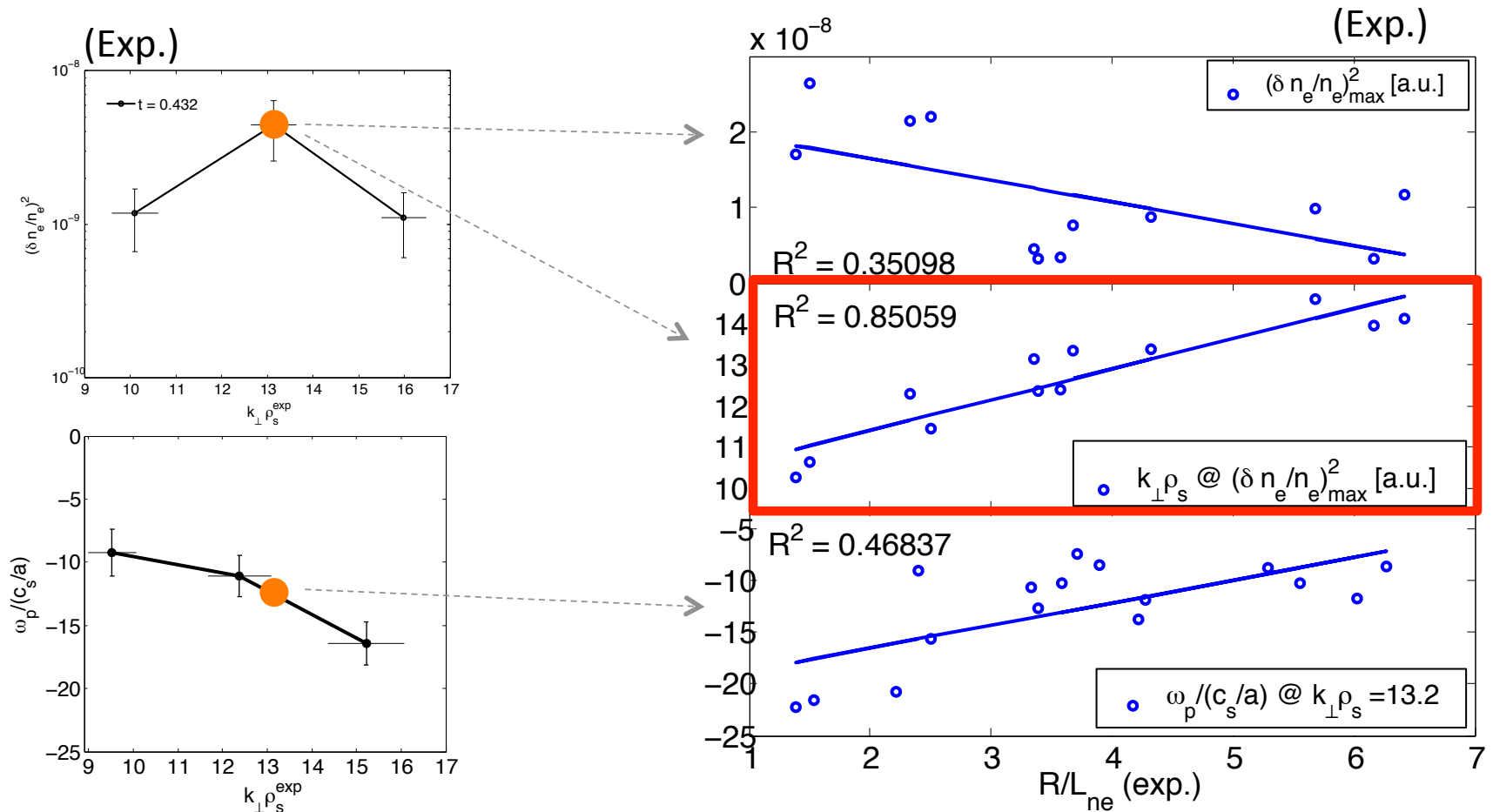


- Experimental  $k$  is linearly stable in GS2.
- $|\omega_p^{exp}|$  and  $|\omega_r^{sim}|$  decrease in time.
- Only 15% change in  $c_s/a$  (normalization).
- Note  $R/L_{ne}$  increases in time.



# Correlation Between Experimental Wavenumber at Maximum Fluctuation Amplitude and Density Gradient

Compute  $(\delta n_e/n_e)^2_{\max}$ ,  $k_{\perp}\rho_s$  @  $(\delta n_e/n_e)^2_{\max}$ , and  $\omega_r/(c_s/a)$  @  $(k_{\perp}\rho_s=13.2)$  and compare to  $R/L_{ne}$ .

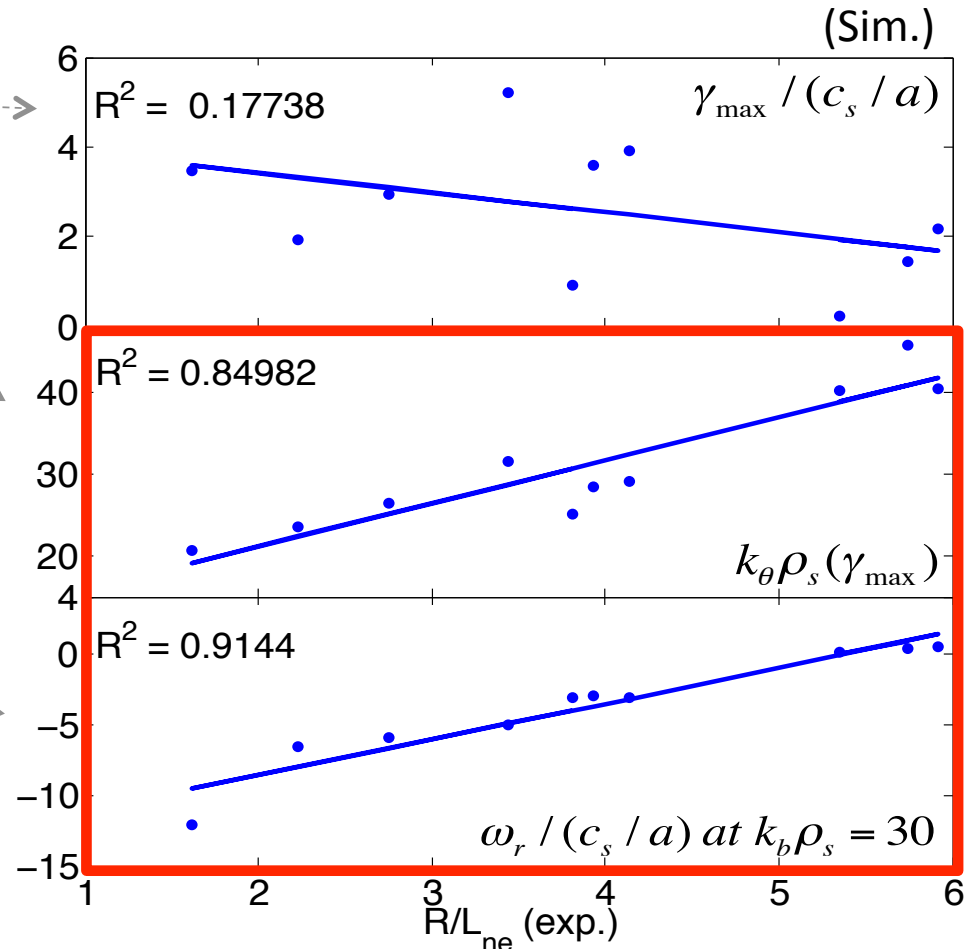
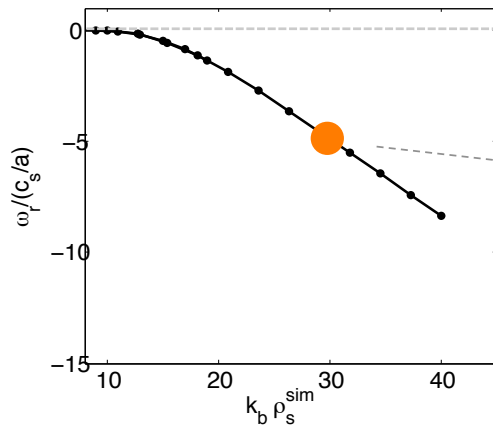
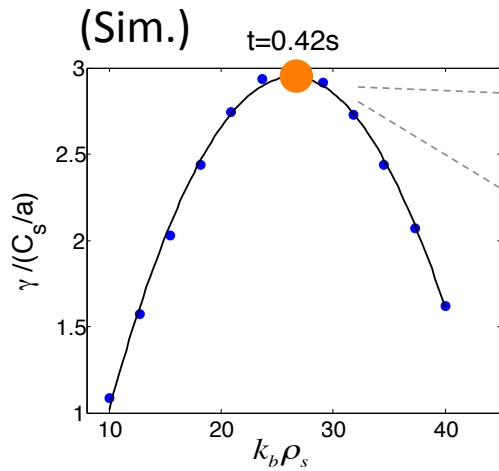


- Low correlation between  $(\delta n_e/n_e)^2_{\max}$ ,  $\omega_r/(c_s/a)$  and experimental  $R/L_{ne}$ , but note **similar trend** as found from GS2 linear simulations (slide 25).
- Correlation between  $k_{\perp}\rho_s$  @  $(\delta n_e/n_e)^2_{\max}$  and  $R/L_{ne}$  (possible beam refraction effects on  $k_{\perp}$ ).



# Correlation Between GS2 Wavenumbers at Maximum Growth Rate, Real Frequency and Electron Density Gradient

Compute  $\gamma_{max}$ ,  $k_b \rho_s(\gamma_{max})$  and  $\omega_r/(c_s/a)$  @ ( $k_b \rho_s=30$ ) in time, compare to  $R/L_{ne}$ .

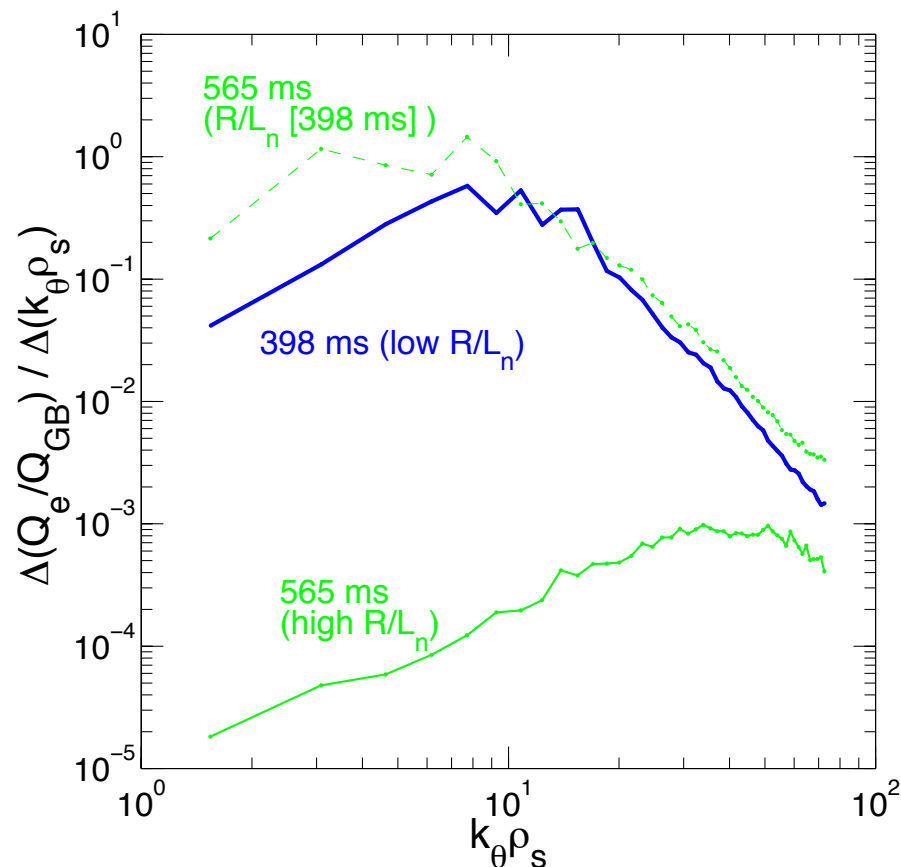
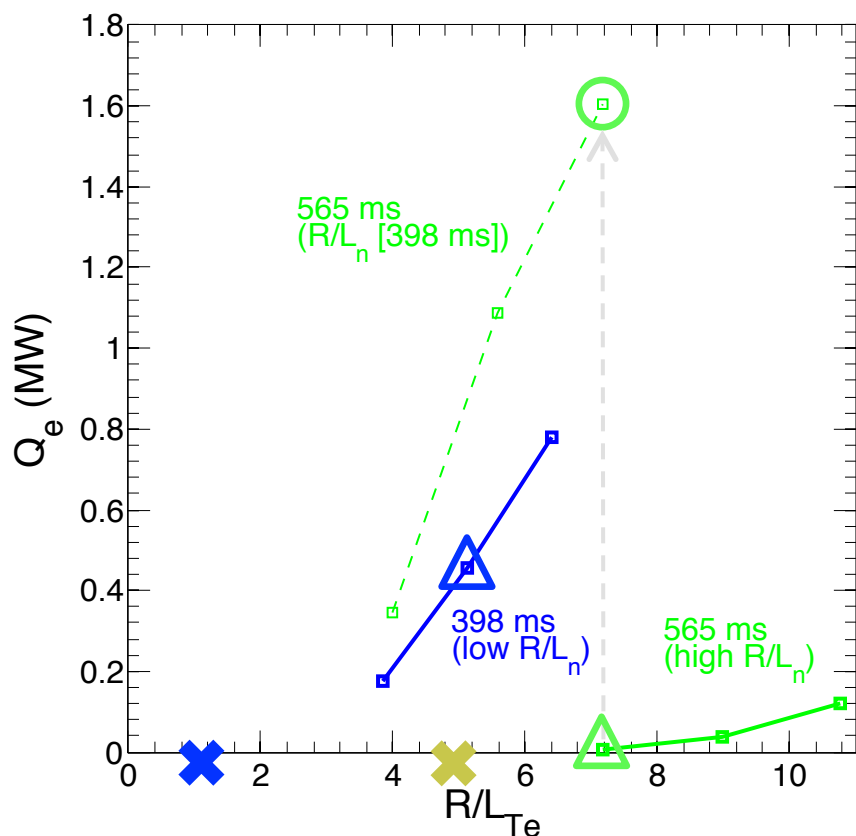


- Low correlation between  $\gamma_{max}$ , and experimental  $R/L_{ne}$ .

- Correlation between  $k_b \rho_s(\gamma_{max})$ ,  $\omega_r/(c_s/a)$  @ ( $k_b \rho_s=30$ ) and  $R/L_{ne}$ .

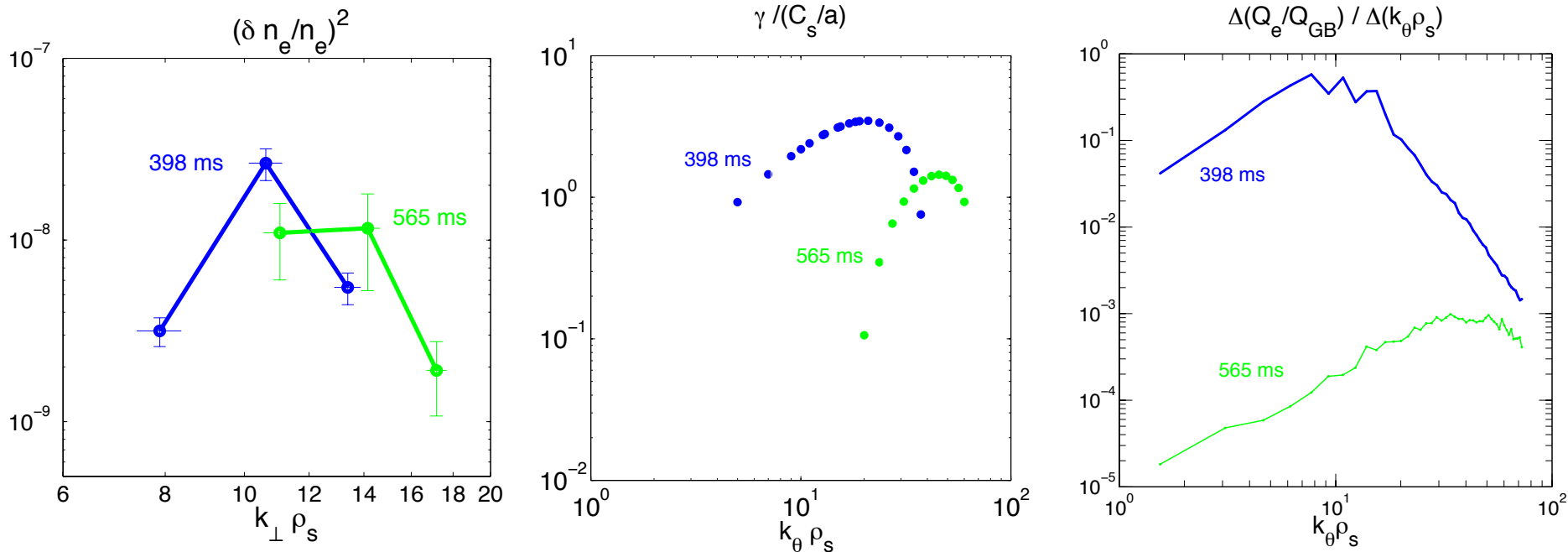
# Comparisons with nonlinear GYRO simulations

# Nonlinear High-k Gyrokinetic Simulations Show High Reduction of Electron Heat Flux at High Density Gradient



- High density gradient increases *nonlinear* critical gradient, and reduces electron heat flux and stiffness.
- $Q_e(t = 565 \text{ ms}) \sim 0 \rightarrow$  unlikely to account for experimental  $Q_e \rightarrow$  study low-k contrib. to  $Q_e$ .
- TRANSP  $Q_e(t=398\text{ms}) \sim 1.54\text{MW}$ ,  $Q_e(t=565\text{ms}) \sim 0.91\text{MW} \rightarrow$  GYRO under predicts  $Q_e$ .
- Nonlinear threshold is the same for blue and green dashed curve (same  $R/L_{ne}$ ).
- Crosses indicate linear critical gradient (Jenko). Nonlinear upshift of critical gradient (cf. Peterson PoP 2013).

# Experimental Density Fluctuations, Linear Growth Rates and Electron Heat Flux Spectrum Behave Similarly with $R/L_n$



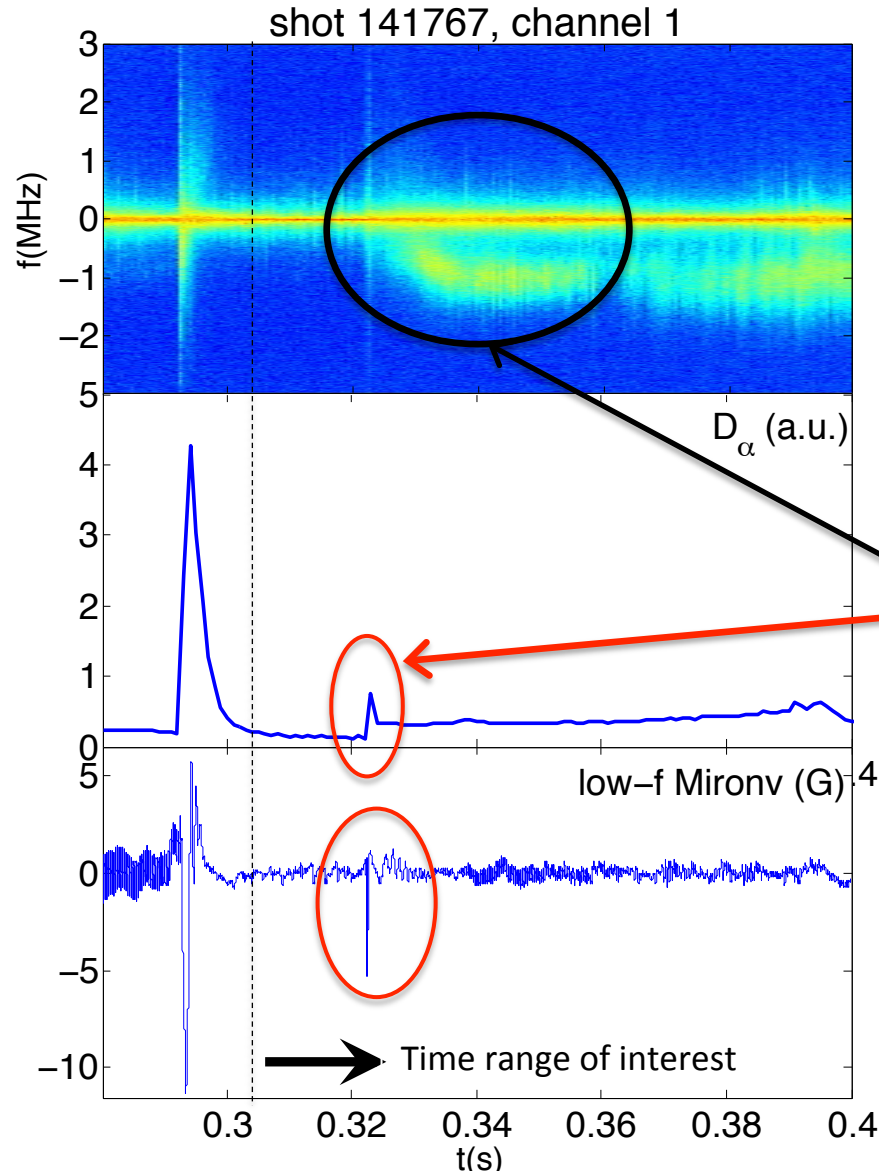
- $t = 398$  ms is low density gradient time,  $t = 565$  ms is high density gradient time.
- $(\delta n_e / n_e)^2$ ,  $\gamma / (c_s / a)$  and  $\Delta(Q_e / Q_{GB})$  are greatest at  $t = 398$  ms.
- The lower- $k$  part of the spectrum of  $(\delta n_e / n_e)^2$ ,  $\gamma / (c_s / a)$  and  $\Delta(Q_e / Q_{GB})$  is most reduced at high density gradient.

## Future Work

- Study low-k and high-k turbulence using nonlinear GYRO in this discharge to determine if they account for experimental electron heat flux levels (TRANSP).
- Develop a synthetic high-k scattering diagnostic for GYRO to perform quantitative comparisons between experiment and gyrokinetic simulations.
- Perform global nonlinear GYRO simulations to understand profile effects in NSTX plasmas ( $\rho_* \sim 1/100$ ).
- Use reduced models such as TGLF.
- Understand how  $I_p$  current ramp down can modify electron density and temperature profiles.

# Back-up slides

# High-k Fluctuations Start after Small Spike in $D_\alpha$ and Mirnov Signal



- Before  $t \sim 290$  ms, MHD activity is high. At  $\sim 290$  ms, an ELM event takes place and MHD activity quiets.
- Between  $t \sim 290$  ms and  $t \sim 320$  ms, high-k fluctuations are absent and MHD activity is quiet.
- **High-k fluctuations** start at  $t \sim 320$  ms, after small ELM event, detected in  **$D_\alpha$  and Mirnov signal**.

# Typical quantities in these NSTX D plasmas

- Measured fluctuation wavenumbers  $k_{\perp} \sim 20 \text{ cm}^{-1} \sim 2000 \text{ m}^{-1}$
- $\omega_{pe} = 2\pi \cdot 90 \text{ GHz} \cdot \sqrt{n_e(10^{20} \text{ m}^{-3})} \sim 3.6 \cdot 10^{11} \text{ s}^{-1}$
- $f_{pe} \sim 57 \text{ GHz}$
- $\omega_{pD} = \omega_{pe} / \sqrt{m_i/m_e} \sim 5.9 \cdot 10^9 \text{ s}^{-1}$
- $f_{pD} \sim 0.94 \text{ GHz}$
- $\Omega_{ce} = 2\pi \cdot (28 \text{ GHz/Tesla}) \sim 8.8 \cdot 10^{10} \text{ s}^{-1}$
- $f_{ce} \sim 14 \text{ GHz}$
- $\omega_{pe} / \omega_{ce} \sim 4 \gg 1$  (no ECH)
- $\omega_{cD} = 2\pi \cdot (7.6 \text{ MHz/Tesla}) \sim 2.4 \cdot 10^7 \text{ s}^{-1}$
- $f_{cD} \sim 3.8 \text{ MHz} \gg$  drift wave fluct (low-f)
- $v_{te} = \sqrt{2} \cdot 4.2 \cdot 10^5 \text{ [m/s]} \cdot \sqrt{T_e \text{ [eV]}} \sim 1.3 \cdot 10^7 \text{ m/s}$
- $c_s = \sqrt{2} \cdot v_{te} / \sqrt{m_i/m_e} \sim 3.03 \cdot 10^5 \text{ m/s}$
- Debye length  $\lambda_{de} = v_{te} / (\sqrt{2} \cdot \omega_{pe}) \sim 2.6 \cdot 10^{-5} \text{ m}$
- e- collisionless skin depth  $\delta_e = c / \omega_{pe} \sim 8.8 \cdot 10^{-4} \text{ m}$
- Alfvén velocity  $v_A = c \cdot (f_{ci} / f_{pi}) / \sqrt{1 + (f_{ci} / f_{pi})^2} \sim c \cdot f_{ci} / f_{pi} \sim 1.21 \cdot 10^6 \text{ m/s}$
- Tor. Rotation vel.  $v_t$  (CHERS)  $\sim 70 \text{ km/s}$
- $\beta = c_s^2 / v_A^2 \sim 0.06$
- $\rho_e = v_{te} / (\sqrt{2} \cdot \omega_{ce}) \sim 0.1 \text{ mm}$
- $\rho_s = c_s / (\sqrt{2} \cdot \omega_{ci}) \sim \rho_e \cdot \sqrt{m_i/m_e} \sim 0.6\text{-}0.7 \text{ cm}$
- $\Omega_d =$



# Spatial Localization and Wavenumber Resolution

- Volume overlap of incident and scattered beams leads to poor spatial localization.
- Theory [cf. Horton *Rev. Mod. Phys.* 1999] predicts  $k_{||} \sim 1/qR \ll k_{\perp} \Rightarrow \vec{k} \cdot \vec{B} \approx 0$
- Plasma fluctuations must satisfy:
 

$k \cdot B \approx 0$	(1) Perpendicular fluctuations.
$k = 2k_i \sin(\theta_s / 2)$	(2) Bragg Condition
- When incident beam forms a small angle with  $\vec{B}$ , (1) and (2) become highly dependent on **toroidal curvature** of magnetic field (cf. scattered beams at  $P_1$  and  $P_2$  in the figure). **Oblique propagation** (outside the midplane) of incident beam exploits this phenomenon and enhances *longitudinal localization* of fluctuations [cf. Mazucatto *Phys. Plasmas* 2003].

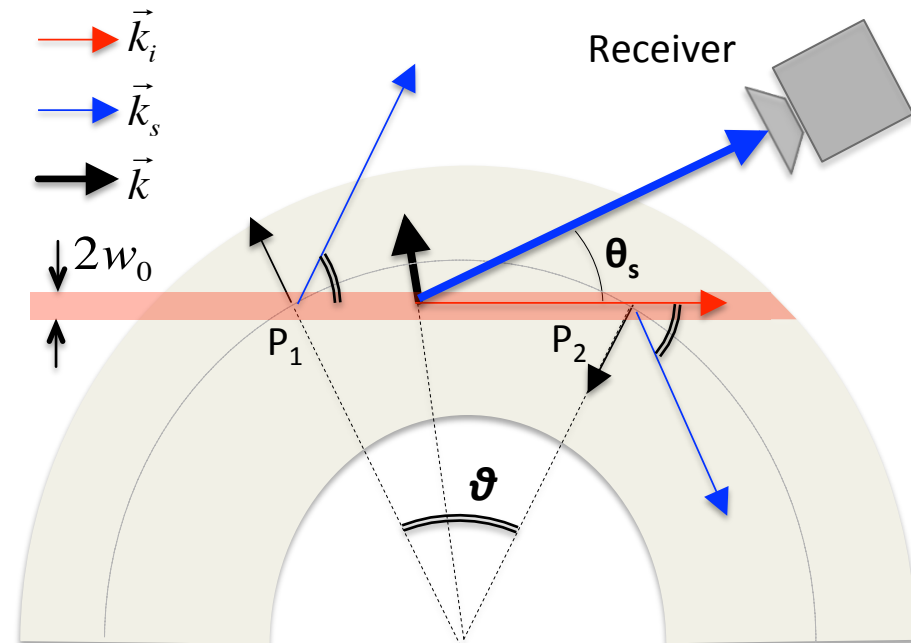
- For midplane propagation, (1) and (2) are only satisfied at  $P_1$  and  $P_2$  and fluctuation wavenumber is purely in the **radial direction**.
- In practice, beam propagation is out of midplane, but oblique angle is small ( $\sim 5^\circ$ ).  $k$  is *mostly* radial.

- Gaussian beam width dictates  $k$  and R-resolution

$$A(r_{\perp}) = \exp(-r_{\perp}^2 / w_0^2)$$

$$G(k_{\perp}) = \exp(-k_{\perp}^2 / \Delta k^2)$$

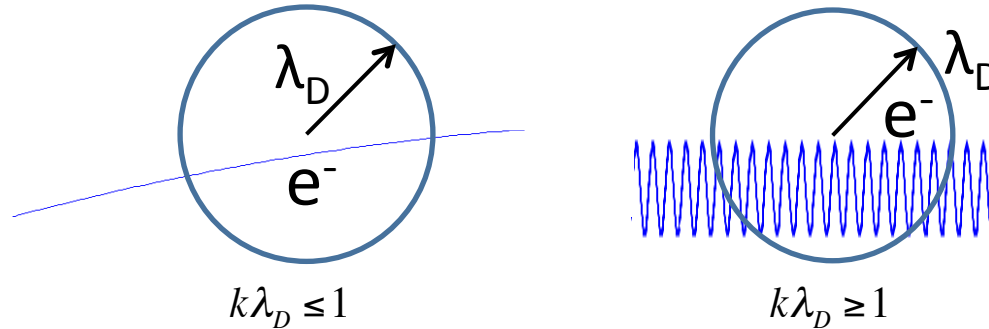
$$\Delta k = 2 / w_0$$



View from top (not to scale)

# Collective Thomson Scattering Theory is used to measure ETG-scale turbulence

- Collective/coherent and incoherent scattering



- Typical values (NSTX)  $\lambda_D \sim 10^{-5}$  m,  $k \sim k_{\perp} < 10^4$  m $^{-1}$  (*high-k*)  
➔  $k\lambda_D < 1$  (**collective scattering**)

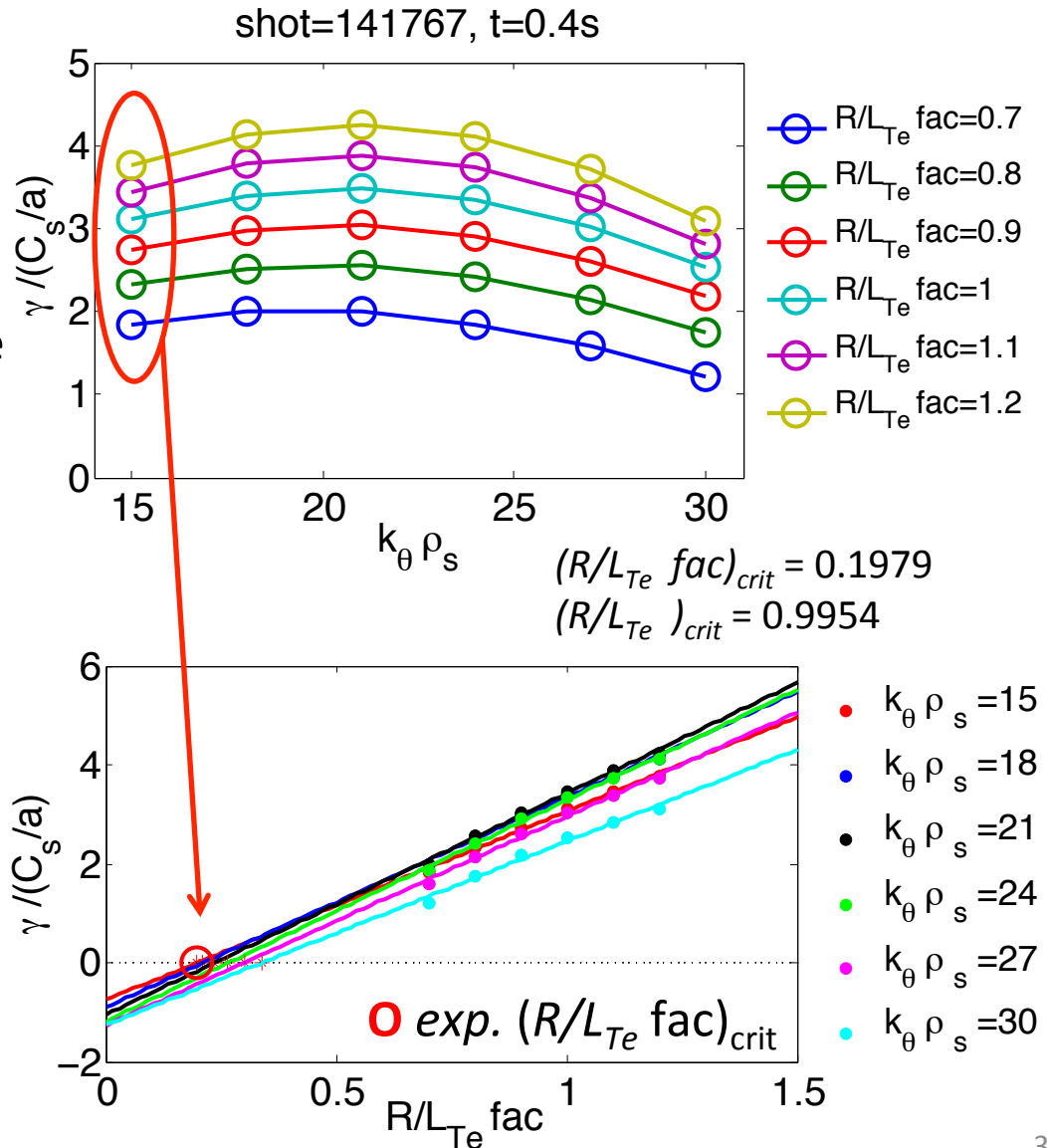
- Scattered power density

$$\frac{d^2 P}{d\Omega d\nu} = P_i r_e^2 L_z |\Pi \cdot \hat{e}|^2 \frac{|\tilde{n}_e(k, \omega)|^2}{VT}$$

$r_e$  classical electron radius  
 $V, L_z$  volume and length of scattering volume  
 $\Pi$  polarization tensor  
 $\hat{e}$  direction of incident electric field  
 $T$  observation time

# A Scan on $R/L_{Te}$ is Performed to Compute a Critical Gradient with GS2 Linear Runs

- $R/L_{Te}$  is varied keeping all other quantities constant. The factor is called ( $R/L_{Te}$  fac).
- High-k linear growth rates saturate with decreasing ( $R/L_{Te}$ ).
- $(R/L_{Te})_{crit}$  is found to be the minimum  $R/L_{Te}$  to satisfy  $\gamma = 0$ .



# Nonlinear GYRO Simulations details

Used numerical parameters based on previous convergence studies [Guttenfelder & Candy, PoP (2011), Ren et al., PoP (2012)]

- 3 kinetic species, D, C, e ( $Z_{\text{eff}} \sim 1.85-1.95$ )
- Electromagnetic:  $A_{||} + B_{||}$ ,  $\beta_e \sim 0.3\%$ .
- Collisions ( $\nu_{ei} \sim 1 c_s/a$ ).
- ExB shear ( $\gamma_E \sim 0.13-0.16 c_s/a$ ), used fixed boundary conditions with  $\Delta^b = 1 \rho_s$  buffer widths.
- Resolution parameters
  - $L_x \times L_y = 6 \times 4 \rho_s$  ( $360 \times 240 \rho_e$ ).
  - $n_x \times n_y = 192 \times 48$ .
  - $k_\theta \rho_s$  [min, max] = [1.5, 73]
  - $k_r \rho_s$  [min, max] = [1.0, 50]
  - $[n_{||}, n_\lambda, n_e] = [14, 12, 12]$

# TRANSP comparison

shot 141767, channel 1

