

Plasmoids formation in a laboratory and large-volume flux closure during simulations of Coaxial Helicity Injection in NSTX-U

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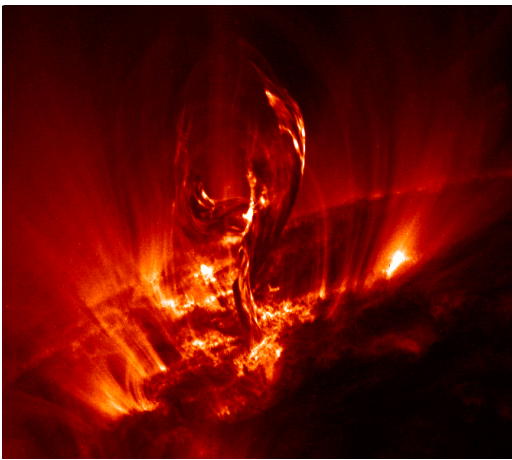
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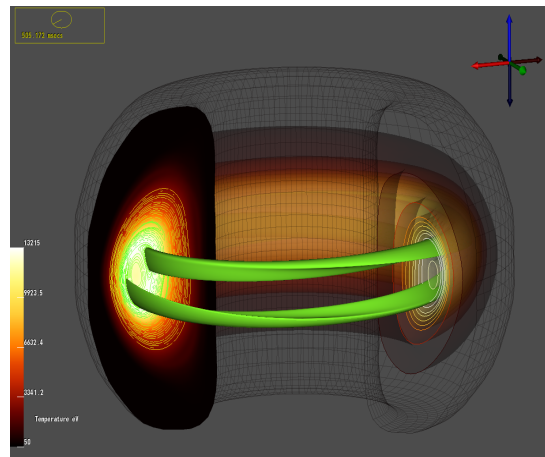


Reconnection plays an important role in the nonlinear dynamics of many processes in laboratory and astrophysical plasmas

Solar flares

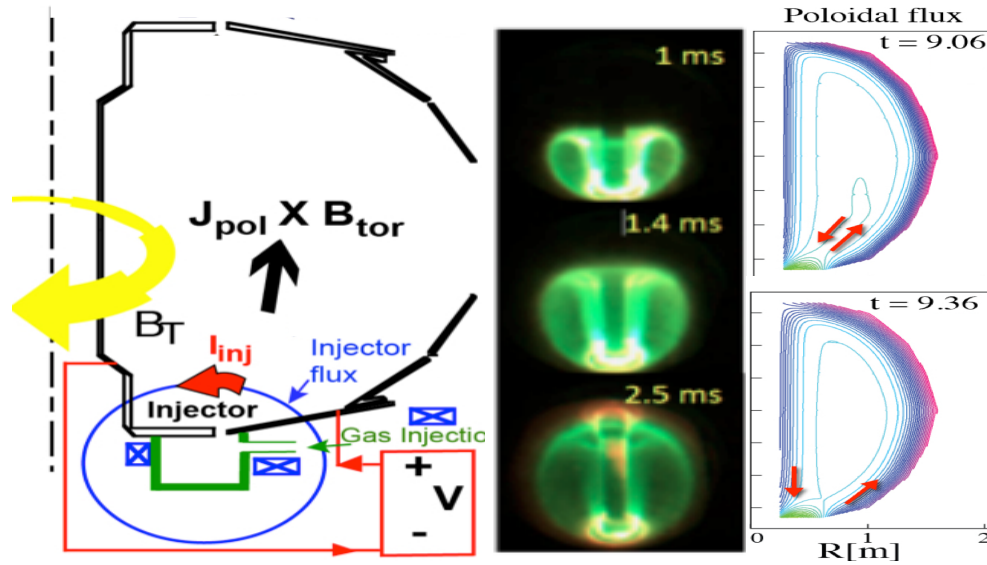


Fusion/lab plasmas



There are two types of trigger mechanism for magnetic reconnection
I- forced (driven) or II- spontaneous reconnection.

Many fundamentals of reconnection physics can be explored during helicity injection



Transient Coaxial Helicity Injection (CHI), the primary candidate for solenoid- free current start-up in the NSTX-U.

Outline

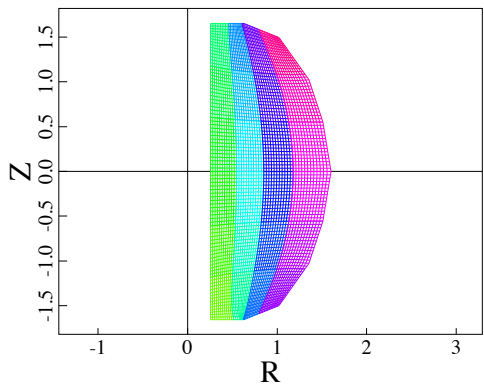
- Flux surface closure during helicity injection
 - Forced (driven) S-P reconnection
 - Spontaneous reconnection

- 3D effects
 - dynamo-driven reconnection plasmoids formation
 - Flux closure and CHI-generated current in NSTX-U

Nonlinear resistive MHD simulations using NIMROD code will be used to study reconnection physics.

The computational model (Nonlinear NIMROD simulations)

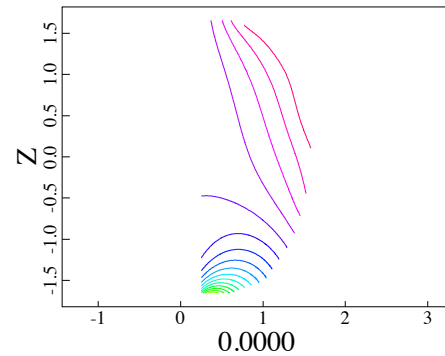
Finite Element Mesh



- Similar geometry to the experiment with a narrow slot.

- Poloidal grid 45×90 fifth order finite elements, 2-D ($n=0$) and 3-D (up to $n=22$ toroidal modes) simulations
- Voltage is applied across the injector gap (V_{inj})
- $E \times B$ normal flows at the gaps

Poloidal flux

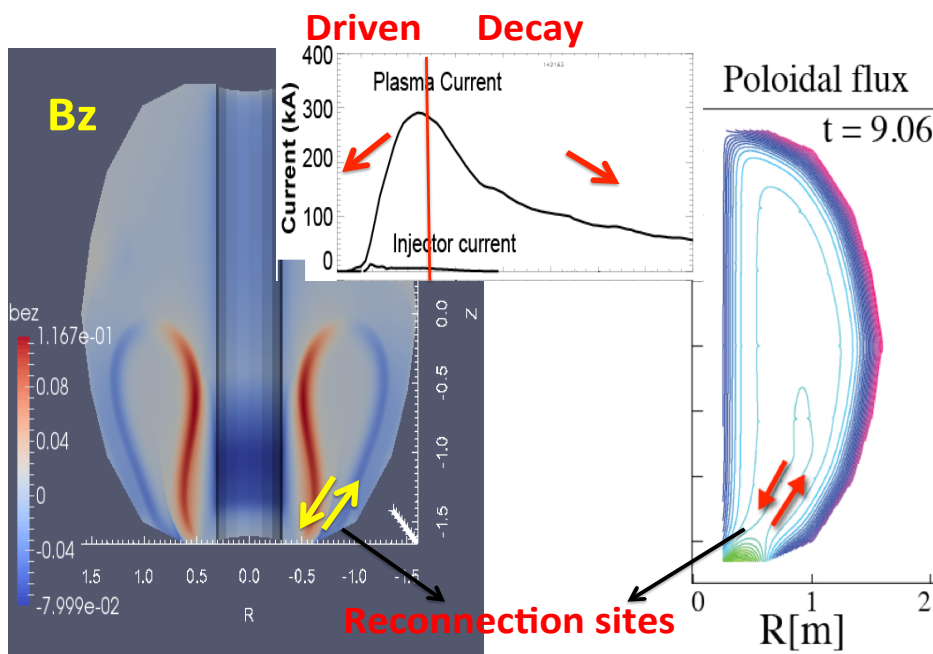


- Initial ψ_{inj} generated by including NSTX poloidal coil currents (with fixed boundary field)

Reconnection could occur during both stages of helicity injection

During injection (V_{inj} on)

During decay ($V_{inj} = 0$)



Outline

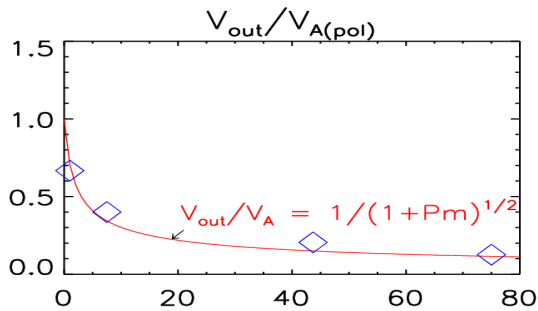
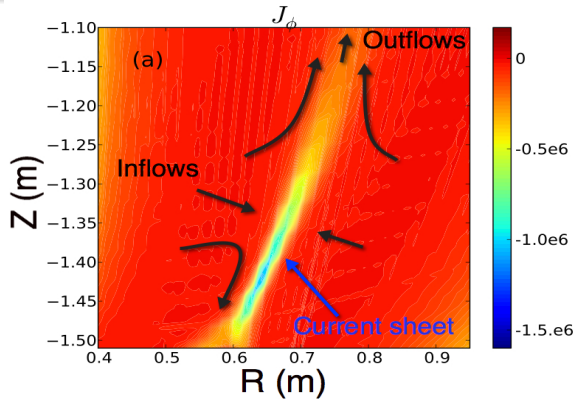
- How are the closed flux surfaces formed?
 - ▶ \implies Forced (driven) S-P reconnection

Outline

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 - ▶ \implies **Forced (driven) S-P reconnection**

How are the closed flux surfaces formed?

I - Forced S-P reconnection



A local 2-D Sweet-Parker type reconnection is triggered in the injection region. Key signatures:

- **I - Elongated current sheets, $L > \delta$.**
- **II - Scaling of the current sheet width**
 $\delta/L \sim (1 + P_m)^{1/4} S^{-1/2} \sim V_{in}/V_{out}$
- **III - Pinch inflow and Alfvénic outflow**

$S = L V_A / \eta$ (Alfvén velocity based on the reconnecting B, L is the current sheet length) F. Ebrahimi, et al. PoP 2013, 2014

Outline

- How are the closed flux surfaces formed?
 - ▶ **Forced (driven) S-P reconnection**
 - ▶ \implies **Spontaneous reconnection**
Could the elongated current sheet become unstable?

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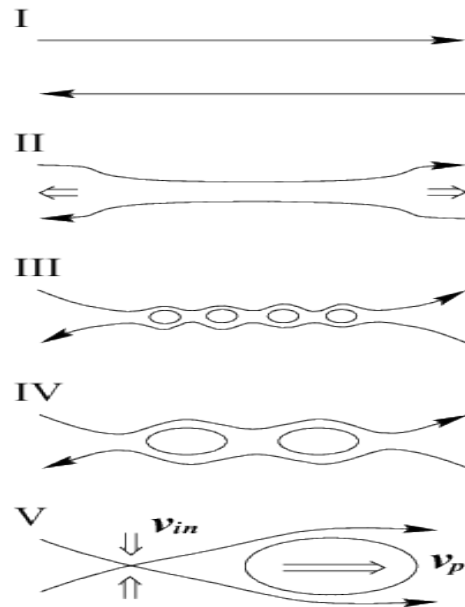
- ▶ **Forced (driven) S-P reconnection**

- ▶ \implies **Spontaneous reconnection**

Could the elongated current sheet become unstable?

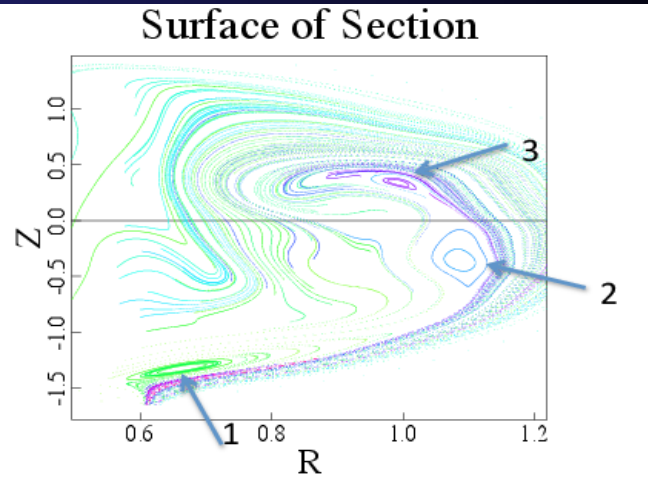
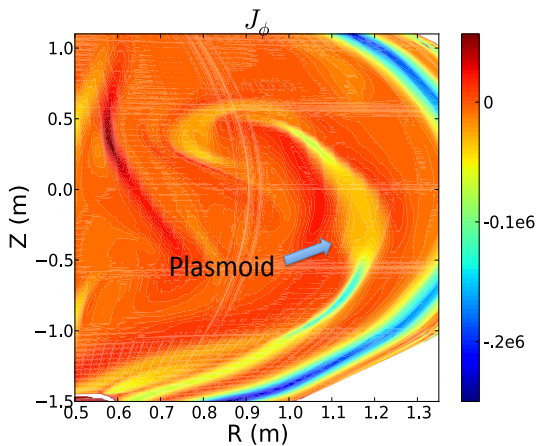
In many fast MHD dynamical processes, plasmoids are essential features

- Elongated current sheet can become tearing unstable at high S . [Biskamp 1986, Shibata & Tanuma 2001]
- The scaling properties of a classical linear tearing changes, ($\gamma \sim S^{1/4}$).
Numerical development: [Shibata & Tanuma 2001, Loureiro et al. 2007; Lapenta 2008; Daughton et al. 2009.; Bhattacharjee et al. 2009]
show fast reconnection.



Shibata & Tanuma 2001

How are the closed flux surfaces formed? II - Spontaneous reconnection

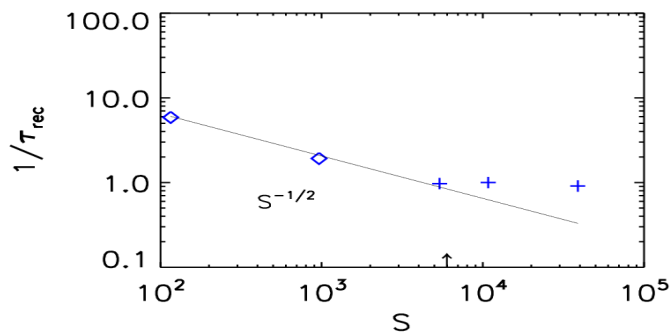
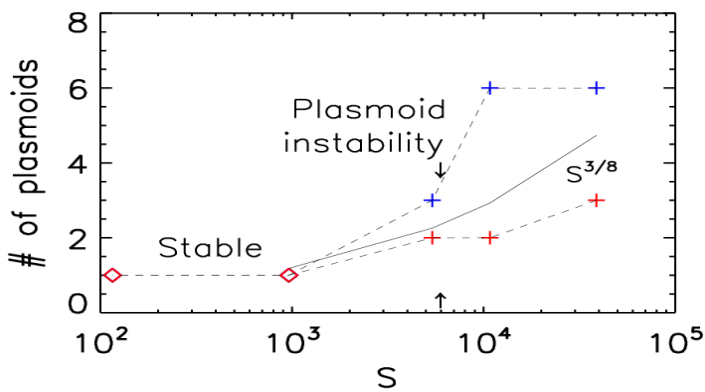


At high S , a transition to a plasmoid instability is demonstrated in the simulations.

Both small sized transient plasmoids and large system-size plasmoids are formed and co-exist. ($S=39000$)

Plasmoids merge to form closed flux surfaces.

At high S, a transition to a plasmoid instability occurs



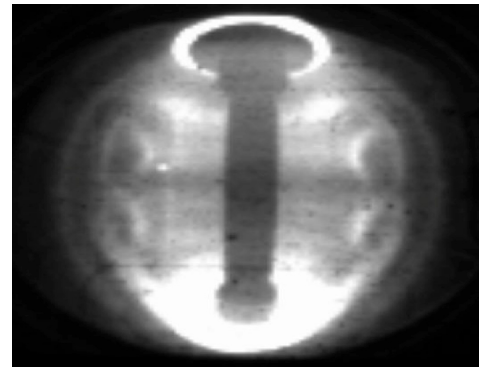
- Number of plasmoids is an increasing function of S. **Blue: small sized transient plasmoids** **Red: large scale and persistent plasmoids**
- As the current sheet evolves in time, static linear theory doesn't apply [L. Comisso et al. PoP 2016]

- Reconnection rate becomes nearly independent of S.

F. Ebrahimi and R.Raman PRL 2015

First documentation of plasmoid formation in MHD regime in laboratory.

(Loading poloidalflux.mp4)



Camera images from NSTX do show the formation plasmoids that then merge into a larger plasma

[Ebrahimi&Raman PRL 2015]

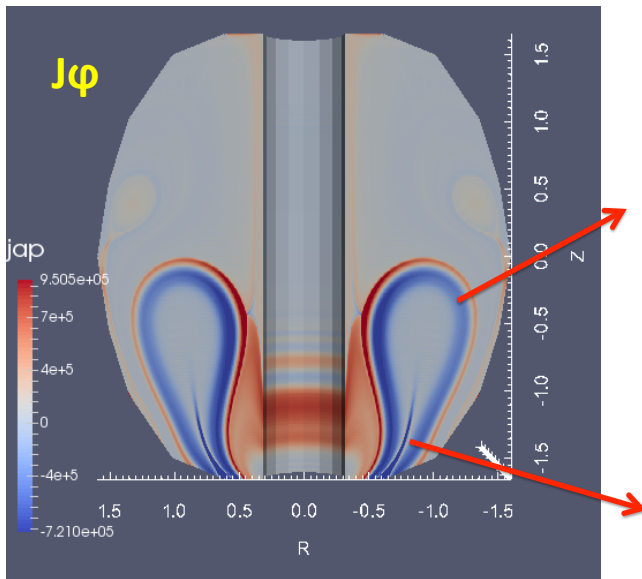
Outline

- ▶ 3D effects
 - ▶ **dynamo-driven reconnection plasmoids formation**
 - ▶ Could large-scale dynamo from 3-D fluctuations trigger reconnecting plasmoids?

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As helicity is injected, two types of current sheets are formed

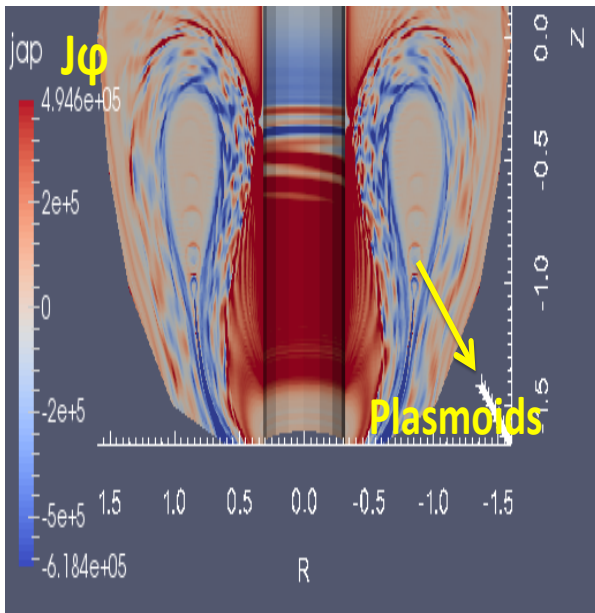


Consider an idealized case at increased helicity injection (increased poloidal flux)

- **1- Edge current sheet** from the poloidal flux compression near the plasma edge.
- **2- Primary reconnecting current sheet** from the oppositely directed field lines in the injector region.

Here in 2-D no plasmoid formation (stable $S < 10^4$)

Edge current-sheet instabilities are triggered in 3-D, and break the current-sheet



- 1 I- Edge-localized modes arising from the current-sheet instabilities
- 2 II- With 3-D fluctuations, axisymmetric plasmoids are formed, **local S increased to $S \sim 15000$.** [Ebrahimi submitted, arXiv:1610.09050]

I- Edge modes grow on the poloidal Alfvén time scales

1 These modes grow fast, on the poloidal Alfvén time scales, but they have tearing-parity structures.

2 These modes saturate by modifying and relaxing the edge current sheet $\gamma_{TA(n=1)} = 0.16$,

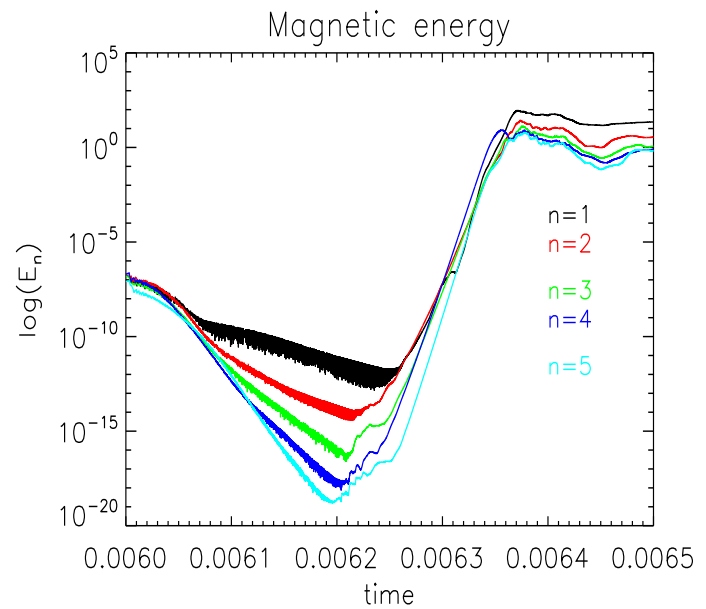
$$\gamma_{TA(n=2)} = 0.18,$$

$$\gamma_{TA(n=3)} = 0.2,$$

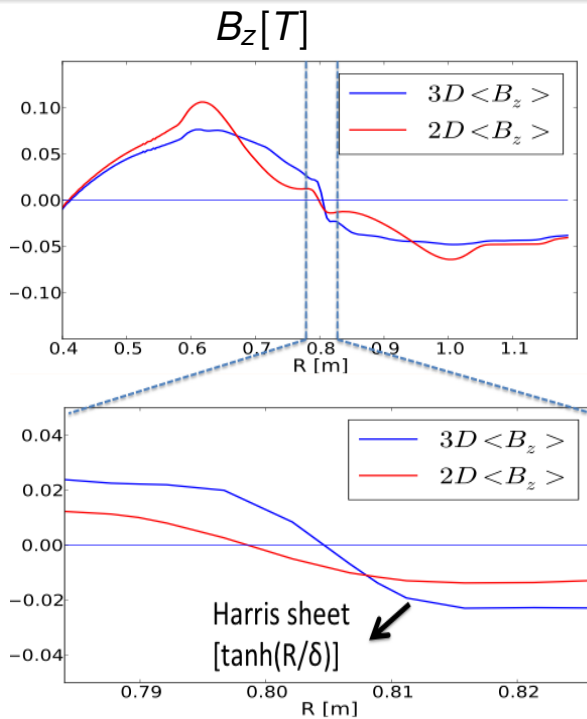
$$\gamma_{TA(n=4)} = 0.23,$$

$$\gamma_{TA(n=5)} = 0.26.$$

$$S = 2 \times 10^5$$



II- A 3-D dynamo poloidal flux amplification is observed to trigger axisymmetric reconnecting plasmoids formation



- For the first time a dynamo poloidal flux amplification is observed
- This **fluctuation-induced flux amplification increases the local S** \implies **triggers a plasmoid instability** \implies breaks the primary current sheet.

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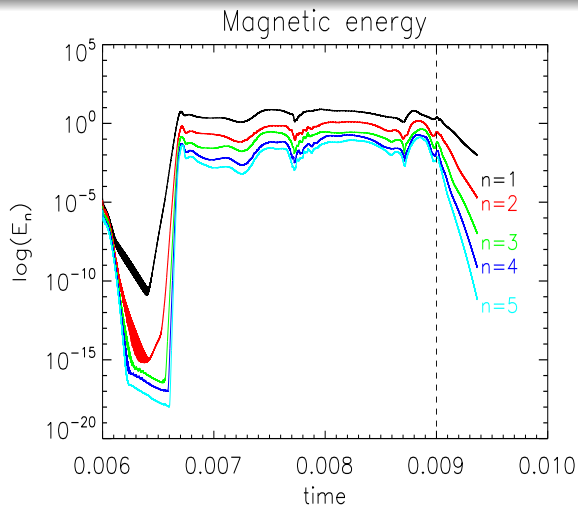
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 - ▶ \implies Flux closure and CHI-generated current in NSTX-U

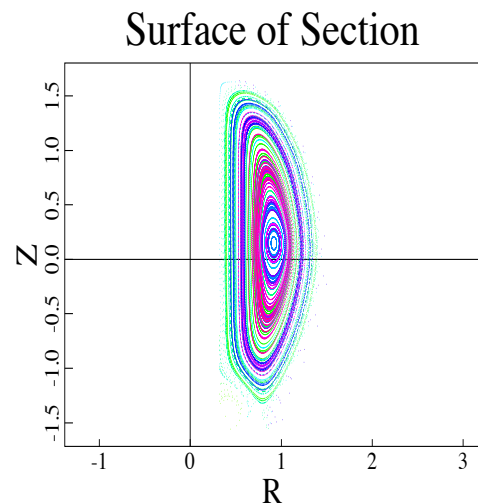
Outline

- Flux surface closure during helicity injection
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 - ▶ \implies **Flux closure and CHI-generated current in NSTX-U**

For NSTX-U relevant parameters, large volume flux closure is formed in the presence of 3-D fluctuations



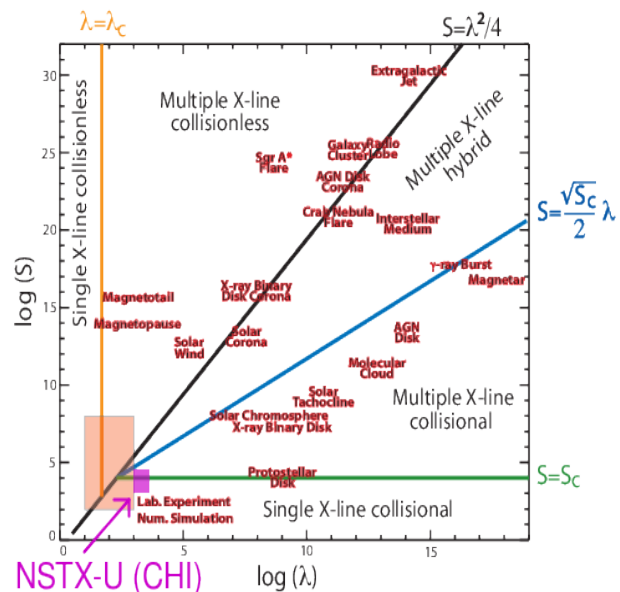
For NSTX-U relevant parameters, the nonaxisymmetric fluctuations saturate at low amplitudes (3-D). These are peeling-type modes with tearing parity



Similar to 2-D ([F. Ebrahimi & Raman NF Letter 2016]), nearly 70% of the injected open flux converts to closed flux in NSTX-U (3-D simulations)

Future directions: In addition to fusion applications, NSTX-U is a rich platform for fundamental reconnection studies

- How transient CHI scales as it is extrapolated to future (larger) devices, such as the ST-FNSF
- The implications of these new results of global plasmoids for 3-D fast reconnection and global tokamak MHD activity including disruption (VDEs and ELMs).



Phase diagram [H. Ji on fundamental reconnection studies]



Summary

I- Forced magnetic reconnection in NSTX/NSTX-U

- A local 2-D S-P reconnection is triggered. [F. Ebrahimi et al 2013, 2014]

II- Possibility of spontaneous reconnection

- a transition to plasmoid instability has for the first time been predicted by simulations in a large-scale toroidal fusion plasma. [Ebrahimi&Raman PRL 2015]
- Motivated by the simulations, experimental camera images exhibit the existence of reconnecting plasmoids in NSTX.

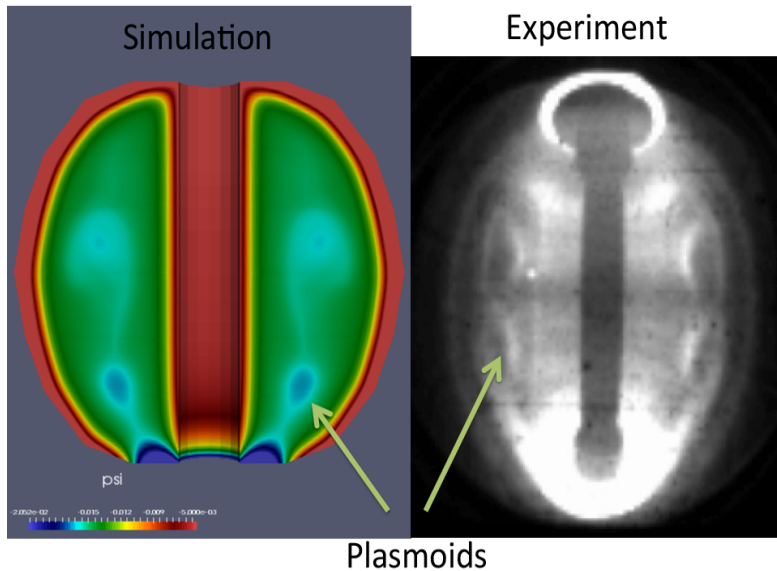
III- For the first time, a 3-D dynamo poloidal flux amplification is observed to trigger axisymmetric reconnecting plasmoids.) [Ebrahimi submitted, arXiv:1610.09050]

IV- Large-volume flux closure is obtained in the NSTX-U simulations in 2-D

[Ebrahimi&Raman NF 2016] **and in 3-D** [to be submitted]

extra slides

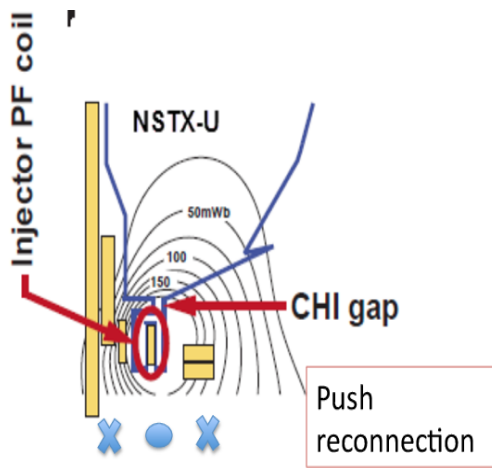
Camera images from NSTX do show the formation, and subsequent separation, of smaller plasmoids that then merge into a larger pre-existing plasma



- $T_e \sim 5eV$,
 $n \sim 4 - 5 \times 10^{18} m^{-3}$
- $B_{rec} \sim 100G$, $L \sim 1m$
- estimated local S
 $=2000-4000$.
- $V_p \sim 25$ km/s

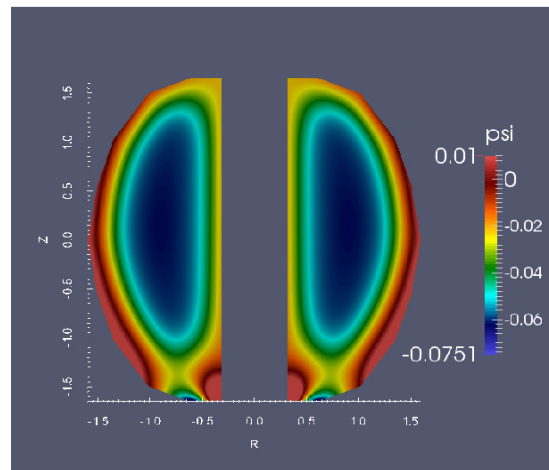
[Ebrahimi&Raman PRL 2015]

Flux closure increases with narrow flux footprints in MHD simulations of NSTX-U.



Because of improved positioning of injector flux and shaping coils in NSTX-U, the volume of flux closure is large.

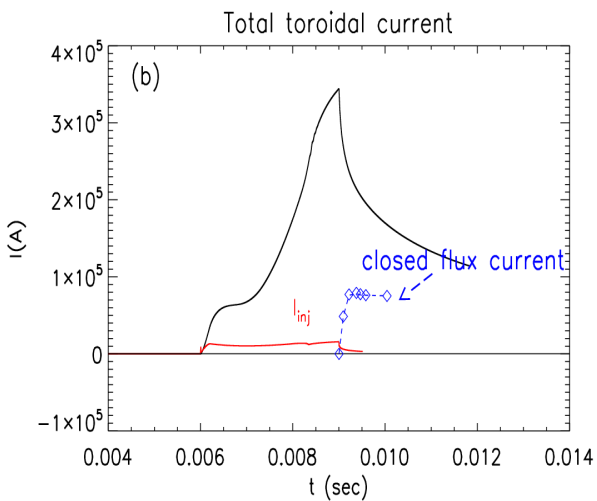
Narrow flux footprint W/ Flux Shaping coils



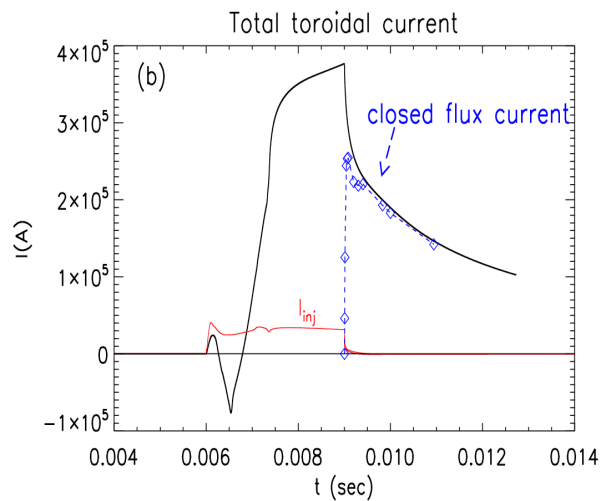
Poloidal flux within the last closed flux surface is 40 mWb and, is about 60-70% of the injector flux.

Nearly all of the CHI-generated current is closed-flux current.

Wide flux footprint



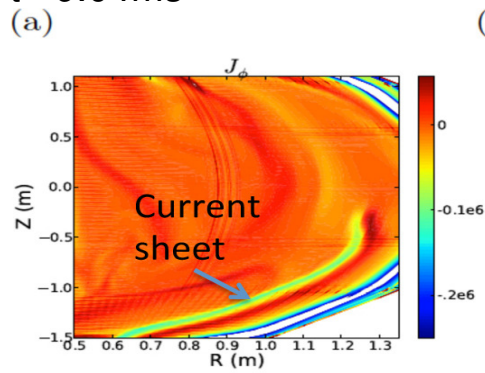
Narrow flux footprint



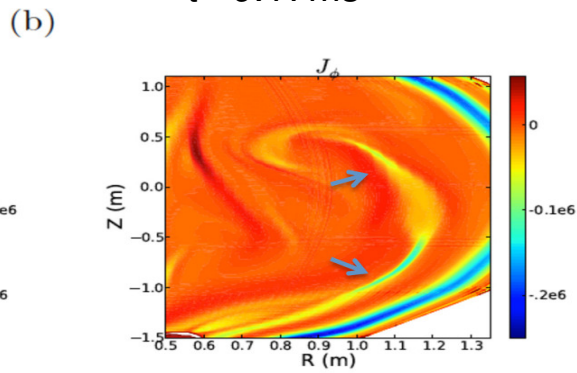
With narrow flux footprint almost all of the total current is in the closed flux region (with a large closed-flux current of about 240kA).

Both small sized transient plasmoids and large scale plasmoids are formed (S=39000)

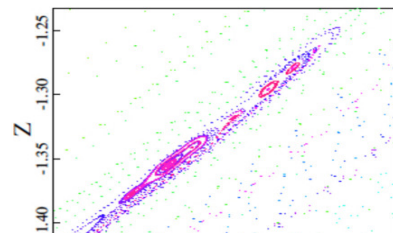
t= 9.04ms



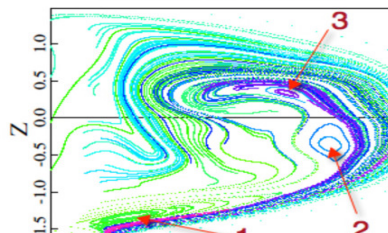
t= 9.17ms



(c) Surface of Section

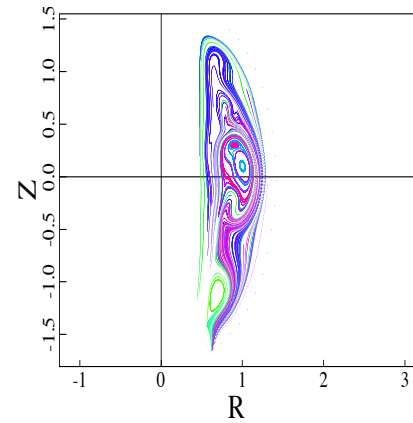


(d) Surface of Section

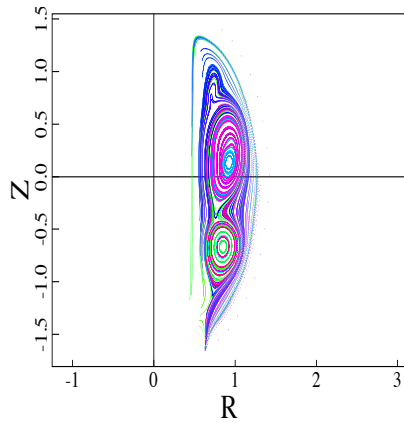


Large scale islands go through a dynamical process and merge to form closed flux surfaces

Surface of Section



Surface of Section



Surface of Section

