

# MHD stability of high- $\beta$ and long-pulse NSTX spherical torus plasmas

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for the NSTX Research Team

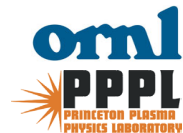
April Meeting of the American Physical Society

Saturday, April 5, 2003

Washington B, Loews Philadelphia Hotel  
Philadelphia, PA

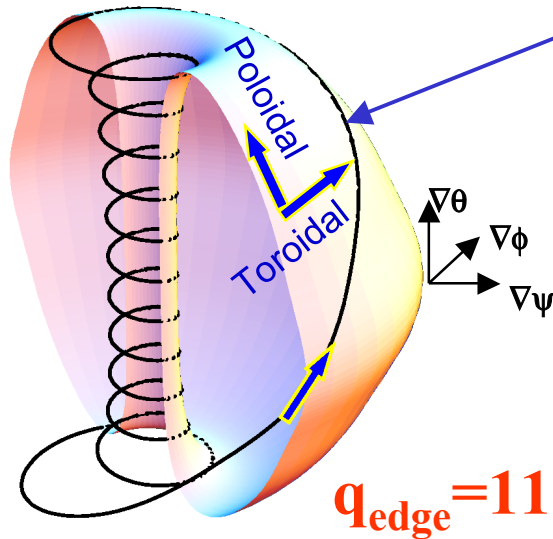


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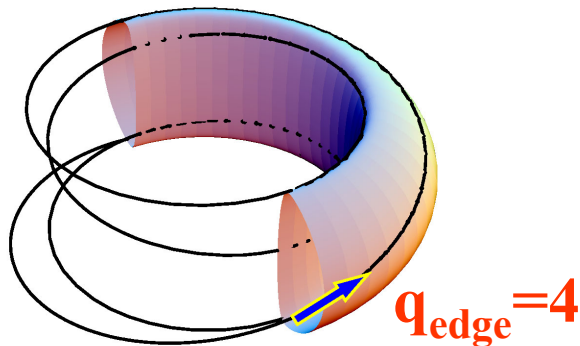
# How does low aspect ratio change stability?

## Spherical torus ( $A = R_0/a \leq 1.6$ )



- **Primarily through safety factor “q”**
  - $q$  is # *toroidal transits* / # *poloidal transits* of a magnetic field-line on a magnetic surface
  - **MHD instabilities try to satisfy  $\mathbf{k} \cdot \mathbf{B} = 0$** 
    - Instability wave-vector  $\mathbf{k} = n\nabla\phi - m\nabla\theta$
    - **Occurs on resonant rational  $q = m/n$  surface**
    - $n$  = toroidal mode number (integer)
    - $m$  = poloidal mode number (integer)
  - **High  $q$  and shear in  $q$  profile stabilizing**
    - Field-line bending energetically unfavorable

## Tokamak ( $A > 2.5$ )



- $q \propto (1 + \kappa^2) a B_T / I_p A$ 
  - **Higher  $q$  for given  $I_p / a B_T$  at low  $A$**
  - **Increased shear in  $q$  profile at low  $A$**

# MHD stability improved at low A



- Efficient reactor needs high toroidal  $\beta$

$\beta$  = kinetic pressure / magnetic pressure

Generating toroidal field (TF) is costly

- Theory and experiment show  $\longrightarrow$

–  $\text{MAX}(\beta_T) \propto I_p / a B_T$

–  $\beta_N \equiv \beta_T(\%) a B_T / I_p(\text{MA}) \leq C \approx 3-6$

- $\beta_N$  increases at low A

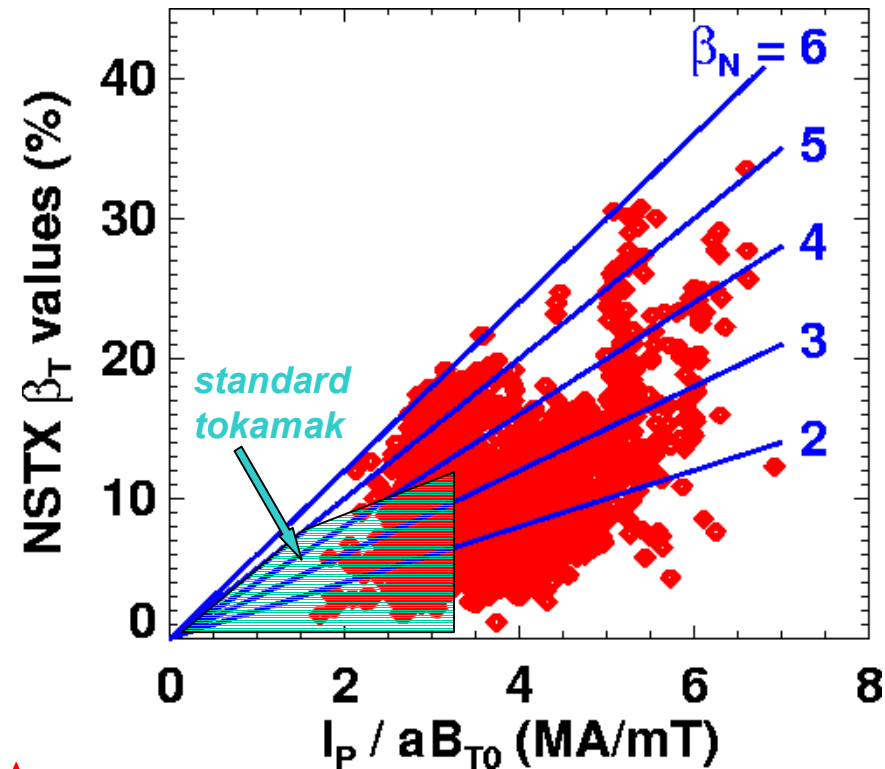
–  $\beta_N$  up to 6 at low A

–  $\beta_N = 3-4$  in standard tokamak

- $I_p / a B_T \propto (1 + \kappa^2) / A q \Rightarrow$

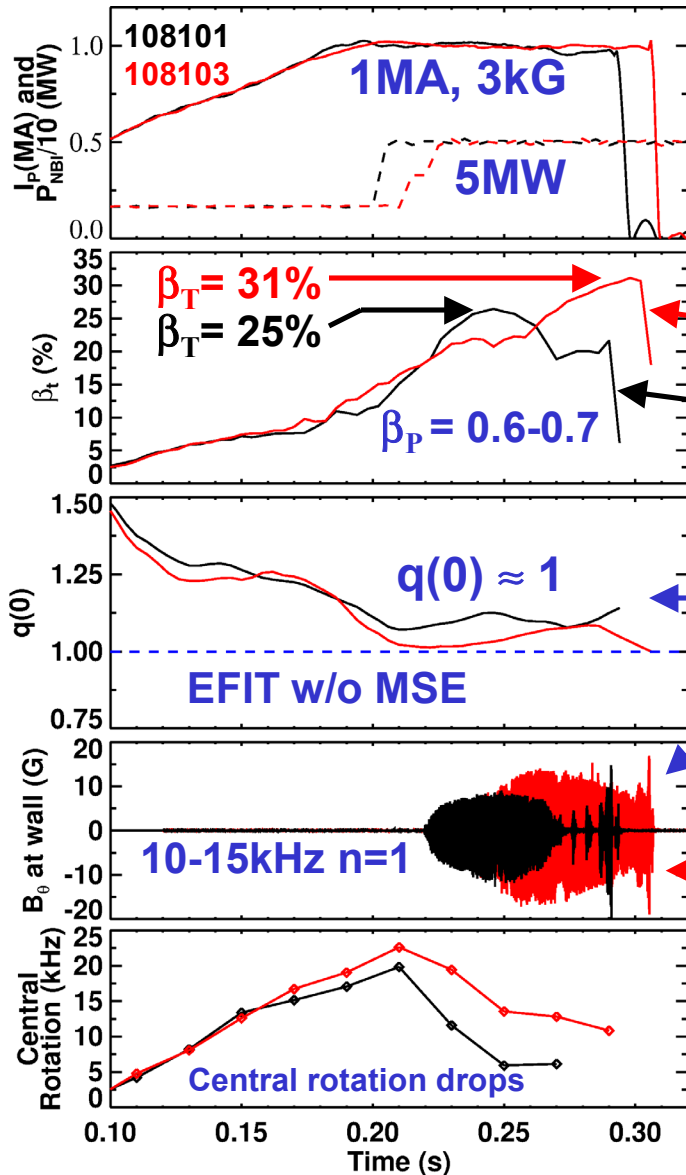
– Stable  $I_p / a B_T$  increased at low A

$$\beta_T \equiv 2\mu_0 \langle p \rangle / B_T^2$$



**Higher  $\beta_N$  and  $I_p / a B_T$  at low A result in  $\beta_T$  up to 35%**

# Highest $\beta_T$ discharges limited by m/n=1/1 modes



- $I_p=1\text{MA}$ ,  $B_T=0.3\text{T}$ ,  $P_{\text{NBI}}=5\text{MW}$ 
  - Both discharges terminate rapidly

- Before rapid termination....

Sometimes,  $\beta$  rises throughout discharge

Most times,  $\beta$  saturates, then drops

When  $q(0)$  is near 1 and  $\beta_T > 20\%$ ,  
10-15kHz n=1 instability appears

n=1 mode larger in high  $\beta$  shot (!)

How is drop in  $\beta$  avoided?

Difference appears to be *sustained rotation*

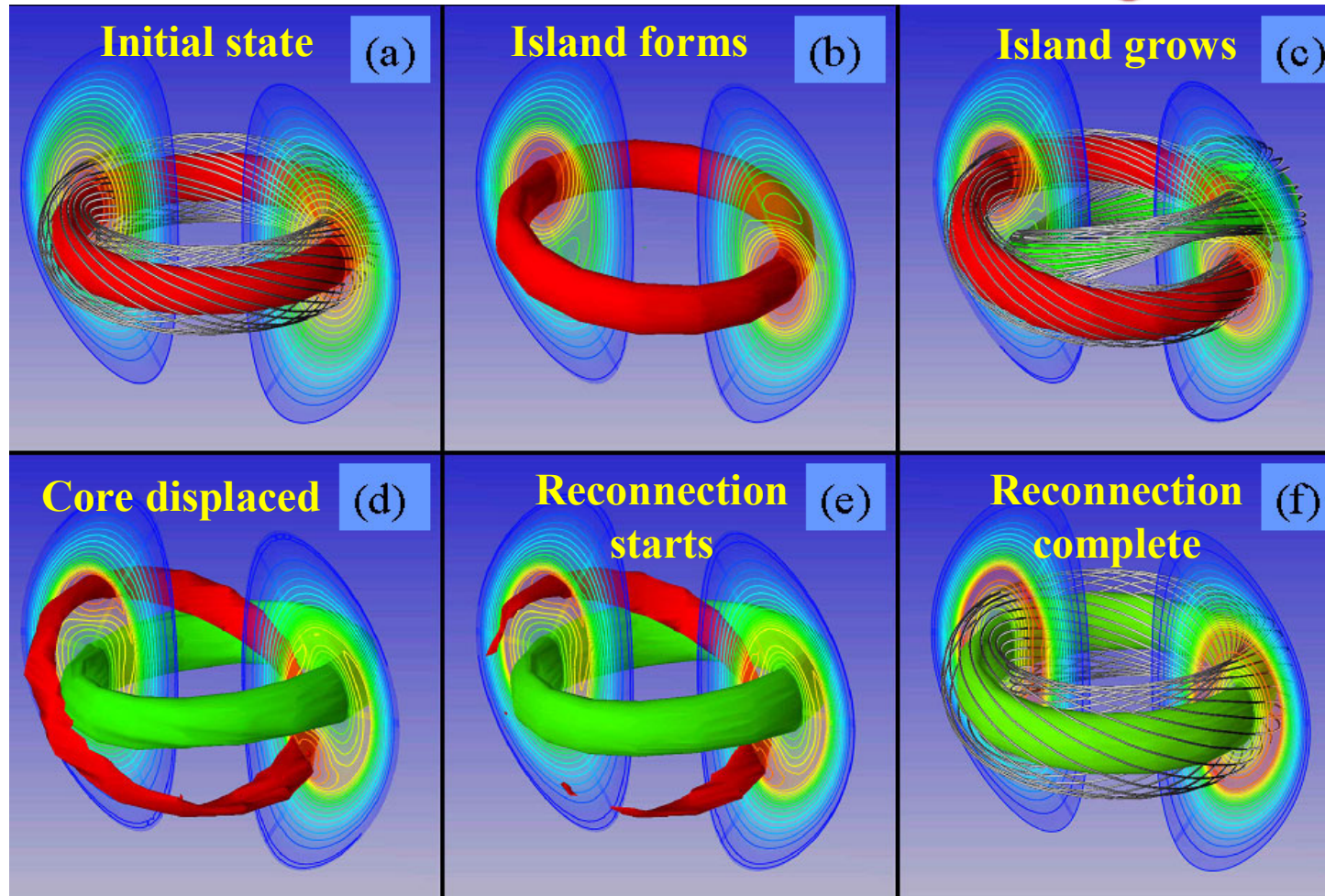
# Instability dynamics from non-linear simulations

(from Wonchull Park, M3D code, PPPL)



Simulation  
without  
rotation  $\Rightarrow$

B-field lines  
**Hot core**  
**Cold island**



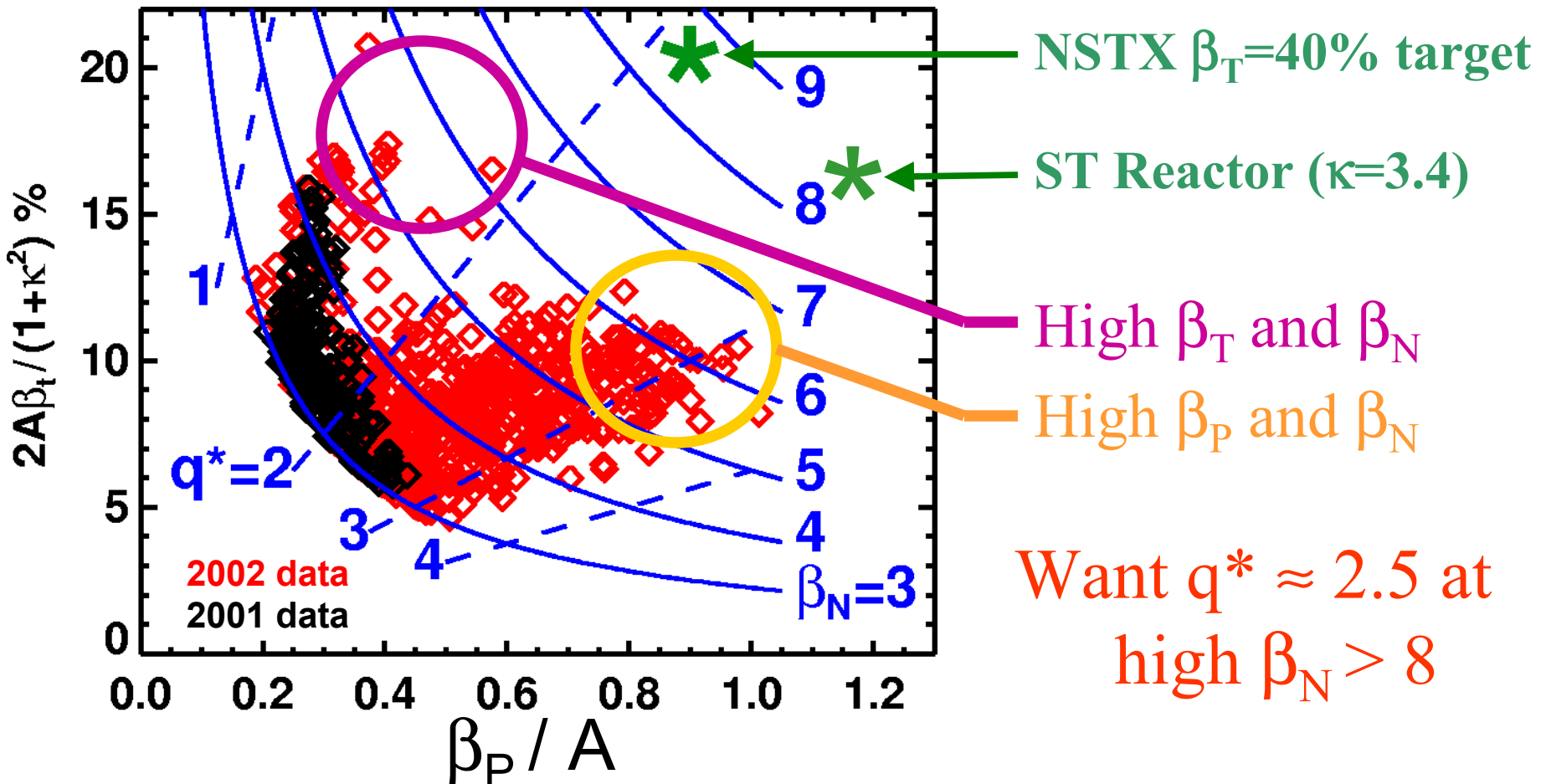
With sufficient rotational flow and shear, reconnection can be interrupted

**May explain long-lived 1/1 modes in high  $\beta_T$  NSTX discharges**

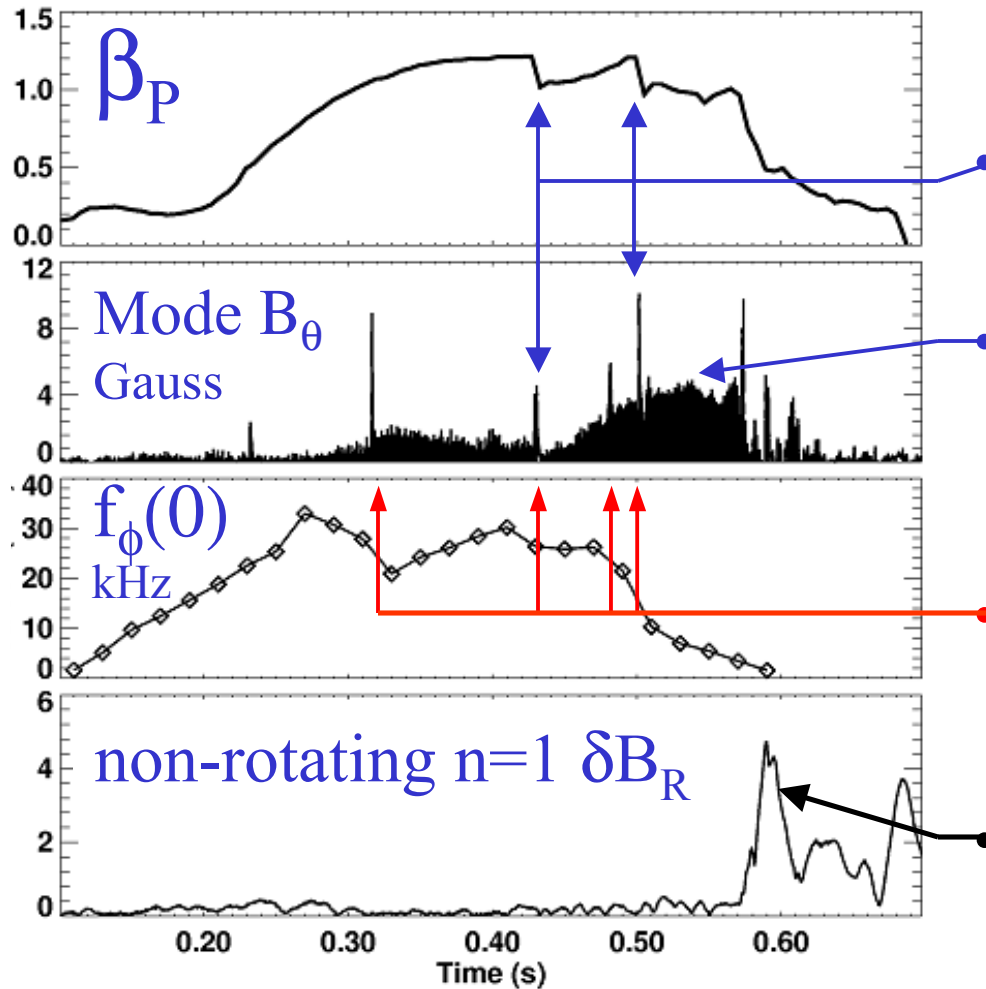
# Steady-state ST also requires high $\beta_P$



- Self-driven current fraction  $\propto \beta_P \equiv 2\mu_0 \langle p \rangle / B_P^2$
- $\beta_T \propto \beta_N^2 / \beta_P \Rightarrow$  Need very high  $\beta_N$  for steady state



# High $\beta_p$ discharges limited by “bursting” $n=1$



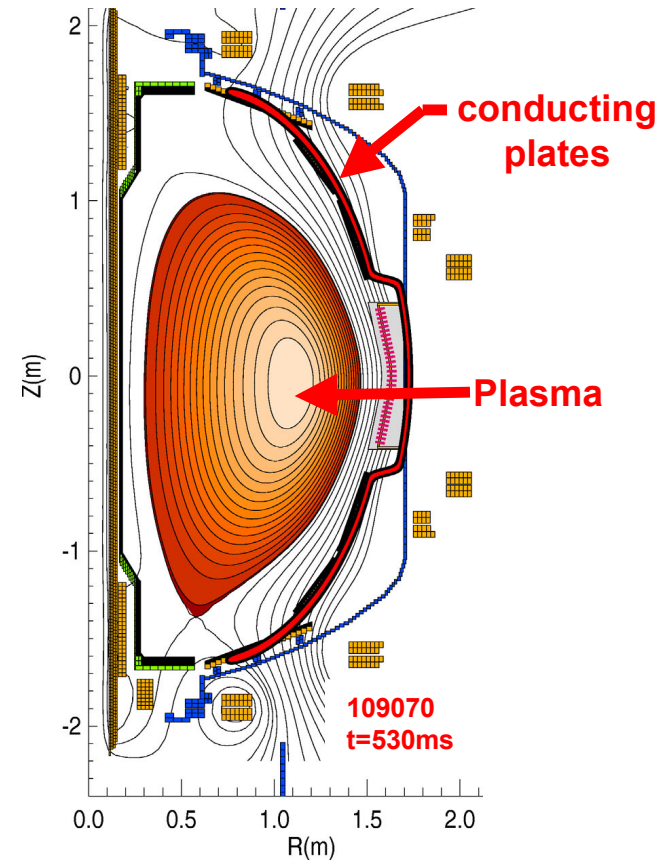
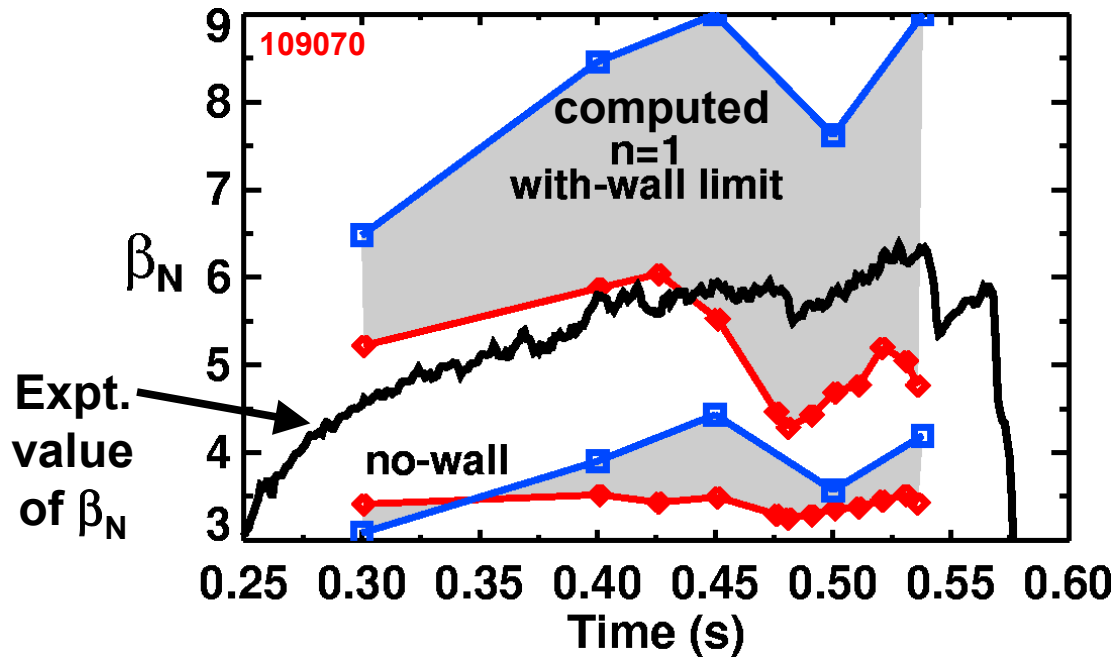
Rapid  $n=1$  bursts cause  $\beta$  drops  
–  $\beta$  can recover between bursts

Continuous modes degrade  $\beta$ ?

Each  $n=1$  burst reduces rotation  
– Also triggers continuous modes?

**Non-rotating**  $n=1$  becomes unstable once rotation is low  
– Causes final collapse of plasma  $\beta$

# High $\beta_p$ shots operate above no-wall limit



Theory and other experiments (DIII-D):  
 $\Rightarrow$  resistive wall and plasma rotation  
can stabilize “resistive wall mode”

*NSTX high  $\beta_p$  shots may be hitting with-wall limit*



# Stability increases at low A



- Low  $A \Rightarrow$  high  $I_p / a B_T$  and  $\beta_N \Rightarrow$  high  $\beta_T \leq 35\%$ 
  - High  $\beta_T$  discharges limited by *long-lived*  $n=1$  modes
- High  $\beta_p$  and high  $\beta_N$  needed for steady-state ST
  - High  $\beta_p$  discharges limited by *bursting*  $n=1$  modes
- Highest  $\beta_p$  discharges are above stability limits w/o wall
  - Rotational stabilization of resistive wall mode