

# Full wave simulations for fast wave heating and power losses in the scrape off layer of tokamak plasmas

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## Introduction

### Experimental:

- In NSTX, up to ~ 60% of the power possibly lost to SOL
  - Larger SOL losses for high density in front of the antenna
- Field-line mapping models flow of lost HHFW power from the antenna region to the divertor [Perkins *et al.*, PRL 2012]
  - Computed strike points form a spiral that matches the observed RF spiral
- In DIII-D, observed increase in edge losses as the SOL density increases in ELMy H-mode plasmas [Petty *et al.*, NF 39, 1421 (1999)]

### Modelling:

- SOL region now included in full wave code AORSA [Green *et al.*, PRL 2011]
- Simulation results show:
  - Indication of cavity mode in AORSA simulations on NSTX
  - Negligible power absorbed in the SOL with only standard kinetic model

**Need to understand and minimize these RF power losses for improving FW performance and compare different devices with different geometry and heating regime**

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## Phys. mech. behind the damping in the SOL not yet been identified

- "Collisional" damping included in AORSA as a **proxy** to represent the real SOL damping processes

- It is essentially a "Krook model":

$$\text{dielectric tensor} \quad \text{related to} \quad Z \left( \frac{\omega - n\Omega_c}{k_{\parallel} v_{th}} \right) \xrightarrow{\text{Damping}} Z \left( \frac{\omega - n\Omega_c + i\nu}{k_{\parallel} v_{th}} \right)$$

Plasma Dispersion Function

- $\nu/\omega$  "collision" parameter is input into AORSA

- absorption obtained through the anti-hermitian part of

$$\alpha = \frac{\omega}{4\pi} \frac{\mathbf{E}^* \cdot \mathbf{e}^{\text{a}} \cdot \mathbf{E}}{|\mathbf{S}|}$$

Poynting vector

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## Larger evanescent region at higher $B_T$ achievable in NSTX-U

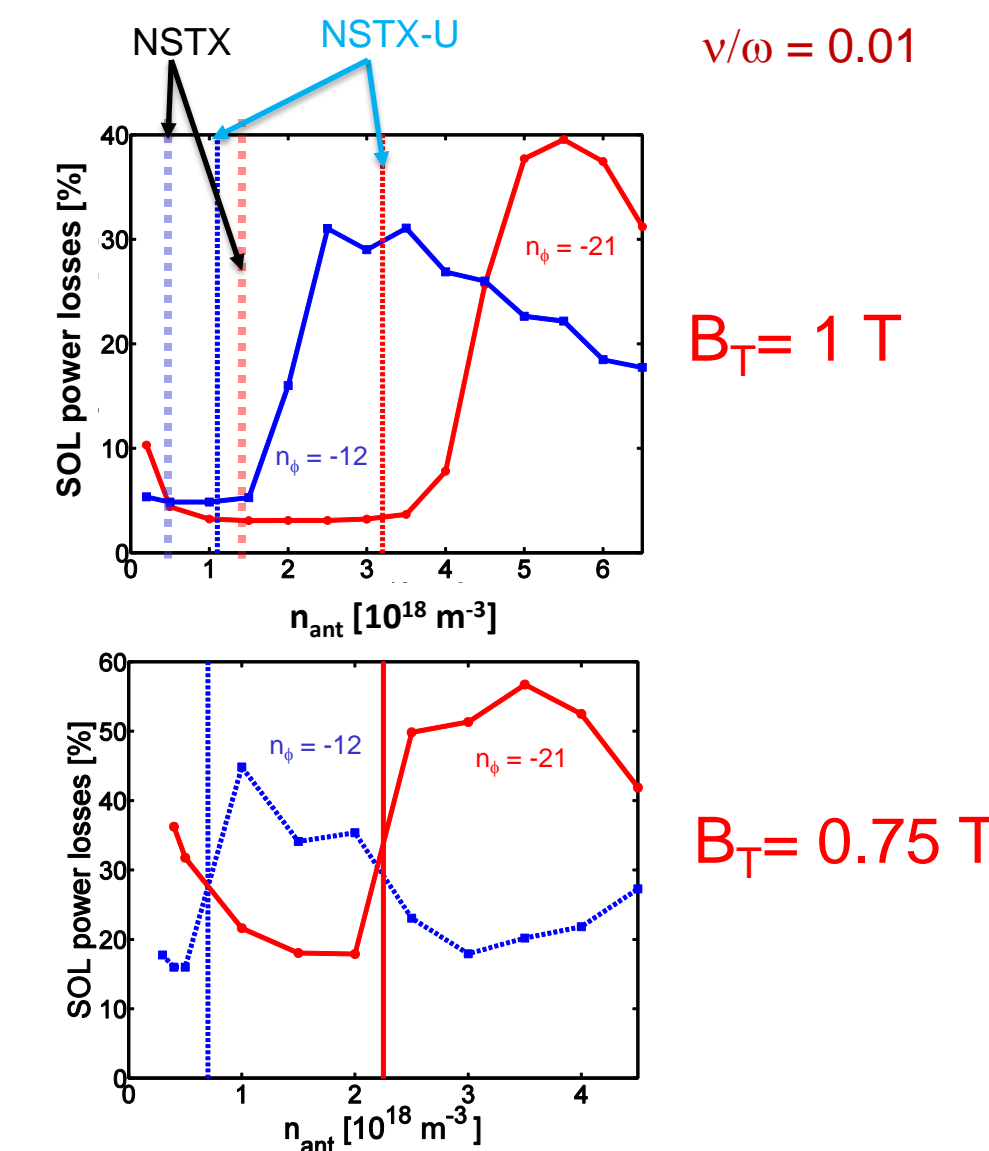
- H-mode scenario planned for NSTX-U, obtained by TRANSP simulations [S. P. Gerhardt, Nucl. Fusion 52, 083020 (2012)]

$$n_{e,FW\text{cut-off}} \propto \frac{k_{\parallel}^2 B}{\omega}$$

- NSTX-U ( $B_T = 1$  T) vs. NSTX ( $B_T = 0.55$  T): transition occurs for ~2x higher in  $n_{ant}$ 
  - wider evanescent region with lower SOL power losses
  - favorable for future experiments

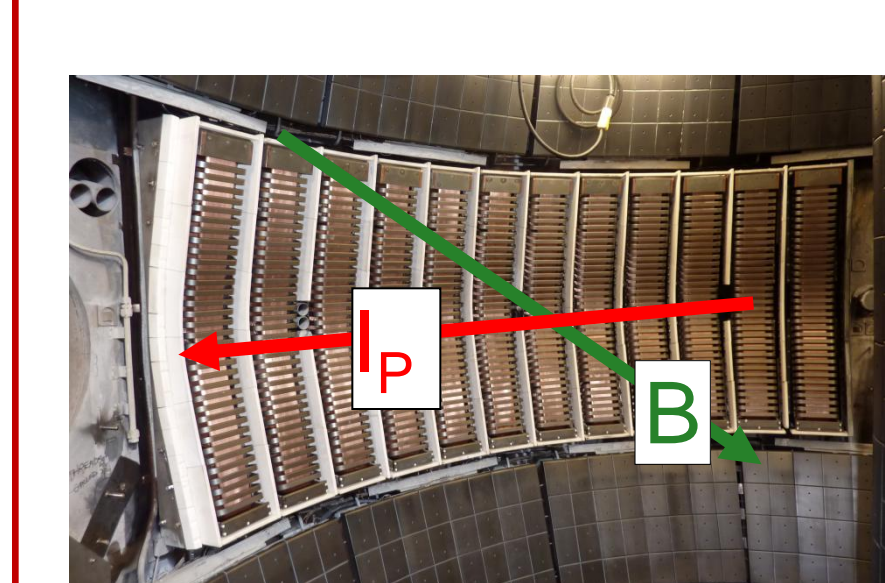
- NSTX-U: optimization of edge density to minimize the SOL losses and maximize the coupling to the core

- Same relative behavior found for NSTX-U case with  $B = 0.76$  T

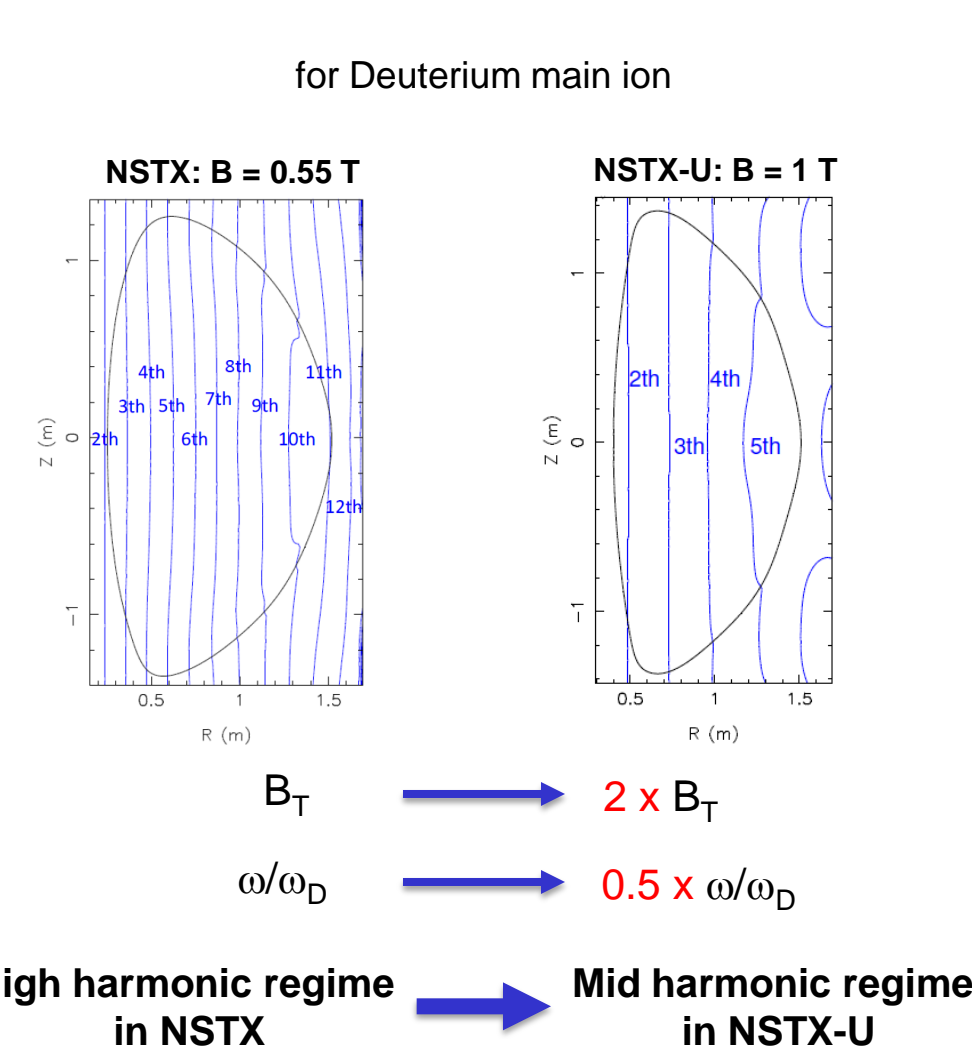


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## Same 30 MHz FW system in NSTX & NSTX-U

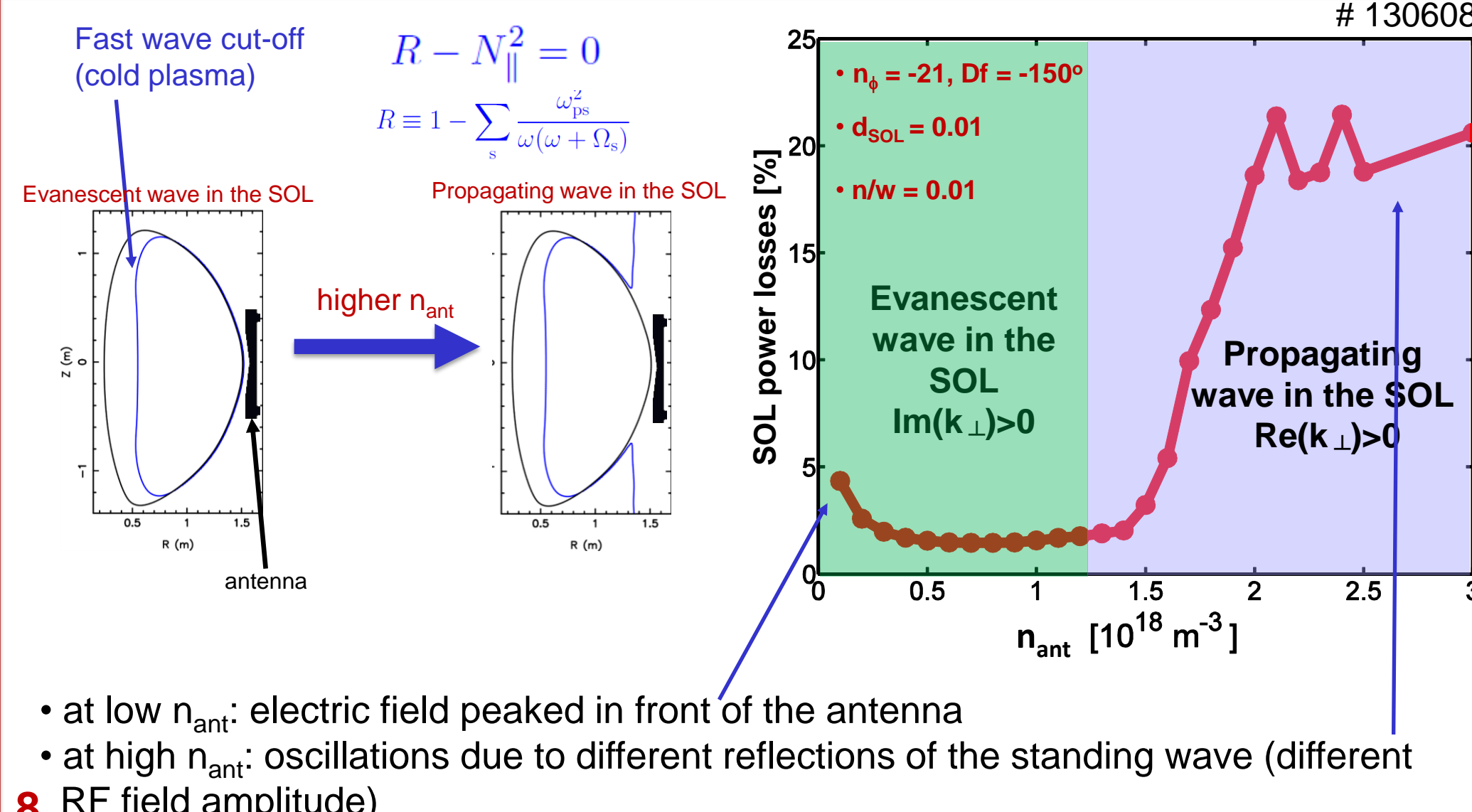


- 12-strap antenna on the outboard midplane, extends 90° toroidally
- Well-defined spectrum
- $|k_{\parallel}| = \pm 3, \pm 8, \text{ and } \pm 13 \text{ m}^{-1}$  when  $\Delta\phi = \pm 30^\circ, \pm 90^\circ, \text{ and } \pm 150^\circ$



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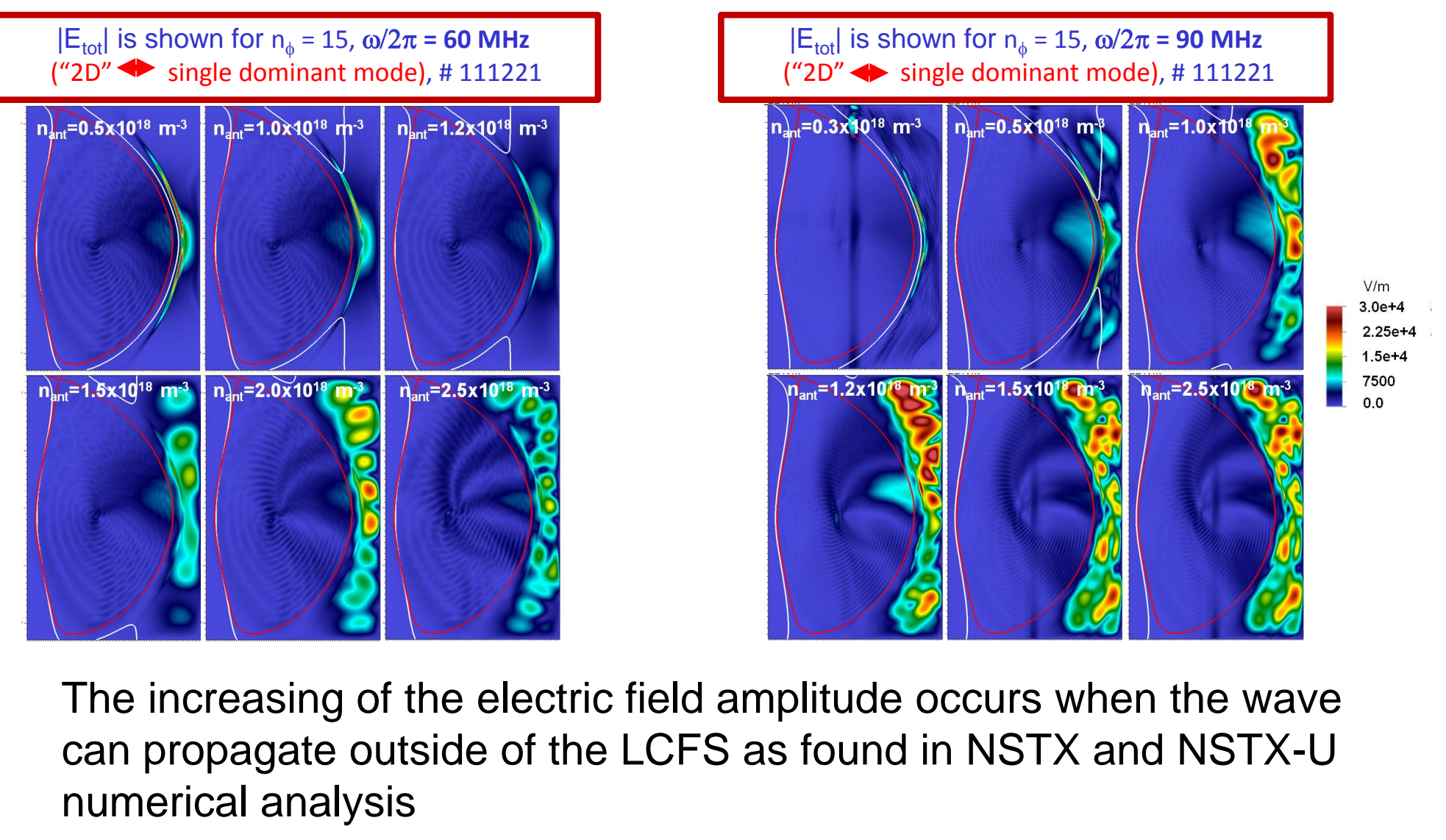
## Sharp increase in SOL pow. loss. when FW propagates within SOL



- at low  $n_{ant}$ : electric field peaked in front of the antenna
- at high  $n_{ant}$ : oscillations due to different reflections of the standing wave (different RF field amplitude)

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## DIII-D: $|E_{\text{sol}}|$ in the SOL increases when FW propagates within SOL



The increasing of the electric field amplitude occurs when the wave can propagate outside of the LCFS as found in NSTX and NSTX-U numerical analysis

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## Full wave code AORSA

Wave fields modeled with linear wave equation (inhomogeneous plasma):

$$\nabla \times \nabla \times \mathbf{E} - \frac{\omega^2}{c^2} \left( \mathbf{E} + \frac{4\pi i}{\omega} \mathbf{J}_p \right) = \frac{4\pi i}{\omega} \mathbf{J}_A$$

where  $\mathbf{e}$  is the dielectric tensor,  $\mathbf{J}_A$  is the antenna current

$$\mathbf{J}_p(\mathbf{r}) = \int d\mathbf{r}' \sigma(\mathbf{r}, \mathbf{r}') \cdot \mathbf{E}(\mathbf{r}')$$

- AORSA includes the complete non-local, integral operator for  $\mathbf{e}$  that is valid for "all orders" ( $k_{\perp} \rho_s > 1$ ) [E. F. Jaeger *et al.*, Phys. Plasmas 9, 1873 (2002)]

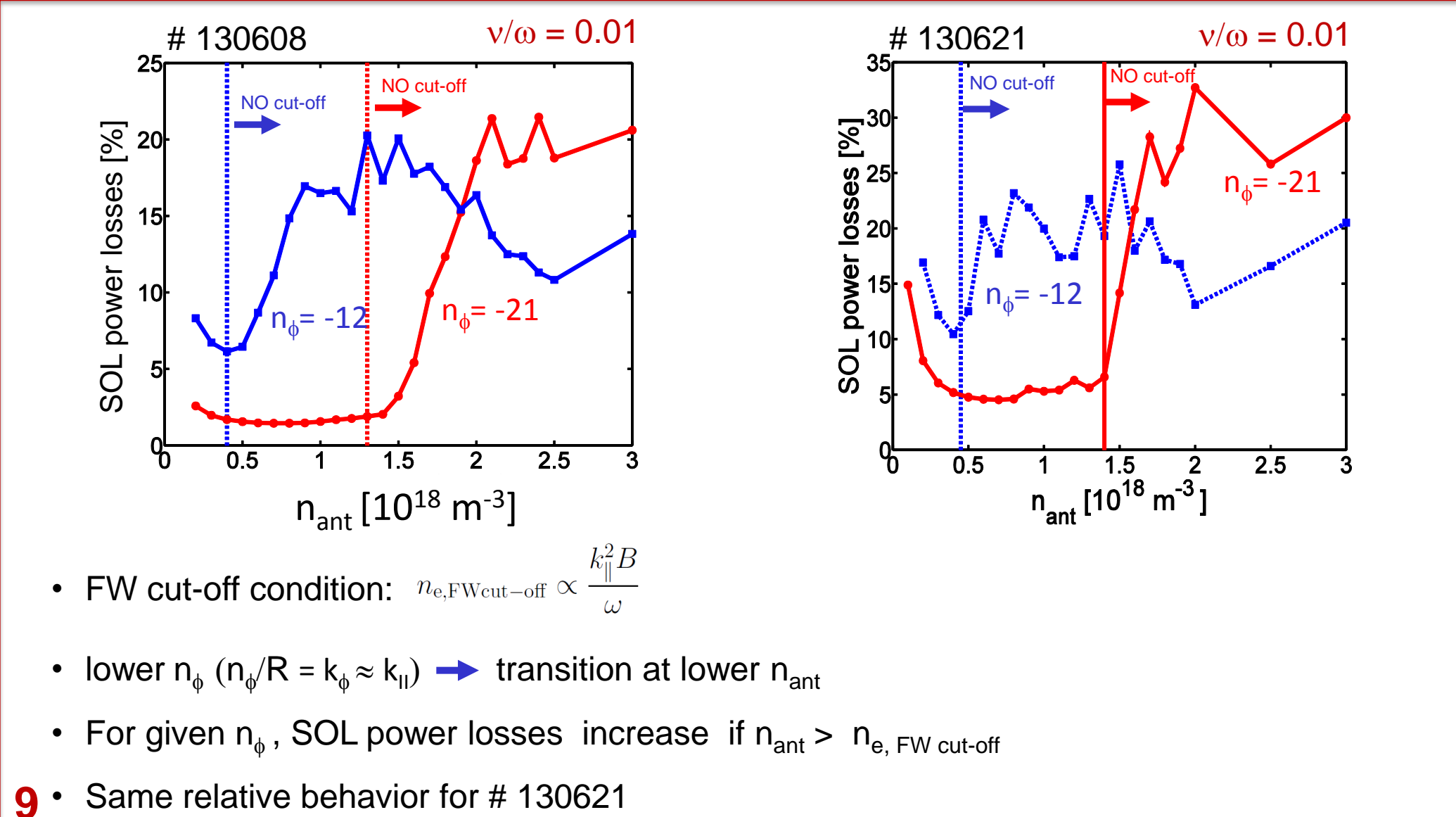
- AORSA utilizes a local Cartesian grid in the poloidal plane and a Fourier decomposition in the toroidal direction of symmetry:

$$\mathbf{E}(x, y, \phi) = \sum_{n, \phi} \sum_{m, \phi} \mathbf{E}_{n, \phi, m, \phi} e^{i(k_{\parallel} x + k_{\perp} y)}$$

where  $x = R - R_0$  and  $y$  is the distance from the midplane

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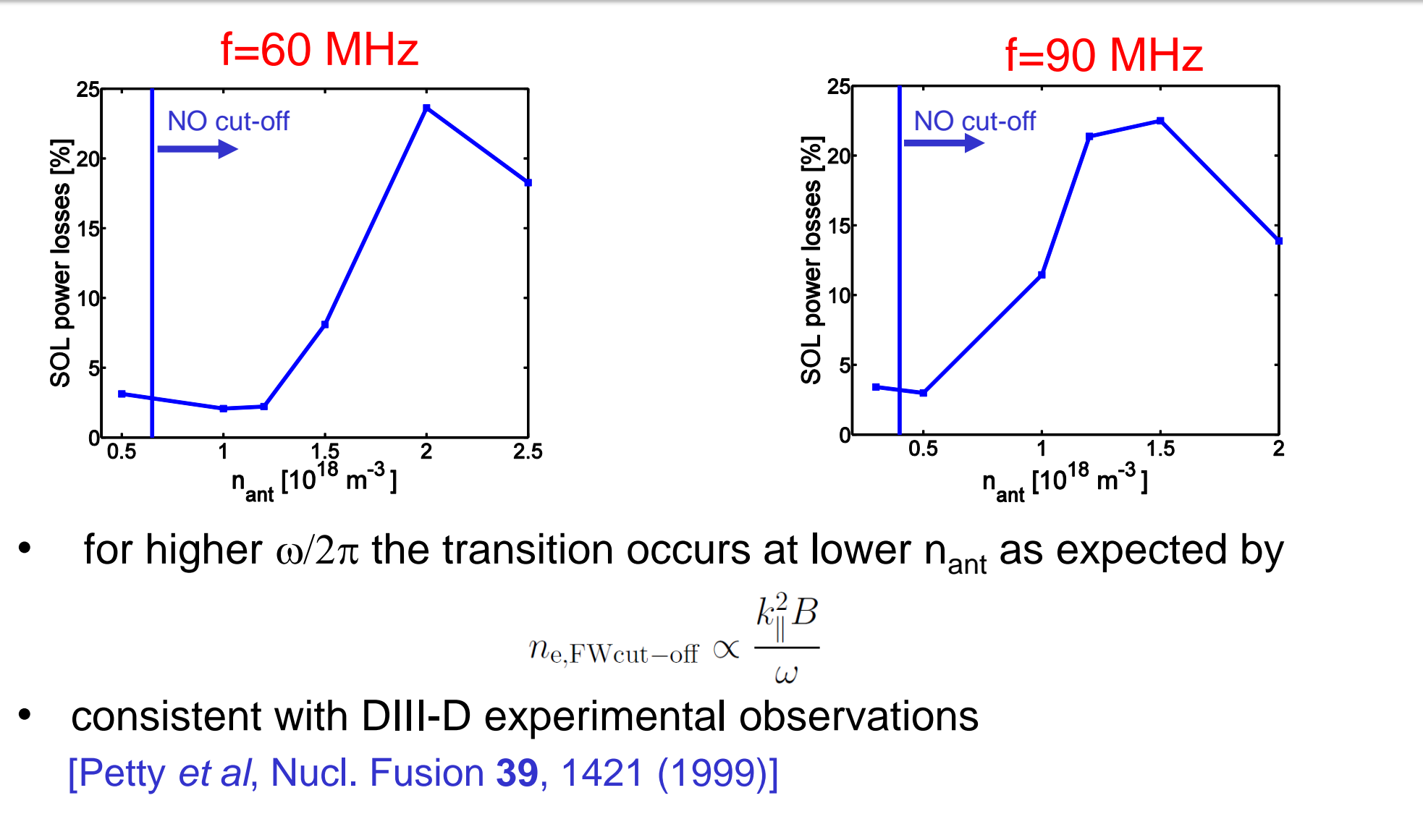
## Same transition is found for lower antenna phases



- FW cut-off condition:  $n_{e,FW\text{cut-off}} \propto \frac{k_{\parallel}^2 B}{\omega}$
- lower  $n_b$  ( $n_b/R = k_{\parallel} \approx k_{\perp}$ )  $\rightarrow$  transition at lower  $n_{ant}$
- For given  $n_b$ , SOL power losses increase if  $n_{ant} > n_{e,FW\text{cut-off}}$

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## Sharp increase in SOL pow. loss. in DIII-D as found in NSTX/NSTX-U



- for higher  $\omega/2\pi$  the transition occurs at lower  $n_{ant}$  as expected by  $n_{e,FW\text{cut-off}} \propto \frac{k_{\parallel}^2 B}{\omega}$
- consistent with DIII-D experimental observations [Petty *et al.*, Nucl. Fusion 39, 1421 (1999)]

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## Exponential decay for the SOL density to fit the exp. density data

- Extended AORSA to include SOL plasma:

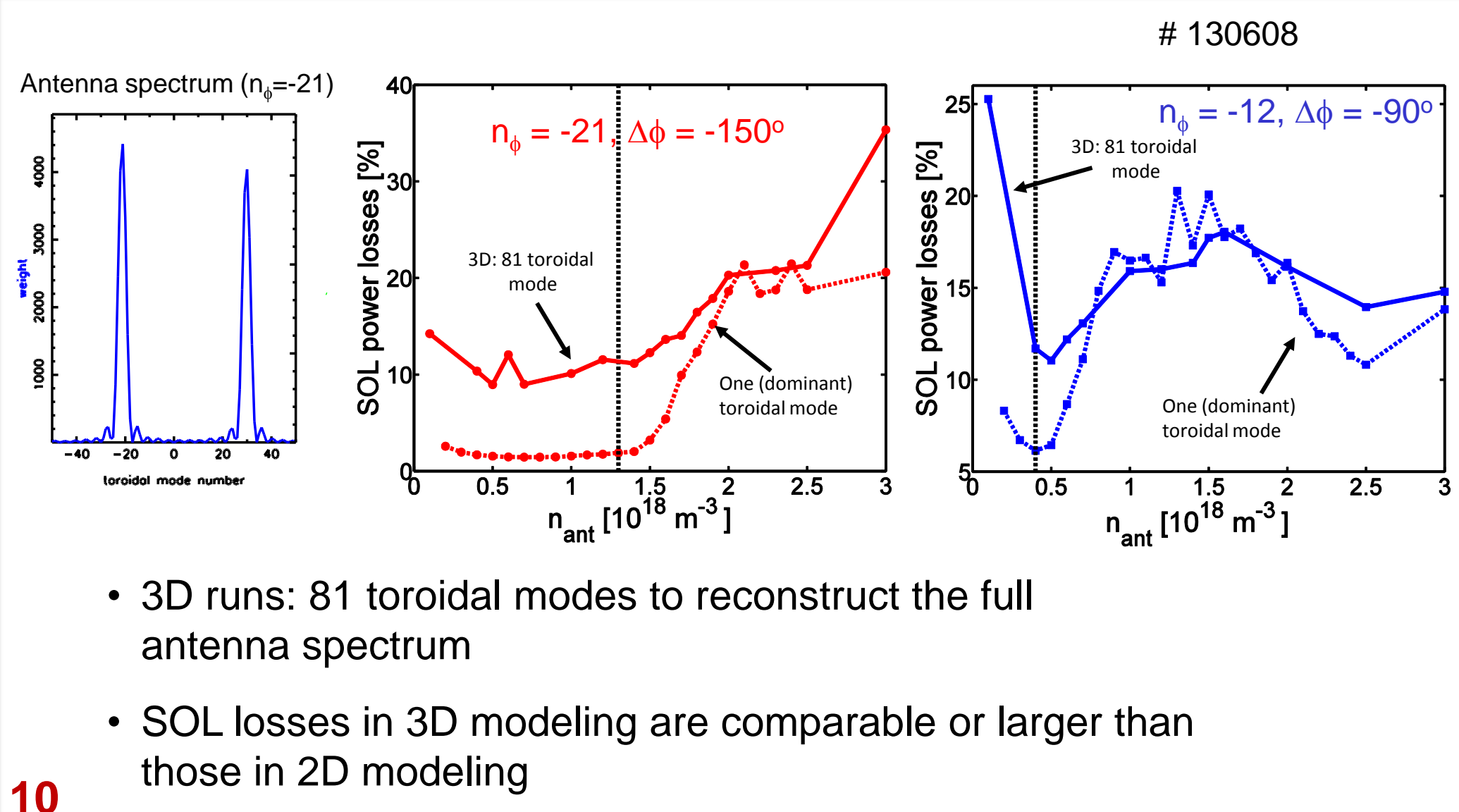
$$n_e(\rho) = n_{e,ant} + [n_e(\rho=1) - n_{e,ant}] \exp\left[-\frac{\rho-1}{d_{SOL}}\right], \quad \rho > 1$$

density at the separatrix      density in front of the antenna      decay length

- Density at separatrix is set by experimental measurements
- $T_e$  and  $T_i$  assumed to be constant across the SOL
- Will vary values of the density in front of the antenna
- Will vary decay length in order to cover the range of possible density profiles

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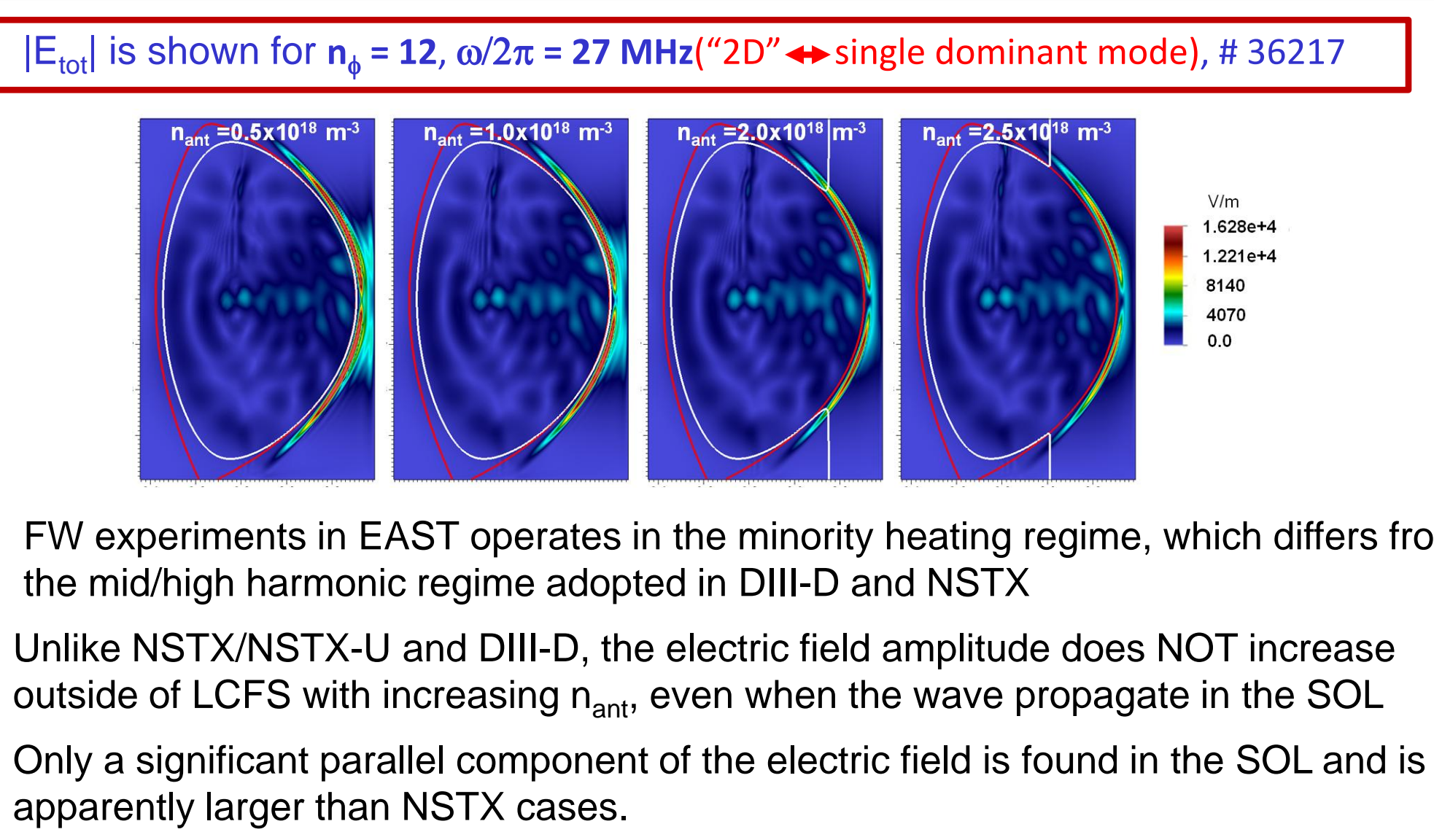
## 3D results show SOL losses enhancement similar to 2D single $n_b$



- 3D runs: 81 toroidal modes to reconstruct the full antenna spectrum
- SOL losses in 3D modeling are comparable or larger than those in 2D modeling

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## NO large electric field amplitude in SOL for EAST case

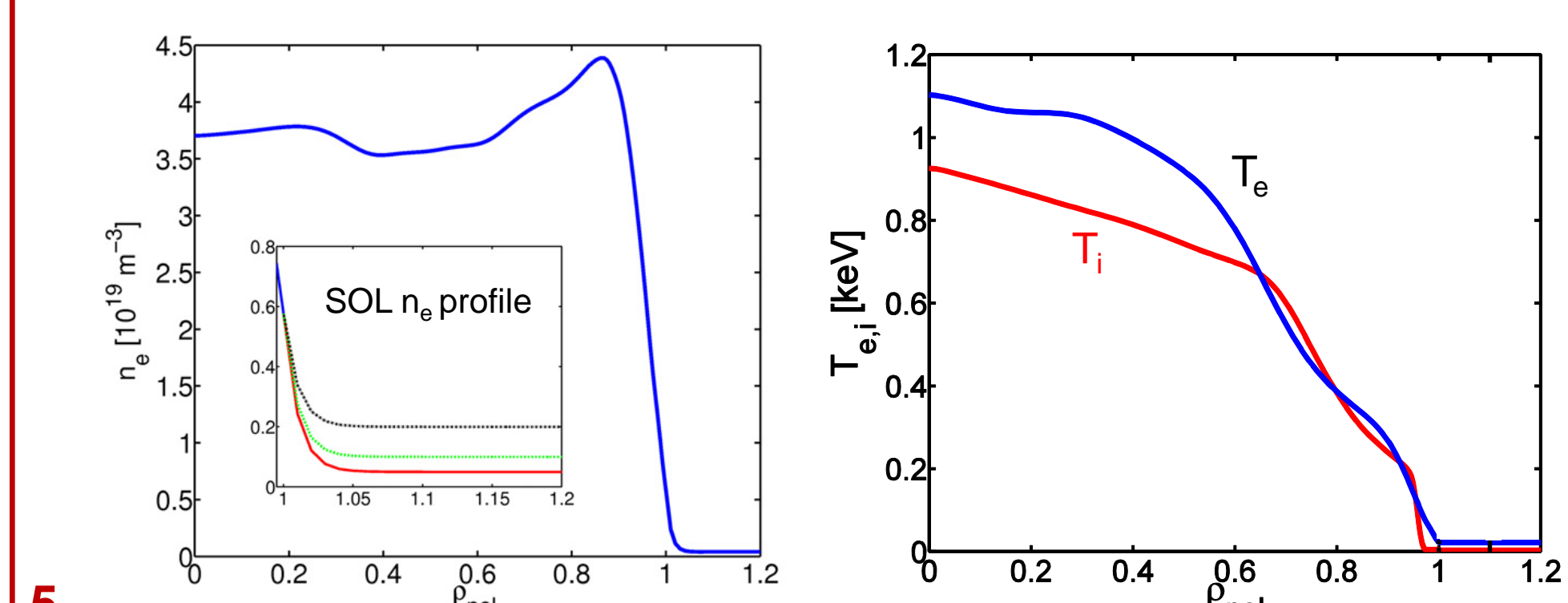


- FW experiments in EAST operates in the minority heating regime, which differs from the mid/high harmonic regime adopted in DIII-D and NSTX
- Unlike NSTX/NSTX-U and DIII-D, the electric field amplitude does NOT increase outside of LCFS with increasing  $n_{ant}$ , even when the wave propagate in the SOL
- Only a significant parallel component of the electric field is found in the SOL and is apparently larger than NSTX cases.

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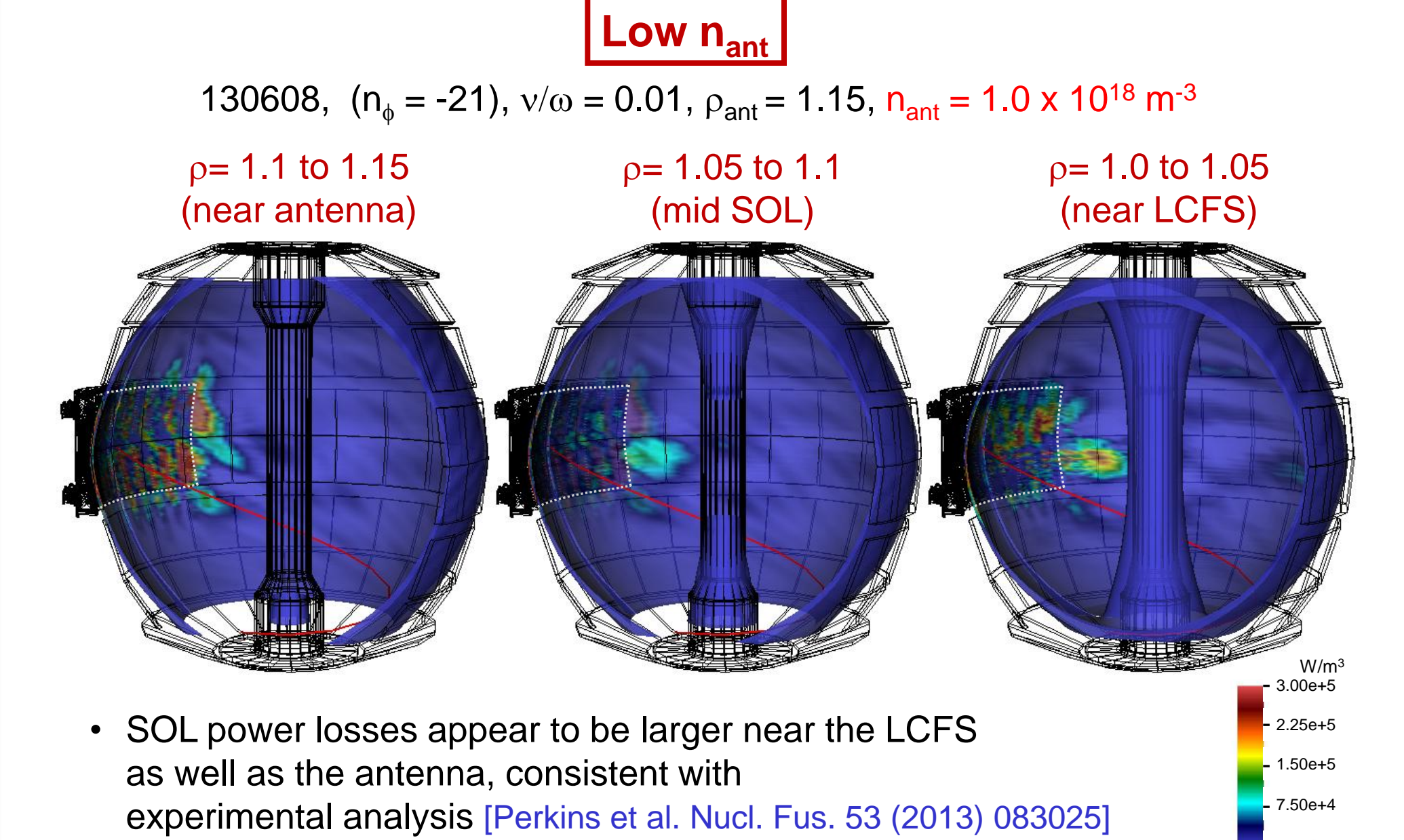
## NSTX discharge 130608: temperature and density profiles

- $B(R_0) = 0.55$  T
- $n_b(0) = 3.7 \times 10^{19} \text{ m}^{-3}$ ,  $n_b(1) \sim 5.8 \times 10^{18} \text{ m}^{-3}$ ,  $T_e(0) \sim 1.1$  keV
- Density (and SOL density), temperature profiles used in the simulations:



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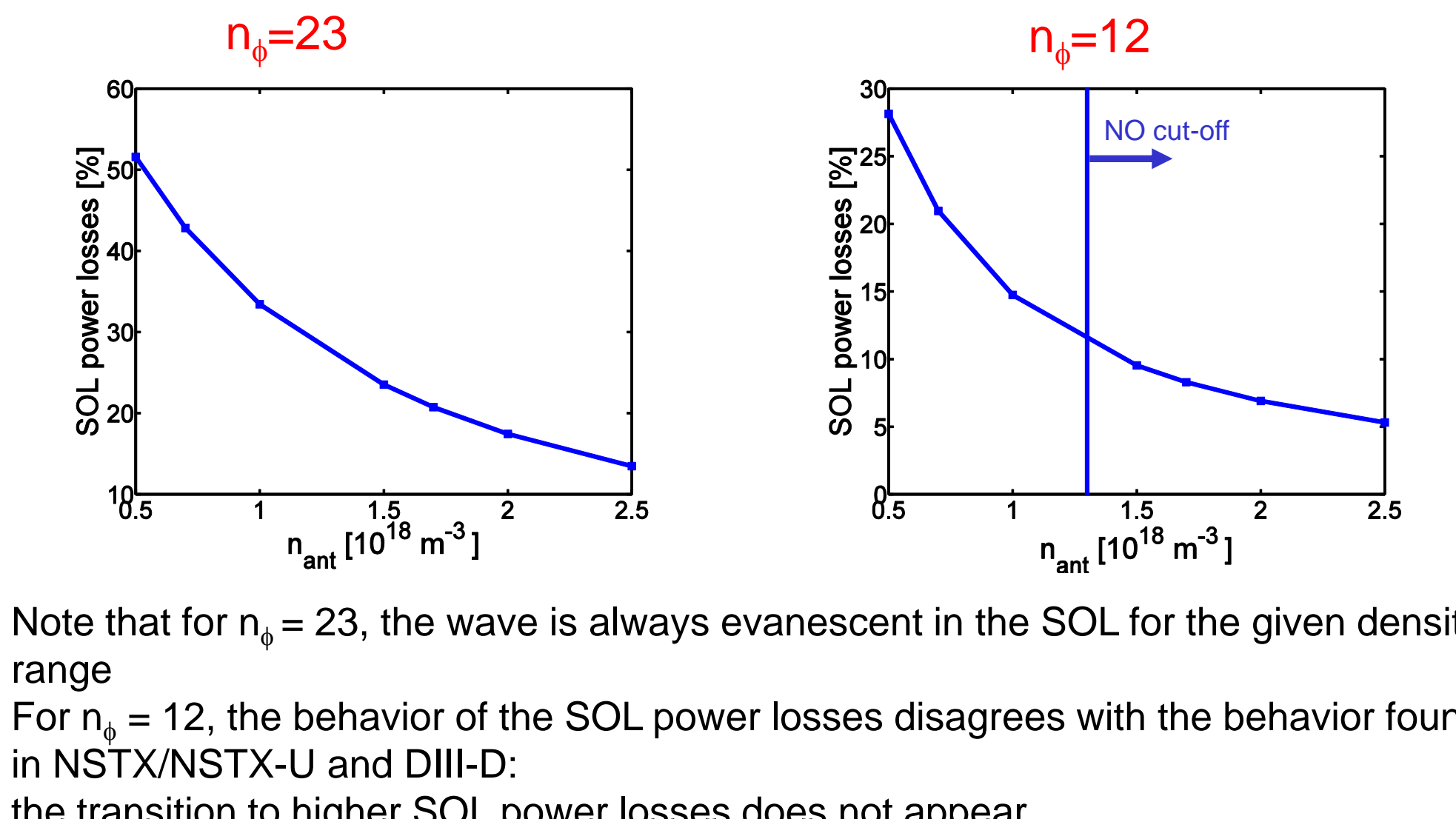
## RF power is deposited near the antenna and near the LCFS



- SOL power losses appear to be larger near the LCFS as well as the antenna, consistent with experimental analysis [Perkins *et al.*, Nucl. Fus. 53 (2013) 083025]

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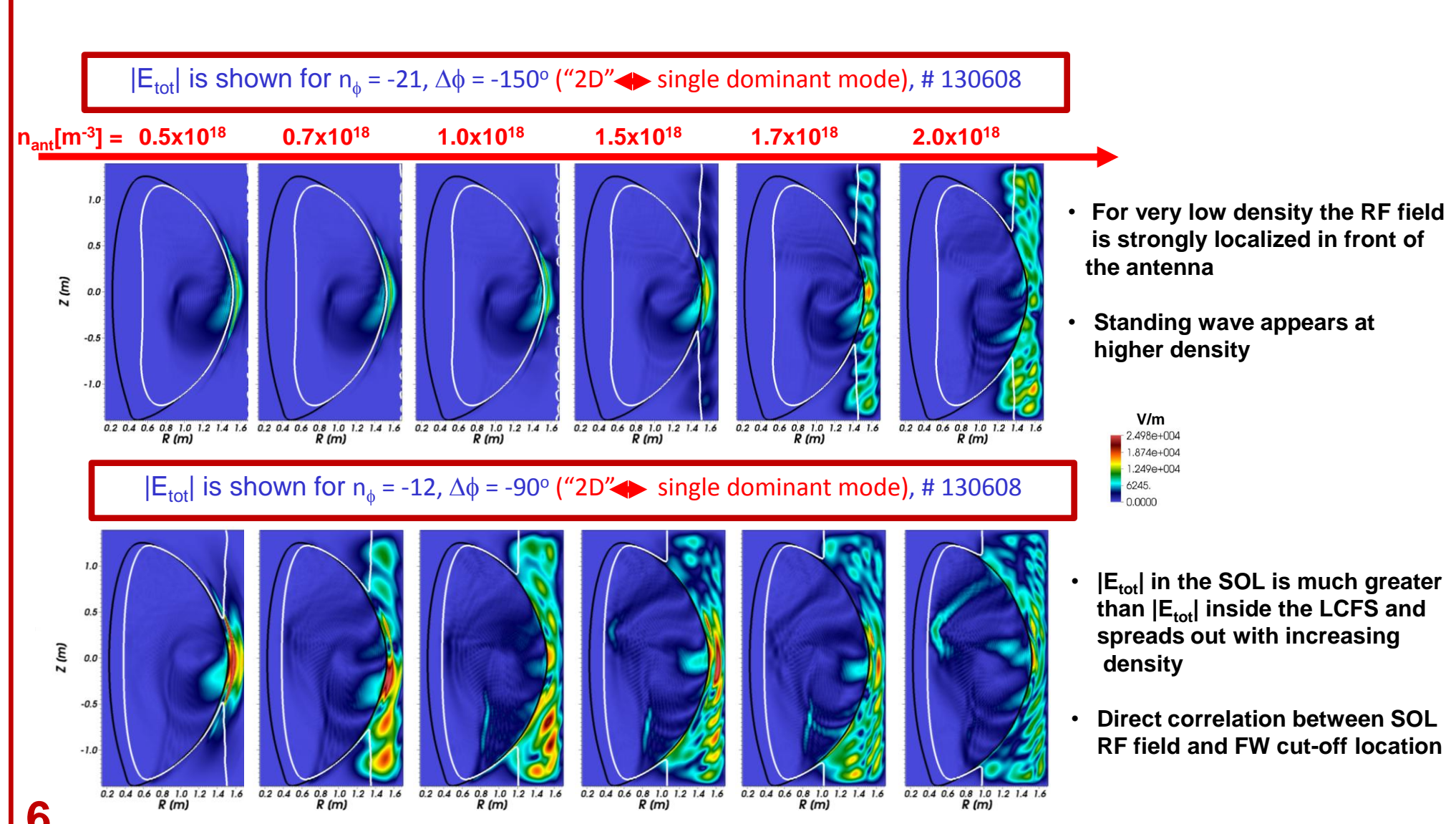
## NO increasing SOL power losses for higher $n_{ant}$ for EAST case



- Note that for  $n_b = 23$ , the wave is always evanescent in the SOL for the given density range
- For  $n_b = 12$ , the behavior of the SOL power losses disagrees with the behavior found in NSTX/NSTX-U and DIII-D: the transition to higher SOL power losses does not appear

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## $|E_{\text{sol}}|$ in the SOL increases when FW can propagate within the SOL



For very low density the RF field is strongly localized in front of the antenna

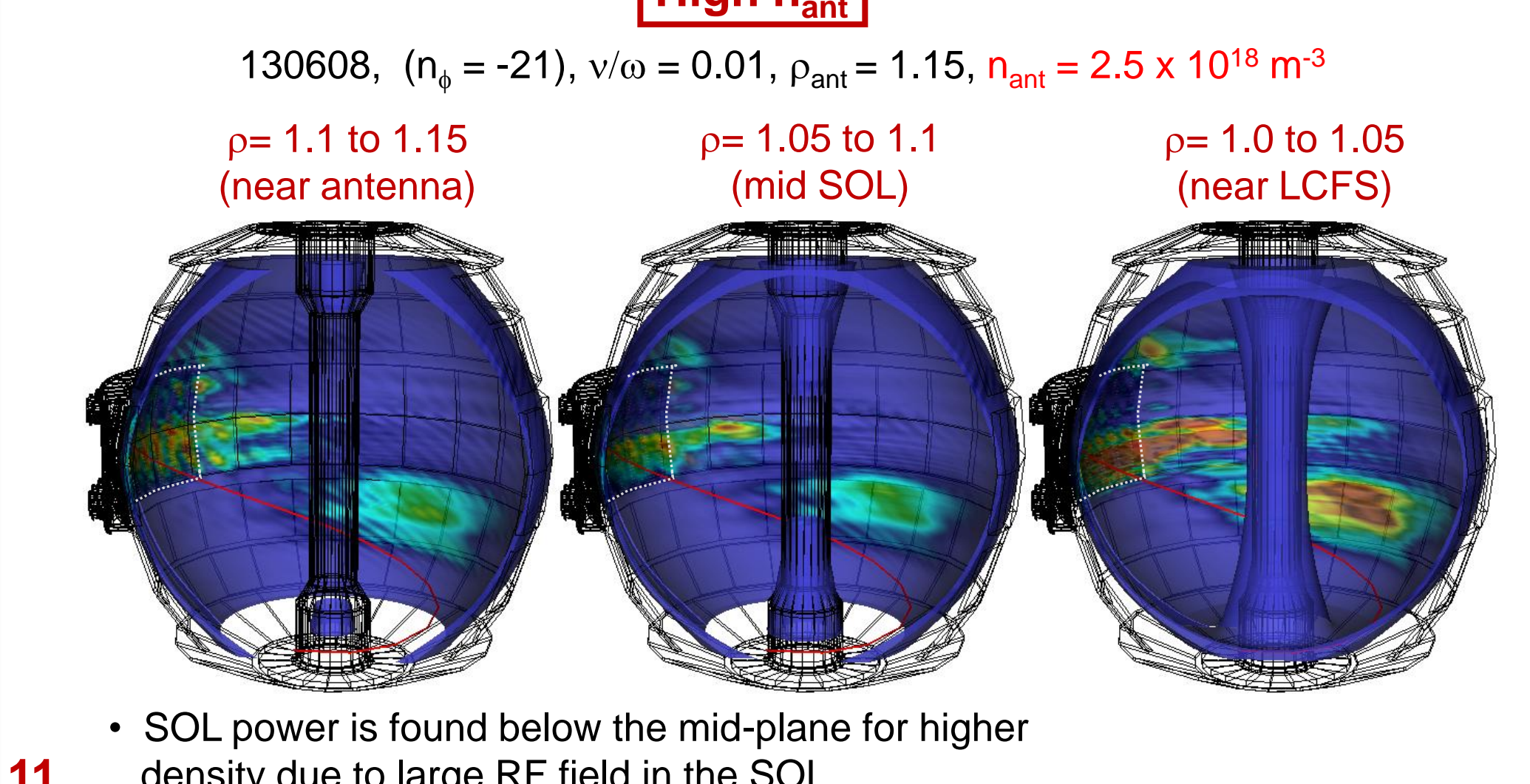
Standing wave appears at higher density

$|E_{\text{sol}}|$  in the SOL is much greater than  $|E_{\text{sol}}|$  inside the LCFS and spreads out with increasing density

Direct correlation between SOL RF field and FW cut-off location

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## High $n_{ant}$



- SOL power is found below the mid-plane for higher density due to large RF field in the SOL

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## Conclusions

- Full wave simulations for NSTX/NSTX-U and DIII-D show a direct correlation between the large SOL RF field, the location of the FW cut-off, and the SOL losses (in 2D and 3D)
  - ✓ SOL losses increase significantly as the wave transitions from evanescent to propagating as the density in the SOL increases
  - ✓ higher SOL losses with lower antenna phase, consistent with experiment
- 3D simulations show larger SOL losses near the antenna and the LCFS
  - ✓ consistent with the NSTX experiment
- NSTX-U simulations predict a wider evanescent region with lower SOL losses relative to NSTX
  - ✓ Optimization of the edge density in order to minimize the RF SOL losses
- The behavior of SOL power losses in NSTX/NSTX-U and DIII-D is really similar and is consistent with the experimental observations for both devices
- EAST (with minority heating regime) results strongly differ from NSTX/NSTX-U and DIII-D results: **under investigation**
- Future direction:
  - ✓ Further studies in order to identify the physical mechanism(s) behind the SOL losses [Sheath effects, PDI, non-linear processes, etc.]
  - ✓ Implementation of a more realistic boundary condition (limiter boundary) in the full wave code AORSA is underway
  - ✓ new experimental observations/measurements in the upcoming NSTX-U experimental campaign

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See for NSTX/NSTX-U also Bertelli *et al.*, Nucl. Fusion 54 (2014) 083004