

V1.1



Stability, Transport, and Active MHD Mode Control **Analysis of KSTAR High Performance Plasmas Supporting Disruption Avoidance***

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Abstract of paper EX/P7-16

H-mode plasma operation in KSTAR has surpassed the n = 1 ideal MHD no-wall beta limit computed to occur at $\beta_N = 2.5$ with $I_i = 0.7$. High β_N operation produced β_N of 3.3 sustained for 3 s, limited by tearing instabilities rather than resistive wall modes (RWMs). High fidelity kinetic equilibrium reconstructions (EFIT) have been developed to include Thomson scattering, charge exchange spectroscopy data, and an allowance for fast particle pressure. In addition, motional Stark effect data are used to produce reliable evaluation of the safety factor, q, profile. The reconstructed equilibria can exhibit significant variation of the q-profile dependent upon the broadness of the bootstrap current profile as computed by TRANSP analysis. TRANSP analysis indicates that the non-inductive current fraction can reach up to 75% while its profile can vary significantly depending upon plasma operational scenario. The classical stability of the m/n = 2/1tearing mode that limited high β_N operation is examined using the resistive DCON code and by the M3D-C¹ code using the kinetic EFIT reconstructions as input. For equilibria at high $\beta_N > 3$, the tearing stability index, Δ' , is more unstable compared to that of equilibria at reduced β_{N} . However calculations of instability of stable lower β_{N} equilibria indicate that the neoclassical components of tearing stability need to be invoked to produce consistency with experiment. MISK code analysis which examines global MHD stability modified by kinetic effects shows significant passive kinetic stabilization of the RWM in these plasmas. Predict-first TRANSP analysis was conducted to design KSTAR plasma operation with 100% non-inductive fraction at β_N above 4. Active RWM control has additionally been enabled on KSTAR. To accurately determine the dominant *n*-component produced by RWMs expected to onset at higher beta utilizing the new second NBI system, an algorithm has been developed that includes magnetic sensor compensation of the prompt applied field and the field from the induced current on the passive conductors. This analysis on stability, transport, and control provides the required foundation for disruption prediction and avoidance research on KSTAR.

US-KSTAR collaborative research on stability, transport, and active MHD mode control for disruption avoidance

Motivation

- To achieve high performance plasma operation with high non-inductive current fraction without disruptions, the estimation of stability and unique equilibrium and transport properties is crucial
- These analyses will provide the required foundation for disruption prediction and avoidance on KSTAR

S.A. Sabbagh, et al., EX/P6-26 in this conference

Outline

KSTAR

- Equilibrium reconstructions using measured kinetic profiles and MSE data for reliable stability and transport analysis
- Transport analysis using TRANSP code to investigate high noninductive current operation and to optimize the stability
- Ideal and resistive stability analysis using the resistive-DCON, MISK and M3D-C¹ code
- Active RWM feedback algorithm implemented in the KSTAR PCS

KSTAR H-mode equilibria have reached and exceeded the computed n = 1 ideal no-wall stability limit



Normalized beta vs. internal inductance from EFIT reconstruction** containing ~9,000 equilibria produced in the 2016 device campaign

*O. Katsuro-Hopkins, *et al.*, Nucl. Fusion **50** (2010) 025019 ** Y.S. Park, *et al.*, Nucl. Fusion **53** (2013) 083029

Parameters for high β_N

□ B_T in range 0.9-1.3 T

u Highest $\beta_N = 4.3$, $\beta_N/l_i > 6$

- $\ \ \, \square \ \ \, MHD \ \, stability \ \, at \ high \ \, \beta_N$
 - Many equilibria operate above the published ideal n = 1 no-wall stability limit (DCON)
 - Plasma is subject to RWM instability, depending on plasma rotation profile
 - High β_N > β_N^{no-wall} operation mostly limited by 2/1 tearing mode onset

High $\beta_N > 3$ equilibria limited by MHD - shot 16295



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Comparative equilibria stable to disruptive MHD instabilities - shot 16325



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Kinetic EFIT with MSE constraints used for accurate stability and transport analysis on KSTAR



(a) Reconstructed pressure and (b) safety factor profile from kinetic EFIT using internal profile constraints

Equilibrium reconstruction using measured internal profile constraints provides reliable pressure profile and internal magnetic geometry

S.A. Sabbagh, Nucl. Fusion 41 (2001) 1601

- KSTAR is equipped with key internal profile diagnostics
 - **Charge exchange** (T_i, V_{ϕ}) : 32 CHs
 - Thomson scattering (T_e, n_e) : 27 CHs
 - Motional Stark Effect : 25 CHs
- MSE background polychrometer installed in the device will improve the measurement
- Reconstructed *q*-profile shows region of low shear appear at varied *q*-values in different operational regimes due to broadness of bootstrap current profile

Kinetic EFIT reconstruction shows evolution of low-sheared *q*-profile region at varied *q*-values



Interpretive TRANSP analysis computes high non-inductive current fraction in the experimental equilibria



- Transport calculations using TRANSP to investigate a disruption-free path to device target high beta operation with high non-inductive current fraction, f_{NI}
- □ Interpretive TRANSP runs indicate f_{NI} can reach up to 75% in the experiment
- While f_{NI} can be high in different operating conditions, the bootstrap and total non-inductive current profile can vary significantly

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High $\beta_N > 4$ can be achieved with near fully non-inductive current by using the planned P_{NBI} increase in KSTAR



A new off-axis 2nd NBI system is being installed for the 2018 device operation

- Experimental $f_{GW} = 0.5$ and $H_{98y2} = 1.25$ are maintained, and total $P_{NBI} = 6.5$ MW ($P_{NBI-1} = 5$ MW, $P_{NBI-2} = 1.5$ MW) is used in the predictive TRANSP runs
- □ Predictive runs indicate $\beta_N > 4$ can be achieved with fully non-inductive current by using a fraction of the planned P_{NBI} increase ($P_{NBI} > 11$ MW from 2019)

DCON stability calculation shows high β_N equilibria are subject to n = 1 ideal instability



- **□** Equilibria at lower $\beta_N \sim 2$ is consistently stable to n = 1 ideal modes in DCON
- $\label{eq:stable} \Box \quad \mbox{Unlike the lower } \beta_N \mbox{ case, DCON calculates unstable } n = 1 \mbox{ mode with no-wall } (\beta_N > \beta_N^{no-wall}) \mbox{ at the achieved high } \beta_N > 3$

A.H. Glasser, Phys. Plasmas 23 (2016) 072505

Tearing stability is examined by the resistive DCON code



- □ The tearing stability index, Δ' , at high $\beta_N > 3$ is more unstable compared to that at reduced β_N later in the discharge
- The intermediate β_N ~ 2 equilibria that are experimentally stable to tearing modes but show an unstable value of Δ' indicate that stabilizing components of the NTM stability combined with sheared toroidal rotation need to be invoked

A.H. Glasser, et al., Phys. Plasmas 23 (2016) 112506

Ideal and resistive stability of high β_N equilibria examined by using M3D-C¹



- The high β_N equilibrium is computed to be unstable to n = 1 ideal mode by linear M3D-C¹ analysis consistent with DCON
- □ Kinetic RWM stability can explain the observed RWM stable operation at $\beta_N > \beta_N^{no-wall}$
- A marginally stable 2/1 tearing mode is computed in M3D-C¹ using a simplified input plasma resistivity estimated from measured T_e profile

(a) The perturbed poloidal flux of unstable ideal mode and (b) the radial velocity eigenfunction of a 2/1 tearing mode from M3D-C¹ S.C. Jar

S.C. Jardin, et al., J. Comput. Phys. 226 (2007) 2146

Kinetic modification of RWM stability is evaluated with including energetic particle effects



- MISK calculations find the equilibrium is stable to RWM as is consistent with experiment (rotation profile is scaled from 0.1 to 2 times in the analysis)
- Energetic particles are predicted to give a strong stabilizing effect to RWMs

J.W. Berkery, *et al.*, PRL **104** (2010) 035003 Y.S. Park, *et al.*, Nucl. Fusion **51** (2011) 053001

RWM active feedback control system installed in KSTAR



D For plasma operation at $\beta_N > \beta_N^{\text{no-wall}}$, RWM control system is prepared in KSTAR

- The middle in-vessel control coils (IVCCs) minimizes the inductive shielding by the copper passive stabilizing plates during RWM feedback
- □ Three sets of RWM B_p sensors at 8 toroidal positions above and below the midplane have been installed at the inside of the passive plates \Rightarrow total 20 independent B_p measurements for RWM identification (f_{sample} = 20 kHz)

Magnetic fields from induced currents are compensated from sensor measurement for accurate RWM identification



Define the DC-compensated δB :

$$\delta B_{\rm DC} = \delta B - \sum_{j}^{N_{\rm coil}} M_j I_j$$

- □ Significant inductive component still remains in δB_{DC} due to high sensor mutual coupling to the passive plates
- Similar to NSTX, use low-pass-filtered (LPF) RWM coil current derivative as the source of the inductive δB
- **Define the AC-compensated** δB :

$$\delta B_{\rm AC} = \delta B_{\rm DC} - \sum_{j}^{N_{\rm coil}} \sum_{k}^{N_{\tau}} p_{j,k} LPF\left(\frac{dI_{RWM,j}}{dt};\tau_{AC,k}\right)$$

■ Tested three τ_{AC} values well compensates the inductive effect with remaining $|\delta B| < 2$ G after applying DC/AC compensations

Utilization of multiple RWM sensor arrays is enabled by examining the phase of the perturbed δB of RWMs



<u>The perturbed δB_{p} for an unstable n = 1 eigenfunction</u>

sensor arrays w.r.t the phase at the midplane

- Utilization of the entire available sensor arrays is beneficial for a better control since the RWM eigenfunction can have varying helicities of δB in the outboard
- □ The phase difference between the midplane and off-midplane region changes with β_N as the response of the wall increases with faster mode growth
- Assumed sensor toroidal angles in mode decomposition are modified accordingly

Measured amplitude and phase of slowly rotating MHD mode



- Since RWMs have yet to be measured on KSTAR, mode identification has been tested for slowly rotating tearing modes (w/o applied feedback)
- Used 10 B_p sensor differences (180° opposing sensor pairs) for n = 1 identification
- Mode identification well measures the evolution of n = 1 locking tearing mode

Conclusions and Next Steps

- Stability and transport of high performance KSTAR plasmas are analyzed
 - Kinetic EFIT with MSE is used for accurate stability and transport analysis
 - TRANSP analysis computes high $\beta_N > 4$ with fully non-inductive current can be achieved with newly increased P_{NBI} in KSTAR
 - Achieved high β_N equilibria are subject to ideal n = 1 instability (DCON, M3D-C¹)
 - Resistive DCON analysis emphasizes the role of the pressure driven effects in the observed tearing stability
 - □ MISK analysis can explain the observed RWM stable operation at high $\beta_N > \beta_N^{\text{no-wall}}$
- Development of algorithm for RWM identification and feedback
 - Significant δB produced by induced currents in conducting structures is compensated from RWM sensors, and utilization of multiple sensor arrays is enabled by examining the perturbed δB structure of RWMs

Next Steps

- □ Improve stability analysis by employing the pressure driven terms of NTM stability
- **Test the developed RWM feedback algorithm in high** β_N experiment
- **Experiments are scheduled in 2018 to improve sustained high** β_N , and probe stability

BACKUP SLIDES

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Onset of strong 2/1 tearing mode terminated high β_N



- At high β_N phase, a benign n = 2 mode (presumably 3/2 mode) exists with strong sawteeth
 - No indication of W_{tot} reduction due to the n = 2 mode having |B_p| ~ 2 G
- High β_N operation was limited by strong 2/1 tearing mode onset
 - Measured mode amplitude > 20 G
 - Both W_{tot} and β_N were reduced by ~35% but maintained H-mode
 - Similar discharges exhibited different 2/1 tearing mode onset time (expected to be triggered by sawteeth)

Plasma rotation profile significantly reduced by > 20% due to the 2/1 mode onset

RWM stability evaluated with ideal and kinetic components allows for passive stabilization of the RWM



- Particle collisionality
- Plasma is stable when rotation is in resonance
 - \Box l = 0 harmonic : resonance with precession drift frequency $\omega_E + \langle \omega_D \rangle = 0$
 - l = -1 harmonic : resonance with bounce frequency

 $\omega_{\rm F} + < \omega_{\rm D} > - \omega_{\rm h} = 0$

New RWM sensors will give superior control performance over the previous device sensors



□ The design advantages of the new RWM sensors result in greater mode control - almost up to the ideal with-wall limit ($C_{\beta} = 98\%$, $\beta_{N} = 4.8$)

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Strong coupling between RWM feedback sensors and passive plates could be detrimental to RWM feedback



Transient effects of the induced currents on the passive plates have been examined by VALEN-3D code

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KSTAR passive plates are found to give a significant inductive effect on RWM sensor measurement when RWM coils are activated

AC sensor compensation can eliminate the remaining δB components after the DC compensations



- Similar to NSTX, use low-passfiltered (LPF) RWM coil currents as the source of the inductive δB
- **Define the AC compensated** δB :

$$\delta B_{\rm AC} = \delta B_{\rm DC} - \sum_{j}^{N_{\rm coil}} \sum_{k}^{N_{\tau}} p_{j,k} LPF\left(\frac{dI_{RWM,j}}{dt};\tau_{AC,k}\right)$$

- The coefficient matrix *p* has max. 20 (N_{sensor}) x 12 (N_{coil}) x 3 (N_{τ}) = 720 elements
- LPF with 3 different *τ* values has been tested
- Tested *τ* set well compensates the inductive effect. For the entire RWM sensors, remaining |*δB*| < 2 G after DC+AC compensations

Algorithms for mode identification

• The magnetic perturbation has an amplitude (A_{RWM}) and phase (ϕ_{RWM})

$$B(\phi) = A_{RWM} \cos(\phi - \phi_{RWM})$$

At the *i*-th sensor, the measured mode amplitude is:

$$B_{i} = A_{RWM} \cos(\phi_{i} - \phi_{RWM})$$

$$B_{i} = A_{RWM} \cos(\phi_{RWM}) \cos(\phi_{i}) + A_{RWM} \sin(\phi_{RWM}) \sin(\phi_{i})$$

$$B_{i} = C_{RWM} \cos(\phi_{i}) + S_{RWM} \sin(\phi_{i})$$

Combine signals to form an amplitude and phase of the plasma 3D perturbation

$$\begin{bmatrix} B_{1} \\ B_{2} \\ \vdots \\ B_{N} \end{bmatrix} = \begin{bmatrix} \cos(\phi_{1}) & \sin(\phi_{1}) \\ \cos(\phi_{2}) & \sin(\phi_{2}) \\ \vdots & \vdots \\ \cos(\phi_{N}) & \sin(\phi_{N}) \end{bmatrix} \begin{bmatrix} C_{RWM} \\ S_{RWM} \end{bmatrix} = M \begin{bmatrix} C_{RWM} \\ S_{RWM} \end{bmatrix} \Rightarrow \begin{bmatrix} C_{RWM} \\ B_{1} \\ \vdots \\ B_{N} \end{bmatrix} \Rightarrow \begin{bmatrix} A_{RWM} = \sqrt{C_{RWM}^{2} + S_{RWM}^{2}} \\ \phi_{RWM} = \operatorname{atan}\left(S_{RWM} / C_{RWM}\right) \end{bmatrix}$$

M : mode-ID matrix (20 x 2) for KSTAR

- Convert the sensor fields at each time point to amplitude and phase
- Outputs passed within PCS to RWM feedback algorithms for control current request
- Algorithm is presently being implemented in KSTAR PCS for use in 2018