Measurement of lower-hybrid waves with microwave scattering on TST-2

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1 Introduction

- 2 Microwave scattering
- Full-wave analysis of LH waves

4 Summary



A microwave scattering diagnostic is being developed to measure the Lower-Hybrid (LH) waves directly and test numerical simulations

- Background
 - Non-inductive plasma start-up with LH waves has been studied on TST-2
 - The current drive efficiency in the actual experiment is much lower than what is predicted by GENRAY-CQL3D
- Microwaves in the range of 12-40 GHz have wavenumber similar to LH waves in TST-2 and suited for diagnosing these waves
- The wave measurements will be compared with numerical simulations to quantitatively understand the current drive physics
- Previous direct LH wave measurements
 - Microwave back-scattering with a reflectometer [Baek 2014]
 - ▶ CO₂ laser scattering [Takase 1985]

The TST-2 Machine



• *R* = 36 cm

RF non-inductive start-up

- $B_t < 0.15 \text{ T}$ (pulse length $\sim 80 \text{ ms}$)
- $\bar{n}_e < 10^{18} \ {
 m m}^{-3}$
- $I_p < 25 \text{ kA}$
- RF 400 kW @ 200 MHz

Ohmic (inductive) start-up

- $B_t < 0.3 \text{ T}$ (pulse length $\sim 30 \text{ ms}$)
- $\bar{n}_e < 10^{19} \ \mathrm{m}^{-3}$
- $I_{
 m p} < 110$ kA

Non-inductive plasma start-up using LH waves is studied on TST-2



- Plasma initiation (ECH, LH)
 I_p ramp-up with LH waves
 Further ramp-up with neutral beam injection
- The LH wave has one of the highest current drive efficiency among rf waves
- Similar scenario tested on JT-60U [Shiraiwa 2004]
- The plasma density must be kept low during *I_p* ramp-up
 ← LH high-density limit

LH waves @200 MHz can access 10^{18} m⁻³ for $n_{\parallel} = 5-9$



- Density limit low at low field
- $\omega > 2\omega_{LH}$ to avoid parametric decay instability

The experimentally achieved plasma current (< 25 kA) is an order of magnitude smaller than simulated



- Simulation by GENRAY (ray-tracing) and CQL3D (Fokker-Planck solver)
- The experimental density limit qualitatively consistent with simulation

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A launched microwave beam can be Bragg-scattered by density fluctuations of the LH waves



Optimum scattering geometry was investigated from the simulated typical wave trajectory on TST-2



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LH waves propagating in a wide range of poloidal cross-section can be detected with a microwave launched from the midplane

- Scan the incident beam angle θ
- Scan the beam frequency
- Detect the scattered light at multiple poloidal locations (x_d)
- The source frequency swept at 1 kHz \rightarrow the whole wavenumber spectrum is measured every 1 ms
- Incident beam angle needs to be scanned on a shot-to-shot basis

The O-mode microwave scattering system



- Microwave bands
 - Ku band (WR-62): 12.4-18.0 GHz
 - K band (WR-42): 18.0-26.5 GHz
 - Ka band (WR-28): 26.5-40.0 GHz
- Incident beam from the midplane horizontal port
- 4-antenna array below the plasma
- Small toroidal tilt to compensate wave k_{ϕ}

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5 Backup

Full-wave simulation of LH waves was performed with $\ensuremath{\mathsf{AORSA}}$



The expected signal level is within the detectable range



 $\sim 2.8 \times 10^{-15}\,\mathrm{m} \cdot 0.01\,\mathrm{m} \cdot 2 \times 10^{16}\,\mathrm{m}^{-3} \cdot 0.05\,\mathrm{m} = 0.03$

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The microwave scattering system was designed for directly measuring LH waves on TST-2

- The current drive efficiency on TST-2 is lower than what is predicted by GENRAY-CQL3D
- A microwave scattering diagnostics at 12-40 GHz is designed and being fabricated
- The LH fluctuation level estimated with AORSA is within the detectable range
- Future work
 - Detect LH waves
 - Investigate changes in the location of LH waves and their wavenumber spectrum for various parameters
 - Test GENRAY-CQL3D
 - ► Test full-wave codes (AORSA, TORLH)

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Strong correlation of plasma current and density leads to upper limit in the driven current at a given magnetic field



• Density limit substantially lower than expected at higher field?

LH accessibility condition

 $\mathsf{LH} \leftrightarrow \mathsf{FW}$ mode conversion condition:

$$\begin{split} \frac{\omega_{pi}}{\omega} &= \textit{N}_{\parallel}\textit{Y} \pm \sqrt{1 + \textit{N}_{\parallel}^2(\textit{Y}^2 - 1)}, \\ \textit{Y}^2 &= \frac{\omega^2}{\omega_{ce}\omega_{ci}}. \end{split}$$

In terms of N_{\parallel} ,

$$N_{\parallel} = rac{\omega_{pi}}{\omega}Y \pm \sqrt{1 + rac{\omega_{pi}^2}{\omega^2}(Y^2 - 1)}.$$

Launching the wave from the low-density side, the parameter is accessible if

$$N_{\parallel} > N_{a} = rac{\omega_{pi}}{\omega}Y + \sqrt{1 + rac{\omega_{pi}^{2}}{\omega^{2}}(Y^{2} - 1)}.$$

Full-wave analysis will be performed

TORLH [Wright 2009]

- Finite Larmor radius code (k_⊥ρ_L < 1)
- θ : spectral ansatz
- ψ : finite element
- Coupled with CQL3D [Harvey 1992]
- \rightarrow Non-thermal electron distribution
 - Weak absorption, multi-pass regime
- \rightarrow Full-wave effect may be important

