

Reduced-model (SOLT) validation from comparisons with GPI data on NSTX

D.A. Russell, D.A. D'Ippolito and J.R. Myra

Lodestar Research Corporation

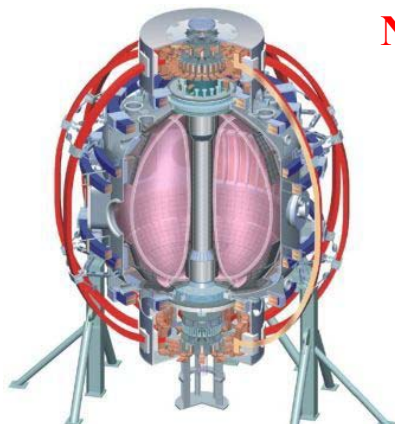
R. Maqueda Nova Photonics

T. Munsat U. of Colorado

D.P. Stotler and S.J. Zweben *PPPL*

and the NSTX Research Team

50th Annual Meeting of the Division of Plasma Physics
Marsalis A/B, Hyatt Regency Hotel,
Dallas, Texas
November 19, 2008



College W&M
Colorado Sch Mines
Columbia U
Comp-X
General Atomics
INEL
Johns Hopkins U
LANL
LLNL
Lodestar
MIT
Nova Photonics
New York U
Old Dominion U
ORNL
PPPL
PSI
Princeton U
Purdue U
SNL
Think Tank, Inc.
UC Davis
UC Irvine
UCLA
UCSD
U Colorado
U Maryland
U Rochester
U Washington
U Wisconsin

Culham Sci Ctr
U St. Andrews
York U
Chubu U
Fukui U
Hiroshima U
Hyogo U
Kyoto U
Kyushu U
Kyushu Tokai U
NIFS
Niigata U
U Tokyo
JAEA
Hebrew U
Ioffe Inst
RRC Kurchatov Inst
TRINITY
KBSI
KAIST
POSTECH
ASIPP
ENEA, Frascati
CEA, Cadarache
IPP, Jülich
IPP, Garching
ASCR, Czech Rep
U Quebec

SOLT 2D Turbulence Code (Scrape-Off Layer Turbulence)

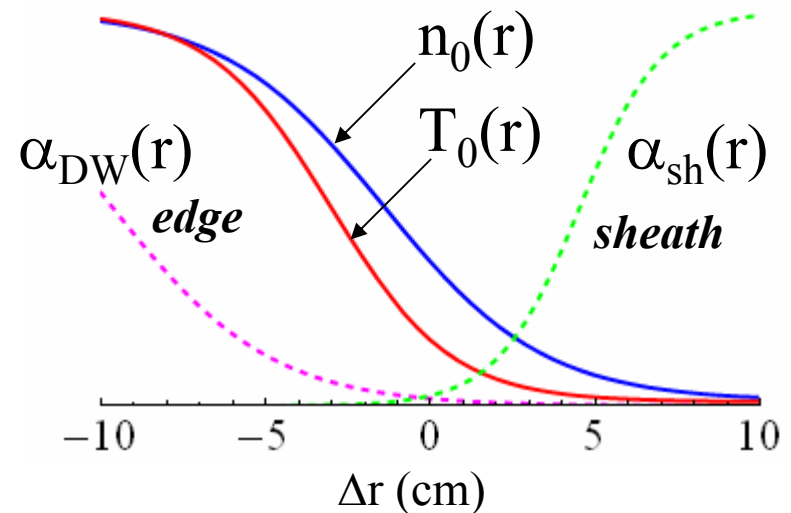
Reduced-Model Equations of Evolution in 2D (OM, $\perp \underline{B}$)
for
Vorticity, Density, Temperature and Zonal Momentum

Two Radially Distinguished Regions:

- Edge Region ($\Delta r < 0$)
 - Sources of Density and Temperature
 - Electron Drift Waves
 - Curvature-Driven (Interchange) Instability (Blob Birth Zone)

- SOL Region ($\Delta r > 0$)
Sheath Absorption: a Sink for all Fields

Initial and Reference Profiles
(normalized)



(no zonal flow initially)

Model Equations

include drift wave, curvature and sheath effects

•Vorticity

$$\partial_t \nabla^2 \tilde{\phi} =$$

$$= \left\{ -\mathbf{v} \cdot \nabla \nabla^2 \phi + \alpha_{\text{DW}} \frac{\bar{T}^{3/2}}{\bar{n}} (\phi - T \ln(n)) + \alpha_{\text{sh}} T^{1/2} (1 - \exp((\phi_B - \phi)/T)) - \frac{\beta}{n} \partial_y (nT) + \mu \nabla^4 \phi \right\}$$

- Electron Adiabatic and Sheath parameters, $\alpha_{\text{DW}}(x)$ and $\alpha_{\text{sh}}(x)$ model an x-dependent connection length $L_{//}$.

- $\beta = 2 \rho_s / R$, the familiar curvature coefficient.

•Density

$$(\partial_t + \mathbf{v} \cdot \nabla) n = \alpha_{\text{DW}} \bar{T}^{3/2} \{\phi - T \ln(n)\} - \alpha_{\text{sh}} n T^{1/2} \exp((\phi_B - \phi)/T) + D \nabla^2 n + S_n$$

•Temperature

- $S_n = v_n(x) (n_0(x) - n)$ and $S_T = v_T(x) (T_0(x) - T)$
($v_{n,T} \sim 1$ as $x \rightarrow 0$ and are ignorably small in the SOL.)

$$(\partial_t + \mathbf{v} \cdot \nabla) T = -\alpha_{\text{sh}} s_E T^{3/2} \exp((\phi_B - \phi)/T) + S_T$$

•Zonal Momentum

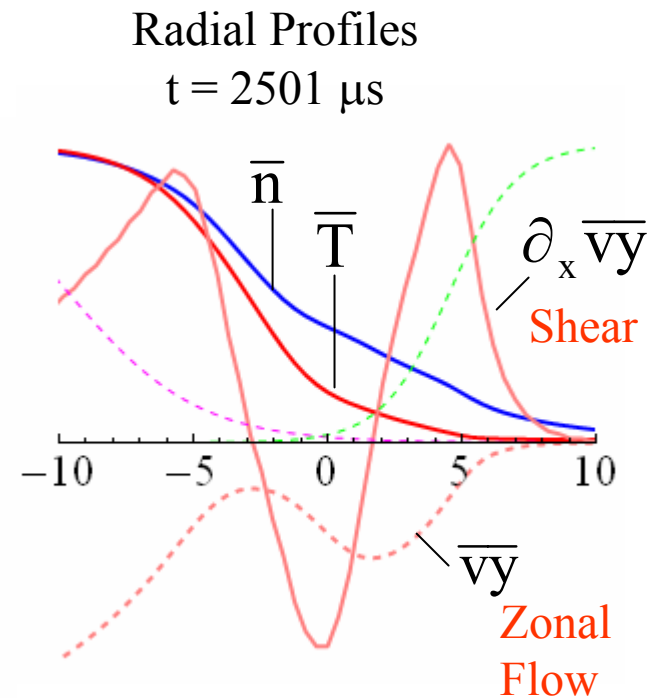
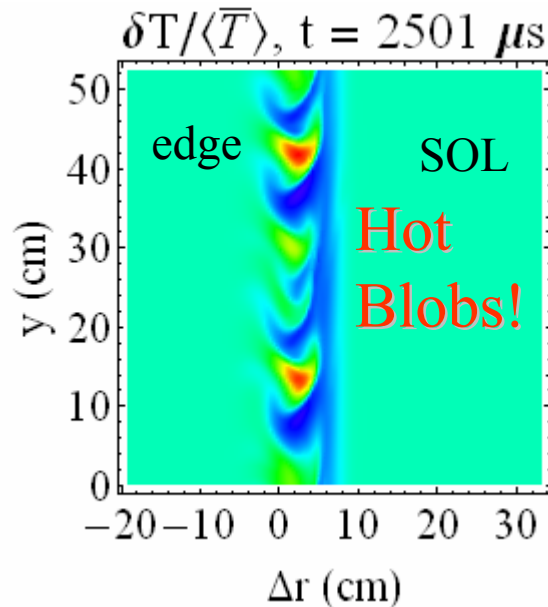
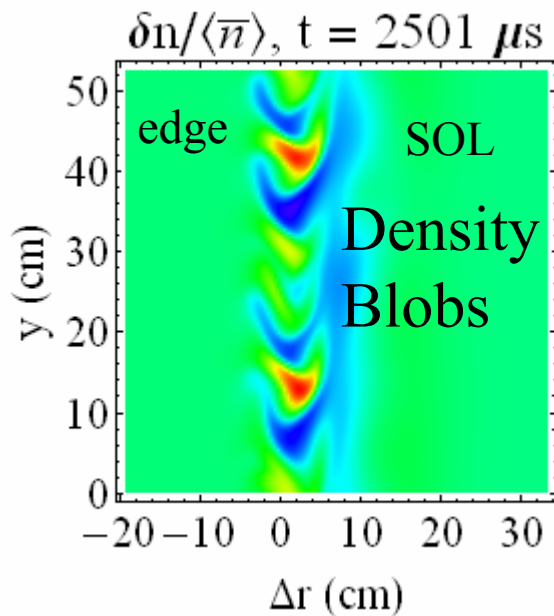
$$\partial_t p_y + \partial_r \overline{n v_r v_y} = \int_r^{L_r} dx \alpha_{\text{sh}} n T^{1/2} (1 - \exp((\phi_B - \phi)/T)) + \bar{\mu} \partial_r^2 \overline{v_y} - \nu_{p_y} p_y$$

$p_y \equiv \overline{n v_y}$, $\overline{v_y} = \partial_r \bar{\phi}$, and where ν_{p_y} is a constant, varied to damp the flow.

Boundary Conditions: $p_y = 0$ (*no slip*) and $\bar{\phi} = \bar{\phi}_{\text{Bohm}}$ at $r = L_r$

The SOLT code resolves strong ($\delta n/n \sim 1$) blob turbulence and zonal flows

- Split-step algorithm
- Convection: Flux-corrected transport (*phenical SHASTA*, J. Boris and D. Book with the 2D flux-limiter by S. Zalesak – J. Comp. Phys. **31**, 331 (1979)).
- Diffusion: Crank-Nicholson (ADI).
- All other terms: evolved explicitly.



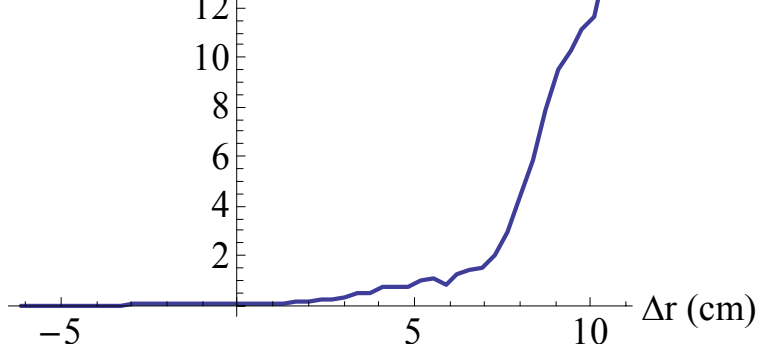
Synthetic diagnostic of GPI intensity, $I = N f_A(n, T)$, allows direct comparison with experiment

NSTX

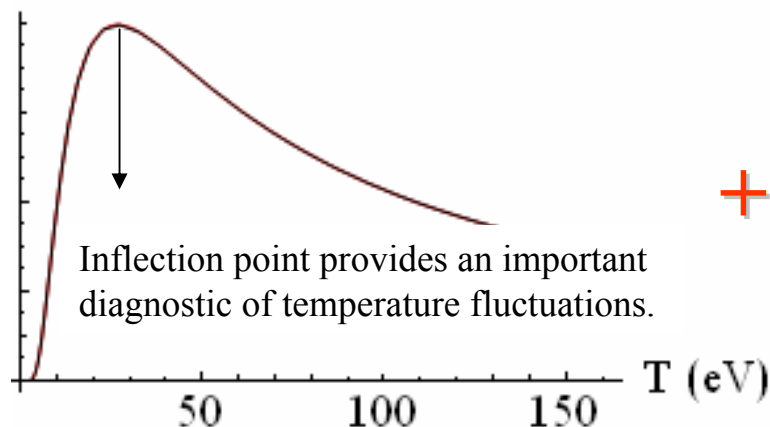
$N/10^{16}$

cm^{-3}

Neutral Density Profile
for Shot #112825



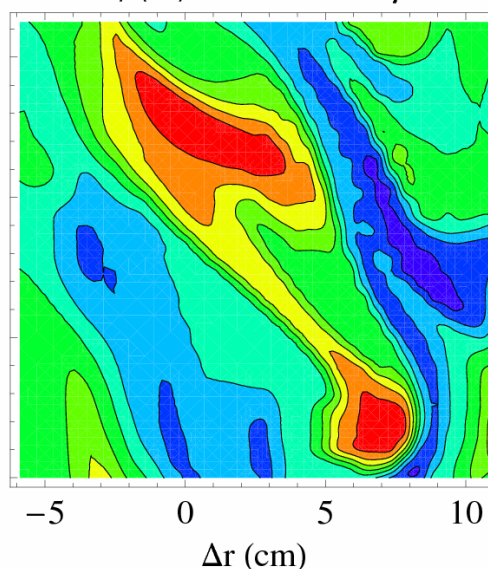
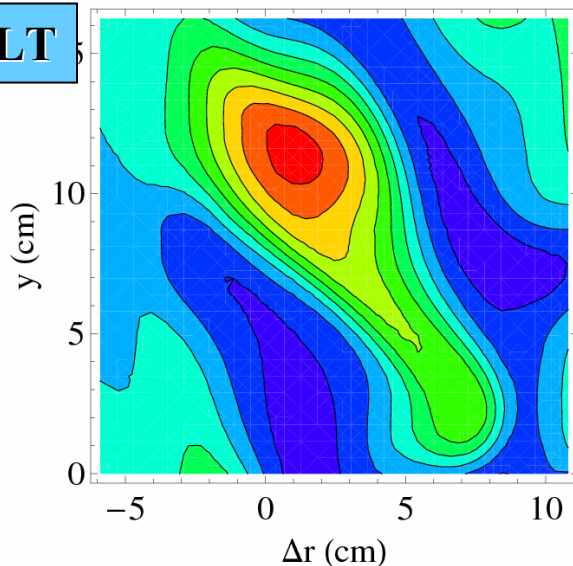
$f_A(n, T)$: atomic emission function [D. Stotler]



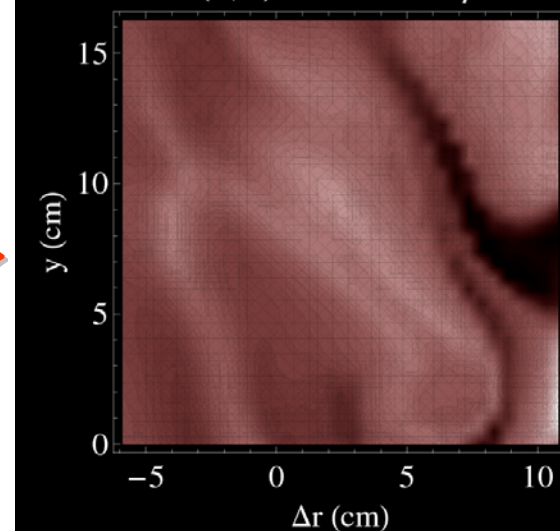
SOLT

$\delta n / \langle \bar{n} \rangle @ t = 4175 \mu\text{s}$

$\delta T / \langle \bar{T} \rangle @ t = 4175 \mu\text{s}$



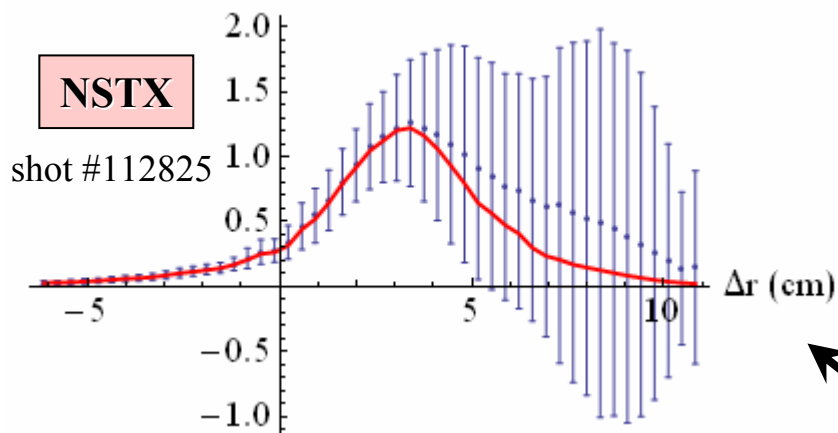
$\hat{I}(n, T) @ t = 4175 \mu\text{s}$



The radial distribution of edge turbulence depends on sheath absorption

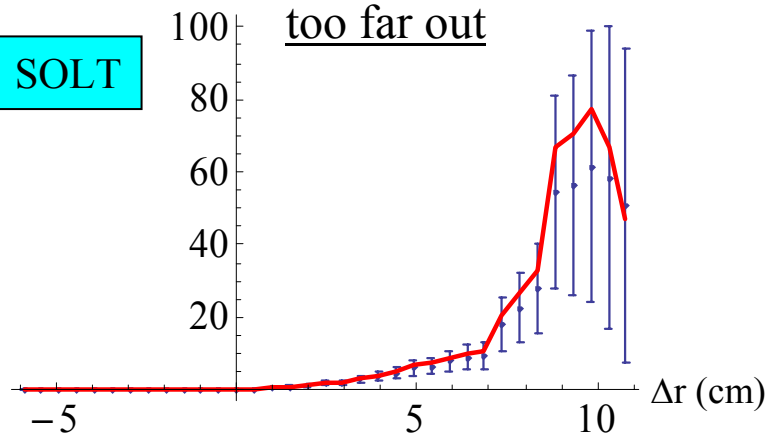
GPI point data

$\langle I \rangle$, σ_I & Median I (red)



SOLT

SHE at $\Delta r = 10.5$ cm
too far out

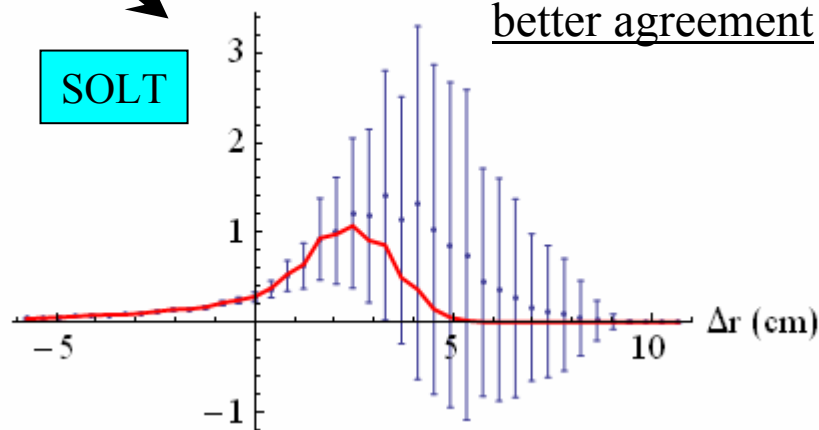


The shape of the GPI intensity curves is a function of the **sheath absorption**, especially the location of the sheath entrance (SHE).

⇒ We can map out the connection length $L_{//}$ (Δr) in the SOL.

SOLT

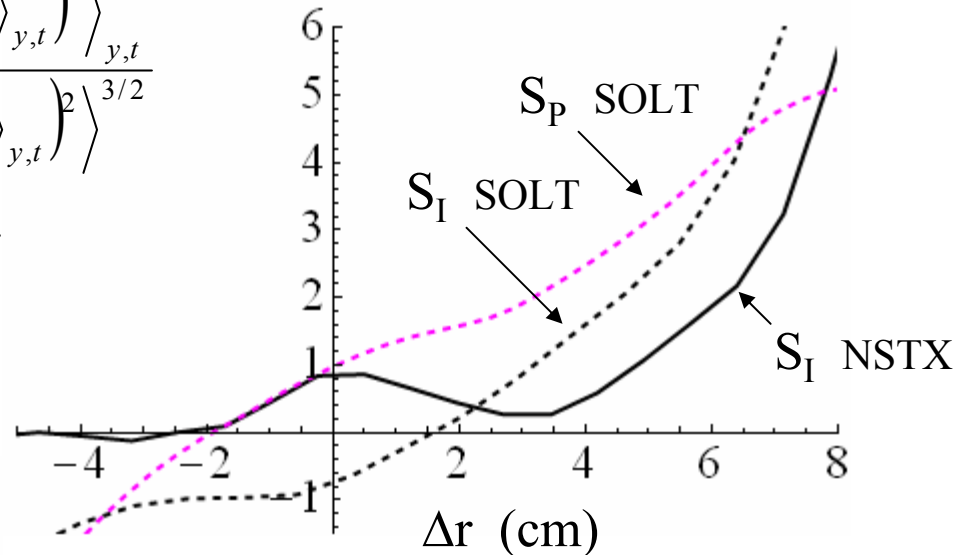
SHE at $\Delta r = 6.7$ cm
better agreement



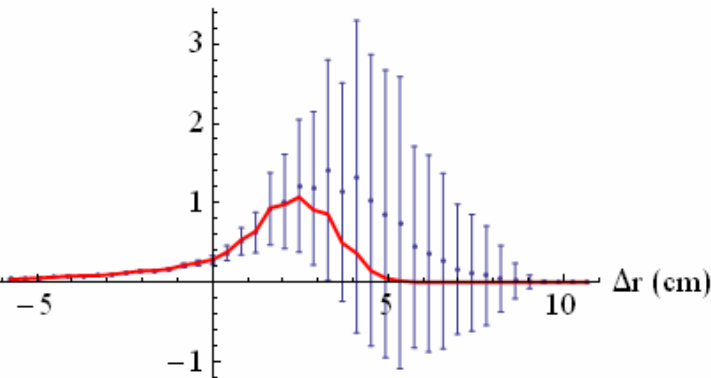
The skewness of intensity fluctuations, S_I , indicates the ballistic (blob) nature of the turbulence

$$S_I(\Delta r) = \frac{\langle (I(\Delta r, y, t) - \langle I \rangle_{y,t})^3 \rangle_{y,t}}{\langle (I(\Delta r, y, t) - \langle I \rangle_{y,t})^2 \rangle_{y,t}^{3/2}}$$

S_P , S_n , S_T - similar



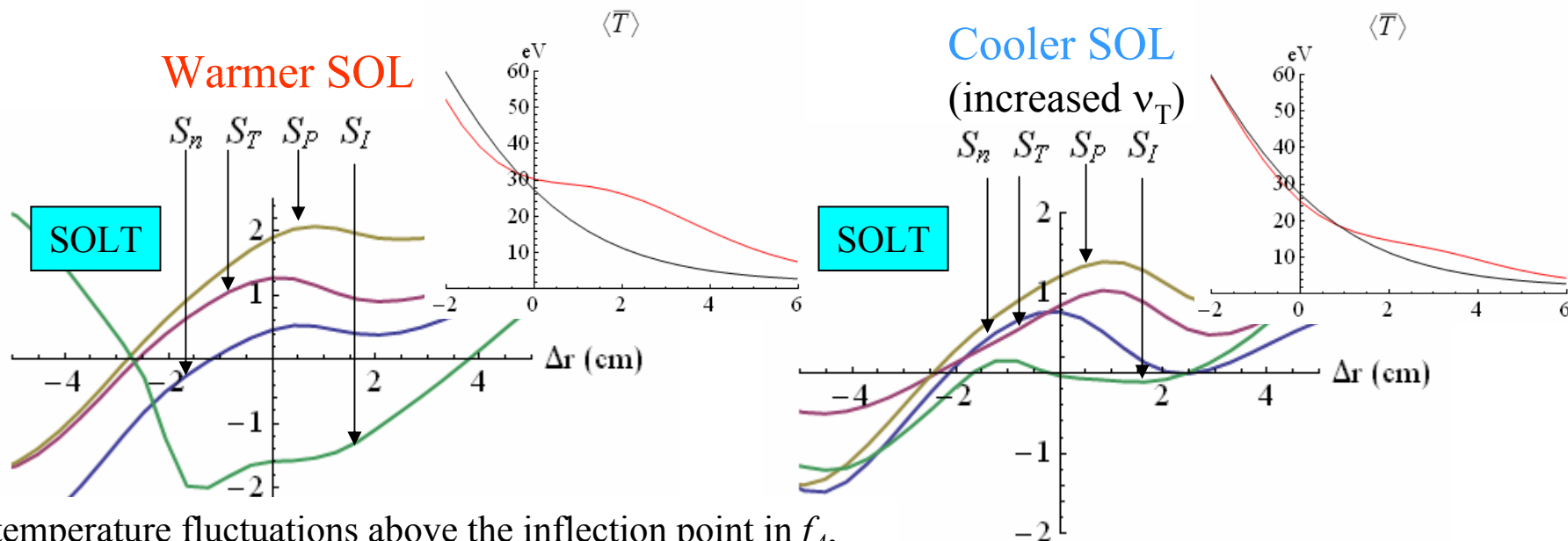
$\langle I \rangle$, σ_I & Median I (red)



Reasonable agreement is observed where the GPI intensity fluctuations are large, $0 < \Delta r < 8$ cm.

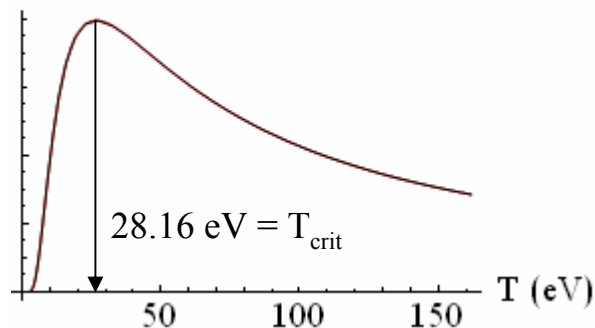
But note $S_I < 0$ near SEP ($\Delta r = 0$) in the SOLT simulations. Cause? (next page)

S_I is a sensitive measure of temperature fluctuations



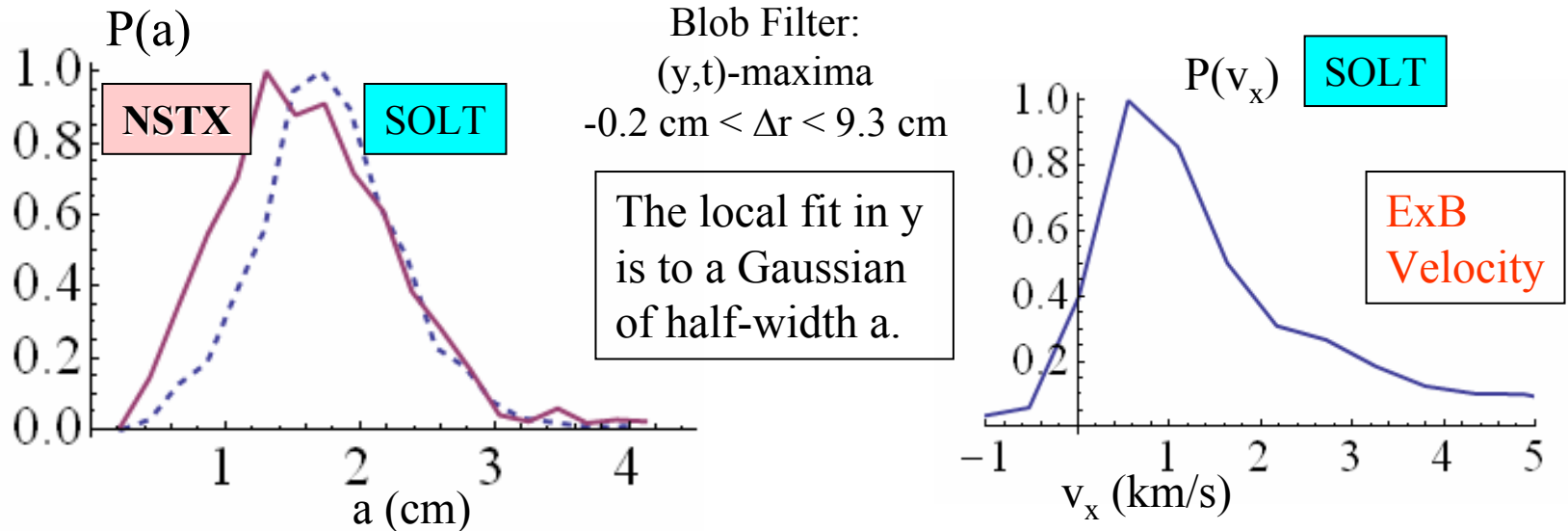
For temperature fluctuations above the inflection point in f_A , δT and δI are anti-correlated, and S_T and S_I tend to have opposite sign. As $\delta T \rightarrow 0$, S_I is dominated by δn .

$$f_A(\langle \bar{n} \rangle |_{\Delta r=2 \text{ cm}}, T)$$

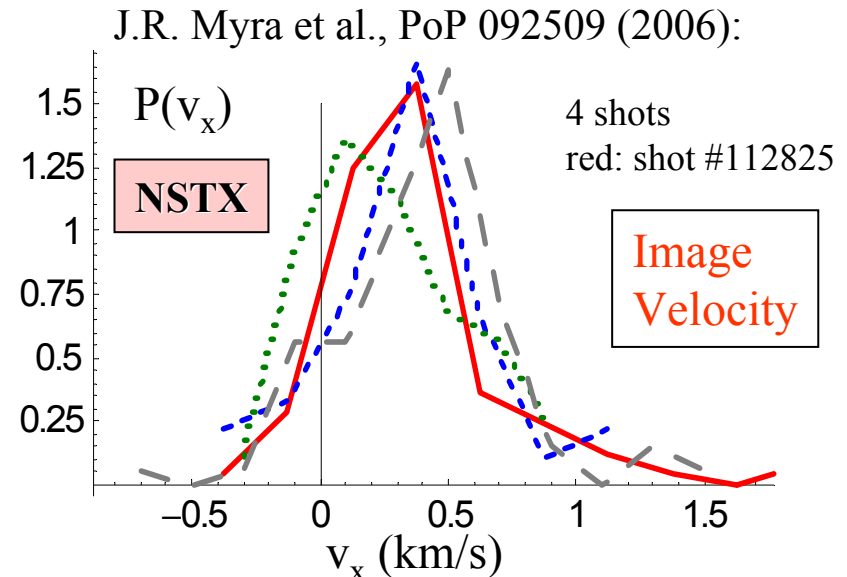


- f_A is max at T_{crit} : $S_I < 0$ when $T_e > T_{\text{crit}}$
- In the simulation, T_e is too large outside SEP ($\Delta r > 0$) due to turbulent heat transport.
- *Note that NSTX data has $S_I > 0$ near the SEP.*

PDFs of blob size and velocity provide a stringent test of theory



- PDFs of blob size are in close agreement.
- SOLT blobs appear to be moving about twice as fast as NSTX blobs.
- We need a less ambiguous measure of image velocity (coming, from Tobin Munsat) and to compare that measure with ExB velocities in the simulations.



SOLT parallel particle and energy fluxes at the midplane have comparable radial widths

Sheath Boundary Conditions on Open Field Lines ($\Delta r > 0$)

$$\Gamma_{//} = \langle c_s n_e \exp[e(\phi_B - \phi) / T_e] \rangle_{y,t}$$

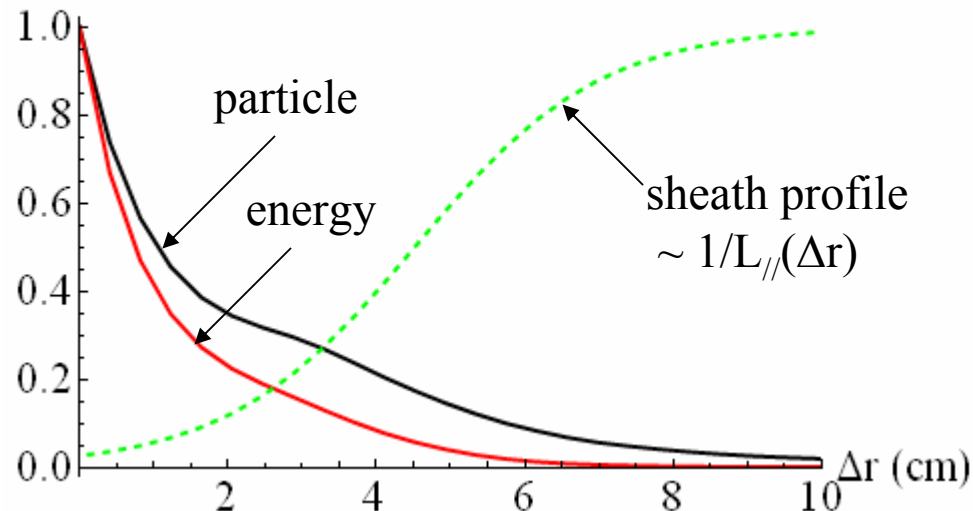
$$q_{//} = \langle s_E c_s n_e T_e \exp[e(\phi_B - \phi) / T_e] \rangle_{y,t}$$

s_E : energy transport
enhancement factor
(= 6)

$\Gamma_{//}/\Gamma_{//}(0)$ & $q_{//}/q_{//}(0)$ (red)

$\Gamma_{//}(0)$:
 $4.6 \times 10^{19} \text{ s}^{-1} \text{ cm}^{-2}$

$q_{//}(0)$:
 $1.4 \text{ kW} / \text{cm}^2$



SOLT simulation:
same as above for
comparison with
NSTX shot
#112825

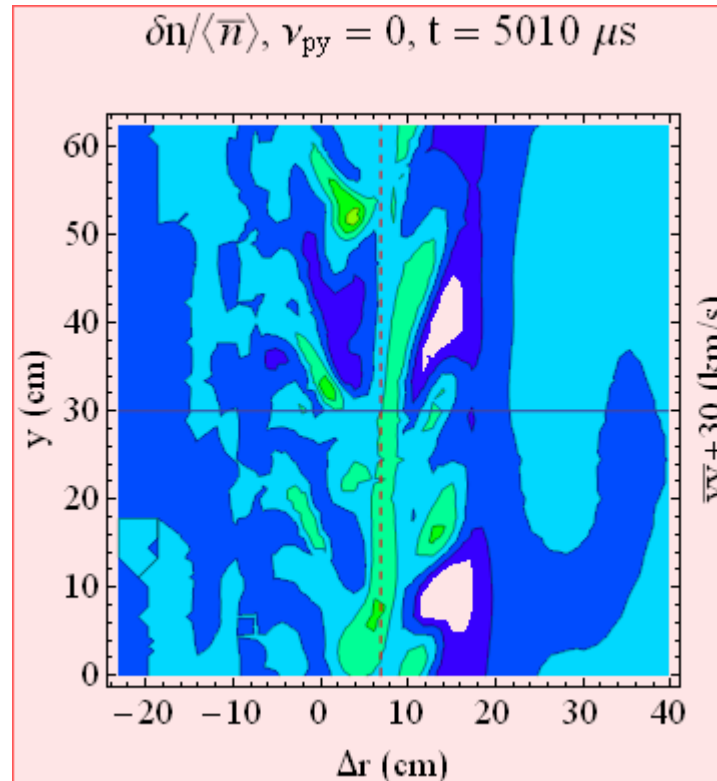
Similar Widths : Both particles and energy are transported ballistically by blobs.

Quantitative comparisons are now feasible between synthetic (SOLT) and experimental GPI data.

- It is essential to analyze the simulation and diagnostic data using the same post-processing algorithms (smoothing and blob filtering).
- The simulated turbulence is sensitive to the parameters that control the stability of the system, e.g. sheath dissipation and sheared flows.
 - The simulated GPI signal is sensitive to the details of temperature and density fluctuations.
 - Still better agreement ($S_I > 0$, slower blobs) will likely be obtained by increasing temperature dissipation in the SOL.
- *Beyond validation*, SOLT predictions can be explored.
 - The role of sheared zonal flow in mediating blob transport across the LCS is under study.

Epilogue

High-flux, intermittent blob bursts and zonal flow reversal (spin-up) observed in a SOLT simulation driven more strongly than the NSTX reference cases above:



Relevant to anticipated higher-gradient NSTX profiles?