

Energy Transport via Compressional Alfvén Eigenmodes Coupling to Kinetic Alfvén Waves

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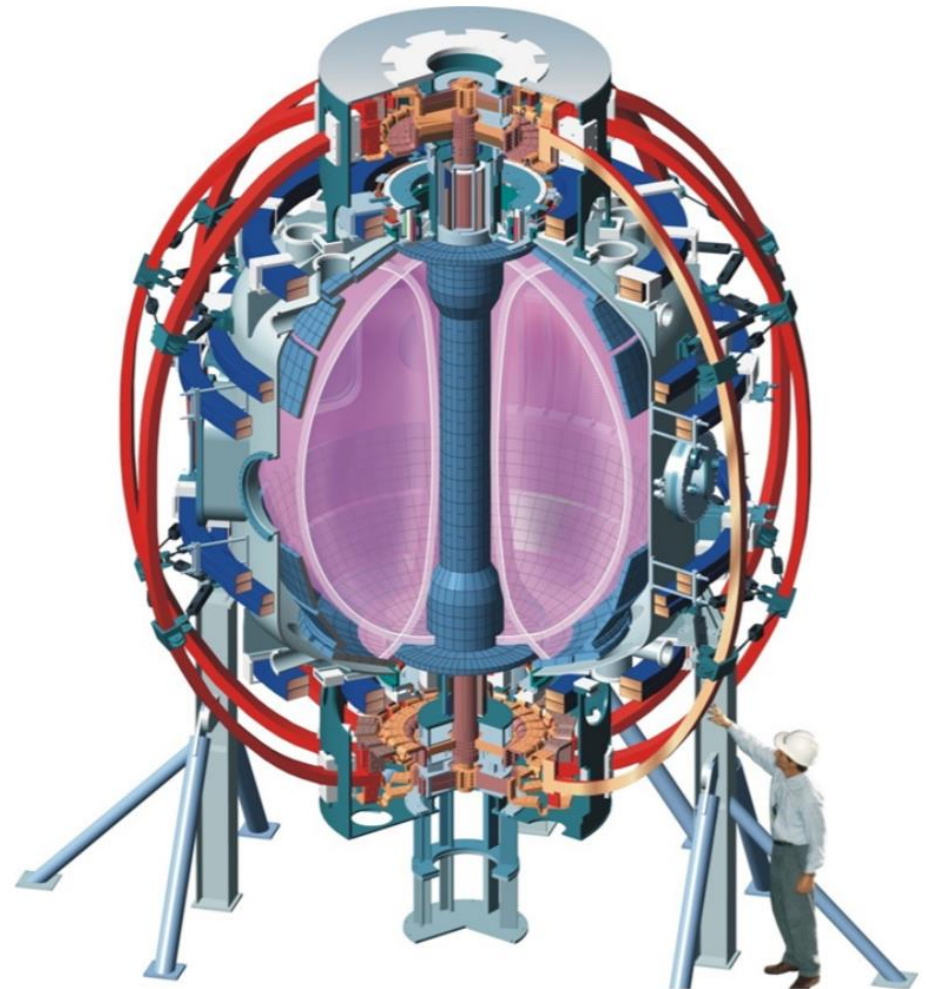


Energy confinement is important in fusion

- Confinement is the challenge
 - Tokamaks are hot (150 million degrees)
 - Need magnetic confinement
 - Heat/Energy is escaping
- Loss mechanisms are an area of intense investigation
- Tokamaks often are heated by beams which excite coherent waves contributing to loss of energy

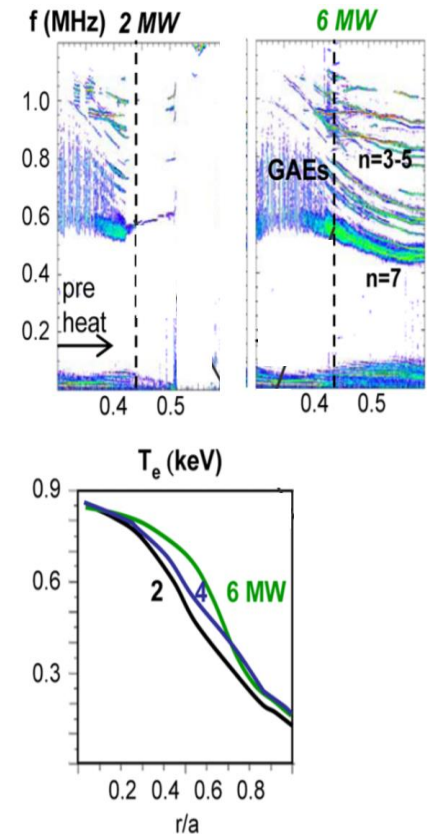
NSTX investigates fusion physics

- Spherical Tokamak
- D+ plasmas
- Temperatures ~ 1 KeV
- $B \sim 0.35 - 0.55$ T
- Current ~ 1 Mamp
- Density $\sim 10^{20} m^{-3}$



Compressional Alfvén Eigenmodes (CAEs) could play a significant role in plasma energy confinement

- Beams excite long wavelength, high frequency ($\omega \lesssim \omega_{ci}$) Alfvén Eigenmodes (AE) including Compressional AE (CAE) in NSTX
- Anomalous temperature profile flattening observed to correlated with high frequency activity in CAEs in the core of National Spherical Torus eXperiment (NSTX) plasma
- Theory suggest CAEs could contribute to the heat loss
 - Multiple mechanisms have been proposed
 - One candidate is CAE-KAW coupling
 - This theory still requires experimental validation



Stutman, PRL 2009

CAE-KAW coupling may explain anomalous flattening

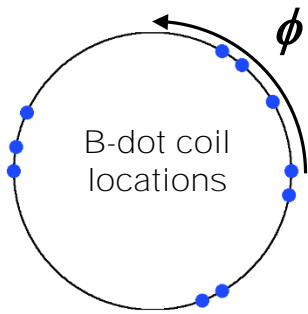
- CAE takes energy from core localized beam ions.
- CAE couple to kinetic Alfven waves (KAW) at resonance location
 - CAE dispersion relationship: $\omega_{cae}^2 = \mathbf{k}^2 v_A^2$
 - Resonance location is where $\mathbf{k} \rightarrow k_{\parallel} \Rightarrow$ KAW excitation
 - Energy dissipated during KAW excitation
 - KAW energy absorbed during propagation
- KAW are short wavelength (in contrast to CAE)
 - $\omega_{kaw}^2 = k_{\parallel}^2 v_A^2 [1 + k_{\perp}^2 \rho_s^2 (1 + \frac{3T_I}{4T_E})]$
- Theory applies for Global Alfven Eigenmodes as well [Kolesnichenko, PRL 2010]

Magnetic diagnostic sensitive to CAEs

- \dot{B} measured from conducting loop at edge NSTX cross-section
- Faraday's law

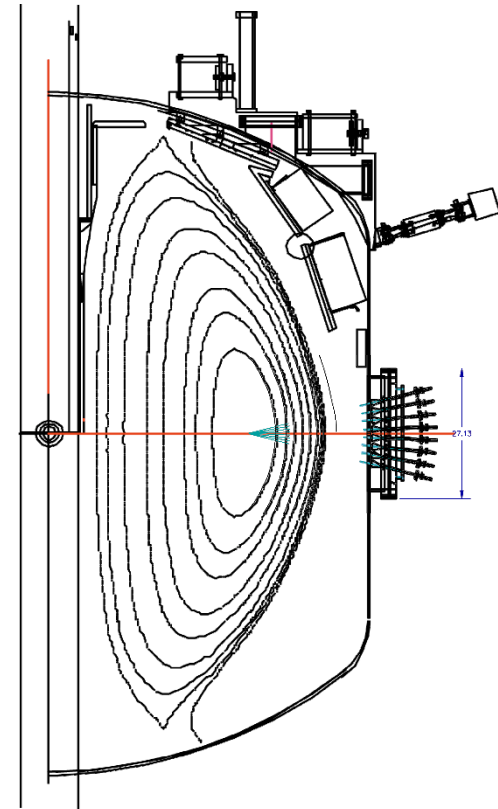
$$V_B = -\frac{d}{dt} \oint \mathbf{B}(t) \cdot d\mathbf{s}$$

- Toroidal mode number (n) via 10 toroidally distributed coils



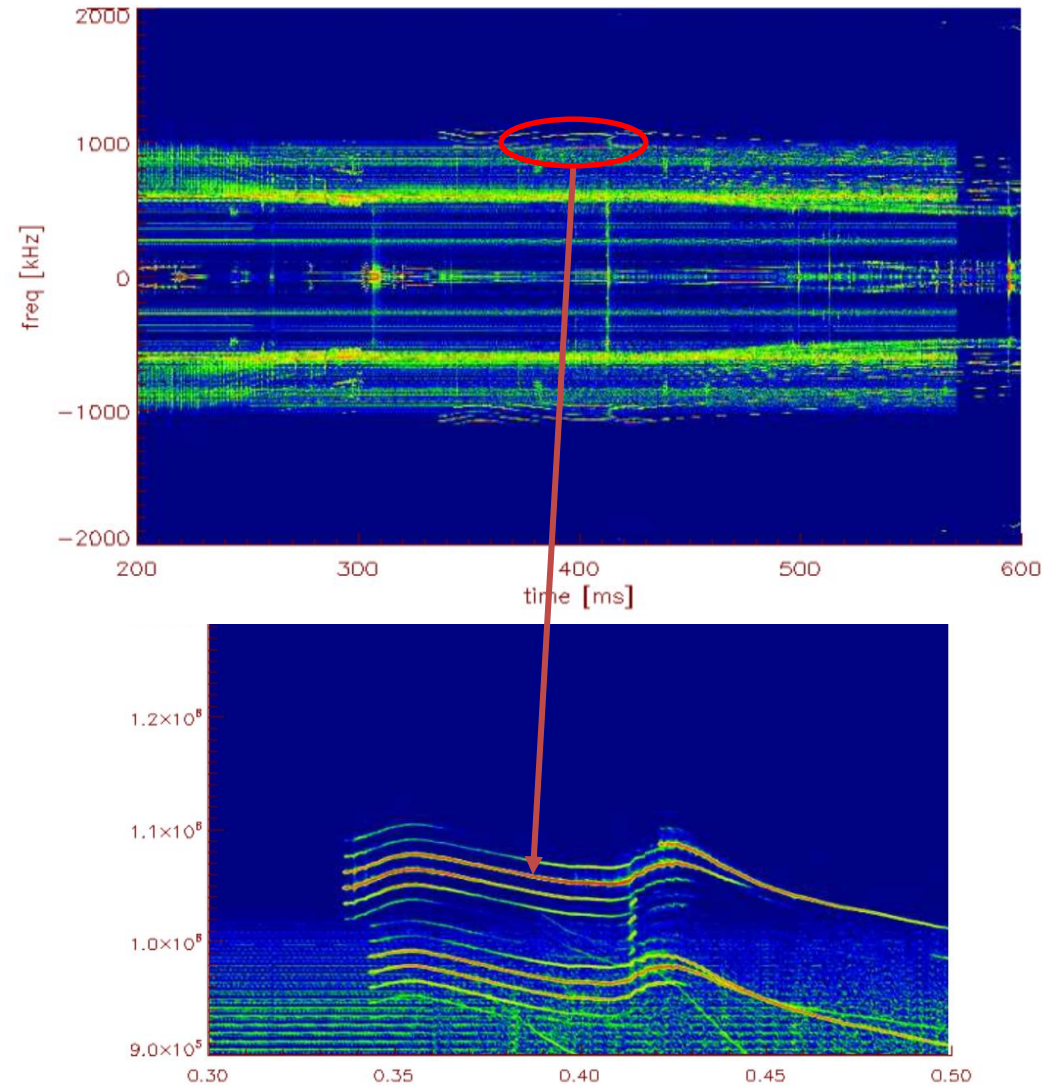
Top Down view of NSTX

$$\chi^2 \equiv 1 - \left| \sum_{\forall \phi} \delta b e^{-in\phi} \right|^2 / \left(N_\phi \sum_{\forall \phi} |\delta b|^2 \right)$$

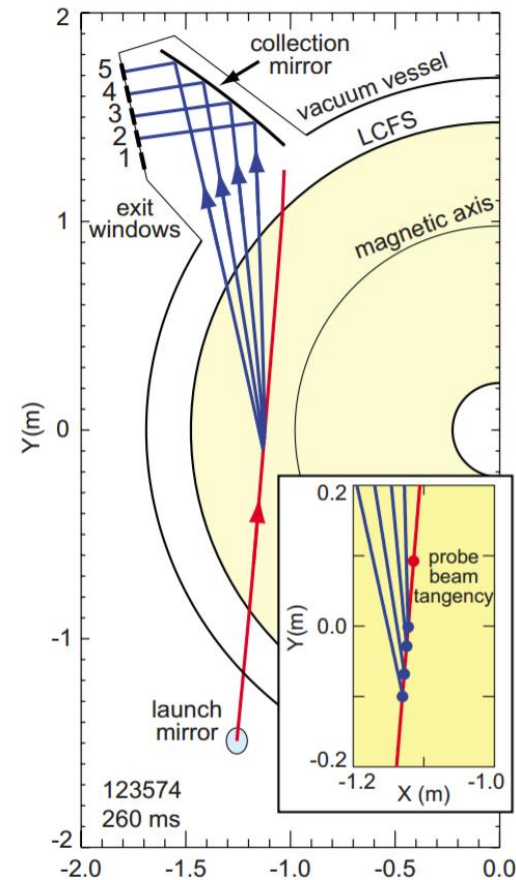


Magnetic diagnostic sensitive to CAEs

- Multiple peaks observed in magnetic spectrum at $f \gtrsim 1\text{MHz}$
- Identified as CAE because f at large fraction of f_{ci}



High K-Scattering diagnostic potentially detects KAWs from CAE-KAW coupling



- 280 GHz millimeter wave probe beam scatters into 5 receivers
- Each channel is sensitive to a different **K**

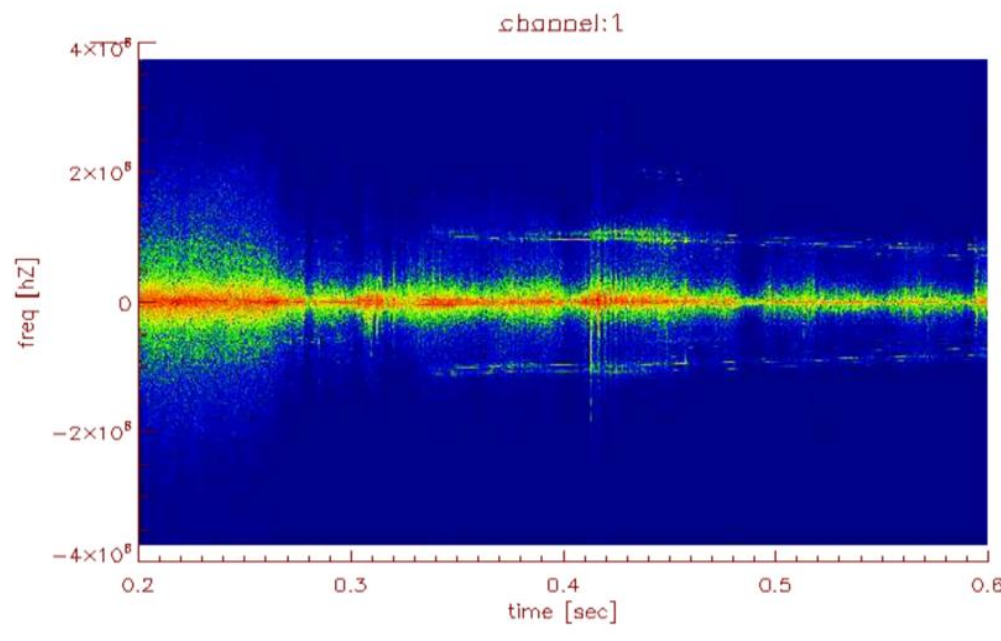
Bragg Scattering

$$\omega_i + \omega_{\delta n} = \omega_{scatter}$$

$$\mathbf{k}_i + \mathbf{k}_{\delta n} = \mathbf{k}_{scatter}$$

- δn is density fluctuation of KAW or turbulence
- beams refracted that changes over time, it has been raytraced, and K's for different channels are determined.

High K-Scattering diagnostic potentially detects KAWs from CAE-KAW coupling



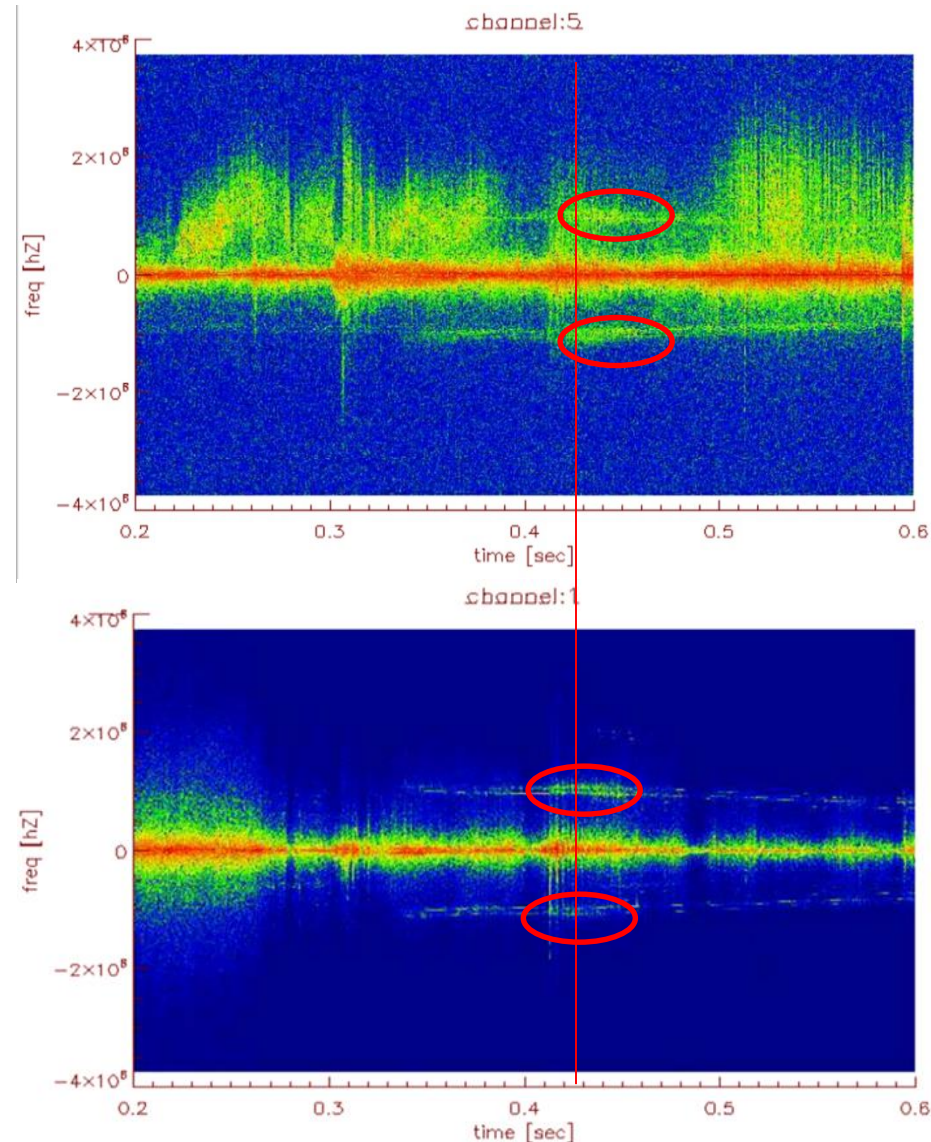
$$E_{hk}(t) \sim e^{i(\omega_s t + \epsilon \sin(\omega_{CAE} t))} + R e^{i(\omega_{KAW} t)}$$

– Two contributions to scattered electric field: **Turbulence** and **KAW**

- ω_s - frequency of turbulence, $\omega_s \ll \omega_{CAE}$
- ω_{CAE} - CAE frequency; large wavelength fluctuations (amplitude ϵ) modulate E-field scattered from turbulence.
- $\omega_{KAW} = \omega_{CAE}$
- R – relative scattering power

Scattering diagnostic probes KAW wavelengths

- $t_e \sim 780 \text{ eV}$
- $t_i \sim 920 \text{ eV}$
- $n_e \sim 5.1e13$
- $B \sim 0.36 \text{ T}$
- Channel 1: $k \sim 7.4 \text{ cm}^{-1}$
- Channel 2: $k \sim 7.6 \text{ cm}^{-1}$
- Channel 3: $k \sim 8.9 \text{ cm}^{-1}$
- Channel 4: $k \sim 11.5 \text{ cm}^{-1}$
- Channel 5: $k \sim 14.6 \text{ cm}^{-1}$
- $\rho_s \sim 1.1 \text{ cm} \rightarrow k\rho_s > 1$
- $T = 0.4483 \text{ sec}$



High K combined with magnetics allows detection of CAE-KAW coupling

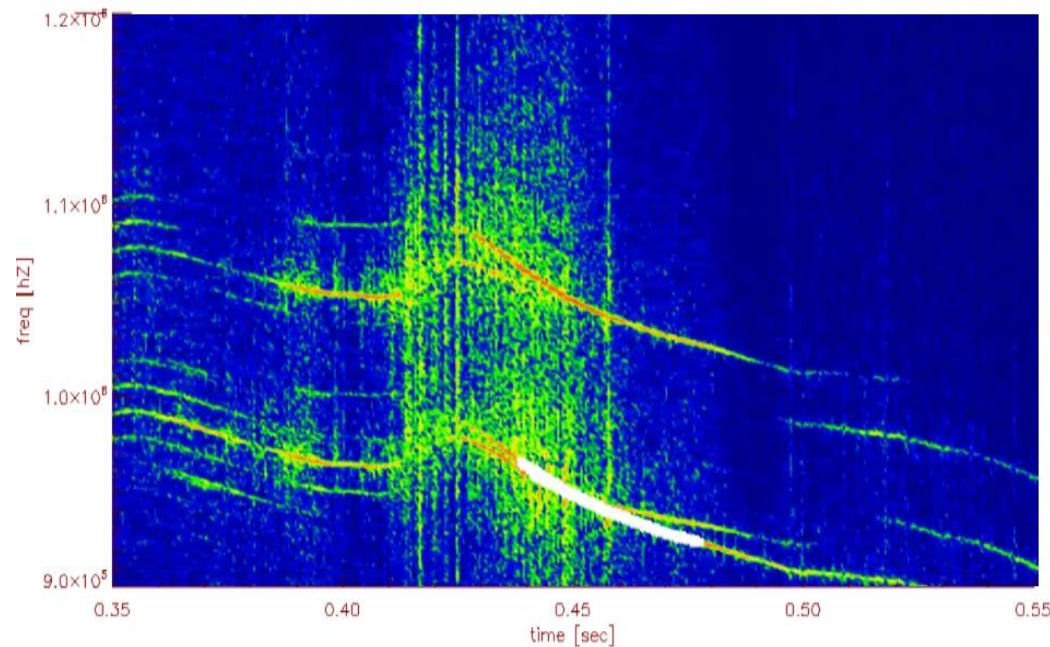
- Both diagnostic show peaks at similar frequency in CAE range → CAE-KAW coupling?
- Expect coupling to produce KAW coherent with CAE → must be verified with analysis
- Verify peak in high k spectrum results from scattering.
- Verify if detected k consistent with KAW theory.
 - Some Channels see peaks, other's don't.

Establishing coherence of high k spectrum peak with magnetic peak is a multistep process

- Isolates mode oscillation from each signal.
 - Determine $f(t)$ for peak \rightarrow filter signals to isolate mode
- Using SVD, isolate global mode oscillation (from mag)
- Calculate coherence of global signal with mode oscillation in scattering signal
 - Test statistical significance for each high-k channel.

Filtering: Determine times and frequencies with image feature analysis

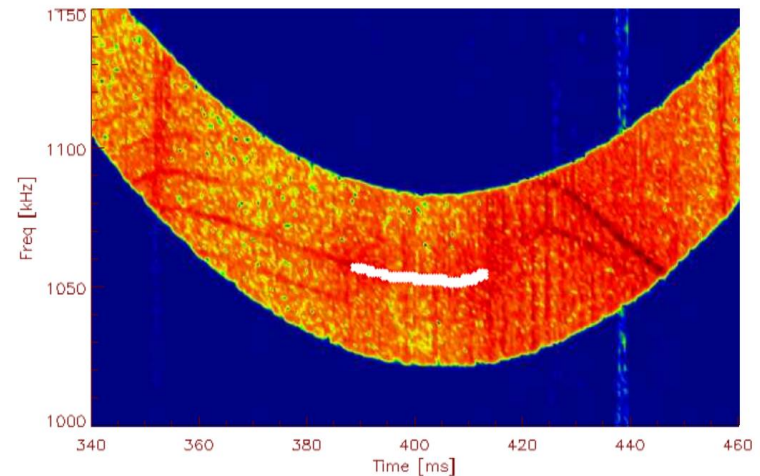
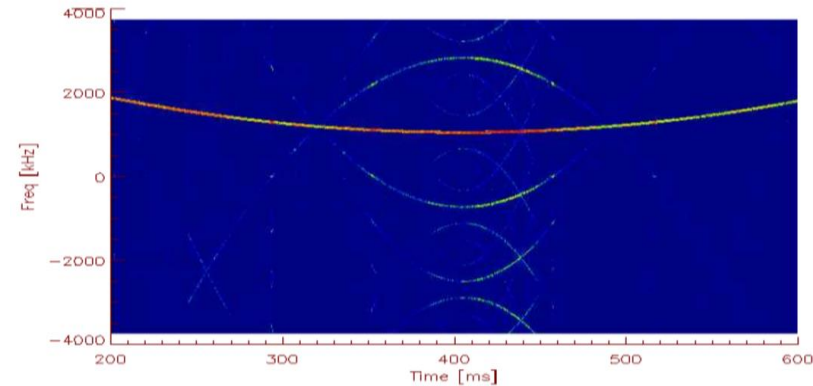
- Track time evolution of mode by choosing a time, frequency range, and a frequency bandwidth.
- Identified $f(t)$ can be used to create filter to isolate mode.



Feature analysis finds mode in spectrum.

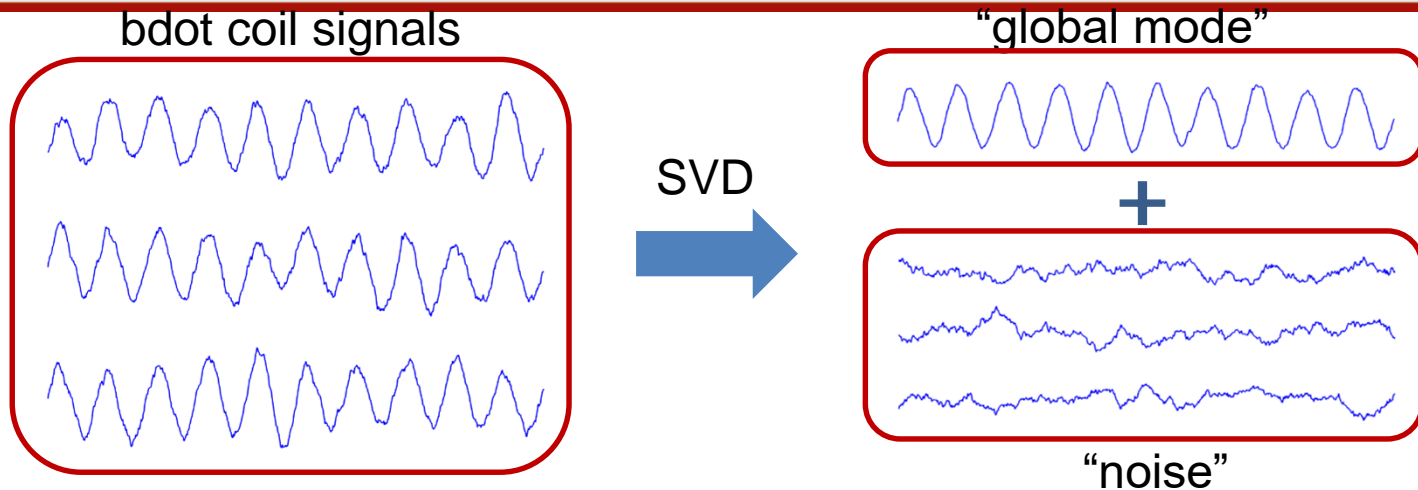
Filtering: Construct time dependent bandpass filter from times and frequencies

- Fit an analytic (polynomial) $f(t)$ to (t,f) from image analysis
- Integrate $f(t)$ to get $\Phi(t)$
- $e^{-i\Phi(t)}$ is used to Dopplershift to DC
- Low pass filter used to isolate mode.
- Un-Dopplershift with $e^{i\Phi(t)}$



$$S(t) = e^{i\phi(t)} \int_{-\Delta\omega}^{\Delta\omega} d\omega e^{i\omega t} \int_t^{t+T} dt' e^{-i\omega t'} e^{-i\phi(t')}$$

SVD: Singular value decomposition gives better “global mode” δb



- global mode observed by 10 bdot coils (HN array)
- “filter” signals with SVD \Rightarrow global mode w/reduced noise
 - SVD factors space & time dependence of signal matrix:

$$b_{jk} = \tilde{b}_j(t_k) \rightarrow \tilde{b}_{0j} \tilde{b}_{global}(t_k) + \epsilon_j(t_k)$$

- Steps *before* SVD ...

- 1) bandpass filter coil signals to isolate mode
- 2) make signals complex \Rightarrow spatial phase (e.g. $n\phi_j$) factors out automatically:

$$\tilde{b}_j(t) = A(t) \cos(\theta(t) + \theta_{0j}) \rightarrow \hat{\tilde{b}}_j(t) = \frac{1}{\sqrt{2}} A(t) e^{i((\theta(t) + \theta_{0j}))} = \frac{1}{\sqrt{2}} \int_0^\infty d\omega e^{i\omega t} \int_{-\infty}^\infty dt' \tilde{b}(t') e^{-i\omega t'}$$

SVD: finds global mode from eigenvector of signal correlation matrix

- SVD solves factoring problem

$$\hat{\tilde{b}}_j(t_k) = \hat{\tilde{b}}_{0j} \hat{\tilde{b}}_{global}(t_k) + \hat{\epsilon}_j(t_k)$$

- by minimizing χ^2 :

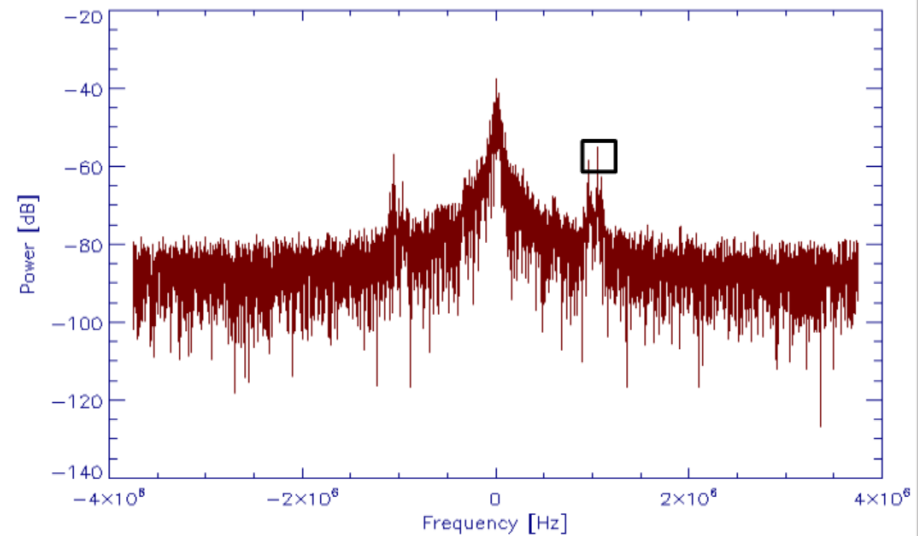
$$\chi^2 = \sum_{j,k} \left| \hat{\tilde{b}}_j(t_k) - \hat{\tilde{b}}_{0j} \hat{\tilde{b}}_{global}(t_k) \right|^2$$

- \Rightarrow spatial coefficients ($\hat{\tilde{b}}_{0j}$) of global mode from ***eigenvector*** of correlation matrix ***with largest eigenvalue***:

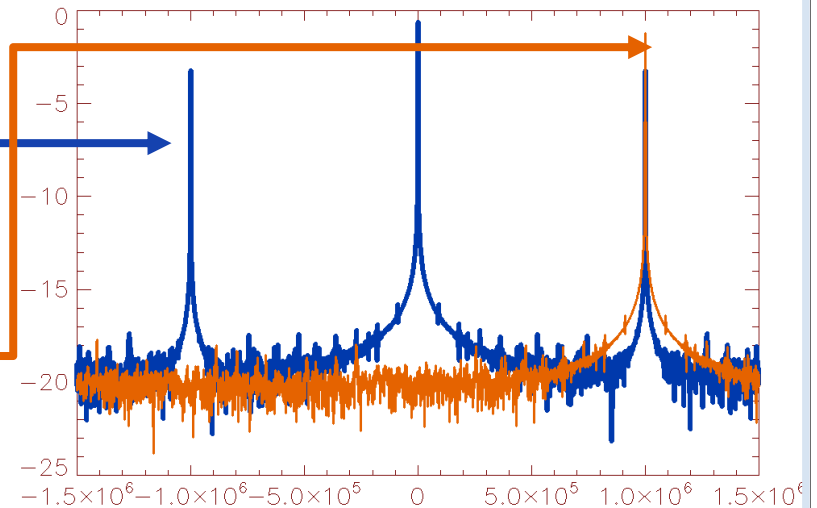
$$\begin{aligned} \mathbf{C} \hat{\tilde{\mathbf{b}}}_0 &= \lambda \hat{\tilde{\mathbf{b}}}_0 \\ [\mathbf{C}]_{ij} &= \left\langle \hat{\tilde{b}}_i(t) \hat{\tilde{b}}_j^*(t) \right\rangle, [\hat{\tilde{\mathbf{b}}}_0]_j = \hat{\tilde{b}}_{0j} \end{aligned}$$

Scattering Verification: Scattering produces an asymmetric spectrum

- Peaks at CAE frequency observed for 2 reasons.
 - Sidebands of scattering from low frequency turbulence
 - Scattering from KAW

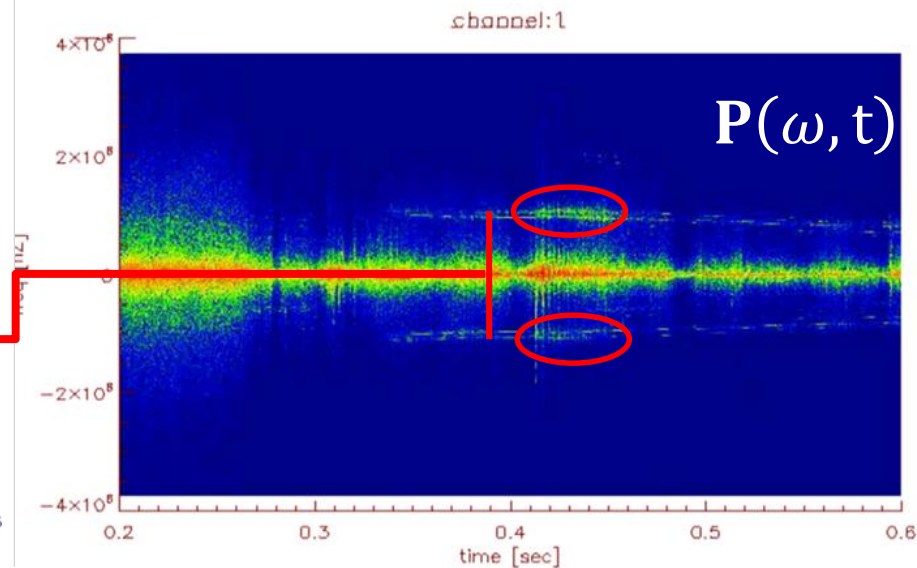
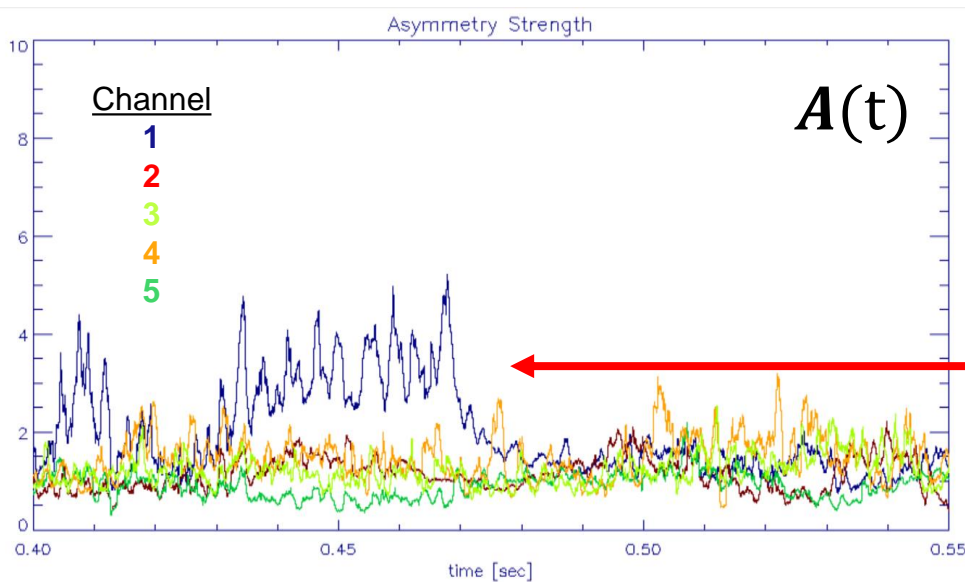


$$E_{hk}(t) \sim e^{i(\omega_s t + \epsilon \sin(\omega_{cae} t))} + R e^{i(\omega_{cae} t)}$$



Scattering Verification: Scattering produces an asymmetric spectrum

- Asymmetry strength across different channels is determined by dividing power at positive mode frequency by power by negative mode frequency
- $A(t) = \mathbf{P}(\omega_{CAE}(t), t) / \mathbf{P}(-\omega_{CAE}(t), t)$

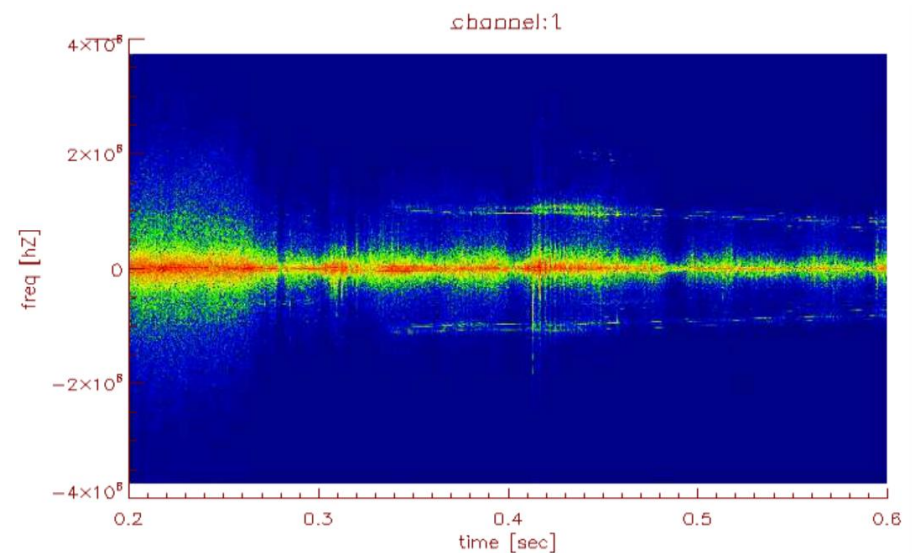
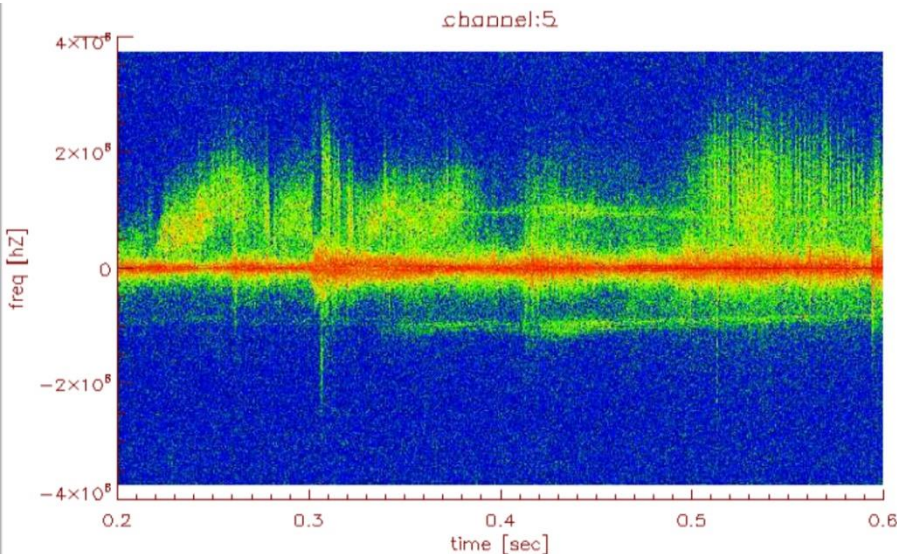


Coherence: Is scattering peak is coherent with magnetic peak as expected?

- Outgoing KAW from mode conversion should be coherent with incoming CAE
- $Coherence = \gamma = \frac{|\langle F_1(t)F_2^*(t) \rangle|}{(\langle |F_1|^2 \rangle \langle |F_2|^2 \rangle)^{1/2}}$
 - $\langle \rangle$ average over time
- Coherence is evaluated with filtered magnetic and scattering spectral peaks
- Resulting coherence is evaluated for statistical significance
 - Monte Carlo Simulation using identically filtered random signals

Comparison of high k channels shows peak at CAE frequency has well defined k

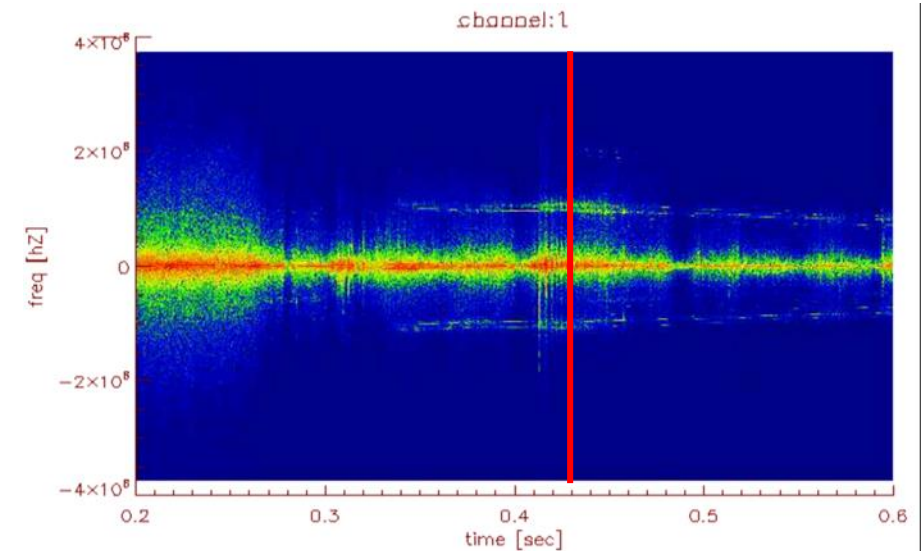
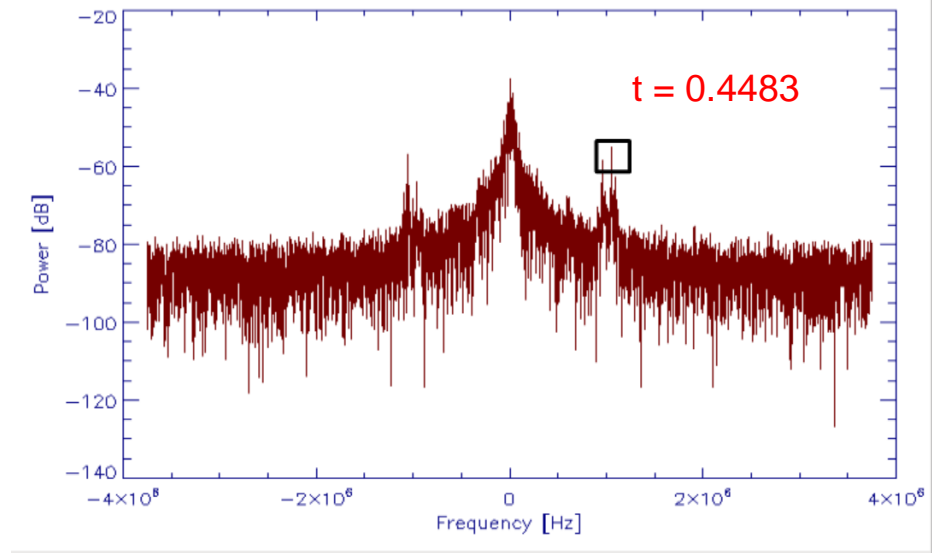
- Expect to see KAW predominantly in only one channel
- Channel sensitive to peak should correspond to k from theory
- Possible to have asymmetry from turbulence with low coherence



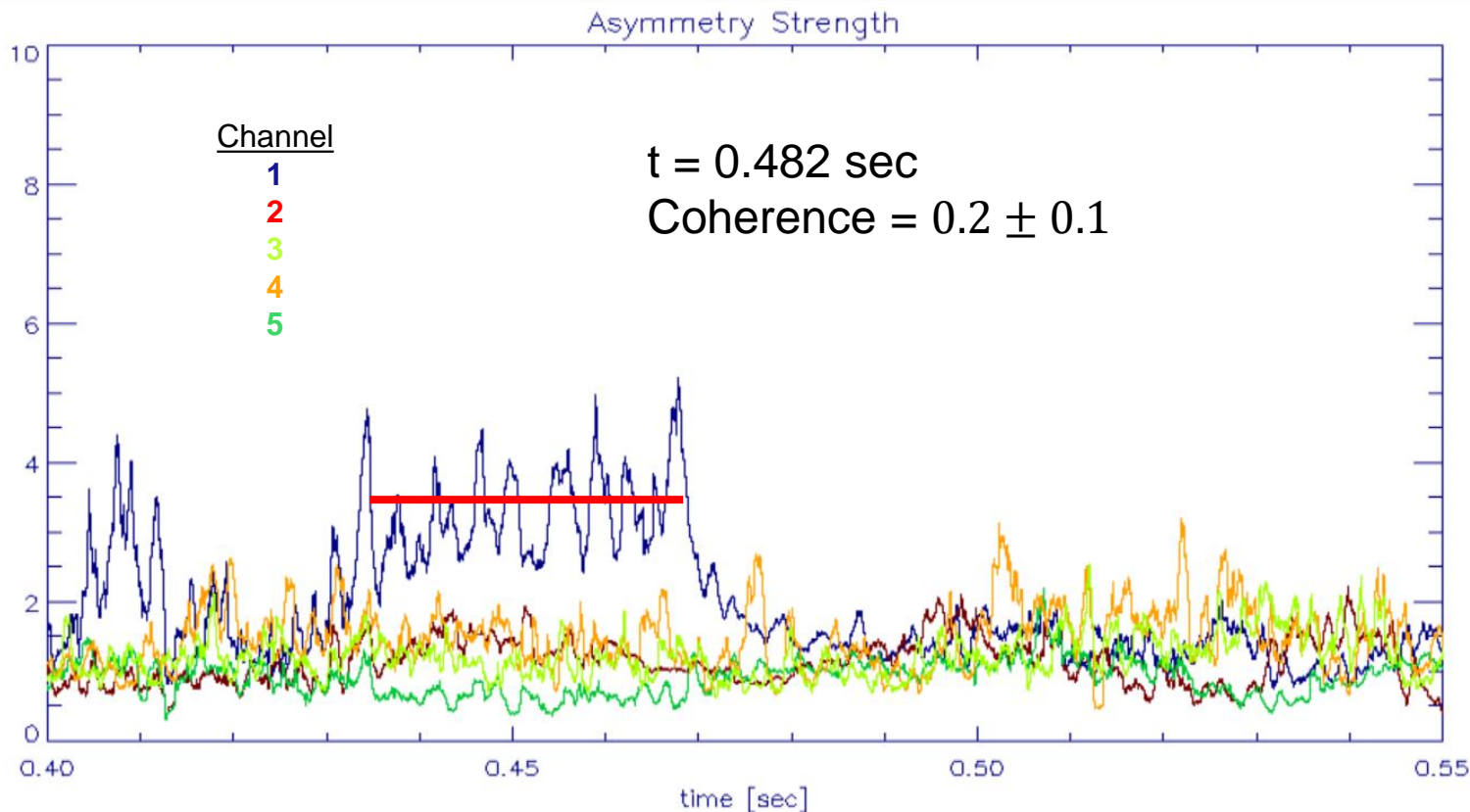
CAEs observed in shot 139395, analyzed for evidence of KAW coupling

- several modes identified
 - 10KHz bandwidth for filtering

t (sec)	f (Hz)	Δt (sec)
0.448	1.07e6	0.046
0.481	1.06e6	0.035
0.487	1.04e6	0.132
0.498	9.11e5	0.047



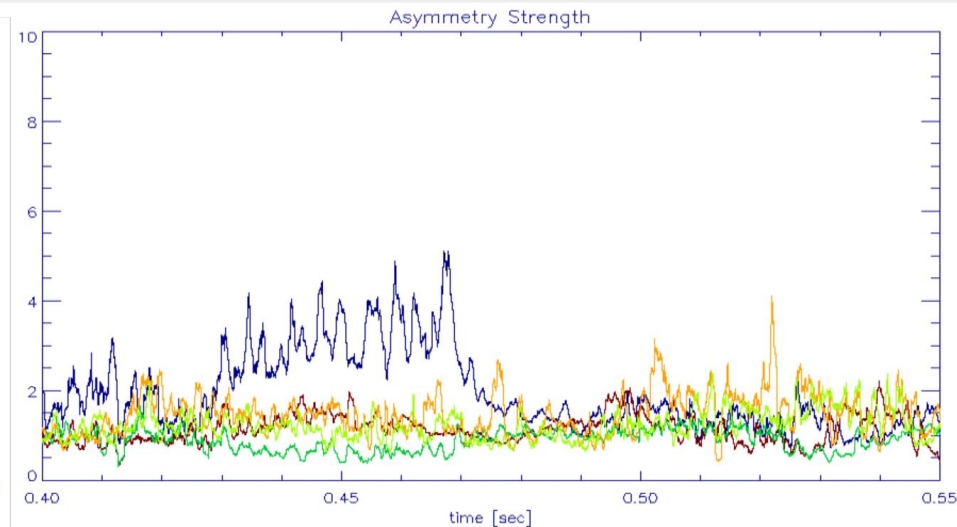
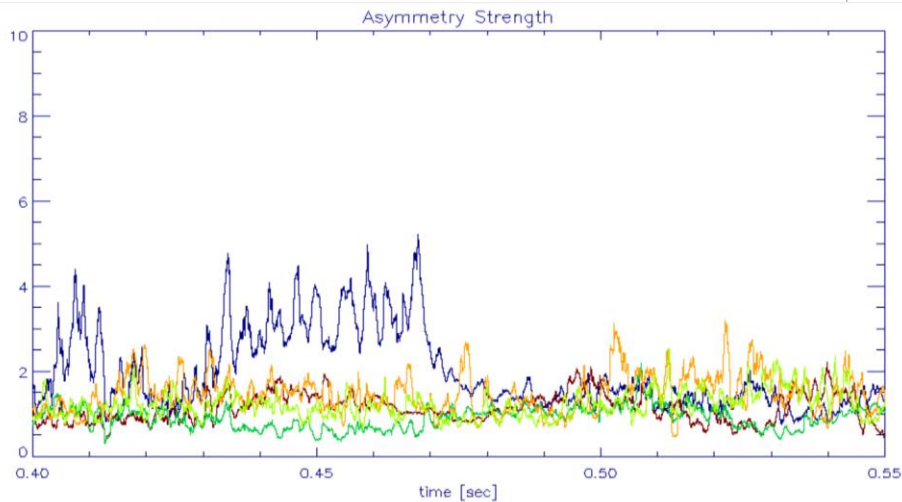
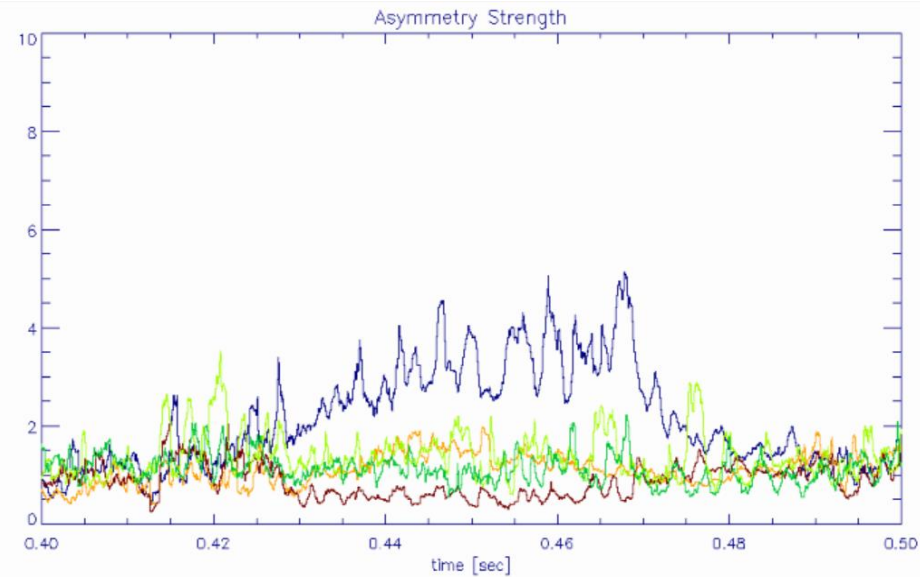
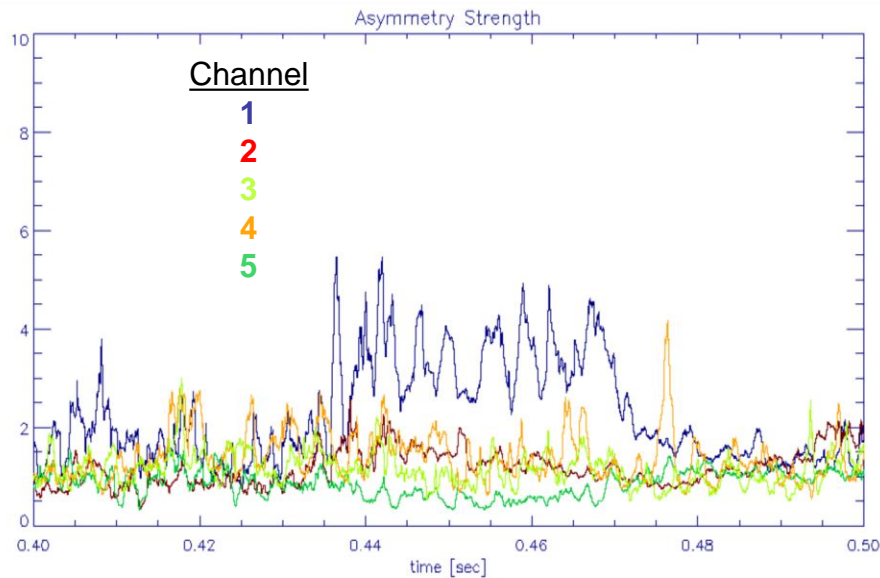
Asymmetry varies with scattered k, as expected



$t(sec)$ $f(Hz)$	A
0.448 1.07e6	6.2
0.481 1.06e6	3.8
0.487 1.04e5	6.3
0.493 9.11e5	4.1

- Channel one ($k = 760 \text{ m}^{-1}$) typically demonstrate a high asymmetry at times where mode is present.

Asymmetry observed in high k channels for multiple modes



Coherence found between magnetics and scattering found for following modes

Coherence	Time
0.1777	0.4483
0.111	0.4817
0.20	0.4871
0.0636	0.4983

Observation consistent with CAE-KAW coupling

- Peaks at same frequency in magnetic and scattering signal are coherent
- Comparison of scattering signals at different k demonstrate well defined k for peaks at CAE freq
- Peaks at lower k spectra demonstrate high asymmetry, in contrast with high k spectra
- Coherence + Asymmetry strongly implies scattering from KAW

Future work

- Compare scattering measurement with expectation from KAW dispersion relationship
 - Predict which channels see KAWs as profiles evolve
 - Assess absorption to determine if significant power can reach measurement location
- Evaluate potential contamination of scattering signal by refraction of probe beam.
 - Beam tracing

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