

Advanced Plasma Shape Control to Enable High-Performance Divertor Operation on NSTX-U

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Abstract

This work presents the development of an advanced framework for control of the global plasma shape and its application to a variety of shape control challenges on NSTX-U. Operations in high-performance plasma scenarios will require highly-accurate and robust control of the plasma poloidal shape to accomplish such tasks as obtaining the strong-shaping required for the avoidance of MHD instabilities and mitigating heat flux through regulation of the divertor magnetic geometry. The new control system employs a high-fidelity model of the toroidal current dynamics in NSTX-U poloidal field coils and conducting structures as well as a first-principles driven calculation of the axisymmetric plasma response. The model-based nature of the control system enables real-time optimization of controller parameters in response to time-varying plasma conditions and control objectives. The new control scheme is shown to enable stable and on-demand plasma operations in complicated magnetic geometries such as the snowflake divertor. A recently-developed code that simulates the nonlinear evolution of the plasma equilibrium is used to demonstrate the capabilities of the designed shape controllers. Plans for future real-time implementations on NSTX-U and elsewhere are also presented.



Introduction

Goals of this work

- The use of advanced divertor configurations, such as the snowflake divertor (SFD), is being considered as a possible means of reducing peak heat flux onto divertor surfaces in NSTX-U.
- Develop an algorithm that is capable of real-time control of all divertor configurations of interest in NSTX-U.
- Address primary limitations of the algorithm as previously implemented on DIII-D:
 - Stable control of the SFD-Plus configuration.

E. Kolemen et al. J. Nucl. Mater. (2015).

- Recovery from high-field-side to low-field-side SFD-Minus.
- Transition to model-based (non-PID) plasma shape control for NSTX-U.

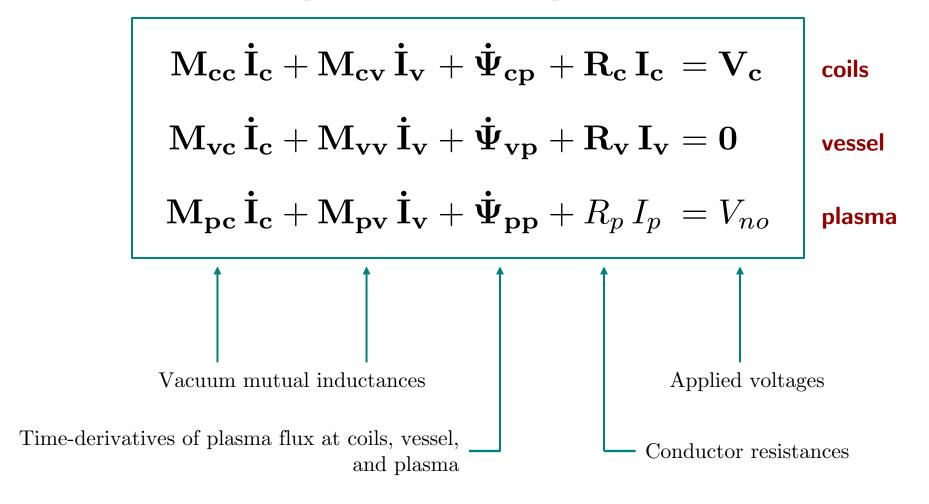
Highlights

- New modeling of X-point position response to PF coil currents.
- PID-based control of the SFD with closed-loop controller tuning using relay feedback.
- Initial development of model-based Linear-Quadratic-Integral control of the plasma shape.



Formalism for plasma shape control modeling

Coupled circuit equations describing dynamics of toroidal currents in coils, passive structures, and plasma.





Linearized circuit equations

$$\begin{aligned} \mathbf{M_{cc}}\,\dot{\mathbf{I_c}} + \mathbf{M_{cv}}\,\dot{\mathbf{I_v}} + & \frac{\partial \Psi_{\mathbf{cp}}}{\partial \mathbf{I_c}}\,\dot{\mathbf{I_c}} + \frac{\partial \Psi_{\mathbf{cp}}}{\partial \mathbf{I_v}}\,\dot{\mathbf{I_v}} + \frac{\partial \Psi_{\mathbf{cp}}}{\partial I_p}\,\dot{I_p} + \mathbf{R_c}\,\mathbf{I_c} = \mathbf{V_c} \\ \mathbf{M_{vc}}\,\dot{\mathbf{I_c}} + \mathbf{M_{vv}}\,\dot{\mathbf{I_v}} + & \frac{\partial \Psi_{\mathbf{vp}}}{\partial \mathbf{I_c}}\,\dot{\mathbf{I_c}} + \frac{\partial \Psi_{\mathbf{vp}}}{\partial \mathbf{I_v}}\,\dot{\mathbf{I_v}} + \frac{\partial \Psi_{\mathbf{vp}}}{\partial I_p}\,\dot{I_p} + \mathbf{R_v}\,\mathbf{I_v} = \mathbf{0} \\ \mathbf{M_{pc}}\,\dot{\mathbf{I_c}} + \mathbf{M_{pv}}\,\dot{\mathbf{I_v}} + & \frac{\partial \Psi_{\mathbf{pp}}}{\partial \mathbf{I_c}}\,\dot{\mathbf{I_c}} + \frac{\partial \Psi_{\mathbf{pp}}}{\partial \mathbf{I_v}}\,\dot{\mathbf{I_v}} + \frac{\partial \Psi_{\mathbf{pp}}}{\partial I_p}\,\dot{I_p} + R_p\,I_p = V_{n.o.} \end{aligned}$$

Linearized response of plasma flux due to changes in external currents and bulk plasma current.

$$\mathbf{X}_{jk} = \frac{\partial \mathbf{\Psi}_{jp}}{\partial \xi_r} \frac{\partial \xi_r}{\partial \mathbf{I}_k} + \frac{\partial \mathbf{\Psi}_{jp}}{\partial \xi_z} \frac{\partial \xi_z}{\partial \mathbf{I}_k} \qquad j, k \in \{c, v, p\}$$

Outputs of a linear plasma response model



State-space representation of the dynamics

Express the linearized circuit equations in state-space form for use with model-based control design tools

$$\delta \dot{\mathbf{x}} = \mathbf{A}(t) \, \delta \mathbf{x} + \mathbf{B}(t) \, \delta \mathbf{u}$$

Perturbed currents

$$\delta \mathbf{x} = egin{bmatrix} \mathbf{I_c} - \mathbf{I_c}_{eq} \ \mathbf{I_v} - \mathbf{I_v}_{eq} \ I_p - I_{p_{eq}} \end{bmatrix}$$

Perturbed voltages

$$\delta \mathbf{v} = \begin{bmatrix} \mathbf{V_c} - \mathbf{V_c}_{eq} \\ 0 \\ 0 \end{bmatrix}$$

$$\mathbf{A}(t) = -\left[\widehat{\mathbf{M}}(t)\right]^{-1}\mathbf{R}$$
 Resistance matrix \mathbf{R}
 $\mathbf{B}(t) = \left[\widehat{\mathbf{M}}(t)\right]^{-1}\mathbf{V}$ Map from voltages \mathbf{V} to coils

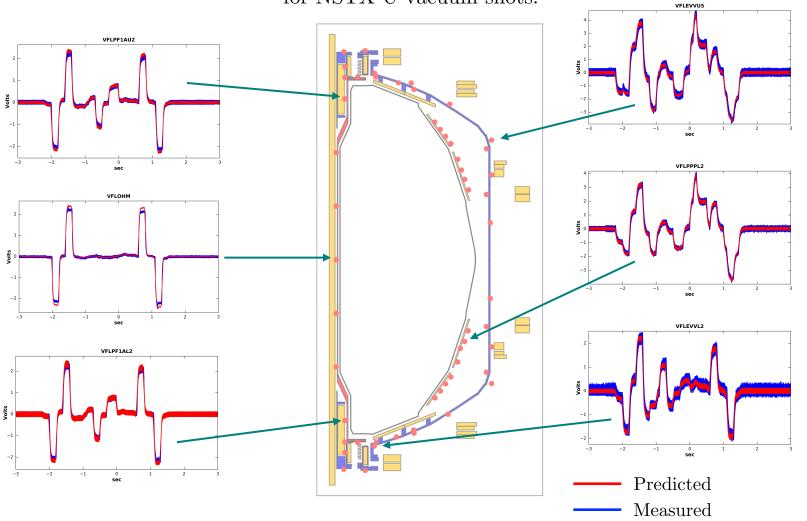
$$\widehat{\mathbf{M}}(t) = \mathbf{M} + \mathbf{X}(t)$$

Vacuum mutual inductances

Effective mutual inductance due to plasma motion

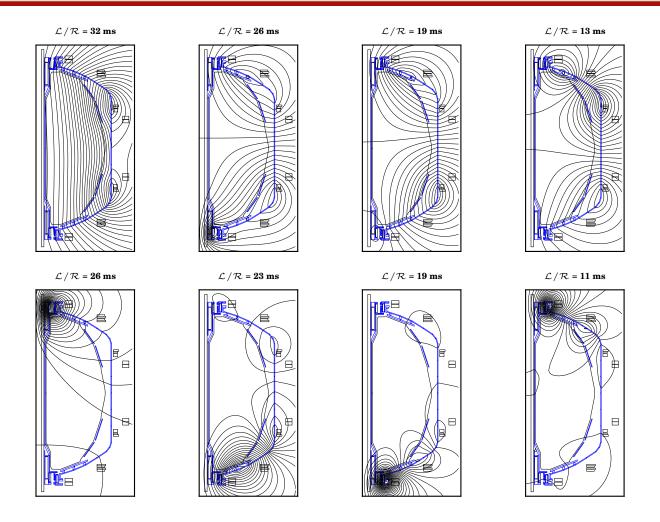
Validation of the no-plasma model

Wall model validated by comparing synthetic and measured magnetic diagnostic signals for NSTX-U vacuum shots.





Dominant eigenmodes of the vacuum vessel



Future validation efforts will seek to identify source of the asymmetry in the vessel model.



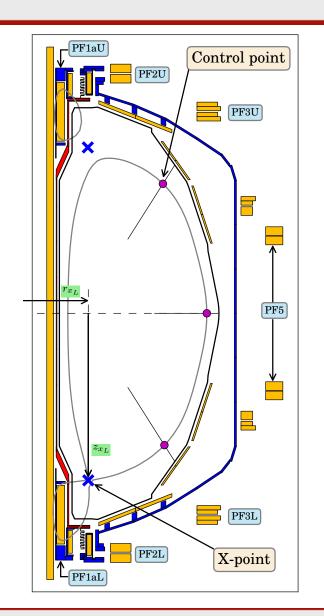
Output equation for ISOFLUX control

State-space dynamics equation paired with output equation relating the inputs (voltages) and states (currents) to quantities of interest for control.

$$\delta \mathbf{y} = \mathbf{C}(t) \, \delta \mathbf{x} + \mathbf{D}(t) \, \delta \mathbf{u}$$

$$\delta \mathbf{y} = \begin{bmatrix} \Delta \psi_1 \\ \Delta \psi_2 \\ \Delta \psi_3 \end{bmatrix} \begin{bmatrix} \text{Control} \\ \text{point fluxes} \end{bmatrix} \mathbf{C}(t) = \begin{bmatrix} \partial_{\mathbf{I}} \psi_1 \\ \partial_{\mathbf{I}} \psi_2 \\ \partial_{\mathbf{I}} \psi_3 \end{bmatrix} \mathbf{C}(t) = \begin{bmatrix} \partial_{\mathbf{I}} \psi_1 \\ \partial_{\mathbf{I}} \psi_2 \\ \partial_{\mathbf{I}} \psi_3 \end{bmatrix} \mathbf{C}(t) = \begin{bmatrix} \partial_{\mathbf{I}} \psi_1 \\ \partial_{\mathbf{I}} \psi_2 \\ \partial_{\mathbf{I}} \psi_3 \\ \partial_{\mathbf{I}} z_{x_U} \\ \partial_{\mathbf{I}} z_{x_U} \\ \partial_{\mathbf{I}} z_{x_L} \end{bmatrix}$$

In general, matrix entries are time-dependent.



Modeling of X-point response

$$\frac{\partial r_x}{\partial \mathbf{I}} = \frac{\partial r_x}{\partial \mathbf{B}} \frac{\partial \mathbf{B}}{\partial \mathbf{I}} = \frac{\partial r_x}{\partial \mathbf{B}} \left(\frac{\partial \mathbf{B}}{\partial \mathbf{I}} \Big|_{\text{vac}} + \frac{\partial \mathbf{B}}{\partial \xi_r} \frac{\partial \xi_r}{\partial \mathbf{I}} + \frac{\partial \mathbf{B}}{\partial \xi_z} \frac{\partial \xi_z}{\partial \mathbf{I}} \right)$$

$$\frac{\partial z_x}{\partial \mathbf{I}} = \frac{\partial z_x}{\partial \mathbf{B}} \frac{\partial \mathbf{B}}{\partial \mathbf{I}} = \frac{\partial z_x}{\partial \mathbf{B}} \left(\frac{\partial \mathbf{B}}{\partial \mathbf{I}} \Big|_{\text{vac}} + \frac{\partial \mathbf{B}}{\partial \xi_r} \frac{\partial \xi_r}{\partial \mathbf{I}} + \frac{\partial \mathbf{B}}{\partial \xi_z} \frac{\partial \xi_z}{\partial \mathbf{I}} \right)$$

Green's functions for the coil-only vacuum fields

Outputs of a linear plasma response model

 $\partial_{\mathbf{B}} r_x$ and $\partial_{\mathbf{B}} z_x$ Computed analytically using saddle-point expansion of the magnetic flux

$$x = r - r_0$$
$$v = z - z_0$$

Flux

$$(r_0 + x) \frac{\partial}{\partial x} \left(\frac{1}{r_0 + x} \frac{\partial \psi}{\partial x} \right) + \frac{\partial^2 \psi}{\partial v^2} = 0$$

Field

$$(r_0 + x) \frac{\partial}{\partial x} \left(\frac{1}{r_0 + x} \frac{\partial \psi}{\partial x} \right) + \frac{\partial^2 \psi}{\partial v^2} = 0 \qquad B_r = -\frac{1}{r_0 + x} \frac{\partial \psi}{\partial v} \quad B_z = \frac{1}{r_0 + x} \frac{\partial \psi}{\partial x}$$

Response of X-point position to B-field

1. Expand the flux function as a series to second-order. Solve for the series coefficients using measurements of B_r and B_z at two points.

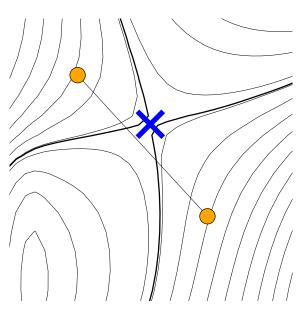
$$\psi(x,v) = l_1 x + l_2 v + q_1 x^2 + 2q_2 x v + q_3 v^2$$

$$B_r = -\frac{1}{r_0 + x} \left(l_2 + 2q_2 x + 2q_3 v \right)$$

$$B_z = \frac{1}{r_0 + x} \left(2(r_0 + x)q_1 + 2q_2 v + 2q_3 r_0 \right)$$

2. Solve for the (r,z) coordinates of the X-point.

$$r = r_0 + \frac{l_2 q_2 - l_1 q_3}{2(q_1 q_3 - q_2^2)}$$
 $z = z_0 + \frac{l_2 q_1 - l_1 q_2}{2(q_2^2 - q_1 q_3)}$

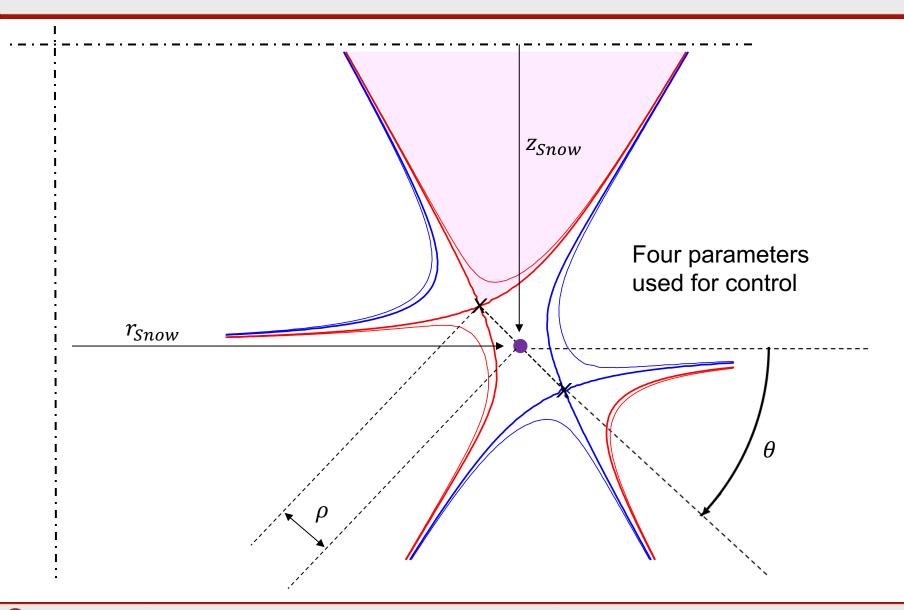


Sample points near the null

3. Compute derivatives of (r_x, z_x) with respect to B_r and B_z . $\partial_{\mathbf{B}} r_x$

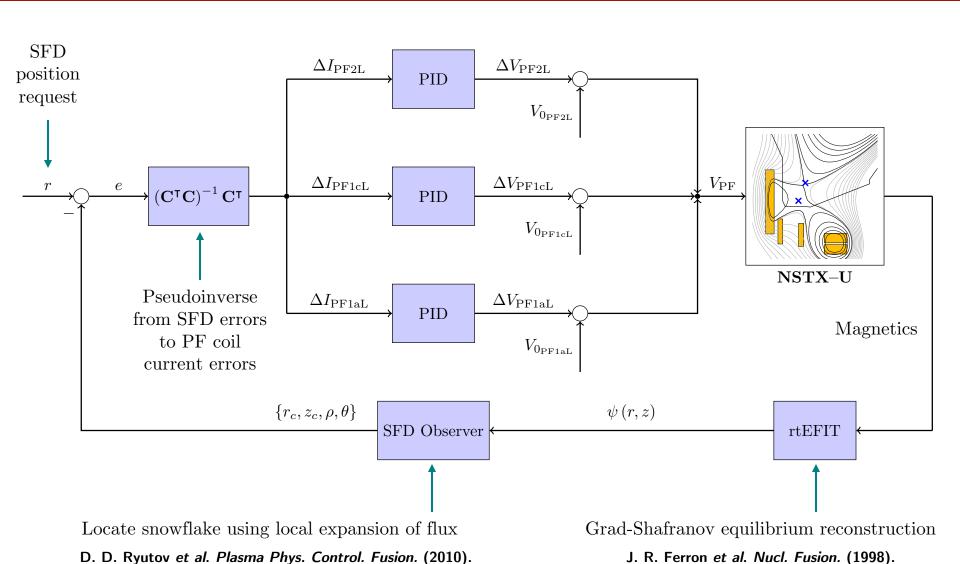
 $\partial_{\mathbf{B}} z_x$

Snowflake Shape Descriptors





SFD control system with PID control





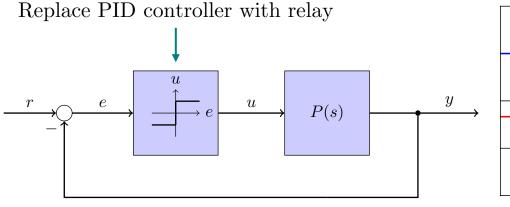
PID Control and Controller Tuning

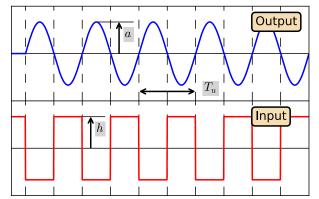
Compute control action from the error between desired and measured signal

$$u(t) = K_{\rm P} \left(e(t) + \frac{1}{T_{\rm I}} \int_0^t e(\tau) d\tau + T_{\rm D} \frac{d}{dt} e(t) \right)$$

Terms proportional to error, integral of error, and derivative of error

Closed-Loop tuning of PID gains using **Relay Feedback**

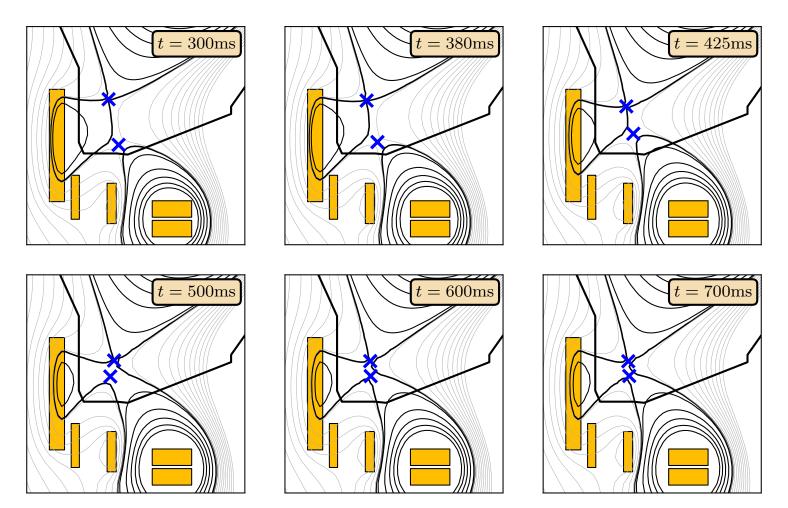




 $K_{\rm u} = \frac{4h}{\pi a}$ Ultimate
Gain

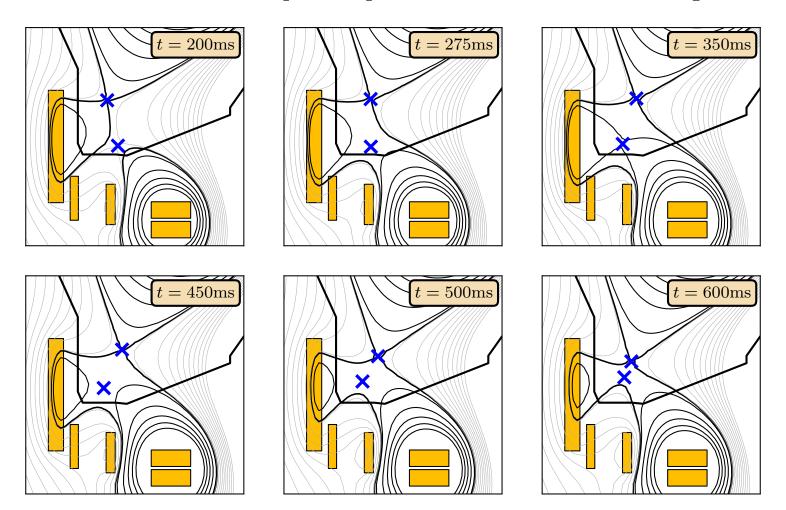
Snowflake radius scan in SFD-Minus

Scenario 1: Scan of the X-point separation in the SFD-Minus configuration.



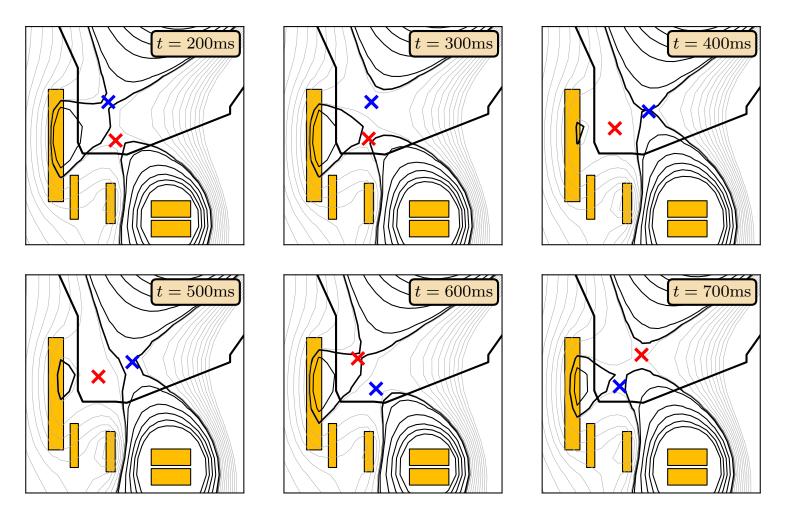
Snowflake radius scan in SFD-Plus

Scenario 2: Scan of the X-point separation in the SFD-Plus configuration.



Snowflake angle scan at constant separation

Scenario 3: Scan of the angular orientation with constant X-point separation.



Future work

- Test the designed controllers for the SFD in-the-loop with free-boundary Grad-Shafranov equilibrium solver (for verification of controller performance).
- Integration of the SFD control into a model-based shape controller designed with LQI for reference tracking of plasma shape parameters.
- Implementation of PID-based and LQI-based controllers for the SFD in the DIII-D and NSTX-U plasma control systems.
- Test the new algorithms in DIII-D SFD scenarios.