
Integrated Tokamak Modeling: when physics informs engineering and research planning

Francesca M. Poli

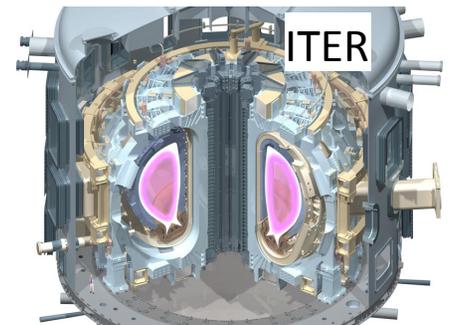
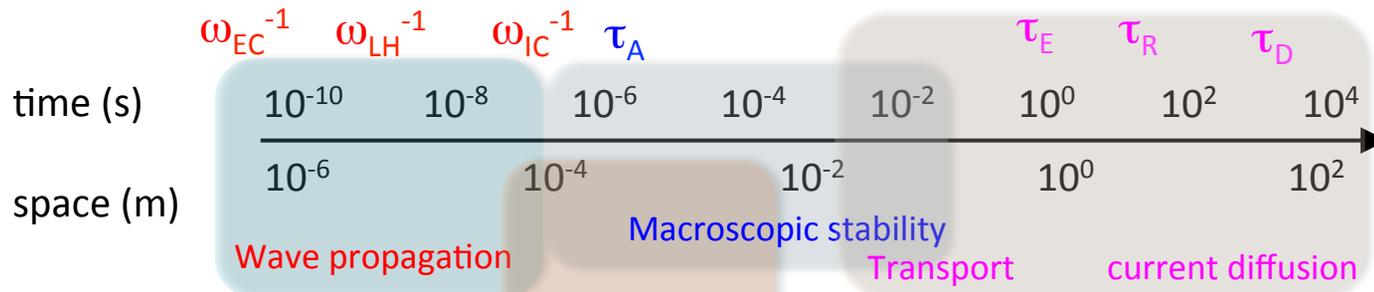
Princeton Plasma Physics Laboratory

For sending material for this tutorial, many thanks go to:

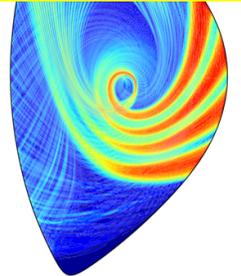
- USA: N. Bertelli, M. Boyer, N. Ferraro, W. Guttenfelder, D. Pfefferle, M. Podesta (PPPL), I. Krebs (PPPL/IPP), V. Soukhanovskii (LLNL), D. Green, J-M Park (ORNL), S. Kruger, T. Jenkins (TechX), P. Bonoli, N. Howard, S. Shiraiwa, J. Wright (PSFC), T. Rafiq (Univ. Lehigh)
- EU: J. Citrin (DiFFer), E. Militello Asp, F. Koechl (CCFE)
- JAPAN: N. Hayashi (QST), A. Fukuyama (Univ. Kyoto)
- ITER Organization: X. Bonnin, Y. Gribov, S-H Kim, R. Pitts



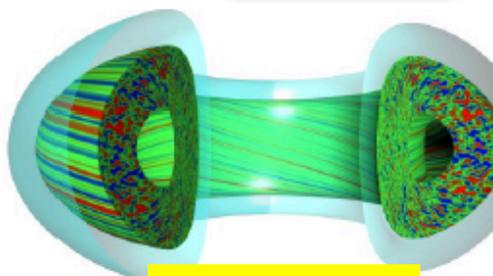
Fusion plasma physics encompasses a wide range of spatial and temporal scales



Full wave solvers
<1M CPU Hours

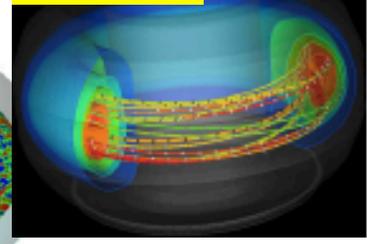


Plasma turbulence

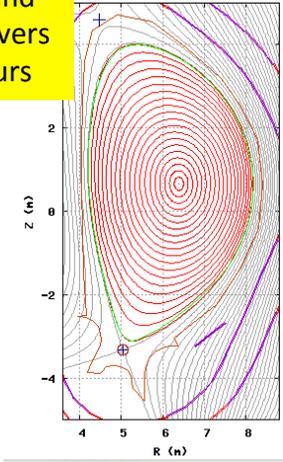


Gyrokinetic codes
<10M CPU hours

MHD codes



Equilibrium and transport solvers
<10² CPU hours



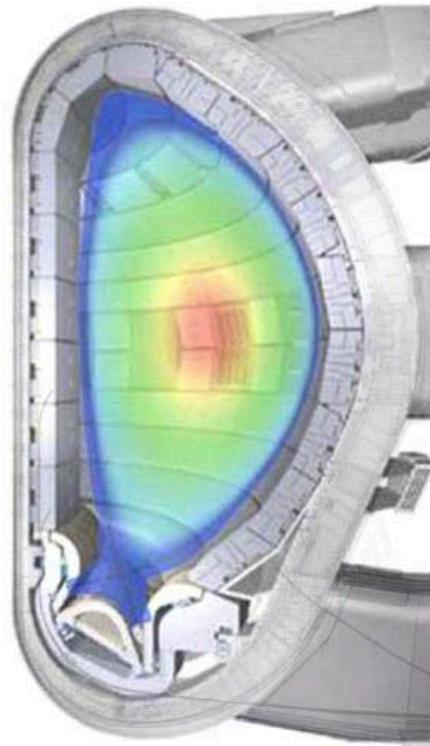
In a tokamak all temporal and spatial scales are coupled

Confined plasma
(closed magnetic field lines)

MHD equilibrium/instabilities,
microturbulence,
energetic particles

Scrape-Off Layer plasma (SOL)
(open field lines)

Microturbulence,
ionization, recombination
radiation



The plasma is surrounded
by solid structures:
Plasma-material interactions

External heating

Radiofrequency waves
Neutral beams

Fueling

Gas injection
Pellets



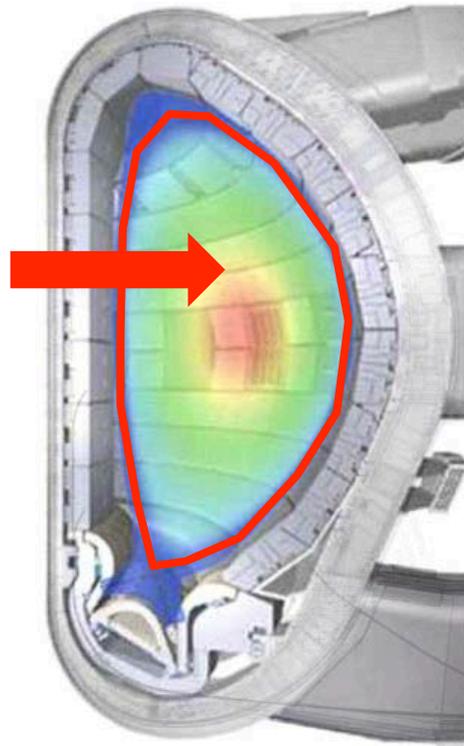
In a tokamak particles and energy are 'confined' by magnetic fields

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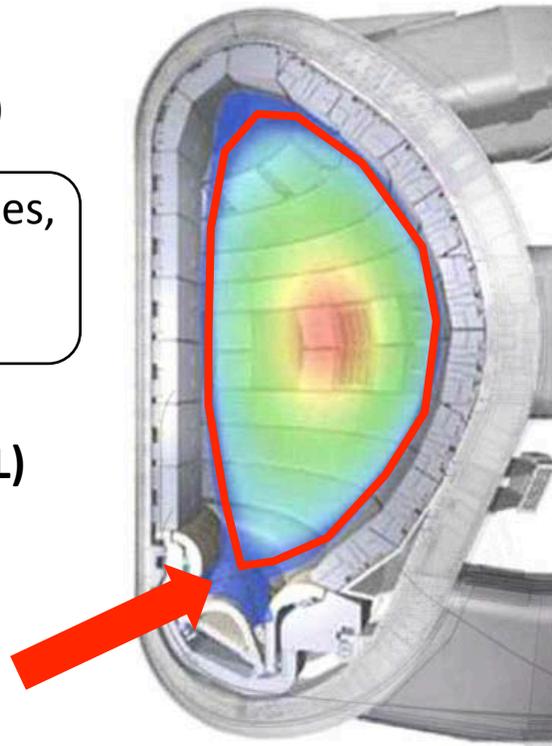
and they can also 'flow' along open magnetic field lines

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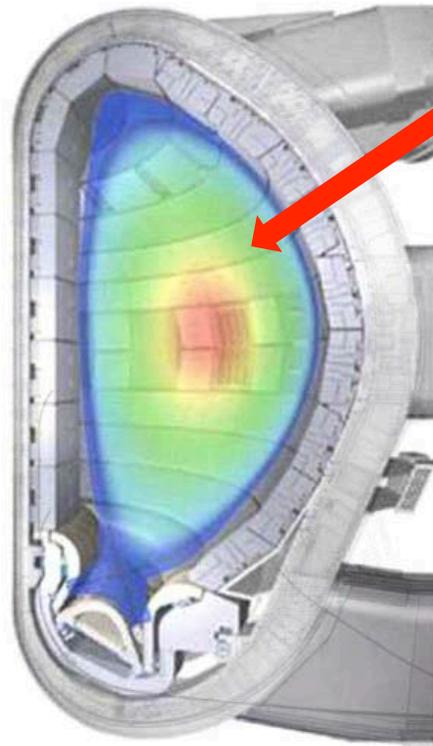
Plasmas in a tokamak are in contact with the 'wall'

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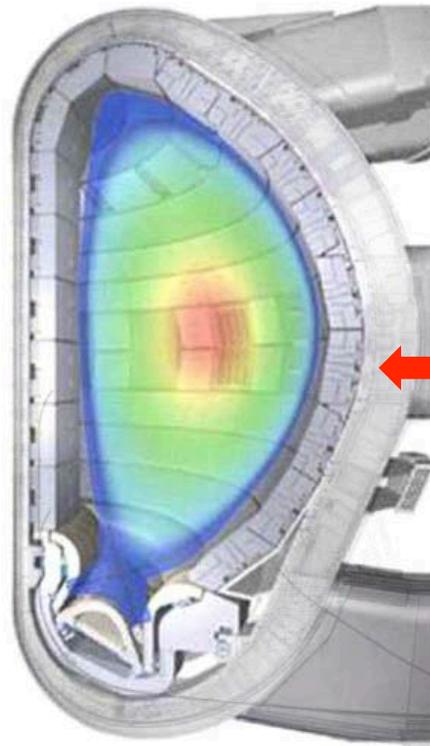
Plasmas in a tokamak need to be 'heated'

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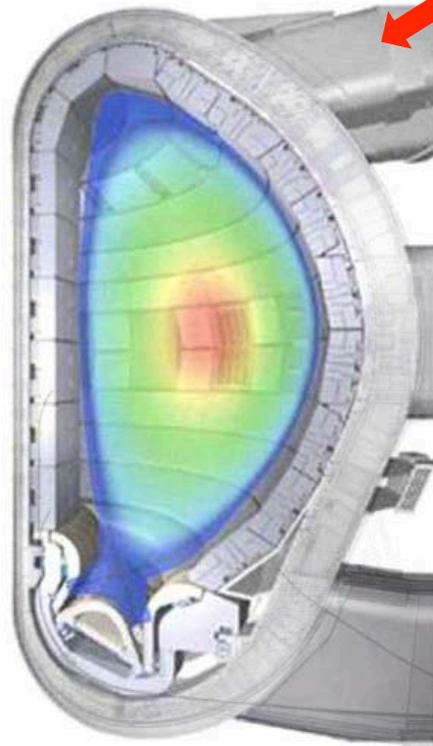
Plasmas need frequent pit stops for 're-fueling'

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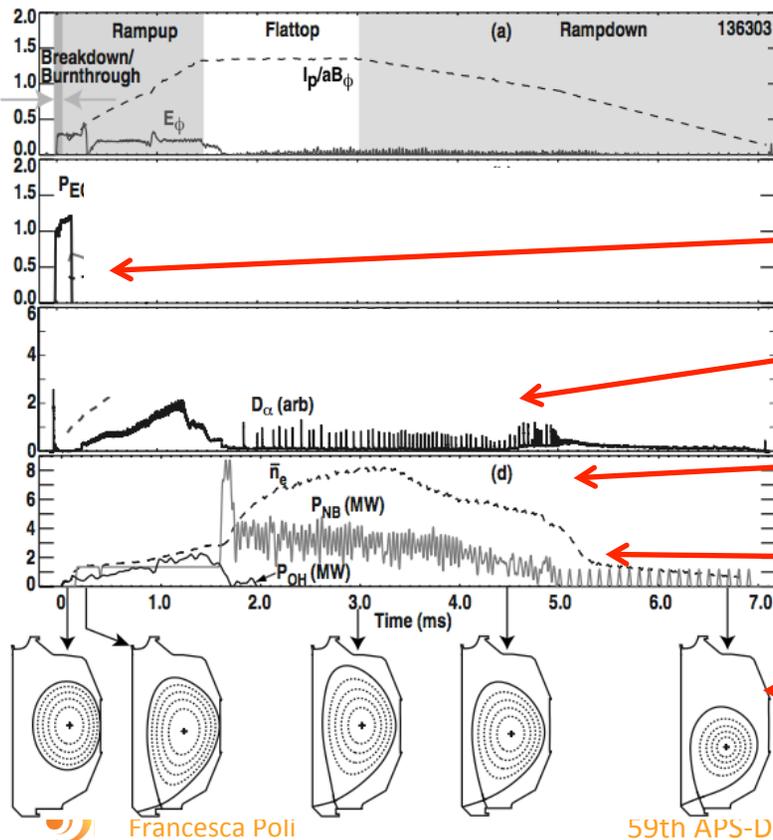
A different perspective on integrated modeling: looking at a tokamak with the eye of an engineer

Goal: simulate an entire plasma discharge from startup to termination
optimizing **ACCURATE PHYSICS and FAST TURNAROUND**

- High fidelity physics at the core of time-slice integrated modeling
- Verified and validated reduced models at the core of (time-dependent) integrated modeling



We model tokamaks to understand experiments and to design stable and reliable reactors



Experiment on DIII-D that simulates an ITER-like plasma

[Fig. 1 - G. Jackson et al, Phys. Plasmas **17** 056116 (2010)]

Heating and current drive (10^{-11} - 10^{-6} s)

MHD instabilities/turbulence (10^{-6} - 10^{-2} s)

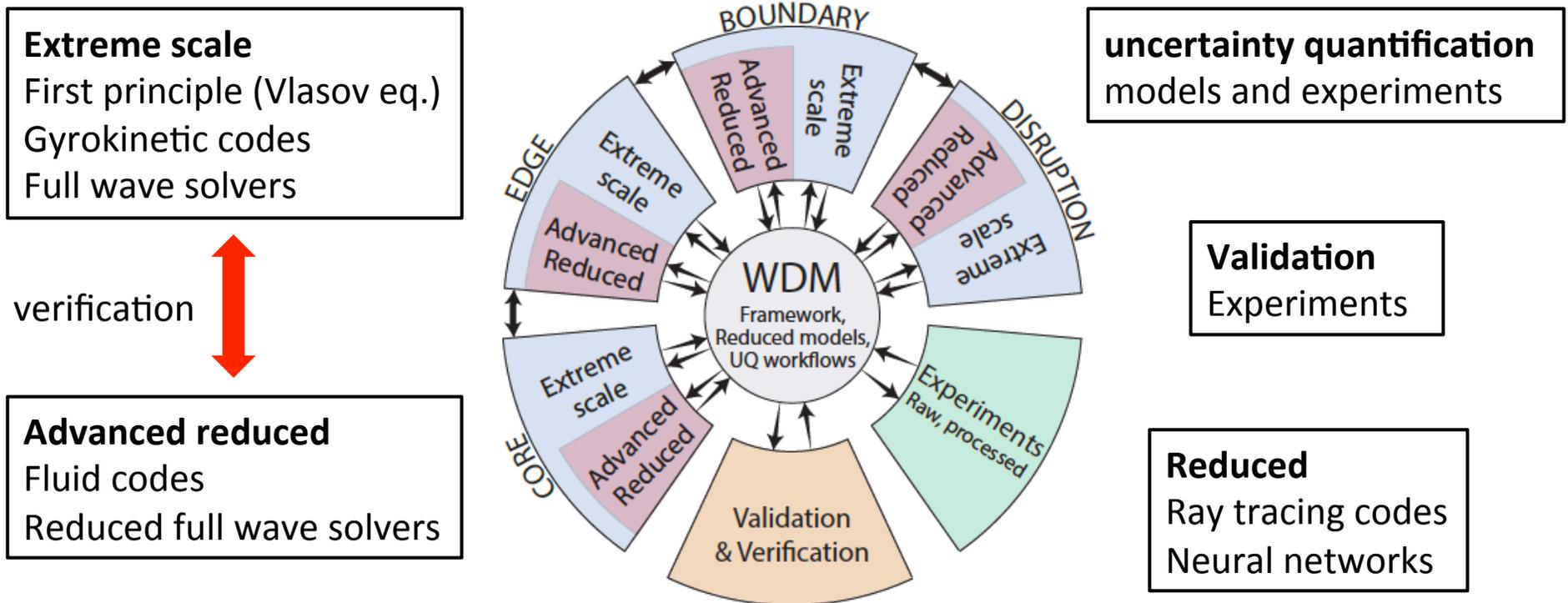
Density evolution (10^{-6} - 10^{-2} s)

Neutral Beams (10^{-6} - 10^{-1} s)

MHD equilibrium (10^{-6} - 10^{-2} s)



A Whole Device Model (WDM) is a comprehensive picture of the complexity involved in modeling a tokamak



A Whole Device Model (WDM) is a comprehensive picture of the complexity involved in modeling a tokamak

Extreme scale

First principle (Vlasov eq.)
Gyrokinetic codes
Full wave solvers

Time-slice
applications

Advanced reduced

Fluid codes
Reduced full wave solvers

Computational time decreases:
millions CPU hours to minutes/seconds
Physics fidelity decreases
Verification/validation needed at each step

Time-dependent
applications

Reduced

Ray tracing codes
Neural networks



- 
-
- From high fidelity physics to reduced models
 - Combining two or more models for an integrated, converged solution on a time slice

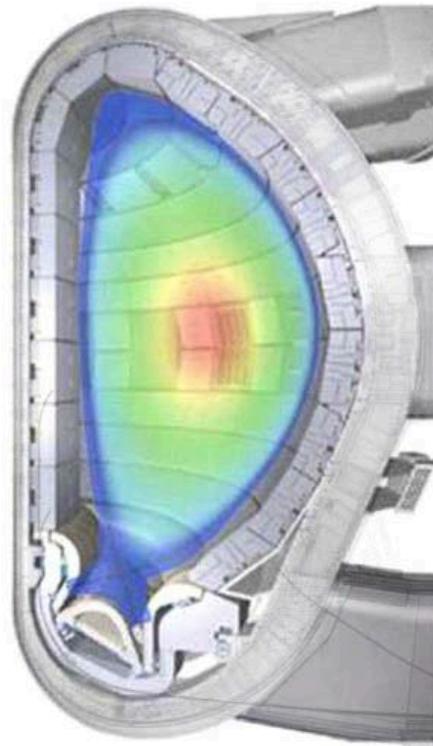


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A plasma in equilibrium can be described by ideal MHD

Equilibrium condition: $\mathbf{j} \times \mathbf{B} = \nabla p$

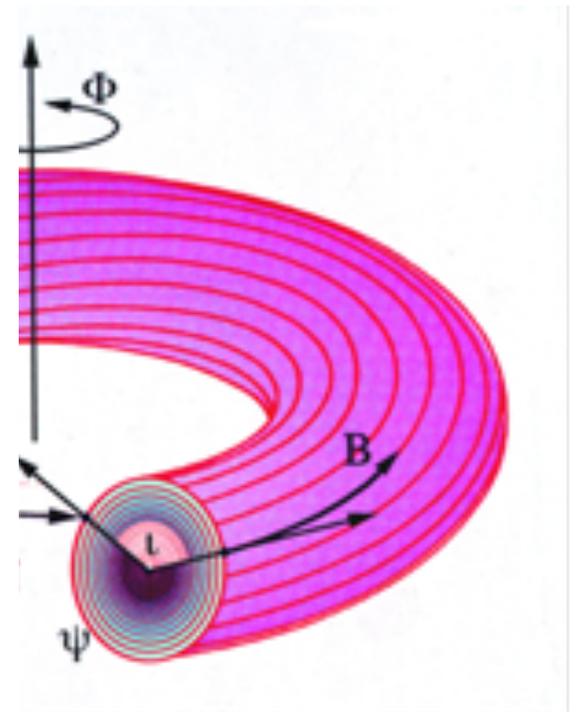
$$\mathbf{B} \cdot \nabla p = 0 \quad \mathbf{j} \cdot \nabla p = 0$$

⇒ \mathbf{j} , \mathbf{B} lie on nested surfaces

⇒ \mathbf{j} , \mathbf{B} , p are described by a flux function ψ

⇒ equilibrium entirely defined by:

$$R \frac{\partial}{\partial R} \left(\frac{1}{R} \frac{\partial \psi}{\partial R} \right) + \frac{\partial^2 \psi}{\partial z^2} = -\mu_0 R J_\phi$$



[Credit, DIFFER website, The Netherlands]

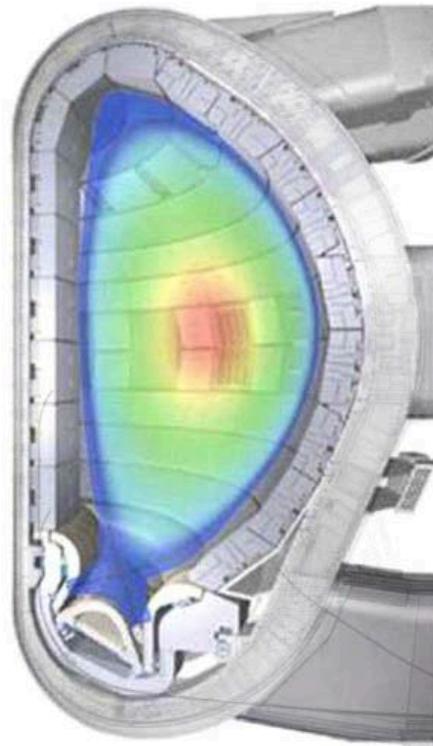


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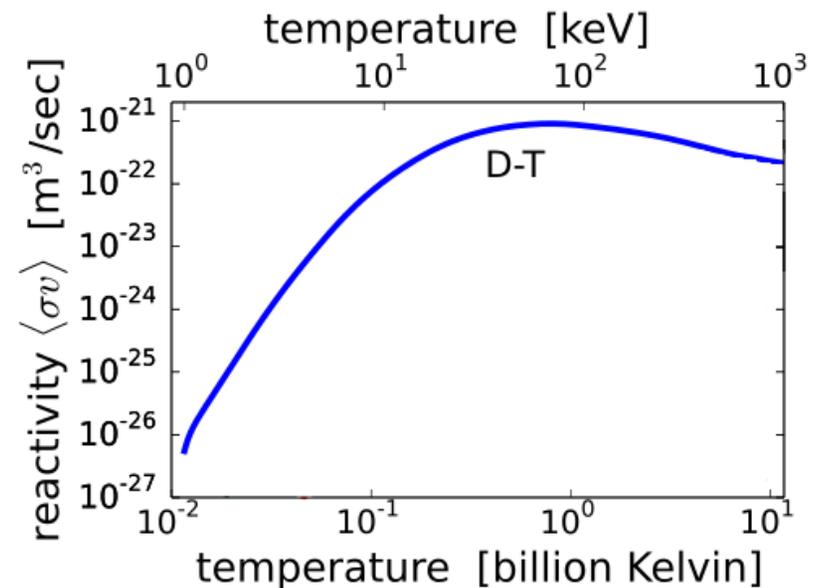
Gas injection
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Ignited (burning) plasmas are dominated by alpha heating ... but they need a boost from external sources before they can burn

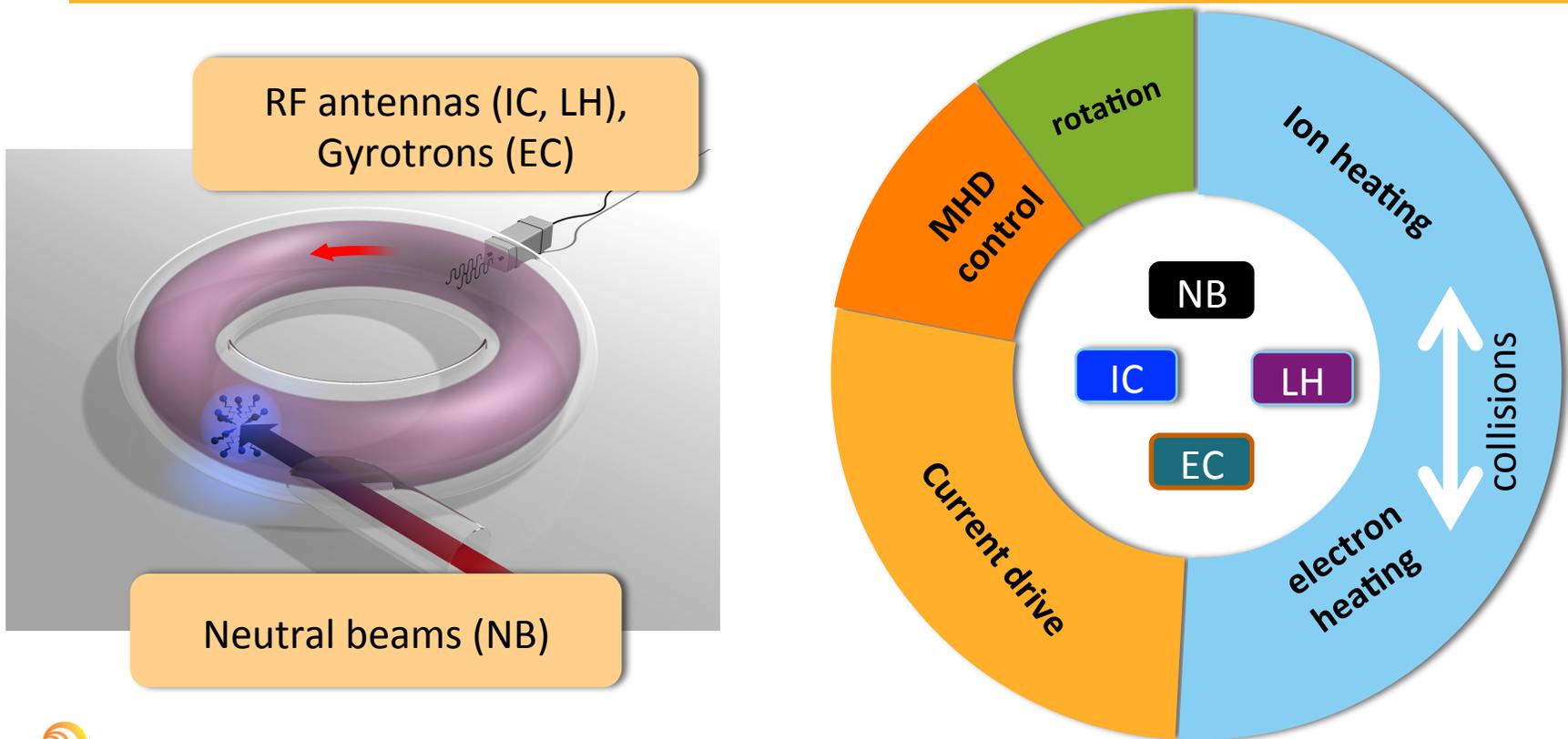
$$S_\alpha = \frac{1}{4} E_\alpha n^2 \langle \sigma v \rangle$$

3.5MeV \nearrow E_α \nwarrow $n_D + n_T$
50-50



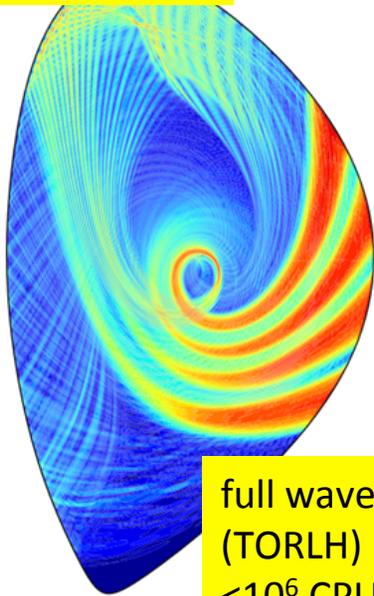
- ⇒ minimum temperature for ignition => 15 keV
- ⇒ alpha heating dominates at temperature >5-7 keV
- ⇒ ohmic heating can heat plasmas only up to 3-5 keV

External sources do more than providing heating and current: they can provide momentum and control of MHD instabilities

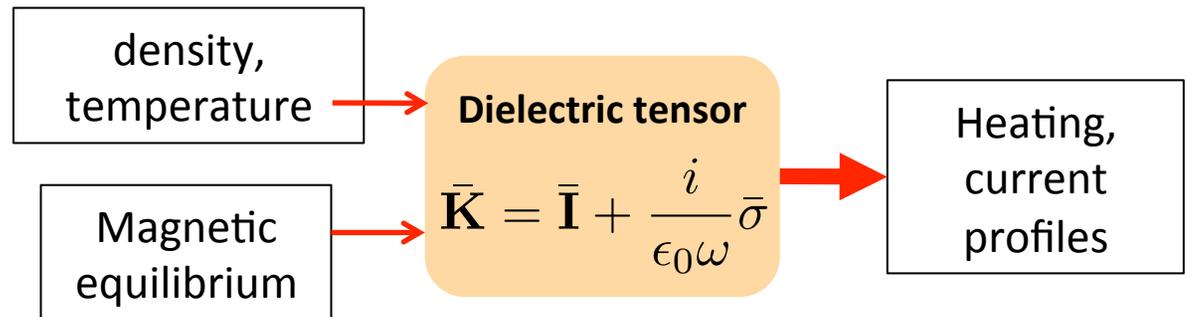


Perturbation of equilibrium enables description of waves propagation by representing the plasma with a dielectric tensor

CMod plasma



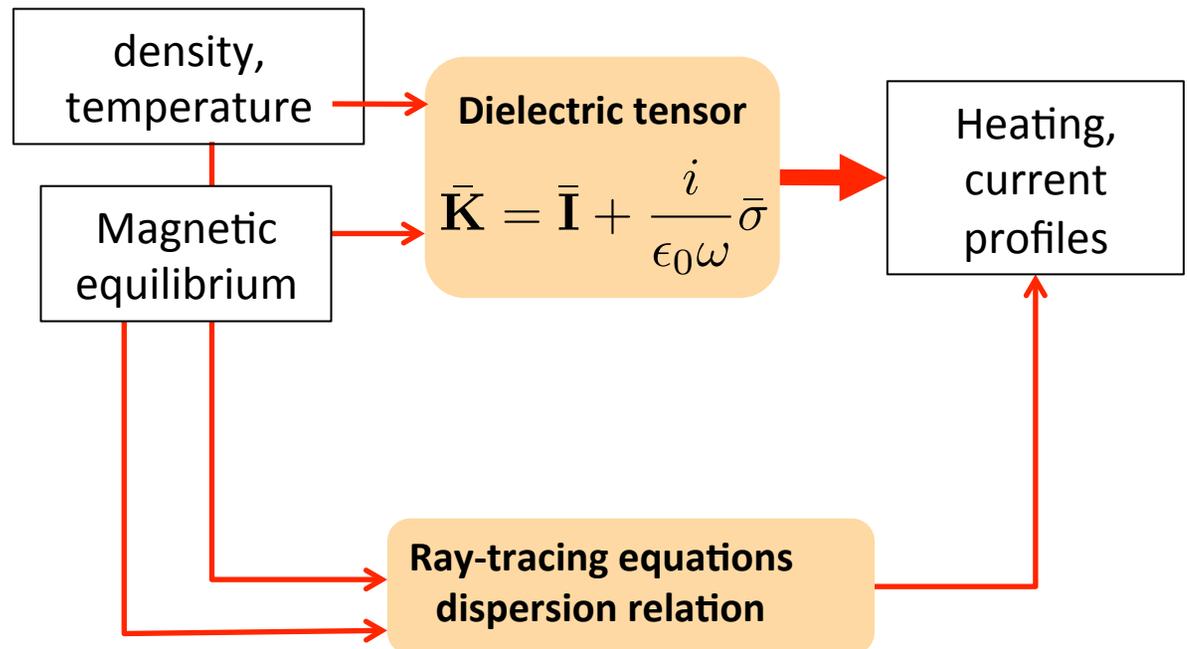
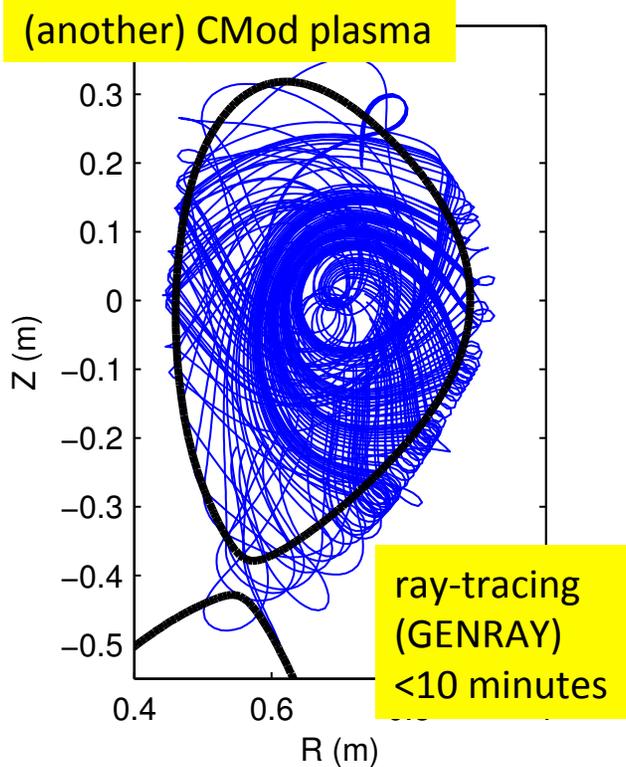
full wave solver
(TORLH)
10^6 CPU Hours



[courtesy of S. Shiraiwa (PSFC), VI2.00003]



Ray-tracing equations are accurate (and fast) approximations of high frequency wave propagation

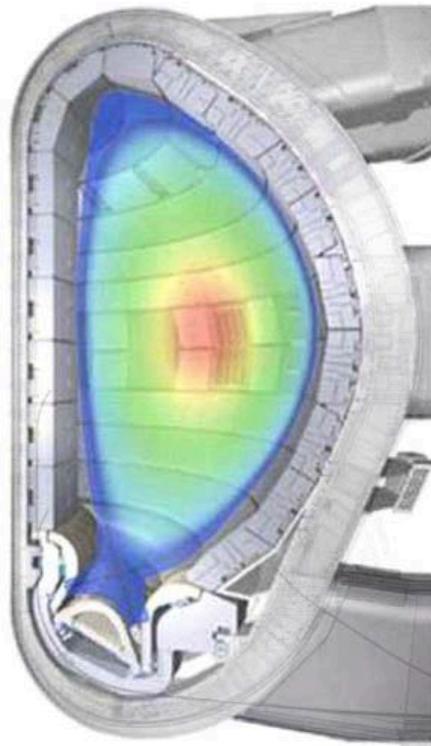


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The plasma is surrounded
by solid structures:
Plasma-material interactions

External heating

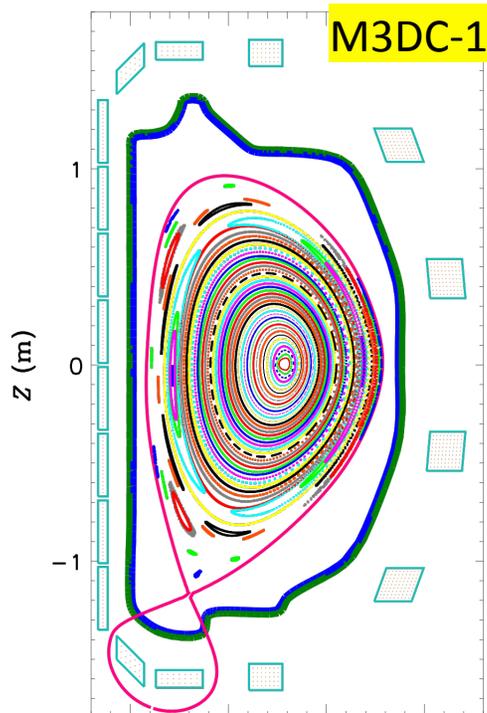
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Plasmas develop instabilities that can degrade performance



M3DC-1

- at specific locations (resonant surfaces)
- they look like magnetic islands
- they can be stabilized by highly focalized beams of high frequency (~ 100 GHz) waves

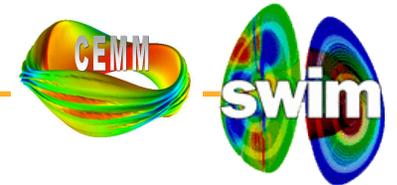
Courtesy of N. Ferraro (PPPL)



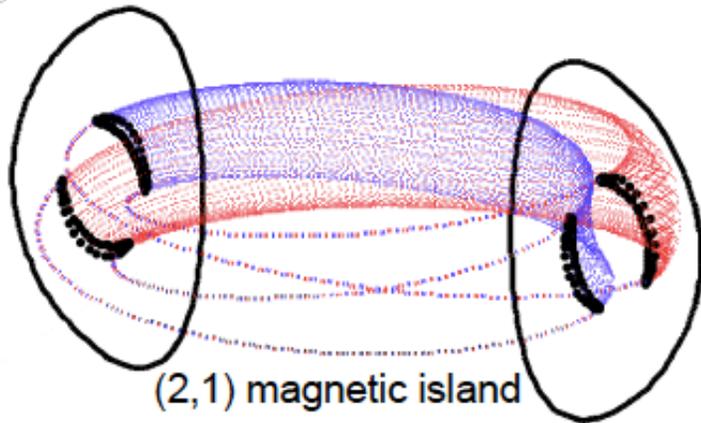
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59th APS-DPP meeting, Milwaukee, WI, October 2017

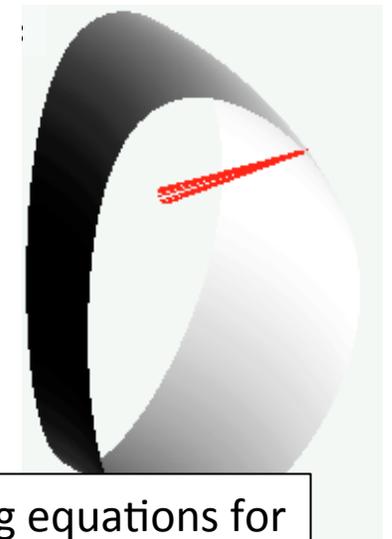
Self-consistent simulations of stabilization of magnetic islands can be done only in the framework of 3D nonlinear MHD



NIMROD/GENRAY



quasilinear corrections
to MHD equilibrium
to account for effect of waves



Ray tracing equations for
wave propagation

T. Jenkins (Tech-X)



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Reduced models are based on a Modified Rutherford Equation

$$\frac{dw}{dt} \propto \sum f_{plasma}(w) + \sum f_{EC}(w)$$

Magnetic island evolution described by:

- analytic expressions for instability threshold and driving
- In a cylindrical approximation
- Analytic expressions and fitting parameters for the stabilizing effect of RF waves

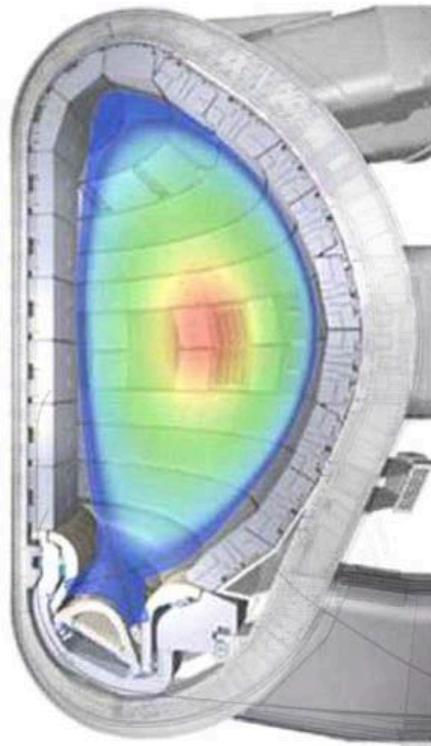


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Separation of scales enables representing (almost) any transport problem as a diffusion/convection-like problem

The goal is to obtain a set of diffusion-like equations of the form:

$$\frac{\partial Q}{\partial t} + \nabla \cdot \Gamma = S(Q, \mathbf{r}, t)$$

⇒ for a physical variable **Q**

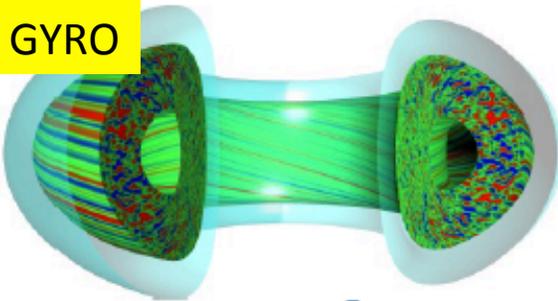
⇒ identify the flux **Γ**

⇒ and the source and sink terms contained in **S**



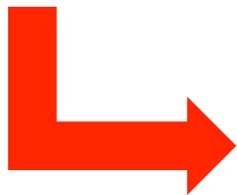
Understanding and modeling tokamak turbulent transport requires theory-based prediction of flux-gradient relationships

GYRO



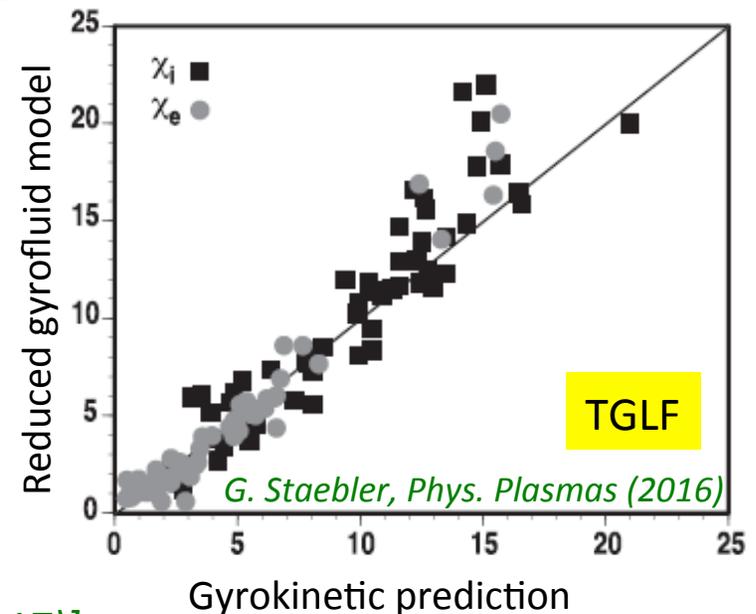
6D Vlasov equations =>
5D nonlinear "gyrokinetic"

State-of-the-art multi-scale ($\rho_i \rightarrow \rho_e$)
~50M CPU-hrs for 3-point scan
[*N. Howard, Nucl. Fusion (2016)*]



fluid eq. with right closure =>
Fit over limited set of GK
(**< 1 sec with NN**)

[*Citrin NF 55 (2015), Meneghini, NF 57 (2017)*]

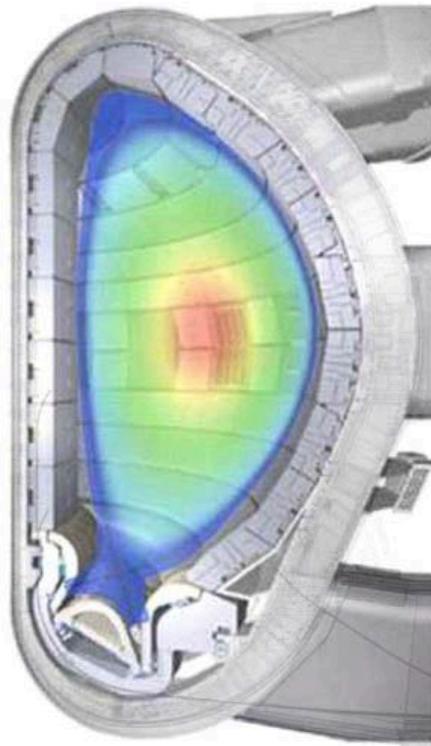


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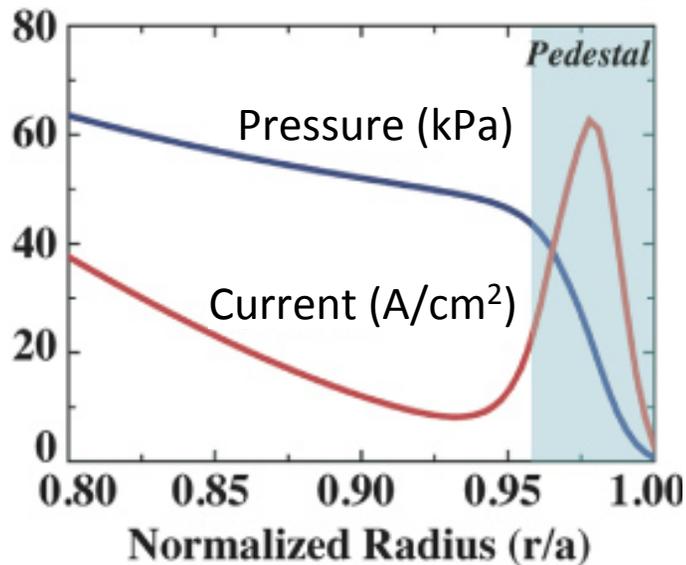
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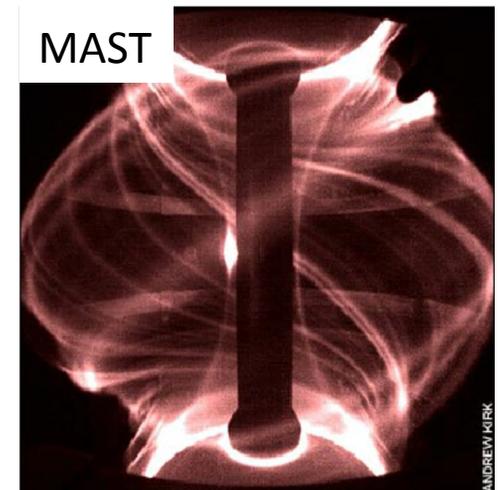


Steep gradient at the plasma edge drives MHD instabilities

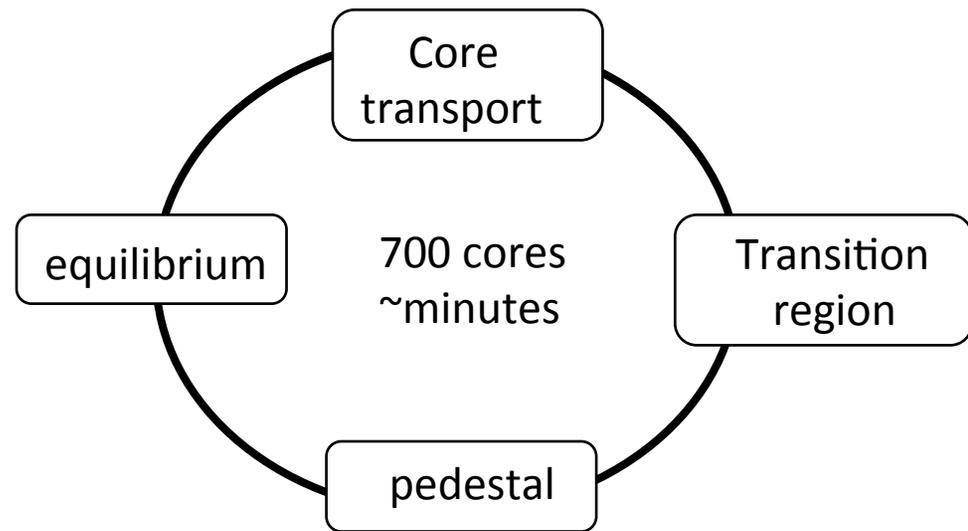
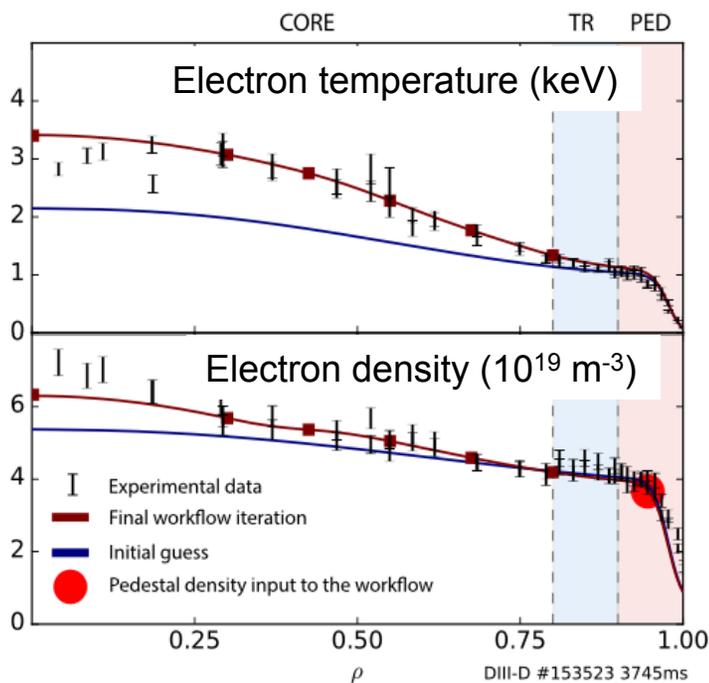


Plasmas spontaneously develop a region (pedestal) that encloses good confinement [*P. Snyder FR1.00001*]

The width and height of the pedestal can be determined univocally by MHD calculations for a given value of density



Efficient workflows allow combined solution of core transport and pedestal MHD stability



Not yet a self-consistent solution:

- model for pedestal density still missing
- model for transition region that includes microturbulence still missing

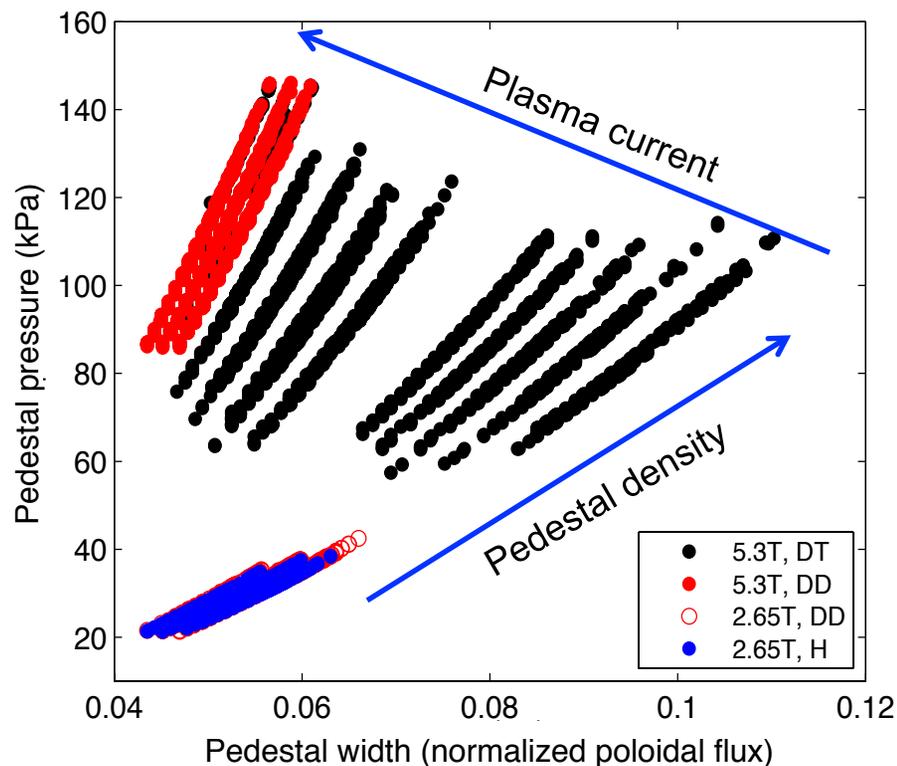
O. Meneghini, *Phys. Plasmas* **23** (2016)



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Replacing MHD stability pedestal calculations with a lookup table advantageous in time-dependent simulations



~6500 points cover the expected ITER operational space

[task under international collaboration]

A neural network has also been developed
[O. Meneghini, NF 57 (2017)]

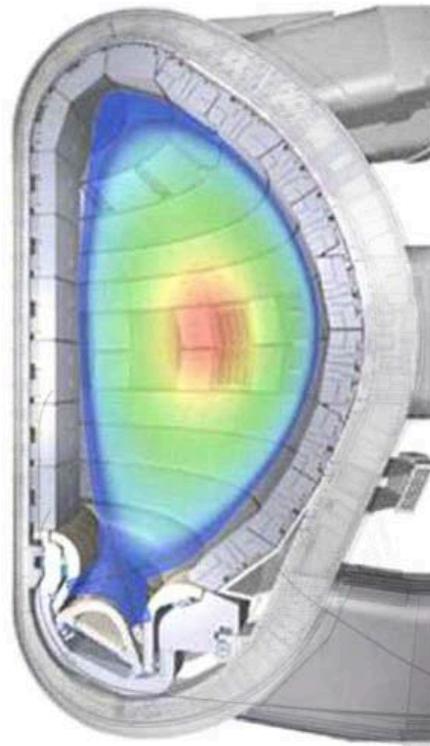


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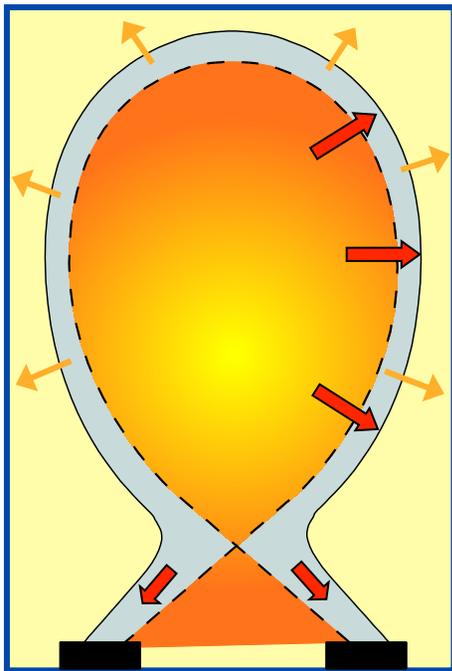
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Where the plasma meets the wall: the Scrape-Off-Layer (SOL)



Heat losses from the plasma can damage plasma facing components

SOL width determined by competition between parallel and perpendicular transport

R. Goldston, NF 52 (2012), PI3.00004

Courtesy of V. Soukhanovskii (LLNL)



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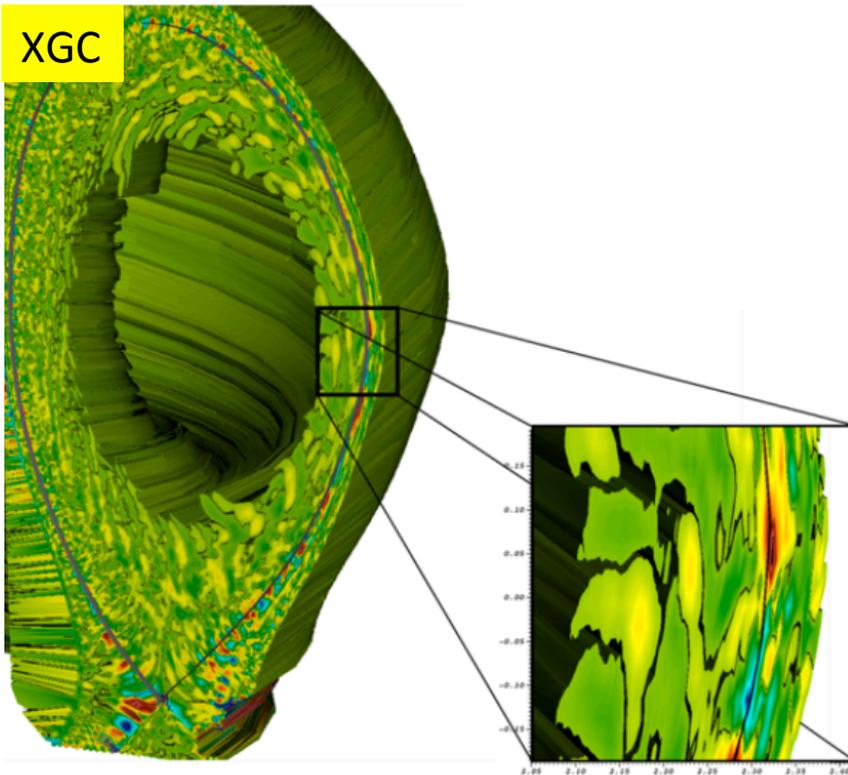
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Gyrokinetics simulations that extend the computation domain to the wall needed to develop reduced models for the edge plasma

XGC

S-H Ku, DI2.00004



We need a better model for the plasma edge, Including the region inside the pedestal,

- Microturbulence
- coherent structures (blobs)
- Separation of ion and electron scales
- Wall sources

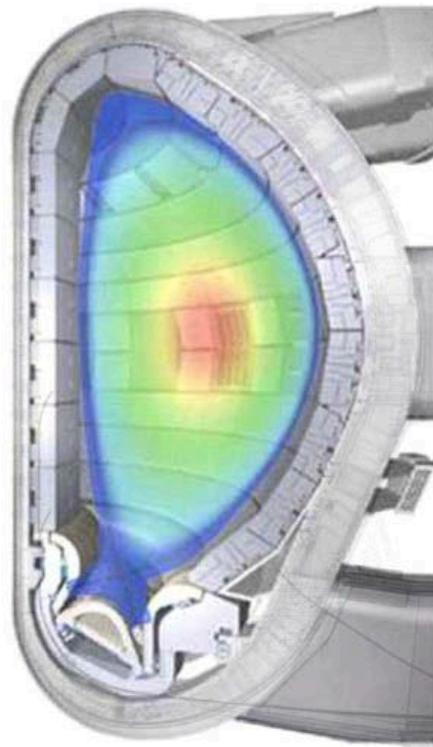


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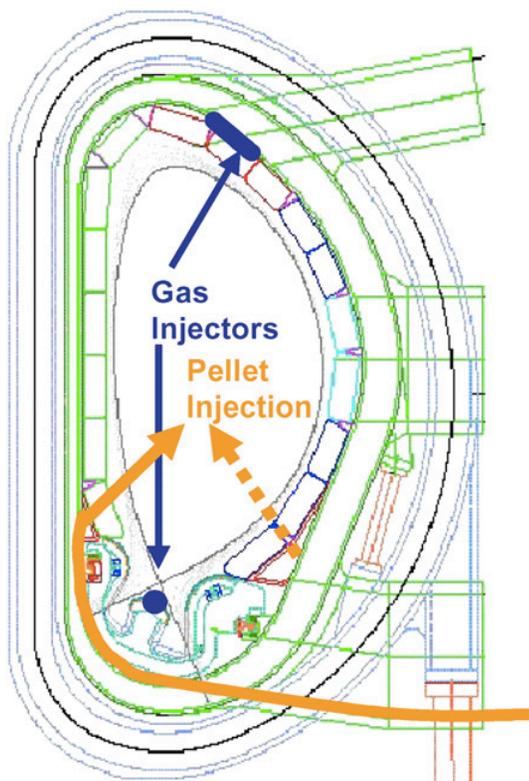
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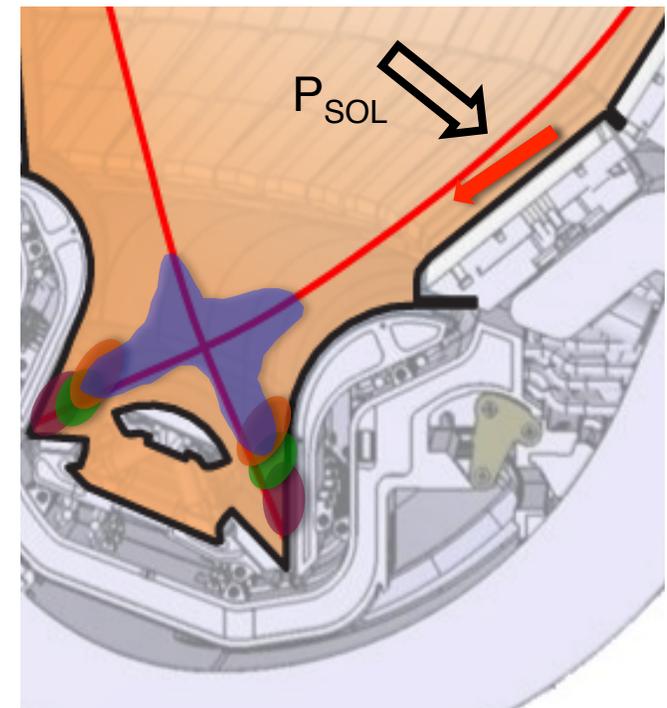


Edge transport and fuelling are critical ingredients to model the plasma evolution in burning plasma conditions



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- Heat conduction zone
- Impurity radiation zone
- $H^0/D^0/T^0$ ionization zone ($T_e > 5eV$)
- Neutral friction zone
- Recombination zone ($T_e < 1eV$)



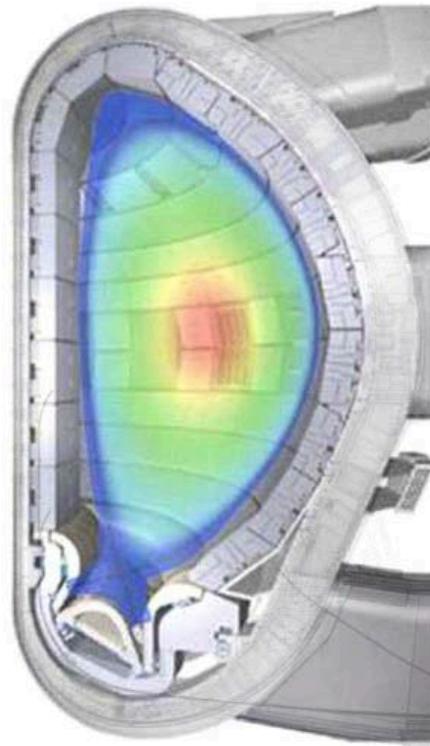
Courtesy of R. Pitts (ITER Organization)

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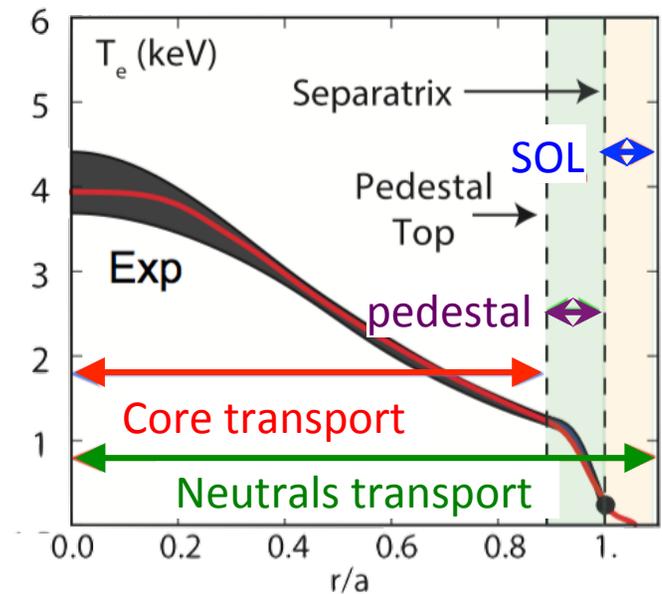
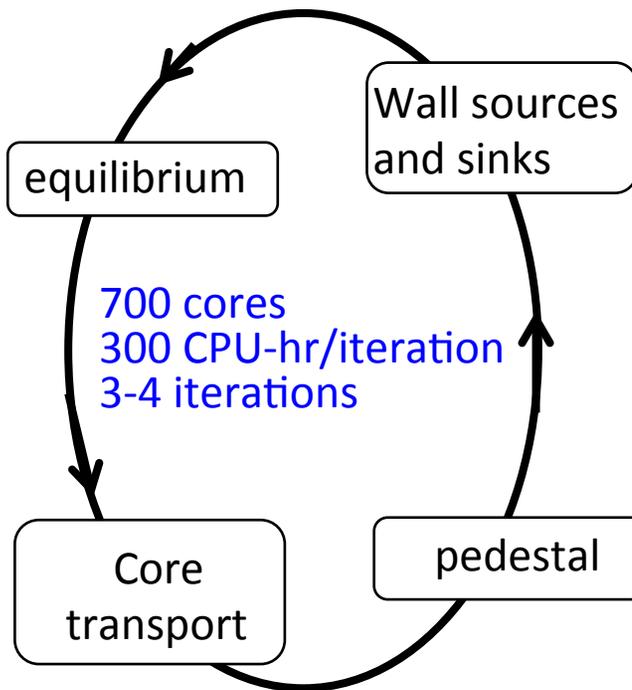
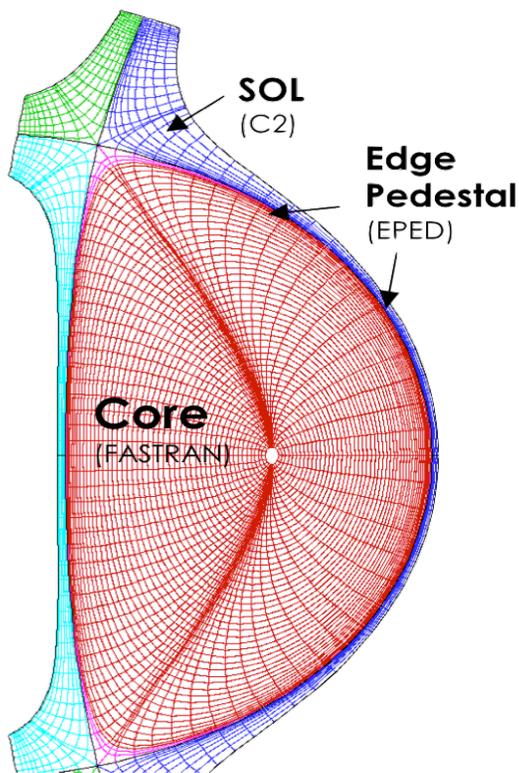
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Modular coupling of 1.5D core and 2D edge transport implemented and validated on DIII-D steady-state plasma

AToM SciDAC



J-M. Park, D. Green (ORNL)



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Integrated modeling on single time-slice has led to great insight

- On small scales dynamics and multi-scale couplings
- On how to plug-in models to describe the plasma from axis to wall

... but plasmas in tokamaks are dynamical, nonlinear systems ...





The other side of integrated tokamak modeling:
an entire plasma discharge from startup to termination



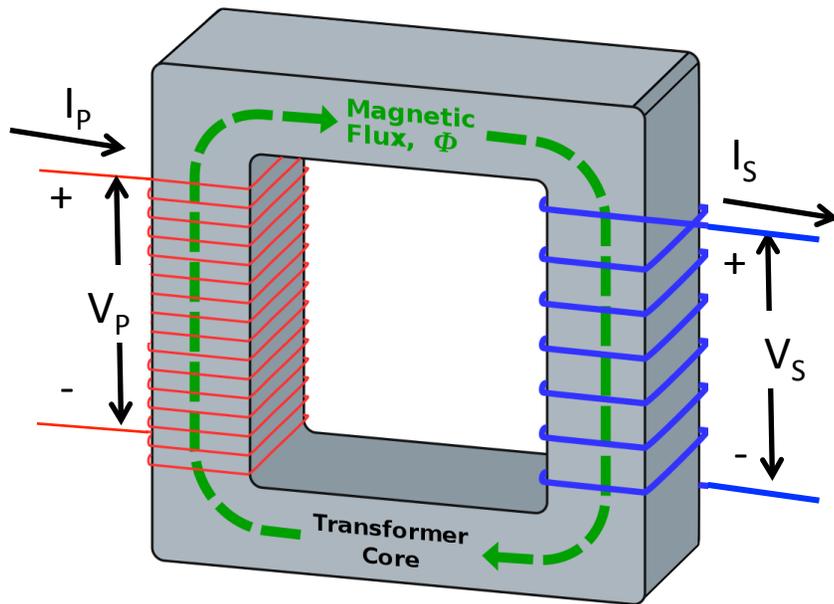


The basics ...

A tokamak is like a transformer, where the plasma is the secondary winding



In a transformer the circuits are magnetically connected



In the primary:

$$I_p(t) \Rightarrow B_p(t)$$

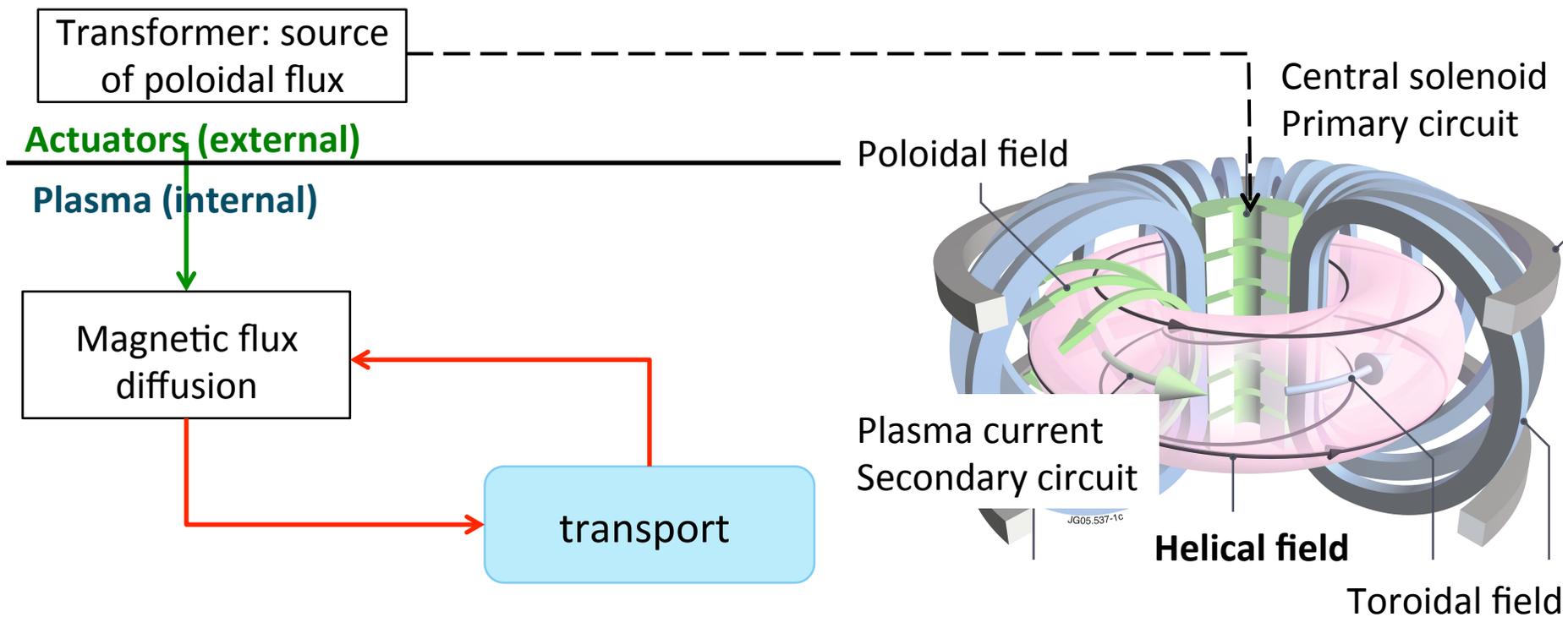
In the secondary:

$$I_s(t) \Rightarrow B_s(t)$$

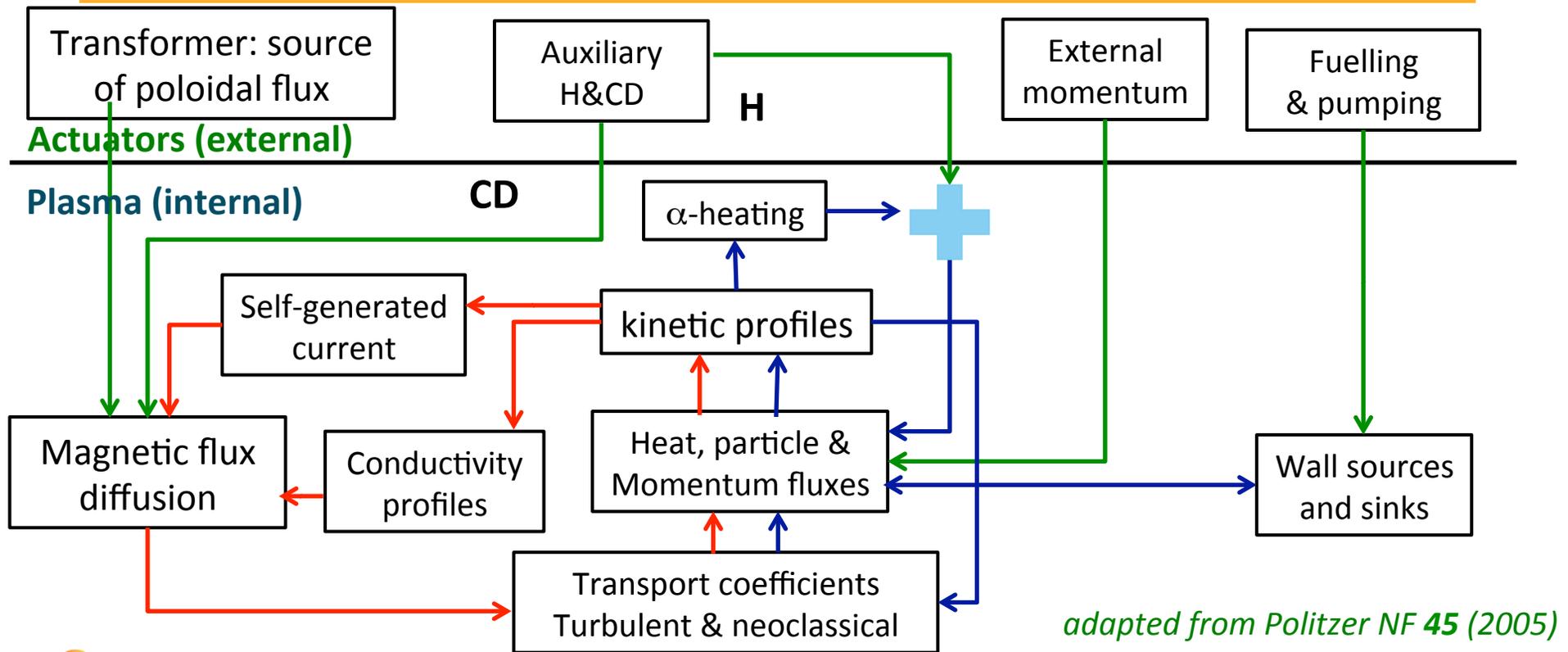
Total magnetic flux density: $B(t) = B_p(t) + B_s(t)$

Limitation: it is pulsed, cannot operate continuously in steady-state

In a tokamak the secondary circuit is a conducting fluid things get complicated



A tokamak simulator needs to connect **fast** (transport) and **slow** (current diffusion) time scales



adapted from Politzer NF 45 (2005)



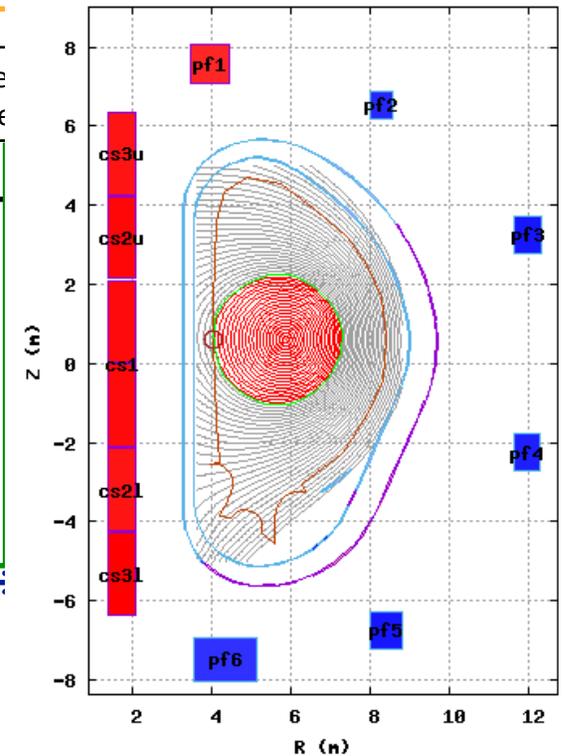
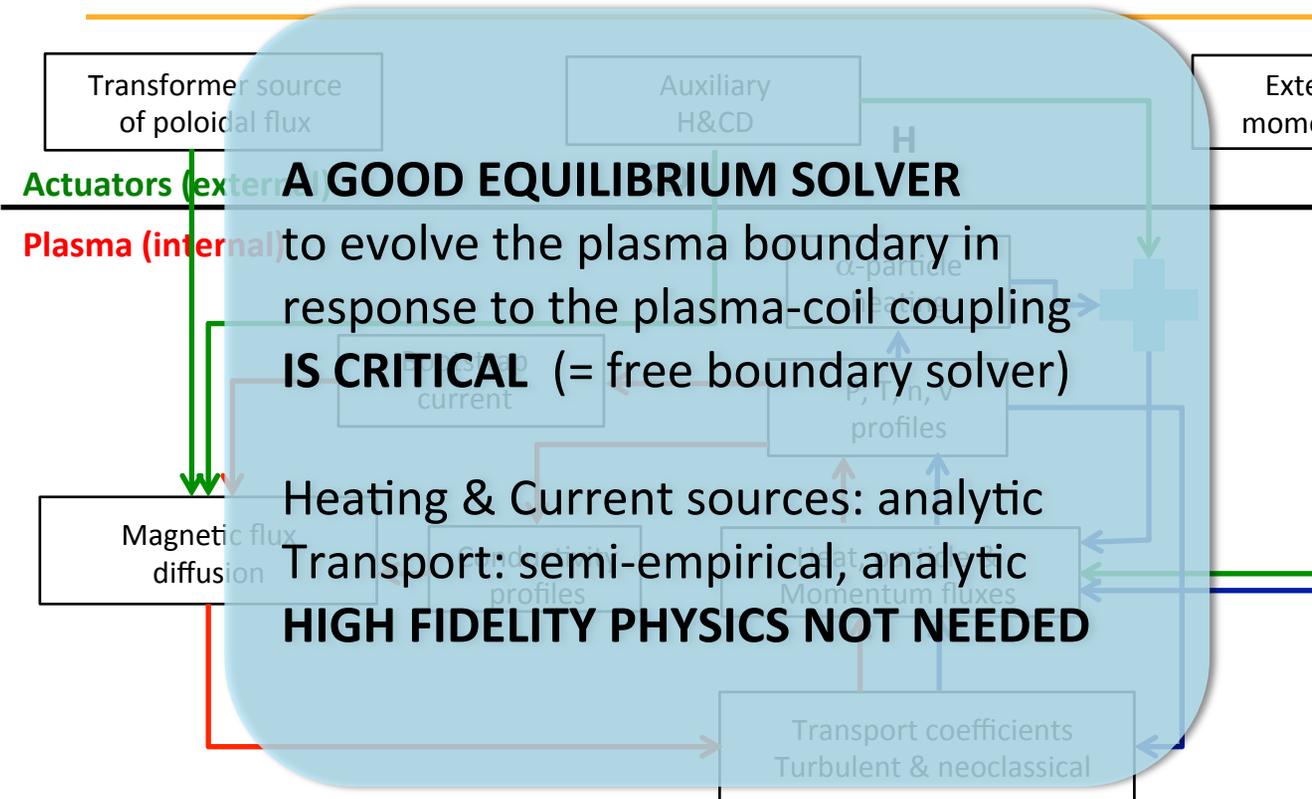
All the steps we take when we model an ITER plasma discharge

The goals of ITER: **I**nternational **T**hermonuclear **E**xperimental **R**eactor

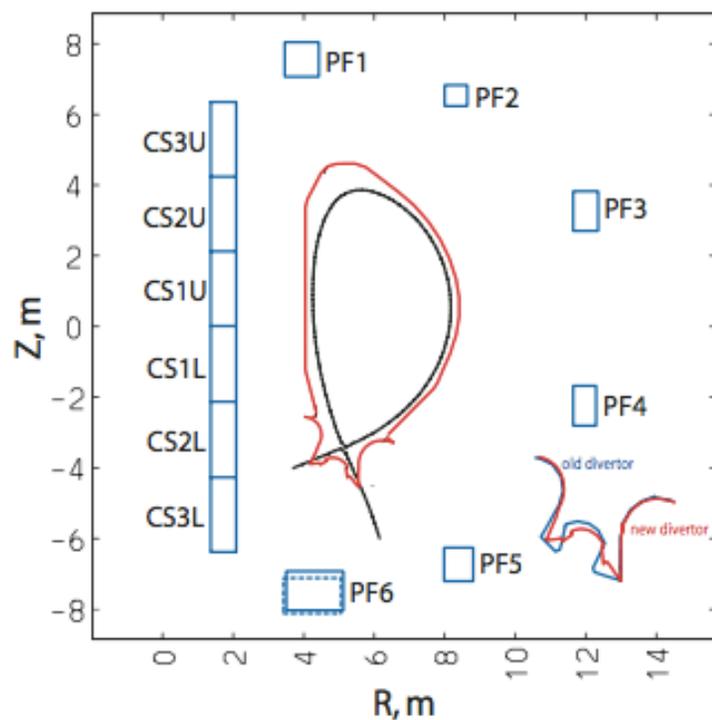
- Produce 500 MW of fusion power
- Demonstrate integrated operation of technologies for a fusion power plant
- Achieve a self-heated deuterium-tritium plasma
- Test tritium breeding
- Demonstrate safety of fusion devices



The first step is to get all coil currents and plasma shape right



These simulations have been valuable to define and revise the ITER coil operational space and the plasma control capabilities



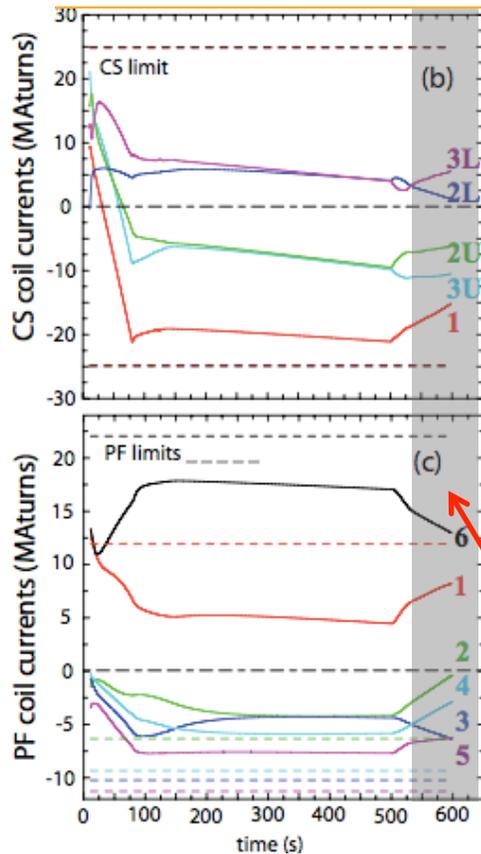
New poloidal field coils layout expands operational space:

- Flattop burn duration
- Operation with broader current profile

[C. Kessel et al, *NF 49* (2009)]



Simulations with simplified transport are designed to test engineering parameters, not to discuss physics



[T. Casper et al, NF 54 (2014)]

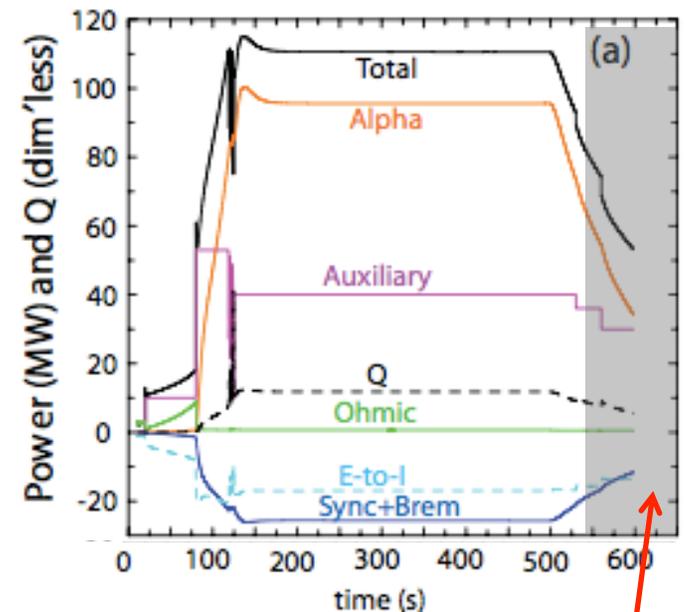
With assumptions on transport, confinement to satisfy target fusion power

⇒ Assessment of coil limits

Plasma termination is a critical area of ongoing research

59th APS-DPP meeting, Milwaukee, WI, October 2017

CORSICA



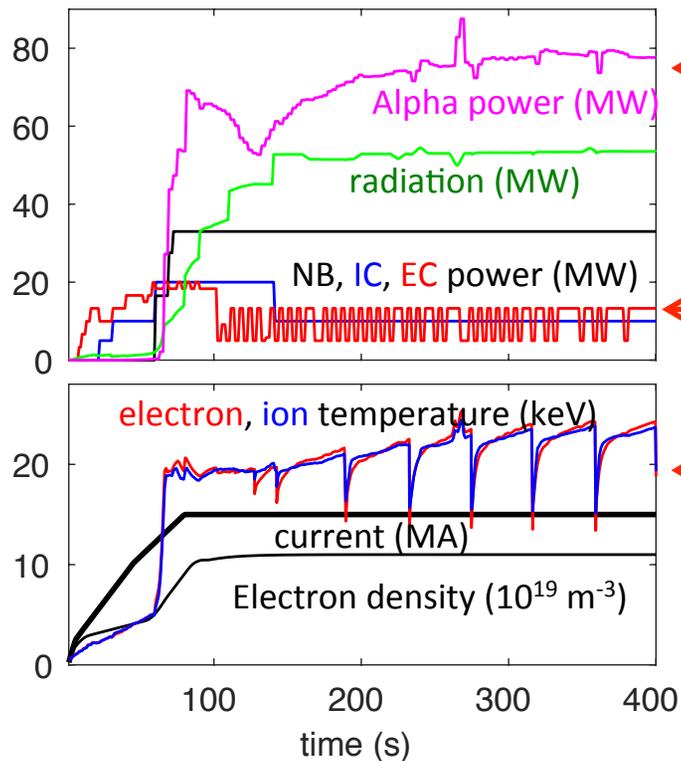
Simulations with simplified transport do not inform on plasma dynamic response to external actuators and internal MHD stability

- Can this plasma really achieve the target?
- How additional physics constraints affect the results?



Physics-based models for thermal transport and current evolution can move the operational point away from target

TRANSP



← alpha power P_α responds nonlinearly to core transport and pedestal width and height

← external power under feedback control

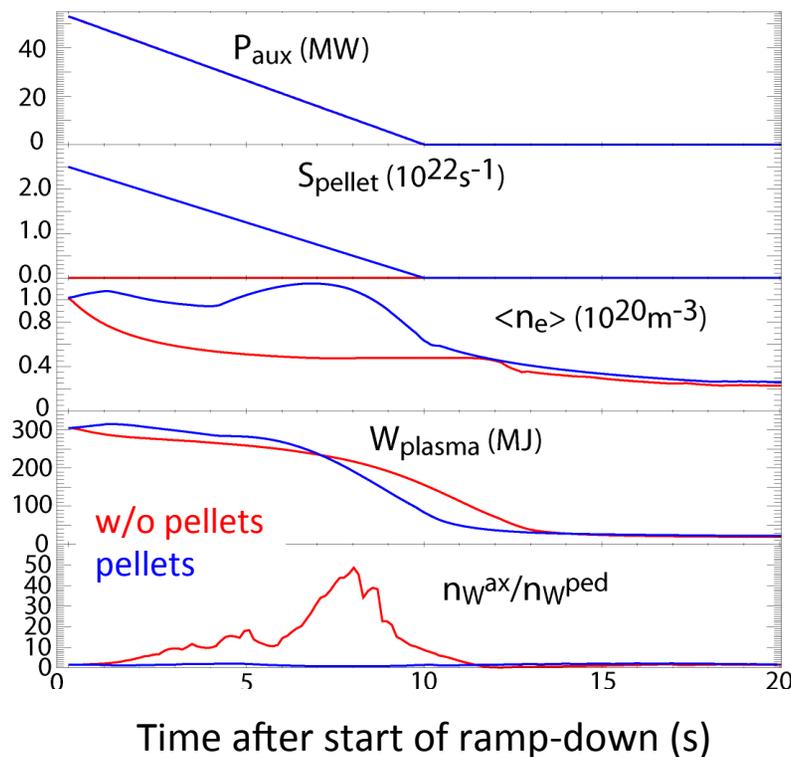
← global MHD stability included in evolution of temperature and current

Plasma does not achieve the target of $P_\alpha=100\text{MW}$



Simulations with fully integrated 2D edge transport plasma have indicated the importance of fueling during ITER plasma termination

JINTRAC



- Pellet fueling at exit from H-mode needed to avoid accumulation of tungsten in the core
- Model has been validated on JET, in dedicated ITER-like experiments.

A. Loarte IAEA-FEC (2016)
F. Koechl, NF 57 (2017)



The path forward to an all-inclusive tokamak simulator

- Reduced models:
 - from nonlinear 3D MHD stability
 - for edge plasma turbulence induced transport
 - for wave propagation in the Scrape-off-Layer
 - for energetic particle induced transport
 - Plasma startup (SNU, Univ. Kyoto)



The path forward to an all-inclusive tokamak simulator

- Reduced models from nonlinear 3D MHD stability
 - **Plasma response to external perturbations**
 - Disruptions (trigger)
 - MHD instabilities (onset, control)
 - global instabilities, like sawtooth cycle [**CP11.00113**]

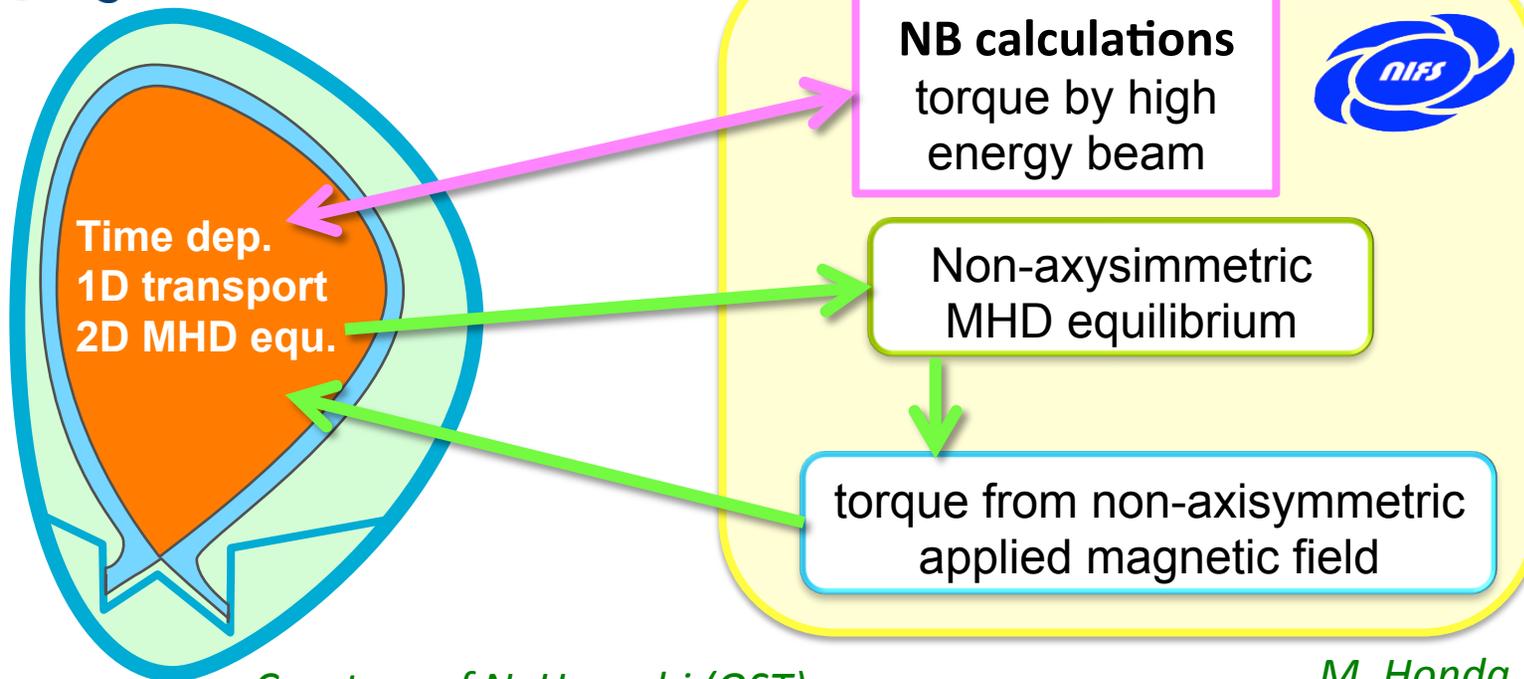
Reduced models for edge plasma turbulence induced transport

Reduced models for wave propagation in the Scrape-off-layer

Reduced models for energetic particle induced transport



Better models for prediction of plasma rotation needed for modeling plasma response to applied magnetic perturbations



Courtesy of N. Hayashi (QST)

M. Honda, NF 57 (2017)



The path forward to an all-inclusive tokamak simulator

Reduced models from nonlinear 3D MHD stability

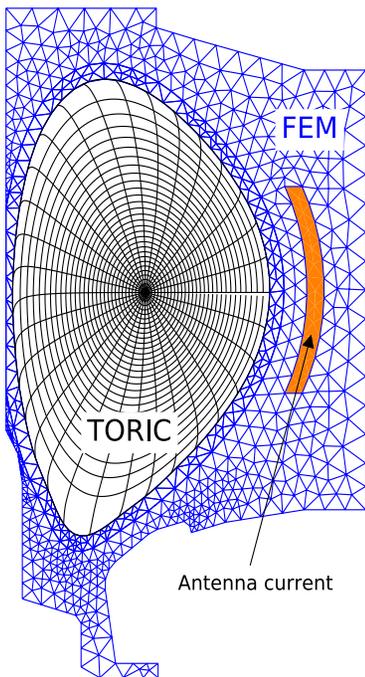
Reduced models for edge plasma turbulence induced transport

- **Reduced models for wave propagation in the Scrape-Off-Layer**

Reduced models for energetic particle induced transport



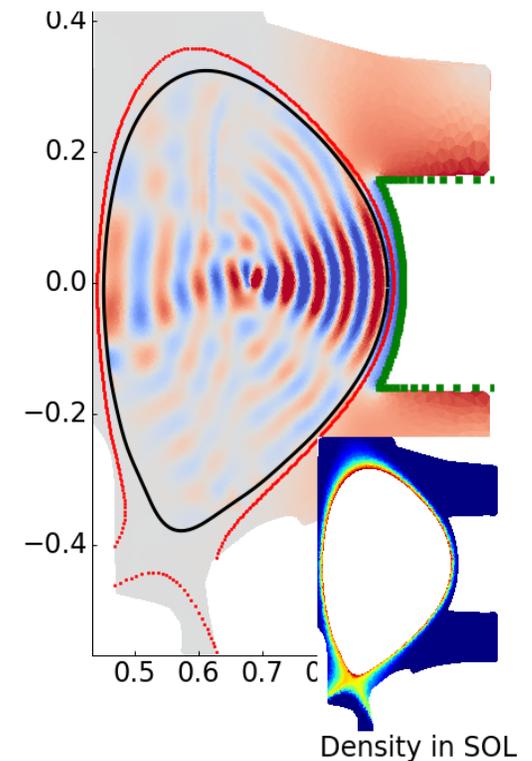
Hybrid approach to modeling of RF wave propagation is a promising avenue towards implementation in tokamak simulator



Core: **Axisymmetric** flux surface grid
Hot plasma conductivity
Dense Matrix Solver
Edge: **Unstructured mesh** with
complicated geometry (either 2D or **3D**)
Cold plasma with collision.

Boundary: **matching technique** to build
integrated solution

Would benefit from realistic model of SOL



S. Shiraiwa (PSFC) VI2.00003



Francesca Poli

59th APS-DPP meeting, Milwaukee, WI, October 2017

The path forward to an all-inclusive tokamak simulator

Reduced models from nonlinear 3D MHD stability

Reduced models for edge plasma turbulence induced transport

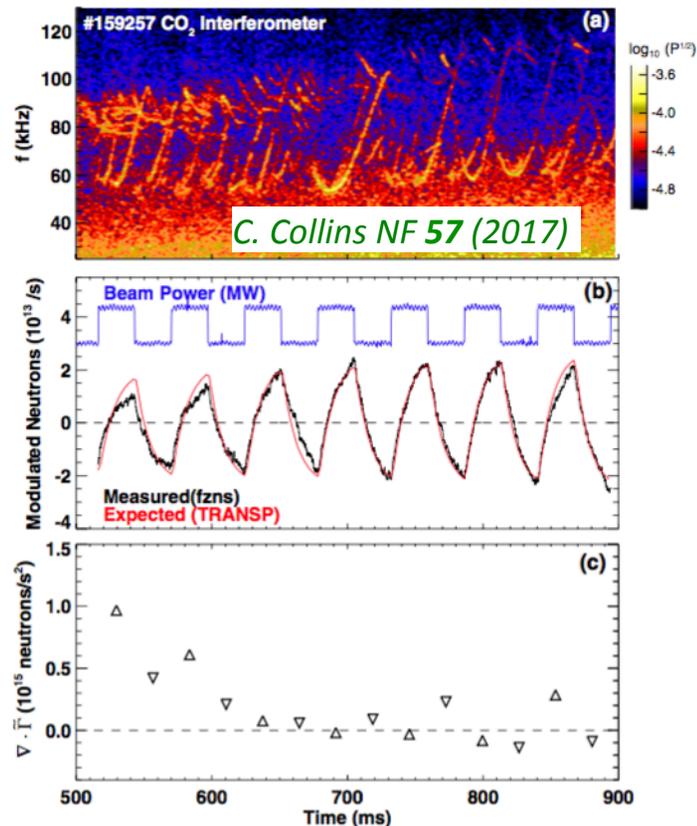
Reduced models for wave propagation in the Scrape-off-layer

- **Reduced models for energetic particle induced transport**



Fast particles can drive instabilities, like Alfvénic modes

[SciDAC project ISEP]



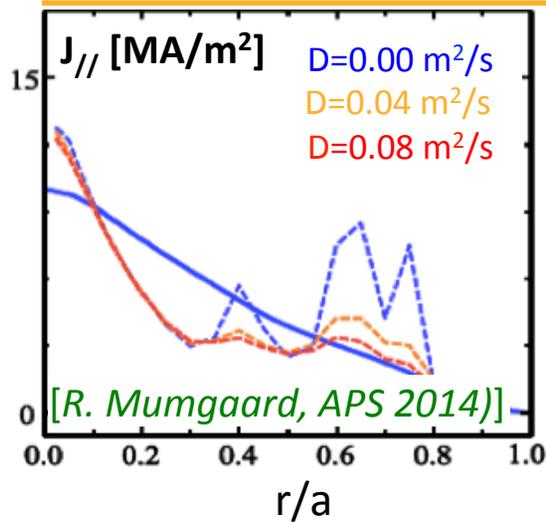
- ⇒ Instabilities eject fast particles (e.g. alphas)
 - ⇒ Decrease performance
 - ⇒ Cause localized losses and damage to wall
- ⇒ Challenges:
 - ⇒ Understand physics to develop scenarios not prone to instabilities
 - ⇒ Develop control tools to mitigate/suppress instabilities



Experimental validation is critical

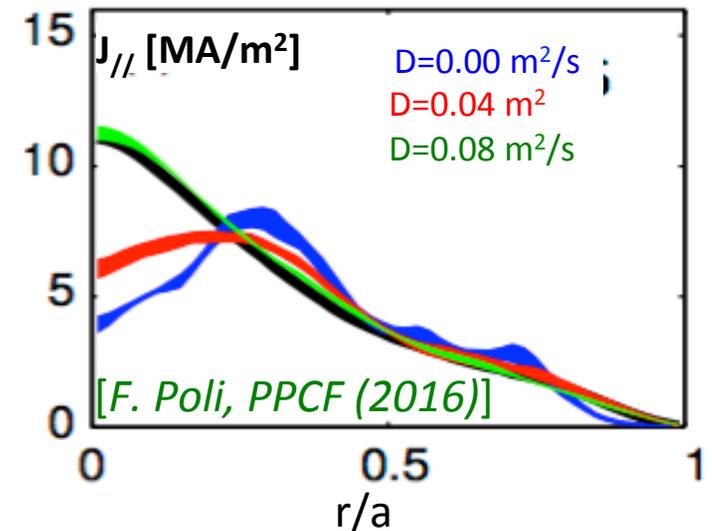
Time-dependent validation is critical for modeling of wave propagation in plasmas

GENRAY/CQL3D



Standalone ray-tracing +Fokker-Planck solver
⇒ large off-axis current, not seen in experiments

Time-dependent simulation
⇒ self-consistent evolution of wave field
and magnetic equilibrium



Concluding remarks

- Increasing computing capabilities allow to solve bigger problems
- Efficient workflows allow to solve complex problems
- The path forward for an all-inclusive tokamak simulator should focus on reduced models
- Validation against experiments and verification against extreme scale calculations are critical to retain fidelity while reducing computational time

